



National Aeronautics and
Space Administration



Unsteady Aerodynamic Modeling for Atmospheric Entry Vehicles in Subsonic and Incompressible Flow: A Frequency Response Approach

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Outline

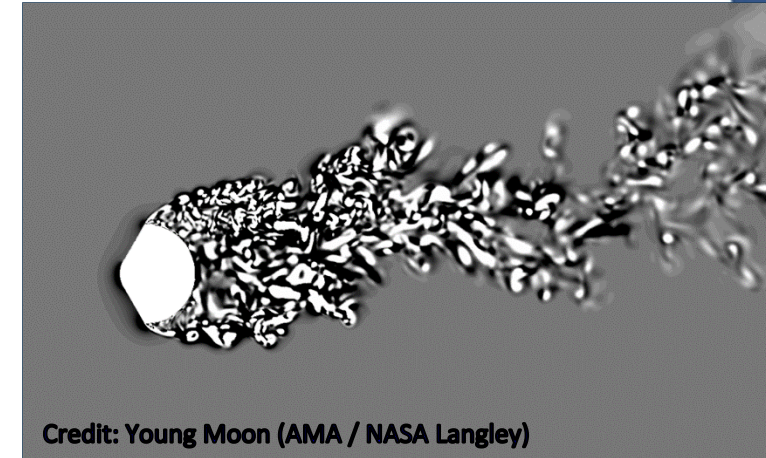
- Background
- Motivation and Objective
- Classical Aerodynamic Theories
- Frequency Response Approach
- Results from the 12 ft Low Speed Tunnel
- Reconstruction from semi-empirical models
- Summary

Motivation

- Planetary missions often require an entry, descent, and landing (EDL) system to successfully deliver a payload to the surface

- *Blunt body entry capsules display complex aerodynamics due to unsteady wake effects*
 - Dynamic instabilities tend to grow with decreasing Mach number
 - Either the vehicle **(i) reaches a stable limit cycle**, or **(ii) the oscillations can grow** to where a stabilizing device may not be safe to deploy, and the vehicle may eventually begin to tumble
 - Flow separation and reattachment shift unpredictably depending on frequency of oscillations, mean angle of attack, and pitch/plunge amplitudes.

- Previous model efforts include:
 - Allen and Tobak, along with Fletcher and Wolhart, laid the foundation for dynamic stability analysis of blunt bodies through early analytical and experimental studies.
 - Schoenenberger and Queen developed differential equations of motion for entry vehicles, emphasizing the dominance of aerodynamic forces over gravitational and centrifugal forces, focusing on the angle of attack's time evolution.
 - Seiff, Whiting, and Ericsson extended Newtonian impact methods to unsteady aerodynamics, refining semi-empirical predictions for slender blunted cones with half-cone angles below 10 degrees and achieving agreement for Mach numbers as low as Mach 3.



Credit: Young Moon (AMA / NASA Langley)

**Understanding unsteady aerodynamics
is an important part of the modeling
and simulation of entry vehicle
trajectories**

[1] Tobak, M., and Allen, H. J., Dynamic stability of vehicles traversing ascending or descending paths through the atmosphere, Vol. 4275, National Advisory Committee for Aeronautics, 1958.

[2] Schoenenberger, Mark, and Eric M. Queen. "Limit cycle analysis applied to the oscillations of decelerating blunt-body entry vehicles." NATO RTO Symposium AVT-152 on Limit-Cycle Oscillations and Other Amplitude-Limited, Self Excited Vibrations. No. RTO-MP-AVT-152. 2008.

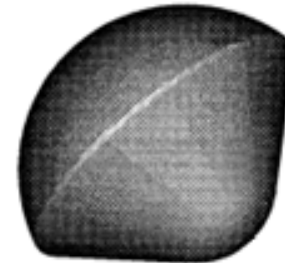
[3] Ericsson, L. E., "Generalized unsteady embedded Newtonian flow," Journal of Spacecraft and Rockets, Vol. 12, No. 12, 1975, pp. 718–726. <https://doi.org/10.2514/3.27870>.

Motivations:

- Develop data driven analysis to model unsteady aerodynamics from forced oscillation wind tunnel tests.
- Traditional aerodynamic databases not capturing vehicle frequency effects on aerodynamic response.
- Investigate a new data reduction technique to extract dynamic aero coefficients
- Develop a system identification like approach that characterizes a transfer function for a blunt body entry vehicle should be able to predict lift and moment time histories as a function of oscillatory inputs to inform controller design

Objective: Utilize a data-driven frequency approach to model unsteady aerodynamics from experimental forced oscillation tests on an entry vehicle.

The chosen test article is a 60-degree spherical cone angle with a hemispherical backshell, similar to Mars Microprobe and Varda's Winnebago-1 (W-1)



Braun, R. D., Mitcheltree, R. A., and Cheatwood, F. M., "Mars microprobe entry analysis," *1997 IEEE Aerospace Conference*, Vol. 1, IEEE, 1997



Credit: Varda Space Industries

Classical Unsteady Aerodynamic Theories

Joukowski's transform: Conformal mapping from a simple geometric shape (circle) into a complex one (e.g., cambered airfoil, blunt bodies). The solution of velocity potential remains accurate [4].

Kutta Condition : Bounded circulation is conserved between the wake and on the body. Flow leaves the trailing edge smoothly. Lift is proportional to flow circulation [5].

How can we develop semi-empirical models to predict unsteady flow for a blunt body ?

Unsteady aerodynamic models: Theodorsen [6] provides a frequency-dependent solution for the unsteady lift and pitch moment resulting from oscillations. Builds on the steady-lift concepts derived from Joukowski and Kutta condition and extends them to unsteady flow via potential flow theory.

Assumptions: Potential flow theory. I.e., incompressible, irrotational, linear and mostly attached

$$C_{L_{Total}} = \underbrace{\frac{\pi b}{U_{\infty}^2} (U_{\infty} \dot{\alpha} - ba \ddot{\alpha})}_{C_{L_{Addedmass}}} + \overbrace{2\pi \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_{\infty}} \right] C(k)}^{C_{L_{US}}}$$

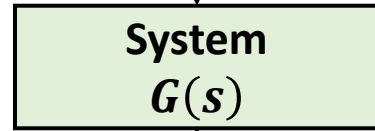
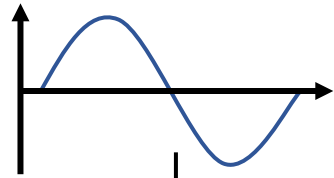
$C_{L_{QS}}$
Thin airfoil theory

Approximated transfer function to account for wake effects [7]

[4] Joukowski, N., "Über die konturen der Tragflächen der Drachenflieger," *Zeitschrift für Flugtechnik und Motorluftschiffahrt*, Vol. 1, No. 22, 1910, pp. 281–285.
 [5] Kutta, W., "Auftriebskräfte in strömenden Flüssigkeiten," *Illustrierte Aeronautische Mitteilungen*, Vol. 6, No. 133, 1902, pp.133–135.
 [6] Theodorsen, T., "General theory of aerodynamic instability and the mechanism of flutter," Tech. Rep. NACA TR-496, National Advisory Committee for Aeronautics, 1949..
 [7] Bisplinghoff, Raymond L, Holt Ashley, and Robert L. Halfman. *Aeroelasticity*. Courier Corporation, 2013, pp 232 - 255

Frequency Response Approach

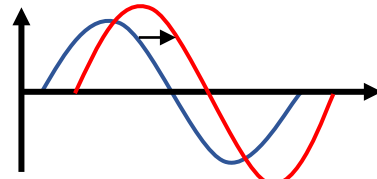
Input: $U(t) = A \sin(\omega t)$



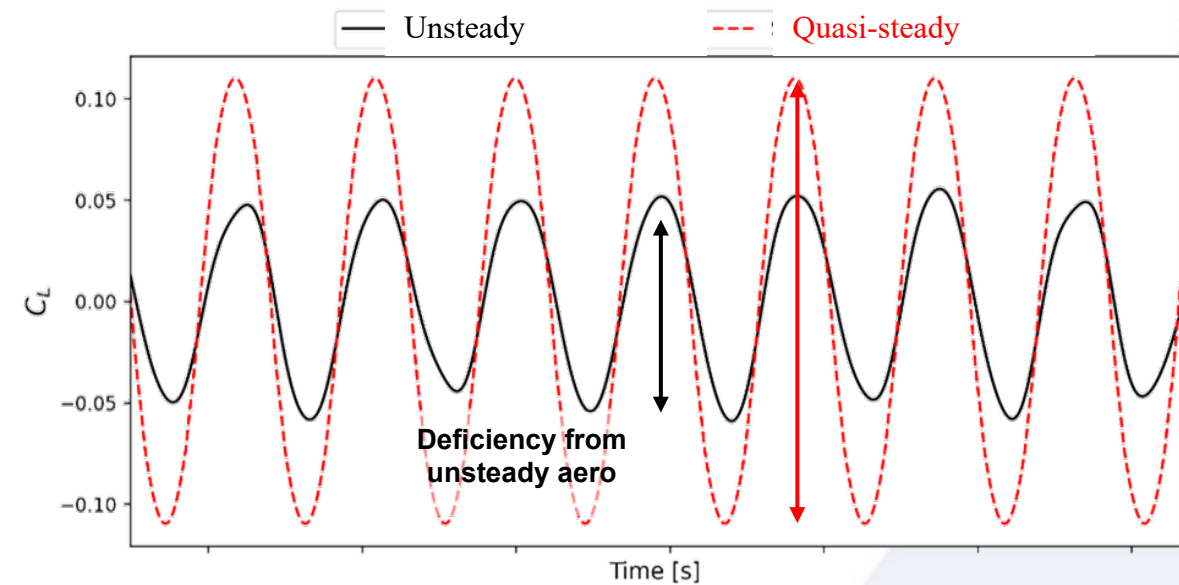
$$|G(s)| = \frac{A'}{A}$$

$$\angle G(s) = \phi$$

Output: $Y(t) = A' \sin(\omega t + \phi)$



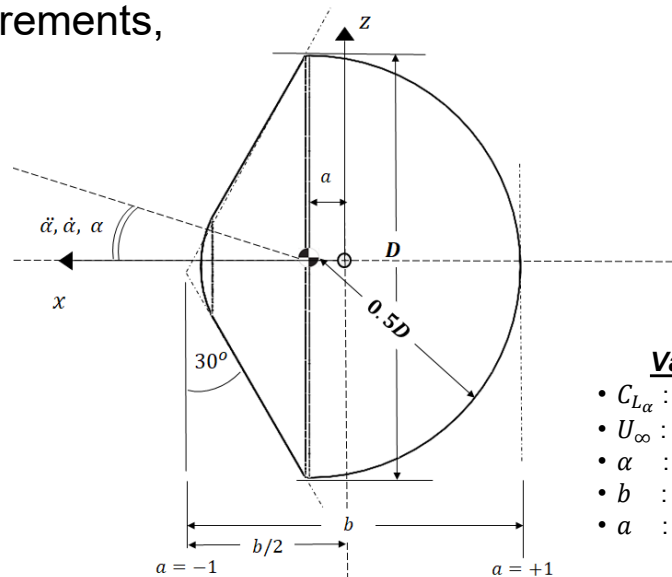
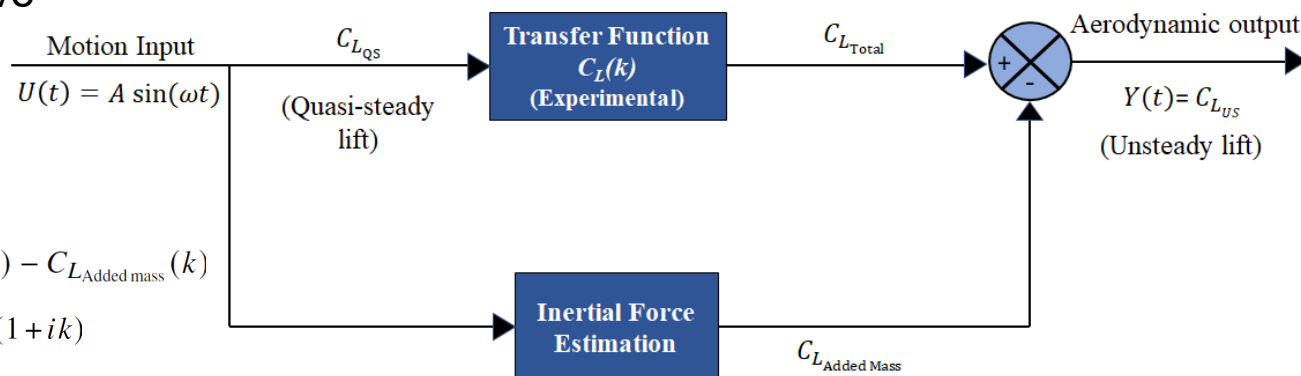
- Model a transfer function, $C(k)$ to describe the system's response (unsteady aerodynamics due to harmonic oscillatory input (quasi-steady aerodynamics))
 - Assumes the system is linear time-invariant (LTI) and that the frequency is preserved
 - Quasi-steady flow follows the motion input (gain of 1 and phase lag of 0°)
 - Manifests as a gain/deficiency, shift in the mean, and/or a phase lag
 - This process can be data-driven and informed by CFD or **dynamic test techniques**



A system identification like approach that **characterizes a transfer function for a blunt body entry vehicle** should be able to **predict lift and moment time histories** as a function of oscillatory inputs.

Frequency Response Approach to Unsteady Aerodynamics

An approximate transfer function from forced oscillation wind tunnel test measurements, addressing unsteady aerodynamic analysis from a dynamical systems perspective



Variable Definitions:

- $C_{L\alpha}$: Lift-curve slope
- U_∞ : Free-stream velocity
- α : Angle of attack
- b : Vehicle span
- a : Normalized position of the axis of rotation from the center of the body
- $C_L(k)$: Transfer function

The effect of unsteady aerodynamics can be captured with a transfer function (transfer of energy from quasi-steady to unsteady aero):

Lift and moment timeseries can be recreated for any given $C(k)$

$$\int_{-\infty}^{\infty} U(t) e^{i\omega t} dt$$

➤ The transfer function is a reduced frequency dependent complex number:

$$C(k) = F(k) + iG(k)$$

➤ $C(k)$ is a complex number that contains information on the gain/deficiency and phase lag resulting from dynamic motion

$$C_L(k) = \frac{C_{LUS}(k)}{C_{LQS}(k)}$$

$$C_L = \underbrace{\frac{C_{L\alpha} b}{2V_\infty^2} [U_\infty \dot{\alpha} - ba\ddot{\alpha}]}_{\text{Added Mass Effects}} + \underbrace{C_{L\alpha} \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right]}_{\text{Term modified by Transfer Function}} C(k)$$

$$C_m = \underbrace{\frac{C_{L\alpha} b}{2U_\infty^2} \left[-U_\infty b \left(\frac{1}{2} - a \right) \dot{\alpha} - b^2 \left(\frac{1}{8} + a^2 \right) \ddot{\alpha} \right]}_{\text{Added Mass Effects}} + \underbrace{C_{L\alpha} b \left(a + \frac{1}{2} \right) \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right]}_{\text{Term modified by Transfer Function}} C_L(k)$$

$k = \frac{\omega D}{2V_\infty}$: reduced frequency, a, b : geometric parameters

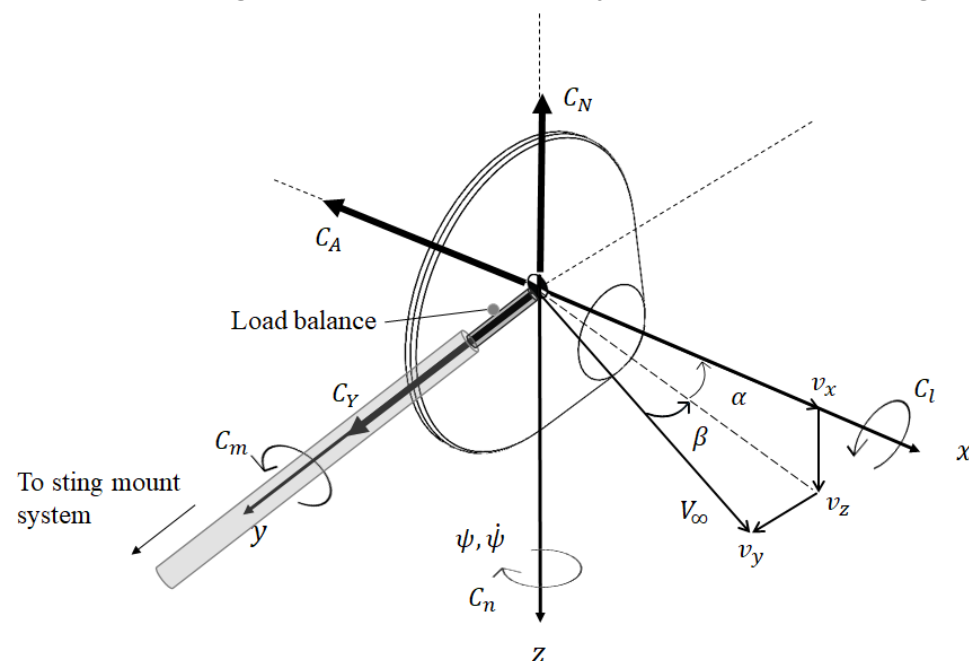
Entry Systems Model - Forced Oscillation Test

Forced oscillation test overview:

- A 60° spherical nose cone and hemispherical backshell vehicle was tested in the *12-Foot Low-Speed Tunnel at NASA Langley*.
- 21 runs were conducted at different reduced frequencies. 250 Hz sampling rate.
- Wind on and Wind off force and moment measurements
- Post-Processing : Cut-off frequency at 4X the driving frequency

Operating conditions

Mean Angle of Attack, α_0	Pitching amplitude, A_α	Reduced frequency, k
0°	10°	0.01 - 0.21



Body coordinate system



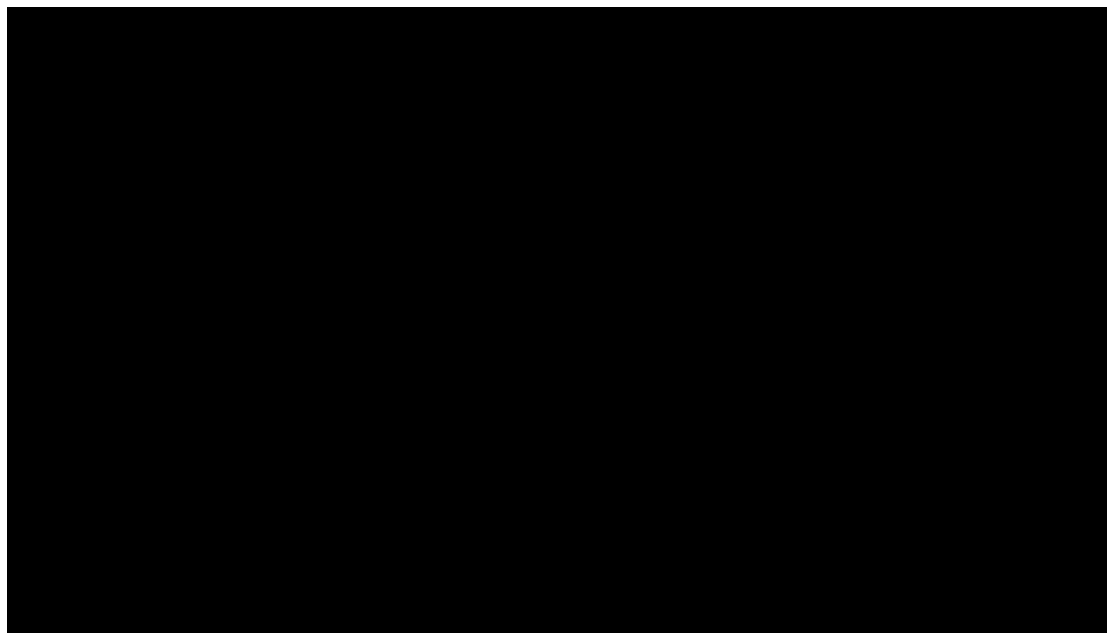
Forced oscillation test set up



Entry Systems Model - Forced Oscillation Test

Operating conditions

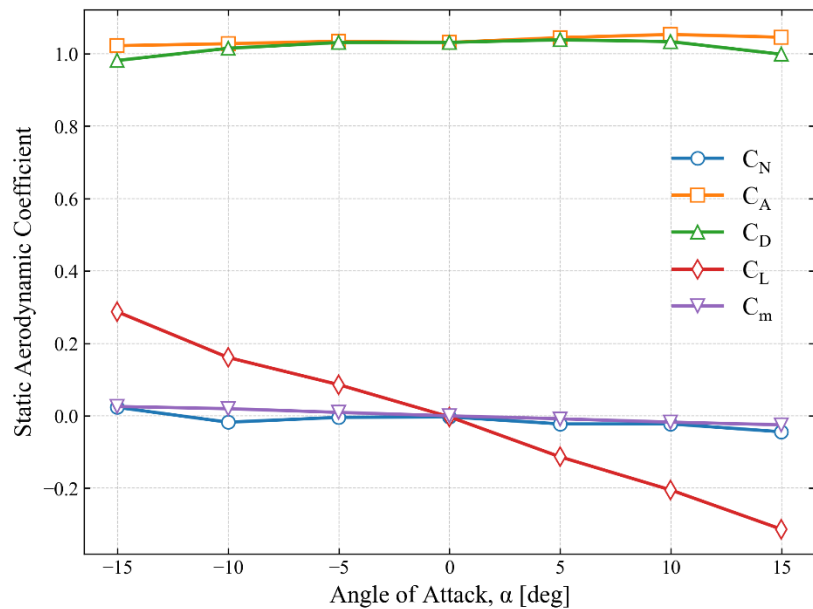
Mean Angle of Attack, α_0	Pitching amplitude, A_α	Reduced frequency, k
0°	10°	0.01 - 0.21



Results: Static and Unsteady Aerodynamics

Static Aerodynamic data:

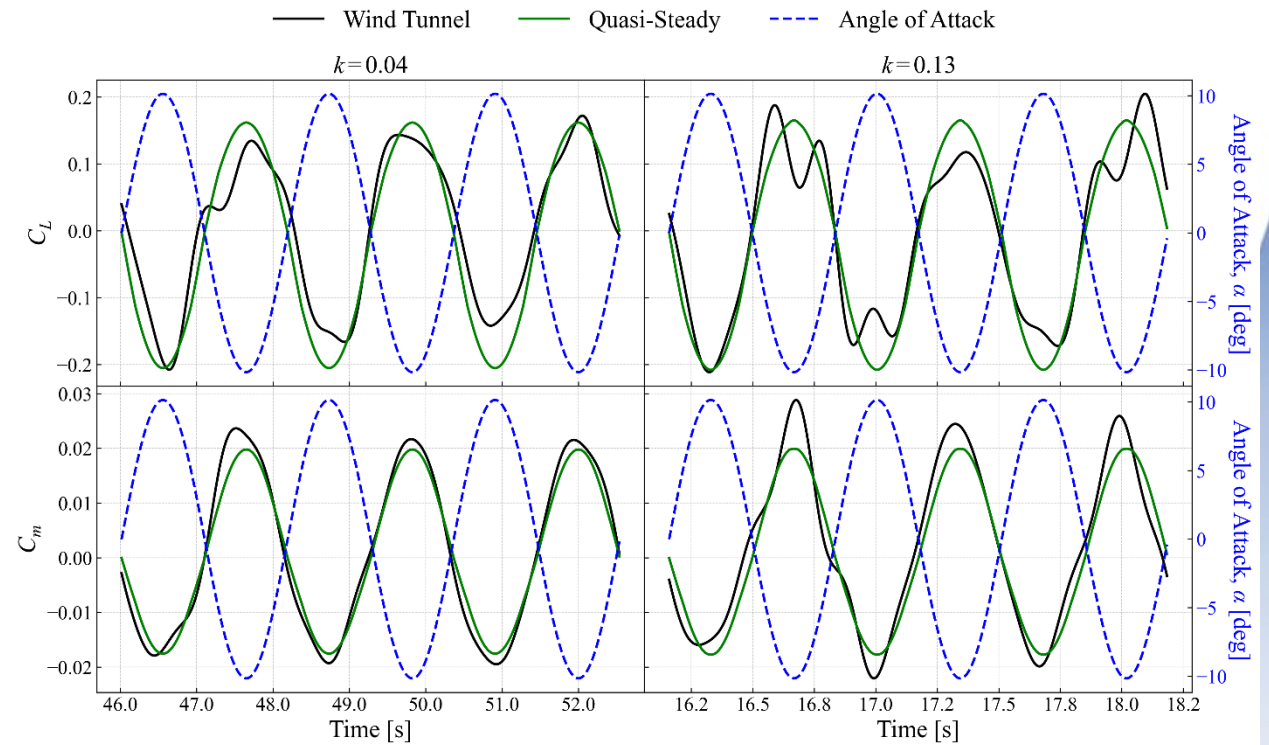
- Static data was also acquired in the 12-Foot tunnel for use in the reconstructions
- $C_{L\alpha}$ is determined from the static run to generate the quasi-steady lift and moment time-series.



Static aerodynamic data

Unsteady Aerodynamic test data

- Quasi-steady departs from the unsteady aerodynamic at larger k



Unsteady lift and moment time-histories



Results: Transfer Function Database

- Values of the transfer function can be stored in databases and queried to calculate C_L and C_m using oscillatory aerodynamic models [8]
 - These values can be stored and queried from aerodynamic databases
 - Typical aerodatabases do not capture frequency dependent and resonant phenomena. These can be captured using the following transfer function database

k	$C_L(k)$	$C_m(k)$	k	$C_L(k)$	$C_m(k)$
0.01	$1.021 - 0.035i$	$0.957 + 0.034i$	0.12	$0.767 - 0.163i$	$1.122 + 0.265i$
0.02	$1.019 - 0.109i$	$0.972 + 0.101i$	0.13	$0.773 - 0.154i$	$1.156 + 0.350i$
0.03	$0.999 - 0.102i$	$0.984 + 0.104i$	0.14	$0.768 - 0.148i$	$1.136 + 0.287i$
0.04	$0.939 - 0.106i$	$1.032 + 0.132i$	0.15	$0.753 - 0.178i$	$1.167 + 0.321i$
0.05	$0.940 - 0.173i$	$1.014 + 0.186i$	0.16	$0.772 - 0.169i$	$1.151 + 0.237i$
0.06	$0.919 - 0.126i$	$1.039 + 0.149i$	0.17	$0.734 - 0.181i$	$1.169 + 0.342i$
0.07	$0.857 - 0.185i$	$1.064 + 0.215i$	0.18	$0.692 - 0.131i$	$1.212 + 0.102i$
0.08	$0.882 - 0.164i$	$1.078 + 0.263i$	0.19	$0.721 - 0.116i$	$1.203 + 0.268i$
0.09	$0.844 - 0.167i$	$1.076 + 0.236i$	0.20	$0.711 - 0.129i$	$1.188 + 0.241i$
0.10	$0.844 - 0.232i$	$1.092 + 0.302i$	0.21	$0.682 - 0.141i$	$1.226 + 0.577i$
0.11	$0.824 - 0.180i$	$1.108 + 0.272i$			

Results: Lift Transfer Function

- Magnitude attenuation and growing phase lag with increasing k
 - Predictions compared with Theodorsen's analytical solution.
- Uncertainty in the phase grows as reduced frequency increases
 - Uncertainties are computed by the Root Sum Square (RSS) of various sources of error: Instrumentation precision, calibration, data interpolation/sampling and data trimming

$$C_L(k) = F \left\{ \frac{L_{wind}(k) - (N/N_{tare})L_{tare}(k)}{qS_{REF}C_{L_{QS}}(k)} \right\}$$

Note: Tare effects are removed through a frequency domain subtraction.

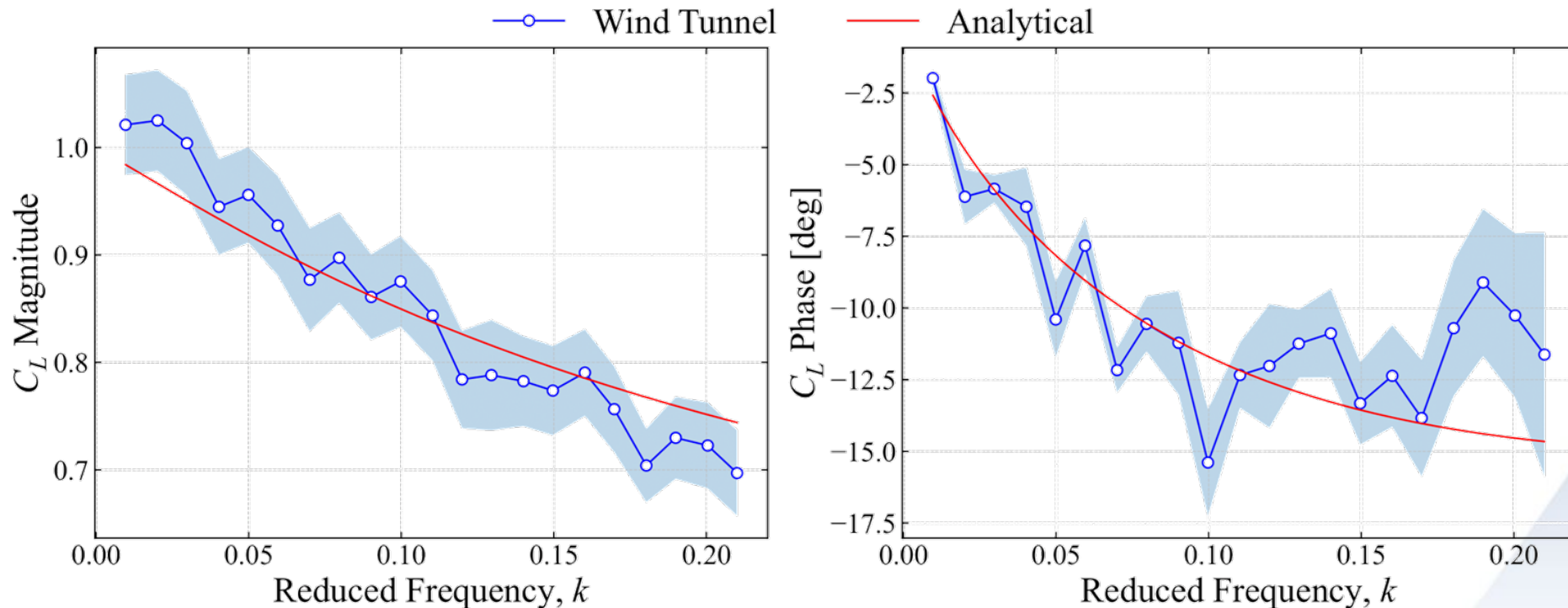


Fig. Bode plot for the lift transfer function and Theodorsen's analytical prediction with 3-sigma uncertainty bands

Results: Pitching Moment Transfer Function

- Magnitude amplification and growing phase shift with increasing k
- Uncertainty in the phase grows significantly as reduced frequency increases
 - Larger uncertainty bounds than were seen in the lift coefficient

$$C_m(k) = \mathcal{F} \left\{ \frac{M_{wind}(k) - (N/N_{tare})M_{tare}(k)}{qS_{REF}DC_{m_{QS}}(k)} \right\}$$

*Tare effects are removed through a frequency domain subtraction. * \mathcal{F} : Fourier transform*

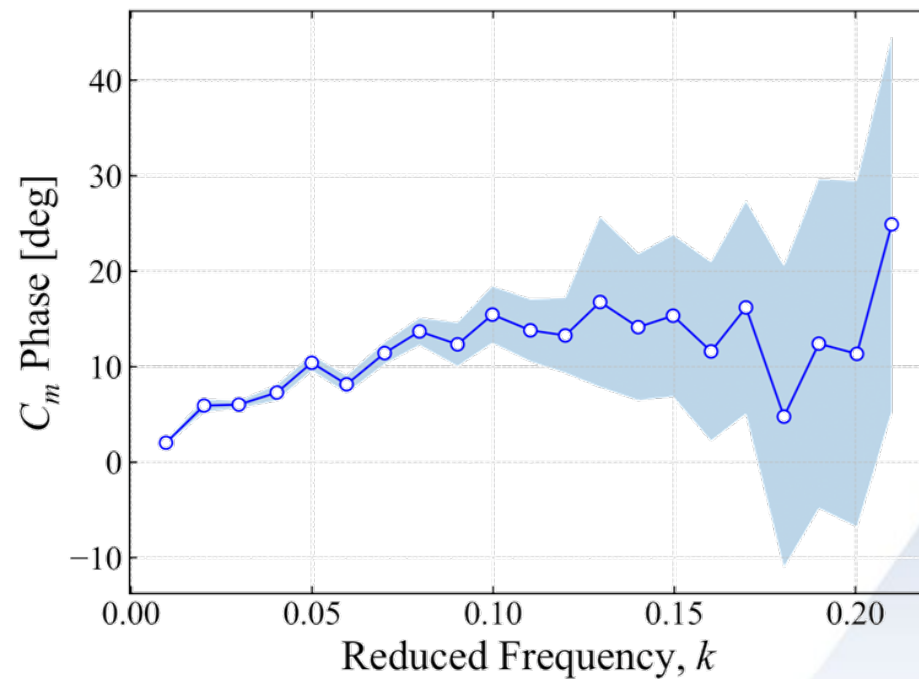
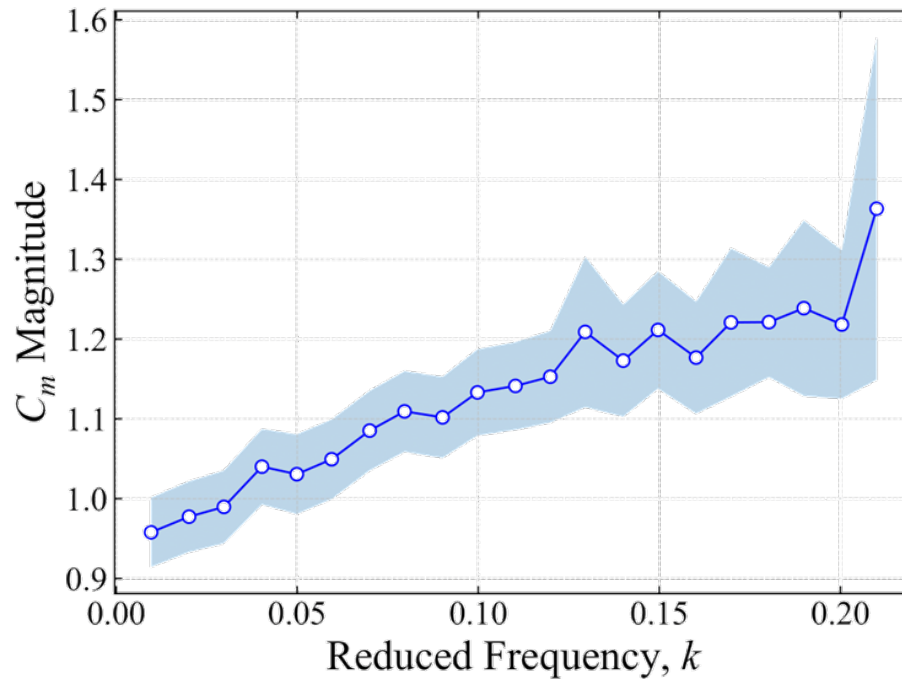


Fig. Bode plot for the pitching moment transfer function where the shaded region represents 3-sigma uncertainty

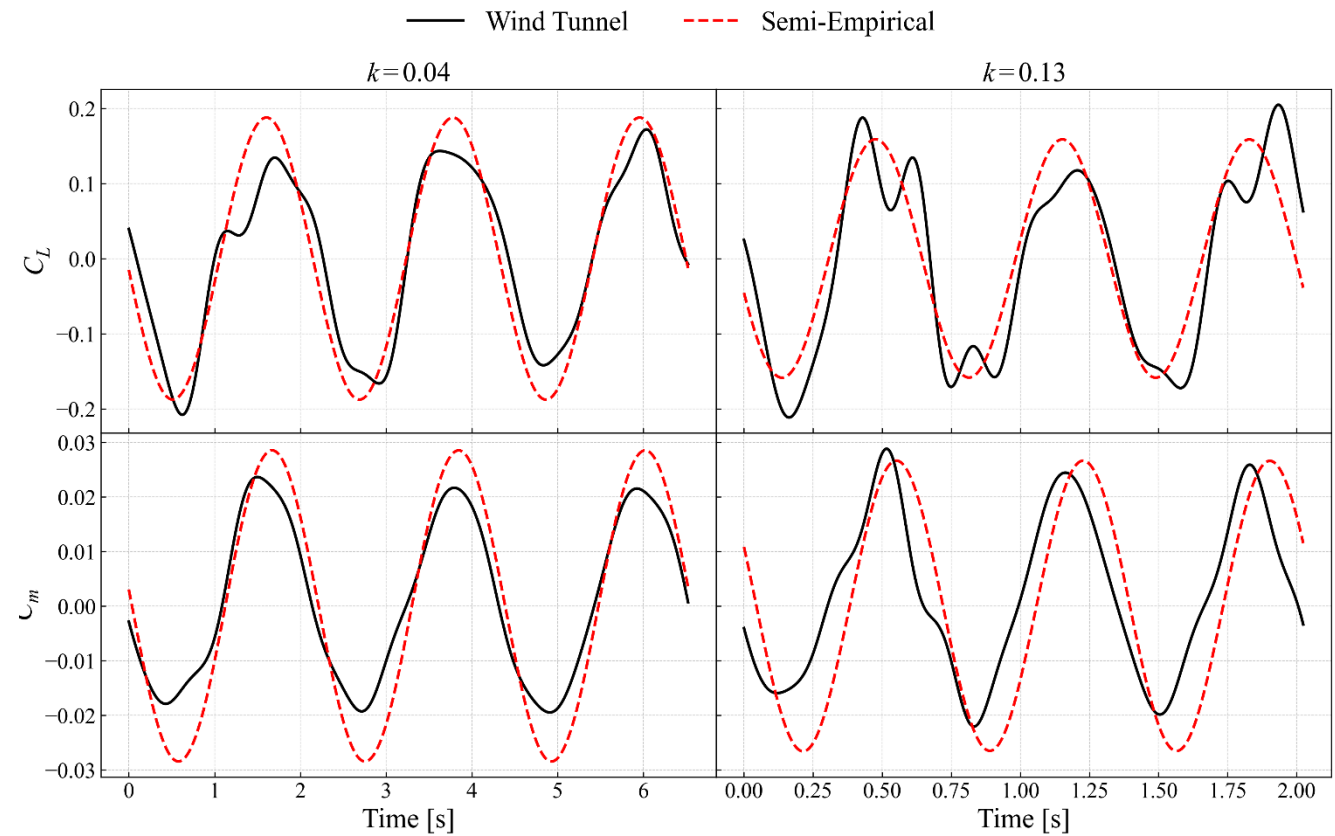
Results: Time-Varying C_L and C_m Reconstruction

➤ Lift and moment time histories can be recreated assuming Theodorsen's equations apply here

- Axis of rotation is assumed to be at the center of mass ($a = -0.329, b = 0.875\text{ft}$)
- Magnitude and phase discrepancies are observed as expected, since Theodorsen's applies to thin airfoils, not blunt-bodies

$$C_L = \frac{C_{L\alpha} b}{2V_\infty^2} [U_\infty \dot{\alpha} - ba\ddot{\alpha}] + C_{L\alpha} \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right] C_L(k)$$

$$C_m = \frac{C_{L\alpha} b}{2U_\infty^2} \left[-U_\infty b \left(\frac{1}{2} - a \right) \dot{\alpha} - b^2 \left(\frac{1}{8} + a^2 \right) \ddot{\alpha} \right] + C_{L\alpha} b \left(a + \frac{1}{2} \right) \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right] C_L(k)$$



Reconstruction of the lift and moment time series for two random reduced frequencies.

Results: Time-Varying C_L and C_m Reconstruction w/ Regression Analysis

- Discrepancy can be reduced using a regression approach: a constant can be added to each term that is solved for using a least squares regression
 - All the time series are used together to solve for the constants

$$C_L = \mathbf{C}_1 \frac{C_{L\alpha} b}{2V_\infty^2} [U_\infty \dot{\alpha} - ba\ddot{\alpha}] + \mathbf{C}_2 C_{L\alpha} \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right] C_L(k)$$

$$C_m = \mathbf{C}_3 \frac{C_{L\alpha} b}{2U_\infty^2} \left[-U_\infty b \left(\frac{1}{2} - a \right) \dot{\alpha} - b^2 \left(\frac{1}{8} + a^2 \right) \ddot{\alpha} \right] + \mathbf{C}_4 C_{L\alpha} b \left(a + \frac{1}{2} \right) \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right] C_m(k)$$

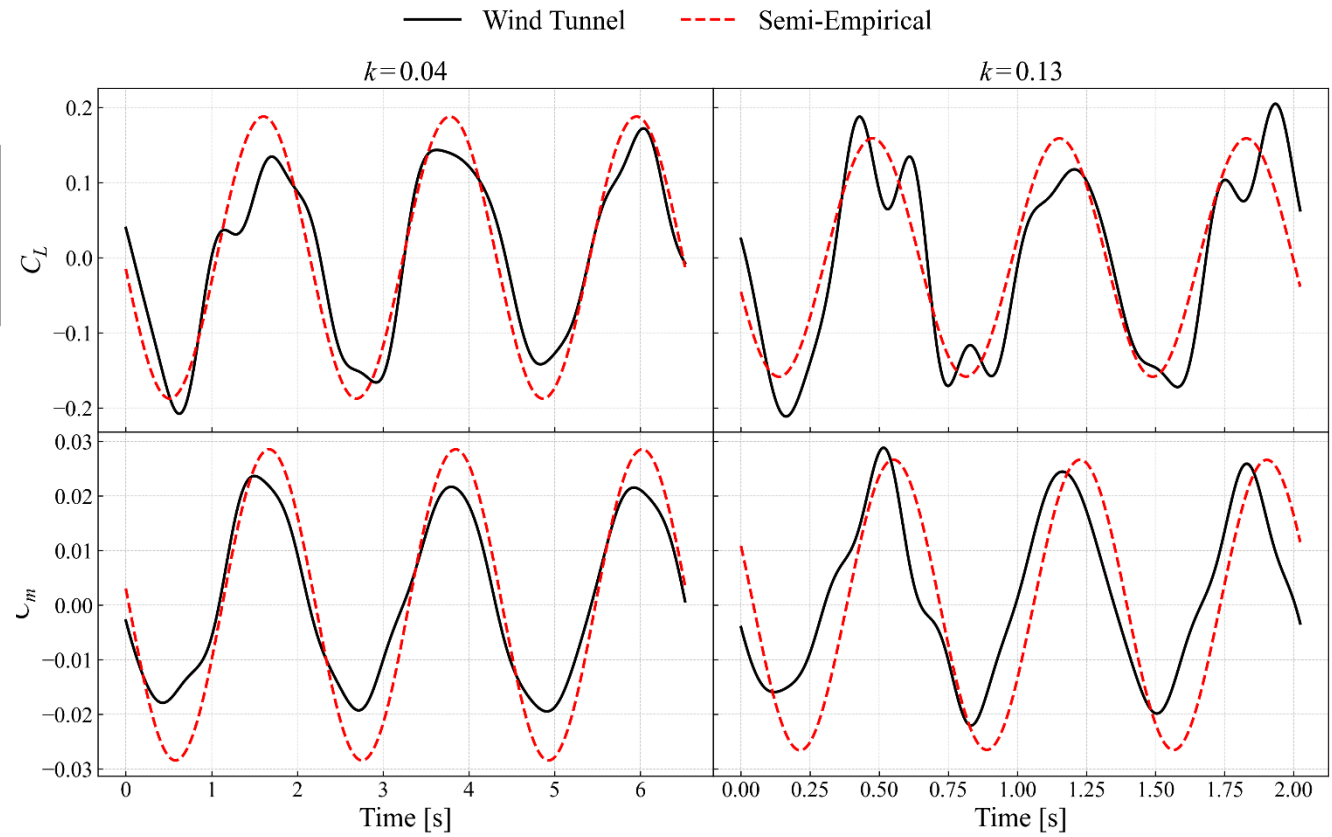


Fig. Reconstruction of the lift and moment time series for two random reduced frequencies. Tare-removed coefficients are shown and only 3 of the 40 cycles are shown.

Results: Time-Varying C_L and C_m Reconstruction w/ Regression Analysis

➤ Discrepancy can be reduced using a regression approach: a constant can be added to each term that is solved for using a least squares regression

- All the time series are used together to solve for the constants
- Results in better predictions (more detailed error analysis is given in the paper)

$$C_L = C_1 \frac{C_{L\alpha} b}{2V_\infty^2} [U_\infty \dot{\alpha} - ba\ddot{\alpha}] + C_2 C_{L\alpha} \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right] C_L(k)$$

$$C_m = C_3 \frac{C_{L\alpha} b}{2U_\infty^2} \left[-U_\infty b \left(\frac{1}{2} - a \right) \dot{\alpha} - b^2 \left(\frac{1}{8} + a^2 \right) \ddot{\alpha} \right] + C_4 C_{L\alpha} b \left(a + \frac{1}{2} \right) \left[\alpha + b \left(\frac{1}{2} - a \right) \frac{\dot{\alpha}}{U_\infty} \right] C_m(k)$$

C_1 (C_L Added Mass)	C_2 (C_L Unsteady)	C_3 (C_m Added Mass)	C_4 (C_m Unsteady)
-2.77	0.924	-0.146	0.598

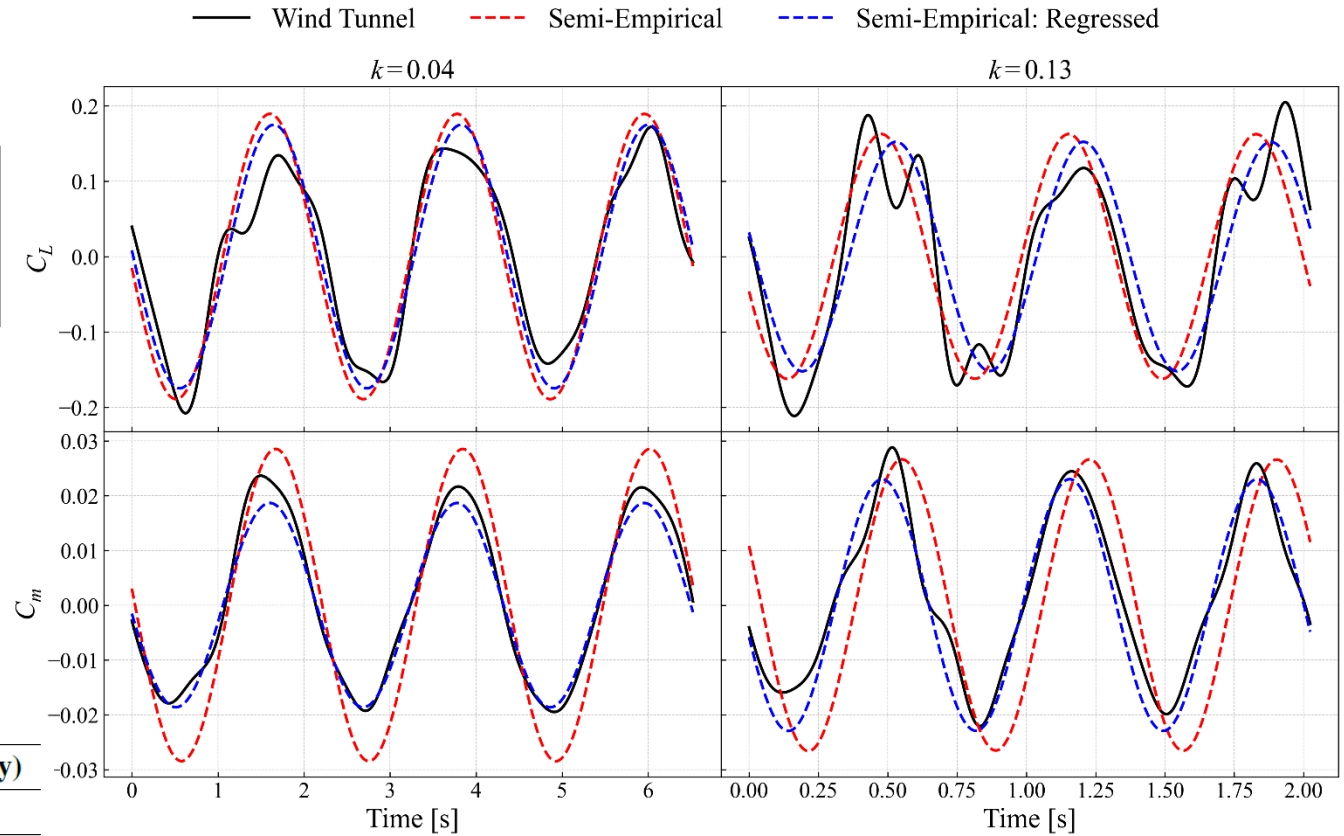
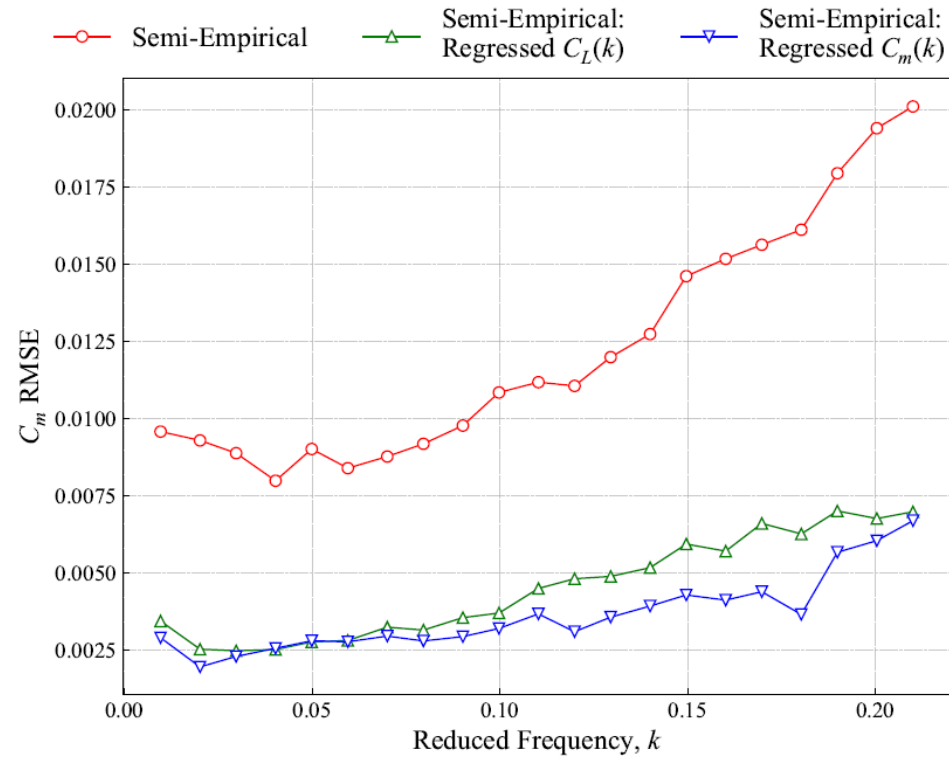


Fig. Improved reconstruction of the lift and moment time series for two random reduced frequencies.

- Discrepancy can be reduced using a regression approach: a constant can be added to each term that is solved for using a least squares regression
 - The Root mean square error (RMSE) of the moment coefficient reconstruction shows that regressed model provides better predictions to the wind tunnel test data
 - $C_m(k)$ captures the phase shift in flow response reasonably better than $C_L(k)$





Summary

Conclusions

- Determined Transfer functions for a blunt body entry vehicle from forced oscillation tests from NASA Langley's 12 ft Low Speed Tunnel.
- Transfer functions describe the gain and the phase of the aerodynamic response due to oscillatory input motion as a function of frequency.
- Modified classical unsteady aerodynamic theories into semi-empirical formulations tailored to blunt body entry vehicle to recreate lift and moment time histories.
- Applied regression analysis with newly determined regressed coefficients to the existing semi-empirical model to better reconstruct the time series of lift and pitch moment.

Future Work:

- Stability coefficients can be extracted from forced oscillation data using a frequency approach as a new data reduction technique. [9]
- Implement the determined transfer functions to develop of frequency dependent dynamic stability coefficients models in aerodynamic databases.

