TIME BEHAVIOR OF SOLAR FLARE PARTICLES TO 5 AU

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Abstract

A simple model of solar flare radiation event particle transport is developed to permit the calculation of fluxes and related quantities as a function of distance from the sun (R). This model assumes the particles spiral around the solar magnetic field lines with a constant pitch angle. The particle angular distributions and onset plus arrival times as functions of energy at 1 AU agree with observations if the pitch angle distribution peaks near 90°. As a consequence the time dependence factor is essentially proportional to R^{-1.7}, (R in AU), and the event flux is proportional to R^{-2}.

I. Introduction

Solar flare particle events constitute one of the important sources of natural nuclear radiation at 1 AU from the sun. Approximately 100 such events have been observed at the earth (see Table I), ranging up to \(10^{10}\) protons/cm²-event above 10 Mev. Various particle flux models have been developed describing the time and energy behavior observed at 1 AU. These models, coupled with prediction mechanisms (usually correlated to sunspot numbers) have been used to estimate the expected solar flare particle environment at 1 AU for future space missions. These techniques have also been used for missions to Mars (1.5 AU) and Venus (0.7 AU).

For missions which involve sending spacecraft far from 1 AU (e.g., the Mercury-Venus and the Outer Planet missions), the solar flare particle model developed from observations at 1 AU is not adequate. Specifically, the dependence of the solar flare particle event model on distance from the sun must be incorporated. This paper presents a simple model for incorporating this spatial dependence which agrees with the observations at 1 AU.

II. Model at 1 AU

The flux of solar flare particles observed at 1 AU is observed to rise quasi-linearly over a period of some hours to a peak value, then decay approximately exponentially over an appreciably longer period. The onset, rise, and characteristic decay times are functions of particle energy. A previously developed mathematical expression which has these characteristics is:

\[
\phi(E_o, t) = \frac{A \cdot e^{-\sigma E_o n t}}{E_o^m} \text{ particles/cm}^2\text{-hr} \tag{1}
\]

where \(E_o\) is the particle energy (MeV), \(t\) is the time from the arrival of the first particles (hrs), \(A\) is a normalization factor, \(\sigma\) approximately constants

\(n\)

If the characteristics of the largest solar flare particle events are used to evaluate \(\sigma\), \(n\), and \(m\), the values obtained are approximately 0.022, 0.4, and 0.75, respectively. An attempt to incorporate the variations of these quantities due to radiation event size led to the relationships:

\[
A = \frac{\phi(>E_o) = 2100 A}{e^{2 \cdot m + 2n}} \frac{\text{particles}}{E_o^m} \tag{2}
\]

\[
\text{rise} = \frac{1}{\sigma E_o^n} = \frac{45}{E_o^{0.4}} \text{ hr} \tag{3}
\]

\[
\text{decay} = 2.15 \frac{E_o^n}{E_o^{1.15}} \text{ hr} \tag{4}
\]

\[
\phi(E_o) = \frac{3200 A}{E_o^{2.55}} \frac{\text{particles}}{cm^2\cdot\text{hr-Mev}} \tag{5}
\]

These relationships have proven useful when applied to large solar flare radiation events at 1 AU. In particular, they have been applied to events for which the data was incomplete--some of the numbers in Table I were obtained this way.

These formulae have no dependence on the distance from the sun, and are therefore strictly
applicable to 1 AU only. In the following sections, modifications of these formulae to account for distance from the sun are discussed.

III. Particle Transport

In order to account for the dependence of the solar flare particle event model on distance from the sun, it is necessary to have a particle transport model. Since particle acceleration and particle transport are related during the acceleration portion of the particle event, it is desirable to go back to the solar flare itself.

A solar flare is an event observed to take place above the photosphere of the sun, generally lasting < 20 minutes. The flare, which is most readily observed at the Lyman-α wavelength of ionized hydrogen (~1215 Å), is believed to derive its energy from the collapse of the magnetic fields associated with sunspots. Various theories have been advanced to account for the acceleration of charged particles during or as a result of a solar flare. Among these theories are betatron acceleration and shock wave acceleration. In betatron acceleration, a charged particle orbiting above a collapsing sunspot gains energy as in a laboratory particle accelerator of the same name. In shock wave acceleration the absorption of the Lyman-α photons in the inner corona produces a supersonic blast wave which accelerates the charged particles. These proton acceleration mechanisms have been investigated by Gold(4), Wentzel,(5) Parker,(6) Weddell,(7) and others. The evidence is not conclusive for any single mechanism, and both betatron and shock wave accelerations probably contribute.

Once the charged particles have been accelerated, their propagation away from the sun is controlled by the interplanetary magnetic fields. These interplanetary fields are complex, having time and spatial dependences which are not well known. However, by making some simplifying assumptions, it is possible to obtain a spatially dependent solar flare particle model.

If it is assumed that the quiet sun interplanetary magnetic field is essentially undisturbed during the propagation of the solar flare protons, it is possible to write:

$$B = \sqrt{\left(\frac{B_r^2}{R^2}\right) + \left(\frac{B_\theta^2}{R^2}\right)}$$

$$Bqv_l = \frac{m v_l}{r}$$

$$J = \frac{m v_l}{r}$$

$$v_l = v \sin \beta$$

$$v_\parallel = v \cos \beta$$

where

- $B$ is the solar magnetic field (webers/m²)
- $B_r$ is the radial component of $B$ (webers/m²)
- $B_\theta$ is the azimuthal component of $B$ (webers/m²)
- $q$ is the charge of the particles (coulombs)
- $v$ is the velocity of the charged particles (meters/sec)
- $v_l$ is the perpendicular component of $v$ (meters/sec)
- $v_\parallel$ is the parallel component of $v$ (meters/sec)

$\beta$ is the pitch angle (angle $v$ makes with the guiding center line of $B$). Since $B_r = B_\theta = 3.5 \times 10^{-9}$ webers/m² (3.5 gammas) at 1 AU, $B$ is ~5 gammas, and the gyroradius is $2.75 \times 10^7$ meters, where $E$ is in Mev. The gyroradius as a function of distance from the sun is approximately

$$r = 3.9 \times 10^7 \frac{E}{m^3} \text{ meters}$$

This function is plotted in Figure 2.

Assuming the particles spiral around the magnetic lines of force, it is expected that their gyroradius will increase as their distance from the sun increases. Based upon the above equations it is expected that their pitch angle $\beta$ will decrease with distance from the sun, according to the well-known relationship.

$$\frac{J q}{2 m E} = \text{constant}$$

However, the pitch angle ($\beta$) cannot exceed 90°.

If the charged particle has this limiting angle close to the sun, its pitch angle might be expected to become quite small by the time it reaches 1 AU. If a particle has a pitch angle of 90° starting at 10 solar radii, the calculated pitch angle will be ~3° by the time it reaches 1 AU (~215 solar radii) according to this model. Since solar flare event particles are observed at 1 AU to be essentially isotropic after the flux peak, this simplified model is clearly unrealistic.

It is possible to obtain reasonable agreement with observations at 1 AU by making the assumption that the pitch angle is independent of distance from the sun. This accounts for the fact that particles of a given energy with small pitch angles will arrive first (as is observed), and the resultant particle angular distributions spread as functions of time. To the extent that this model appears to agree with the observations at 1 AU, it may be useful for calculating the spatial dependence of solar flare particle radiation.

The transit time ($t_{tr}$) as a function of particle energy and pitch angle (both assumed to be constants of the motion) is

$$t_{tr} = \frac{s}{v_\parallel}$$

where $s$ is the length of the line of magnetic flux which serves as the guiding center for the particle gyration. This length(s) may be calculated from the relationship:
But since
\[ \tan^{-1} \frac{r \, d\theta}{dr} = \tan^{-1} \frac{R}{Br} \]

Therefore
\[ r \, d\theta = R \, dr \]

Thus
\[ s = \int \frac{r}{\sqrt{1 + R^2}} \, dr \]

\[ = \frac{1}{2} \left[ R \sqrt{1 + R^2} + \log \left( R + \sqrt{1 + R^2} \right) \right] \quad (7) \]

When the distance from the sun (R) is small the two terms are approximately equal in value, but as the distance from the sun increases, the first term predominates. This function is graphically displayed in Figure 3. A reasonable approximation is:
\[ s \approx -R \sqrt{1 + R^2} \quad (8) \]

If it is assumed that the pitch angle (\( \beta \)) of the particle is a constant of the motion, the transit time as a function of pitch angle and particle energy may be calculated. Based on Equation (7), the result is (for protons):
\[ t_{tr} = 5.5 \times 10^{3} \frac{R \sqrt{1 + R^2} + \log \left( R + \sqrt{1 + R^2} \right)}{E^0 \cos \beta} \quad (9) \]

where R is in AU and E is in Mev. With only a few percent error in the region of interest (0.1 - 10 AU) this may be approximated by
\[ t_{tr} = 1.1 \times 10^{4} \frac{R}{\sqrt{E^0 \cos \beta} \sqrt{1 + R^2}} \quad (10) \]

Values of this function are listed in Tables II-IV, and are plotted in Figures (4) to (6). These figures show the expected dependence on proton energy (E), pitch angle (\( \beta \)), and distance from the sun (R). In principle, this expression may be solved to obtain any of the quantities as functions of the others. The quantity of major interest is the proton flux as a function of these parameters. This is determined by the initial proton distribution in \( \beta \), assuming no protons are lost after being accelerated.

If no particles are lost after being accelerated, the total number per cm\(^2\) per event is expected to be proportional to R\(^{-2}\) (R = distance from the sun in AU). Assuming an R\(^{-2}\) time-integrated spatial dependence effectively makes the solar flare particle event probability spatially independent. Any other assumption introduces the complication of a spatially dependent probability function. Since the time-integrated particle flux for a mission is essentially the product of the flux per event and the probability of that event, little, if anything, is gained by considering a spatially dependent particle event probability.

If a significant fraction of solar flare particles is lost after being accelerated, the flux per event will exhibit an R\(^{-n}\) spatial dependence where n > 2 over the region where particle loss is important. The major probable loss mechanism is collision with objects in space (magnetospheres, planets, moons, etc.) Considering the relative emptiness of space this is not expected to be a major loss mechanism. Even the asteroid belt is only \( \sim 10^{-6} \) opaque, which means that the solar flare particle traversing it will be small. (It may not be completely negligible, however, since the proton helical path length through the asteroid belts will be \( \sim \) two orders of magnitude larger than that of a solar photon). The probability of a solar flare particle encountering a planet, even assuming the particles are essentially confined to an ecliptic plane is \(< 1\%\).

It is interesting to calculate the pitch angle which corresponds to the peak flux rate. At 1 AU the onset + rise time (time from the flare on the sun to the peak flux rate) is approximately:
\[ t_{rise} = \frac{45}{E_0^0.4} \text{ hrs} \]

where \( E_0 \) is the proton energy in Mev. (This relationship was obtained by averaging the values observed for the largest radiation events--smaller events often have shorter onset + rise times). Since the proton transit time (given by Equation 9) is proportional to \( R^{-0.5} \), there is apparently a constant delay time to be added. A little arithmetic shows that this delay time is \( 860/\cos \beta \) seconds. Thus:
\[ t_{rise} = t_{tr} + \frac{860}{\cos \beta} \text{ seconds} \]

This accounts for the acceleration time during which the protons do not migrate away from the sun significantly as well as the time between the arrival of the first protons of a given energy and the peak flux of these protons. It will be noted that a fairly steep proton energy spectrum has been assumed (based on observations at 1 AU), since the onset + rise time for particles above energy \( E_0 \) is taken to be essentially that of particles of energy \( E_0 \).

For 1, 10, and 100 Mev protons at 1 AU, the observed onset + rise times (\( t_{rise} \)) are \( \sim 1.6 \times 10^{5} \) seconds, \( \sim 6.5 \times 10^{4} \) seconds, and \( \sim 2.6 \times 10^{4} \) seconds, respectively. The corresponding transit times are \( \sim 1.5 \times 10^{5} \) seconds, \( \sim 5.2 \times 10^{4} \) seconds, and \( \sim 1.5 \times 10^{4} \) seconds. These are the transit times for protons with pitch angles of \( \sim 90^\circ \). It will be noted that any change of pitch angle as the protons migrate away from the sun will decrease the transit time below that observed. Thus, unless there is some mechanism which acts to increase the
pitch angles of protons as they move away from
the sun, the assumption of constant pitch angle
appears to be superior to any obvious simple
alternative. Also noteworthy, is that beyond
2 AU solar flares on the back side of the sun can
be important sources of particles.

Another interesting conclusion, valid to the
extent that the assumption of a constant pitch
angle is valid is that the bulk of the protons
have pitch angles close to 90°. Based upon
Equation (2), only ~26% of the protons above a
given energy have arrived at 1 AU by the time the
peak flux rate of these protons has been reached.
Thus, ~74% of the protons observed at 1 AU
appear to have pitch angles between 86° and 90°.
Assuming a random distribution of proton velocity
directions prior to acceleration, a sin β pitch
angle distribution would be expected from solid
angle considerations. Apparently, the acceler-
ation mechanism favors particles with pitch angles
close to 90°. This seems to favor the betatron
acceleration mechanism over the shock wave accel-
eration mechanism. However, since the constant
pitch angle assumption is merely a simplified
model, any conclusions based upon it must be con-
sidered uncertain.

The maximum pitch angle determines the limits
of the particle angular distribution as a function
of time and energy. The particle angular dis-
tribution within these time and energy dependent
limits according to this model is determined by
the initial particle pitch angle distribution.
However, not enough angular distribution infor-
mation is available to derive statistically mean-
ingful pitch angle distributions as functions of
time and energy. When such data becomes available,
its incorporation into this model is straight-
forward.

IV. Spatially Dependent Model

It is now possible to use the transport
model to calculate the parameters of the solar
flare particle radiation as functions of distance
from the sun. Since the pitch angle is assumed
to remain constant, so does \( v_g \) (the particle
velocity parallel to the line of magnetic force
which serves as the gyrorotation guiding center).
The time for a particle to reach a given distance
from the sun is simply proportional to the spiral
path length (given by Equation 7). Thus, the
time dependence of the solar flare particle flux
is expected to be

\[
t' \sim \frac{t}{f(R)} + t_0
\]

where \( f(R) \) is given by Equation (7) and \( t_0 \) is ~1.2 \times 10^4 sec. It is, therefore, possible to
rewrite Equations (1) to (5) incorporating \( f(R) \) (the
distance from the sun, as a parameter). The
results are:

\[
\phi(E, t) = \frac{A}{R^2 f^2(R) E_0} e^{-0.022E_0^{0.4} f^{-1}(R)t}
\]

where all symbols have been previously identified.
As mentioned previously, \( f(R) \) can be approxi-
mated by

\[
f(R) \approx R^{1.7}
\]

For \( R > 1 \), this may be further approximated by \( R^{1.7} \)
with \( \leq 15\% \) error at 1 AU and a smaller error for
\( 1 < R < 5 \).

The relative spatial dependences of the time-
integrated particle flux rate \( \sim R^{-2} f^{-2}(R) \)
are shown in Figure 7. These curves show that the
relative solar flare particle environment becomes
severe as the sun is approached. For the largest
such particle event observed at 1 AU (11-12-60),
the peak and time-integrated proton fluxes above
10 Mev at Mercury (\( \sim 0.4 \) AU) would be \( \sim 7.2 \times 10^3 \)
particles/cm^2-sec and \( \sim 6.3 \times 10^4 \)
particles/cm^2-Mev, respectively --

\[
\frac{[0.75 + 0.0088t^{-1}(R)E_0^{0.4}]}{f(R)} \text{ particles/cm}^2\text{-hr-Mev}
\]

which is enough to affect sensitive spacecraft components.
For this same event the corresponding numbers at
Jupiter (\( \sim 5.2 \) AU) would be 120 p/cm^2-sec and
3.7 \times 10^3 p/cm^2-Mev, which should cause no trouble.

It will be noted that particle loss and
change of particle energy after acceleration have
been neglected. If either of these changes takes
place (and there is some evidence that they do),
the solar flare particle environment as the sun
is approached will be even more severe than here
calculated. Conversely, at distances greater than
1 AU, the environment may very well be less severe
than expected.
V. Conclusions

The solar flare particle event model developed to fit the observed data at 1 AU has been modified to account for the distance from the sun. This was accomplished by assuming each flare particle independently spirals around the solar magnetic field (carried frozen-in by the solar wind), while maintaining its pitch angle constant. Among the consequences of this transport model are:

(a) The majority of the solar flare particles have pitch angles close to 90°. Calculated pitch angle particle distributions as a function of time agree fairly well with observations at 1 AU.

(b) The solar flare particle flux rate is quite dependent on distance from the sun \( \sim R^{-3.7} \), while the time integrated particle flux obeys the expected inverse square distance relationship \( R^{-2} \).

(c) The solar flare particle energy spectra are essentially independent of distance from the sun, since particle losses after acceleration are neglected. Some particle losses probably occur, leading to a softening of the energy spectra and larger flux spatial dependences than herein calculated.

This solar flare particle event model is directly applicable to missions which involve sending a spacecraft on missions in the ecliptic plane far from 1 AU from the sun. It is particularly applicable to missions to Mercury and to the outer planets. It is not intended to apply far from the ecliptic plane or past the heliosphere. Within these limits it provides an easy-to-use improvement over spatially independent models.

VI. Nomenclature

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Definition</th>
</tr>
</thead>
<tbody>
<tr>
<td>A</td>
<td>Normalization factor for solar flare particle event model</td>
</tr>
<tr>
<td>AU</td>
<td>Astronomical unit ((1.5 \times 10^8 , \text{km}))</td>
</tr>
<tr>
<td>B</td>
<td>Solar magnetic field strength ((\text{webers/meter}^2))</td>
</tr>
<tr>
<td>B_r</td>
<td>Radial component of solar magnetic field strength ((\text{webers/meter}^2))</td>
</tr>
<tr>
<td>B_t</td>
<td>Tangential component of solar magnetic field strength ((\text{webers/meter}^2))</td>
</tr>
<tr>
<td>E</td>
<td>Particle energy ((\text{MeV}))</td>
</tr>
<tr>
<td>E_0</td>
<td>Particle energy limit ((\text{MeV}))</td>
</tr>
<tr>
<td>J</td>
<td>Angular momentum of particle gyrorotation around solar magnetic field line ((\text{kg-meter}^2\cdot\text{sec}^{-1}))</td>
</tr>
<tr>
<td>m</td>
<td>Parameter in solar flare particle event model</td>
</tr>
<tr>
<td>q</td>
<td>Particle charge ((\text{coulomb}))</td>
</tr>
<tr>
<td>R</td>
<td>Distance from sun ((\text{AU}))</td>
</tr>
<tr>
<td>r</td>
<td>Radius of particle gyrorotation around solar magnetic field line ((\text{meters}))</td>
</tr>
<tr>
<td>s</td>
<td>Length of solar magnetic field flux line ((\text{meters}))</td>
</tr>
<tr>
<td>t</td>
<td>Time ((\text{hrs}))</td>
</tr>
<tr>
<td>t_rise</td>
<td>Onset + rise time ((\text{time from flare on sun to particle flux maximum observed at R AU}) \ ((\text{seconds}))</td>
</tr>
<tr>
<td>t_r</td>
<td>Particle transit time from the sun ((\text{seconds}))</td>
</tr>
<tr>
<td>v</td>
<td>Particle velocity ((\text{meters/sec}))</td>
</tr>
<tr>
<td>v_H</td>
<td>Component of particle velocity parallel to solar magnetic field ((\text{meters/sec}))</td>
</tr>
<tr>
<td>v_L</td>
<td>Component of particle velocity perpendicular to solar magnetic field ((\text{meters/sec}))</td>
</tr>
<tr>
<td>( \alpha )</td>
<td>Parameter in solar flare particle event model</td>
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<tr>
<td>( \beta )</td>
<td>Pitch angle of particle velocity with respect to solar magnetic field ((\text{degrees}))</td>
</tr>
<tr>
<td>( \phi )</td>
<td>Particle (proton) flux ((\text{particles/cm}^2\cdot\text{hr}\cdot\text{Mev}))</td>
</tr>
<tr>
<td>( \phi_t )</td>
<td>Time-integrated particle flux ((\text{particles/cm}^2))</td>
</tr>
<tr>
<td>( \hat{\phi} )</td>
<td>Peak particle flux rate ((\text{particles/cm}^2\cdot\text{hr}))</td>
</tr>
<tr>
<td>( \theta )</td>
<td>Azimuthal angle ((\text{polar coordinate}))</td>
</tr>
</tbody>
</table>

VII. References


### Table I. Annual Totals for Solar Flare Radiation Event Particle Fluxes at 1 AU

<table>
<thead>
<tr>
<th>Year</th>
<th>Number of Events ($f_{&gt;10\text{ Mev}}$)</th>
<th>Annual Totals (particles/cm$^2$)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>$&gt;10^6$</td>
<td>$&gt;10^7$</td>
</tr>
<tr>
<td></td>
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<td></td>
</tr>
<tr>
<td>1956</td>
<td>1 1 1 1</td>
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</tr>
<tr>
<td>1957</td>
<td>1 4 6 3</td>
<td></td>
</tr>
<tr>
<td>1958</td>
<td>1 4 2 4</td>
<td></td>
</tr>
<tr>
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<td>1 1 5</td>
<td></td>
</tr>
<tr>
<td>1960</td>
<td>3 11 1 1 1</td>
<td></td>
</tr>
<tr>
<td>1961</td>
<td>1 5 3 1</td>
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<td>1963</td>
<td>1 1 1 1</td>
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</tr>
<tr>
<td>1964</td>
<td>1 1</td>
<td></td>
</tr>
<tr>
<td>1965</td>
<td>2 1</td>
<td></td>
</tr>
<tr>
<td>1966</td>
<td>3 2</td>
<td></td>
</tr>
<tr>
<td>1967</td>
<td>3 4 1 1</td>
<td></td>
</tr>
<tr>
<td>1968</td>
<td>1 3 4 1</td>
<td></td>
</tr>
</tbody>
</table>

### Table II. Transit Times (Seconds) for Protons as a Function of Pitch Angle at 1 AU

<table>
<thead>
<tr>
<th>Proton Energy (MeV)</th>
<th>Pitch Angle</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>$\beta = 0^\circ$</td>
</tr>
<tr>
<td>1.0</td>
<td>1.23 x $10^4$</td>
</tr>
<tr>
<td>3.0</td>
<td>7.1 x $10^3$</td>
</tr>
<tr>
<td>10</td>
<td>3.9 x $10^3$</td>
</tr>
<tr>
<td>30</td>
<td>2.24 x $10^3$</td>
</tr>
<tr>
<td>100</td>
<td>1.23 x $10^3$</td>
</tr>
<tr>
<td>300</td>
<td>8.18 x $10^2$</td>
</tr>
</tbody>
</table>
Table III. Minimum Transit Times for Protons as a Function of Energy ($\beta = 0$)

<table>
<thead>
<tr>
<th>Distance From Sun (AU)</th>
<th>1</th>
<th>3</th>
<th>10</th>
<th>30</th>
<th>100</th>
<th>300</th>
</tr>
</thead>
<tbody>
<tr>
<td>0.1</td>
<td>1.07 x 10³</td>
<td>6.2 x 10²</td>
<td>3.4 x 10²</td>
<td>1.95 x 10²</td>
<td>1.07 x 10²</td>
<td>7.1 x 10¹</td>
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<tr>
<td>0.4</td>
<td>4.3 x 10³</td>
<td>2.48 x 10³</td>
<td>1.36 x 10³</td>
<td>7.8 x 10²</td>
<td>4.28 x 10³</td>
<td>2.84 x 10²</td>
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<tr>
<td>0.7</td>
<td>8.0 x 10³</td>
<td>4.63 x 10³</td>
<td>2.55 x 10³</td>
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<td>8.0 x 10³</td>
<td>5.33 x 10²</td>
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<tr>
<td>1.0</td>
<td>1.23 x 10⁴</td>
<td>7.1 x 10³</td>
<td>3.9 x 10³</td>
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<td>8.18 x 10²</td>
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<tr>
<td>1.5</td>
<td>2.09 x 10⁴</td>
<td>1.21 x 10⁴</td>
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<td>2.0</td>
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<td>3.0</td>
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<td>4.0</td>
<td>9.95 x 10⁴</td>
<td>5.75 x 10⁴</td>
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<td>9.9 x 10³</td>
<td>6.6 x 10³</td>
</tr>
<tr>
<td>5.0</td>
<td>1.5 x 10⁵</td>
<td>8.65 x 10⁴</td>
<td>4.75 x 10⁴</td>
<td>2.73 x 10⁴</td>
<td>1.5 x 10⁴</td>
<td>9.95 x 10³</td>
</tr>
</tbody>
</table>

Table IV. Transit Times (Seconds) for 1 Mev Proton as a Function of Pitch Angle

<table>
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<tr>
<th>Distance From Sun (AU)</th>
<th>$\beta = 0°$</th>
<th>$\beta = 15°$</th>
<th>$\beta = 30°$</th>
<th>$\beta = 45°$</th>
<th>$\beta = 60°$</th>
<th>$\beta = 75°$</th>
<th>$\beta = 80°$</th>
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<td>1.52 x 10³</td>
<td>2.14 x 10³</td>
<td>4.13 x 10³</td>
<td>6.15 x 10³</td>
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<td>2.48 x 10⁴</td>
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<td>2.46 x 10⁴</td>
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<td>7.08 x 10⁴</td>
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<td>2.42 x 10³</td>
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<td>8.07 x 10⁴</td>
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<td>8.62 x 10⁵</td>
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</tbody>
</table>
Figure 1. Characteristics of the Quiet-Sun Magnetic Field Carried Frozen in by the Solar Wind.

Figure 2. Proton Gyro-Radius as a Function of Distance From the Sun for Undisturbed Interplanetary Magnetic Field.

Figure 3. Spiral Path Length of Solar Magnetic Field as a Function of Distance From the Sun.

Figure 4. Transit Times for Protons at 1 AU as a Function of Pitch Angle (Assumed Constant) and Proton Energy.
Figure 5. Minimum Proton Transit Time ($B = 0$) as a Function of Proton Energy and Distance From the Sun.

Figure 6. Proton Transit Time for 1 Mev Protons as a Function of Distance From the Sun and Pitch Angle.

Figure 7. Ratio of Integral Particle Fluxes ($\phi$) and Peak Particle Flux Rate ($\dot{\phi}$) to the Corresponding Values at 1 AU