FINAL REPORT

For

ANALYSIS OF AURORAL PARTICLE FLUXES

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INTRODUCTION

Three Javelin rocket flights from Fort Churchill, Canada have provided detailed data on electron and proton precipitation in different types of auroral forms. These rockets named TWINS 1, TWINS 2, and p/α were launched in 1967 and 1968. A thorough understanding of auroral phenomena requires a high time resolution study of electron and proton energy spectra and pitch angle distributions with the ability to separate spatial and temporal effects. The TWINS 1 and 2 data furnish this necessary information and in addition study such things as the characteristics of the backscattered electrons and the transit times of the different energy auroral particles. These transit times give a measurement of the auroral source distance. Analysis of these data have given a greater insight into the details of the physics of the auroral phenomenon.
ACCOMPLISHMENTS

The analysis of the TWINS 1 and 2 rocket data began with the writing of a program to decommutate the SPECS data and to sort the data according to energy and pitch angle. Also the TWINS trajectory information given by the radar tracking measurements was combined with the McIlwain BIL program to calculate the values of the magnetic field B and the L-parameter along the flight path.

The next step in the data analysis was the calculation of the instrument geometric factors. This was done through a combination of the original calibration data from Rice University with Ni\(^{63}\) spectral measurements made at Lockheed during October of 1971. The measurements made at Lockheed compared the Ni\(^{63}\) sources used to calibrate the instruments at Rice with a Lockheed Ni\(^{63}\) source having a known spectrum. Once the geometric factors were derived from the calibration data, they were then used to convert the instrument count rates to absolute fluxes.

For purposes of data analysis, we desired the capability of viewing electron spectra at any selected time in the flight, and of viewing the pitch angle variations of the electron spectra. To that end, the rather limited decommutation program written initially was expanded to a general purpose computer program which generated electron spectra plots and pitch-angle distribution plots at all electron energies. The investment of time required to develop this capability proved to be well justified, for we
found it necessary to examine large quantities of the flight data in rather fine detail just to clarify and solidify our impressions of the general behavior of the electron precipitation in a complex auroral situation.

With the data reduced to a workable format, we were able to conduct an extensive analysis including the study of the variation of the incident electron energy spectra and pitch angle distributions throughout the flight, the study of the backscattered electron spectra and pitch angle distributions, and a detailed comparison of the measured and calculated backscattered spectra with regard to the physical processes involved in the interaction of auroral electrons with the atmosphere. This work is discussed in more detail in the enclosed reprints.

Work on the TWINS data to date has resulted in two J.G.R. articles, one which has already been published entitled "Measurements of Highly Collimated Short-Duration Bursts of Auroral Electrons and Comparison with Existing Auroral Models" by O'Brien and Reasoner, and a second which has been submitted to J.G.R. entitled "TWIN Payload Observations of Incident and Backscattered Auroral Electrons." Also, a talk entitled "TWIN Payload Observations of Precipitated and Backscattered Auroral Electrons" was presented at the 1972 spring meeting of the American Geophysical Union by Chappell, Reasoner, and O'Brien.
SCIENTIFIC RESULTS

The TWINS data have revealed several surprising and interesting features of the auroral processes. These features include the measurement of highly collimated bursts of auroral electrons, very high $\sim 100\%$ backscatter ratio of electrons below 1 keV in auroral arcs, and the presence of the "continuum" background spectrum in the auroral region which has peaks in energy superimposed upon it. The peak is seen to shift in energy as the payloads move into and out of the arc regions. The TWIN payload capability of the TWINS was also able to show the motion of the auroral particle precipitation region which corresponds to the motion of the arcs themselves. It is thought that the TWIN 2 observations are probably the first established observation of the same arc at two distinct times during one rocket flight. The detailed results and conclusions of the scientific analysis during the first year of the contract are presented in the two enclosed preprints of the J.G.R. articles.

PROGRAM FOR NEXT YEAR

The analysis of the TWINS data will be continued under a follow-on contract to the present NASW-2212. The future work will carry on the detailed analysis of both the TWINS data and the other pertinent satellite data which bears on the overall problem of auroral and magnetospheric dynamics.
NEW TECHNOLOGY

The instruments discussed here were developed under previous NASA grants. Consequently there has been no new technology developed under this contract.

CONCLUSIONS

This work has produced more data on the complexity of the auroral phenomenon. The details of this complexity are discussed in the enclosed publications. The complicated nature of the auroral processes is evident both in the high space-time variability and complexity of the incident auroral particle spectra and in the unexpected variation of the backscattered particle spectra with energy. Further studies of these data during the coming year will give even more information on these processes.
TWIN PAYLOAD OBSERVATIONS OF
INCIDENT AND BACKSCATTERED AURORAL ELECTRONS

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Energy spectra and pitch angle distributions of auroral electrons have been measured in a pre-midnight multiple arc auroral display by a Javelin rocket containing two identical payloads which separated in flight. The rocket (code-named TWINS II) was launched from Fort Churchill, Canada at 0459 GMT on March 2, 1968, and covered an altitude range up to 800 km. The electron energy spectra between 40 eV and 20 keV show a "continuum" spectrum with a superimposed energetic peak. The center energy of the peak was observed to shift from 10 - 12 keV over the arcs to 2 - 3 keV between the arcs. This spectral structure is shown to be similar to the inverted "V" structure reported by other investigators. The inflight separation of the two payloads allowed investigation of spatial vs. temporal effects in the auroral precipitation. In one interval of the flight an arc was observed to be moving northward with a velocity of 0.6 km./sec. Calculated backscattered spectra are compared with those actually measured. Good agreement was observed for electron energies above about 1.5 keV, but below this energy the backscatter ratio was observed to be \( \sim 100\% \). Several explanations for this unusually high ratio are considered, and a likely possibility is shown to be a parallel, downward-directed electric field between \( \sim 150 \) and 500 km.
INTRODUCTION

Measurements of auroral particle fluxes, their spectra and pitch-angle distributions have been the object of numerous satellite and rocket investigations. An excellent review of pertinent measurements has been given by Paulikas (1971).

The general features of auroral particle precipitation are now well-known, although the details of the origins of the particles and of their interaction with the atmosphere are far from understood. Several investigators have reported "anomalous" results in measurements of auroral particles. For example, Mozer and Bruston (1966a,b) reported energetic protons traveling upward from mirror points well within the atmosphere and anti-correlations between electron and proton fluxes. A recent study by Reme and Bosqued (1971) reported observations of anti-correlated proton and electron fluxes and proton pitch angle distributions peaked along field lines. These observations were interpreted as indicating the presence of electric fields parallel to magnetic field lines in the upper ionosphere with magnitudes ranging up to a few hundred millivolts/meter. Furthermore, Moser and Bruston (1967) reported direct experimental evidence by means of rocket-borne electric field probes of such parallel electric fields in the upper ionosphere.

Choy and Arnoldy, (1971) have observed an anomalously high reflection coefficient, or backscatter ratio for auroral electrons. With energies
below 1 keV. These results are indeed surprising in view of the relatively low altitude of the measurements (< 150 km., R. L. Arnoldy, private communication).

Chappell (1968) has compared, calculated and measured auroral electron backscatter ratios using data from the TWINS-I twin-payload sounding rocket (Westerlund, 1969.) The theoretical backscatter fluxes were calculated by means of a Fokker-Planck treatment of atmospheric Coulomb scattering. Discrepancies between calculated and measured backscatter were suggestive of atmospheric heating, a parallel electric field below the rocket, or a combination of both effects.

In this paper we report observations of precipitating (30° < a < 60°) and backscattered (120° < a < 150°) auroral electron fluxes in the energy range 40 ev - 20 kev over a series of IBC-2 auroral arcs. The measurements were made with the TWINS II payloads, a pair of identical payloads which were launched on an Argo D-4 Javelin sounding rocket and were separated in flight. Five distinct regions of particle precipitation are identified in terms of total energy input, and differing particle spectra characteristics are shown for each region. Variations in spectra near the edge of an arc are shown to resemble the "inverted V" spectral structure (Frank and Ackerson, 1971). Finally, measured and calculated backscatter spectra are compared and an attempt is made to explain the observed discrepancies.
DESCRIPTION OF THE EXPERIMENT

A detailed description of the TWINS-II sounding rocket experiment and the auroral conditions at launch may be found in O'Brien and Reasoner (1971). Briefly, the payloads were launched from Fort Churchill, Canada on March 2, 1968 at 0459 U.T. into an IBC-II northward-moving multiple arc structure. The ground magnetometers recorded a 100 y negative bay at the time of the flight.

The rocket carried two identical payloads, TWINS 2A and TWINS 2B, which were separated in flight with a relative velocity along the flight direction at separation of approximately 4 meters/second. The relative and absolute velocities perpendicular to the magnetic field lines were 1 ± 0.2 meters/second and 585 ± 6 meters/second respectively. The flight azimuth was 37° east of true north. Auroral electron spectra were measured over an altitude range of 200 - 800 kilometers, and the payloads covered L - values from L = 9 to L = 13.

Each payload carried identical detector complements, consisting of two wide-range electron spectrometers, three fixed-energy differential electron detectors, and a large-area geiger tube. Of relevance to this study were the two electron spectrometers, code-named SPECS (O'Brien, et al, 1967). The SPECS detectors were electrostatic deflection spectrometers employing channel electron multipliers and were sensitive to electrons with 40 ev < E < 20 kev. The upward-looking SPECS scanned a pitch angle range
of $30^\circ$ to $60^\circ$ in $\sim 5^\circ$ increments and the downward-looking SPECS simultaneously scanned a range of $120^\circ$ to $150^\circ$. The SPECS required 3.6 seconds to complete a measurement cycle of all electron energies. Payload aspect information was provided by a fluxgate magnetometer.

EXPERIMENTAL RESULTS

An overview of the auroral conditions encountered during the TWINS II flight is best shown by a plot of the total energy deposited, obtained by integrating over the measured spectra. Figures 1 and 2 show the values of the directional electron energy deposited and backscattered as calculated from the measurements of TWINS 2A and 2B respectively for electron energies $>400$ ev and pitch angles of $30^\circ < \alpha < 60^\circ$ and $120^\circ < \alpha < 150^\circ$. The payloads encountered two distinct periods of enhanced precipitation at 240-350 seconds and again at 670-710 seconds. We can associate these periods of enhanced precipitation with the visible auroral arcs. The precipitated fluxes were isotropic to within a factor of 2-3 over the pitch angle range $30^\circ < \alpha < 60^\circ$, and we assume that the isotropy extended over the entire upper hemisphere. With this assumption the total energy deposited was therefore $\sim 20$ ergs/cm$^2$-sec. This agrees well with the photometrically measured brightness of the auroral forms at launch of 10-20 kilorayleighs.

The gross features of Figures 1 and 2 are in good agreement, as would be expected since the maximum payload separation was 700 meters. The
backscattered flux measured by TWINS 2A is apparently less than that measured by TWINS 2B, but this was due to a failure of two of the six channels of the downward-looking SPECS on TWINS 2A. Therefore a portion of the electron spectrum was not measured by this detector.

A closer examination of Figures 1 and 2 reveals that the decrease in precipitation at T = 350 seconds was observed 6 ± 3 seconds earlier by TWINS 2B (the following payload) than by TWINS 2A (the leading payload). This means that the arc observed from 250 - 350 seconds was moving northward more rapidly than the payloads and effectively overtook and ran ahead of them. Considering the payload absolute horizontal velocity of .585 km/sec and the separation at this time of 217 meters, the arc is calculated to be moving northward with a velocity along the payload flight direction (37° east of north) of .61 km/sec. Flux enhancements later in the flight (e.g. T = 630 seconds) do not display any time delays between detection by the two payloads, although the SPECS detectors with their 3.6 second cycle time would not have been able to resolve the time delay of ∼ 2 seconds expected if the payloads encountered a stationary or slowly-moving spatial structure. It is therefore possible that the payloads initially observed an arc moving northward which subsequently slowed and was again encountered by the northward-moving payloads. The unique capability of the TWINS payloads to distinguish temporal and spatial effects demonstrated above was also instrumental in the discovery of short-duration field-aligned electron bursts reported previously by O'Brien and Reasoner (1971).
The next series of figures (3 - 8) shows precipitated and backscattered electron spectra measured at various times in the flight by TWINS 2B. These spectra represent regions of differing precipitation intensity at T = 281, 414, 479, 630, and 670 seconds. These times are indicated by the arrows on Figure 2. Each spectrum represents an average over 3 SPECS cycles, or about 10 seconds and the periods were selected to insure that the data were stationary over the averaging period. Each point is an average of 147 samples of 10 milliseconds each.

Figure 3 shows the incident, or precipitated, spectrum at T = 281 seconds. The fluxes in units of electrons/cm$^2$-sec-stev-eV are plotted as a function of electron energy in eV. This spectrum was measured at an altitude of 585 kilometers and represents pitch angles of $32^\circ \rightarrow 56^\circ$. The loss cone (referenced to the 100 km level) at 585 km is $64^\circ$ and hence all of the observed incident electrons would theoretically mirror below 100 kilometers and thus be subject to severe atmospheric scattering and loss. All down-going electrons measured throughout the flight were in the loss cone; hence no information on the trapped particle population was acquired. The group of data points at each energy represent the spread in flux at a given energy for different pitch angles between $32^\circ$ and $56^\circ$ and indicate the variation of energy spectrum with pitch angle. This variation is a factor of 2 to 3 maximum. The error bars for the points are also indicated. The errors are dominated by uncertainties in the instrument calibration of $\pm 20\%$. Errors due to counting statistics are negligible, typically $< 5\%$. 
The precipitating electrons shown in Figure 3 have a peak in the energy range of 6 - 17 keV which contains the bulk of the total precipitating energy.

Figure 4 shows the electron spectrum which was backscattered from the atmosphere below the payload as measured simultaneously at T = 281 seconds by the downward-viewing SPECS. The data format is the same as for the previous figure. The pitch angle range measured was 124° to 148°. Note that the backscattered spectrum exhibits a small peak in the range 5 - 15 keV which is almost an order of magnitude lower in intensity than the precipitated flux.

As the payloads passed from the region of intense precipitation into one of intermediate precipitation (T = 414 seconds), the spectra changed to those shown in Figure 5. Here we have shown only the sketched incident and backscattered spectra without including the data points showing the spread in pitch angle. This is done here and in subsequent presentations for the sake of brevity and to afford easier comparison of spectra.

The incident spectrum at T = 414 (solid line) shows two peaks, a prominent peak in the range 2 - 9 keV and a less prominent one in the range 500 eV - 1.2 keV. Note that as the payloads passed out of the region of enhanced precipitation (Figure 3) the peak shifted down from 6 - 17 keV to 2 - 9 keV. The backscattered spectrum (dashed line) shows no evidence of a peak. Note also that the flux of backscattered electrons with E < 600 eV is comparable to or greater than the corresponding flux of
precipitating electrons. We will comment further on this surprising backscatter ratio in a later section.

Figure 6 shows the incident and backscattered spectra in the region of low intensity precipitation at $T = 479$ seconds. The peak in the incident spectrum has shifted to the energy range $1.5 - 4$ keV and has diminished in differential flux level by a factor of 2 from the previous spectrum. This precipitated spectrum measured between the intense precipitation regions is similar to the background "continuum" discussed by Westerlund (1969) using data from TWINS I, a similar rocket payload fired one year earlier. Here, the backscattered spectrum shows no peak in energy, unlike the backscattered spectrum in the intense precipitation region ($T = 281$ seconds, Figure 4). Again, the backscatter ratio is $\sim 100\%$ for energies less than 500 eV.

Following the low intensity region, the rocket encountered another region of intermediate intensity precipitation with the incident electron energy spectrum showing a characteristic upward shift in the energy of the peak to the $3 - 12$ keV range. Figure 7 shows representative spectra from this region measured at $T = 630$ seconds. The backscattered spectrum has a slight peak in this same energy range. In this case the backscatter ratio is $\sim 100\%$ for energies less than 1.2 keV.

Finally, the payloads again encountered a region of intense precipitation (Figure 8) at $T = 670$ seconds. The energy of the peak is in the $3 - 16$ keV range, and the incident spectrum is similar to that observed
during the first enhanced period at $T = 281$ seconds. This adds support to the thesis advanced earlier that the two regions of intense precipitation were one and the same arc. The backscatter peak is again prominent, and the backscatter ratio is $\sim 100\%$ for energies below 1 keV.

From this series of spectra it is clear that changes in the total energy deposited were caused by shifts in the energy of an energetic peak which appeared to be superimposed on a much more stable "continuum" spectrum. During this TWINS II flight, unlike the case in the TWINS I flight (Westerlund, 1969) the peak never completely disappeared although it was considerably diminished in both intensity and average energy in the period of $T = 460$ to $T = 500$ seconds.

A higher time resolution study of this peak shift effect is shown in Figure 9. Here three successive incident spectra at $T = 351$, $T = 366$, and $T = 405$ seconds are plotted to display the spectral change over the time period of a sharp decrease in total precipitating intensity (Figure 1). The shift of the energy of the peak to progressively lower values is clearly evident. We will discuss this further in a later section.

Perhaps the most surprising feature of these spectra is the very high, $\sim 100\%$, backscatter ratio for electron energies less than about 1 keV. The expected ratio computed from Coulomb scattering is usually on the order of 10%. In the following section we consider the backscatter phenomena in detail, presenting calculations of expected backscattered fluxes for comparison with those actually observed.
COMPARISON OF EXPERIMENTAL AND THEORETICAL BACKSCATTER RESULTS

To examine the interaction of the precipitated auroral electrons with the atmosphere we have used the measured incident electron spectrum as input for a theoretical computer code which calculates the expected backscattered electron spectrum. This calculated backscattered spectrum can then be directly compared with the backscattered spectrum measured simultaneously by the downward-viewing SPECS detector on the TWINS payload. Any differences between the calculated and measured backscattered spectra can be interpreted in terms of the physical processes occurring in the interaction region.

The theoretical calculations were made using the Fokker-Planck diffusion equation solution of Walt, MacDonald, and Francis (1968). This treatment includes angular scattering, continuous energy loss, and the effect of a converging magnetic field. The solution is in the form of a distribution function for the auroral electrons as a function of energy, pitch angle, and altitude. The calculations were made using differential cross-sections given by the first Born approximation (Mott and Massey, 1965) and the continuous energy loss formula of Bethe (1933). These calculations are reasonably accurate (±10%) for isotropic pitch angle distributions and for electron energies down to about 500 eV. The computations done in this work are carried down to 1 keV.

Figure 10 shows the measured incident (solid line) and backscattered (x's) spectra at T = 281 in the first region of enhanced precipitation.
The incident spectrum which was assumed to be isotropic in pitch angle over the upper hemisphere was used as input to the backscatter code. The blue shaded area represents the calculated backscattered flux in the range of pitch angles of $126^\circ \rightarrow 149^\circ$ corresponding to those pitch angles measured by the downward viewing SPECS. The agreement between the calculated backscattered flux and the measured flux (x's) is reasonably good above about 1.5 keV. However below 1.5 keV the measured backscattered fluxes are roughly equal in magnitude to the incident fluxes (backscatter ratio of $\sim 100\%$). The theoretical calculations, which employ Coulomb scattering only, predict a backscatter ratio more on the order of 10-20%, not 100%.

In other words, the electrons of energy below 1 - 2 keV in this enhanced precipitation region are either kept out of the atmosphere preventing any loss or the incident primary particles are replaced by secondaries produced in atmospheric scattering.

The most obvious explanation for the 100% backscatter below 1 - 2 keV is the presence of a parallel downward directed electric field above the atmosphere with a total potential drop of about 1 keV in magnitude. This field would prevent electrons of less than 1 keV from entering the atmosphere and would reflect them undiminished in flux value. Other potential explanations also exist for this unusually high backscatter ratio. The production of multiple secondary electrons in the atmosphere by the energetic precipitating primaries is a possibility. Also, a severe modulation of the magnetic field by the auroral electrojet currents could possibly affect the particle pitch angles and cause reflection. In anticipation of
the results of a later discussion on the relative merits of these possibilities, we will examine theoretically the effect of a parallel downward directed electric field of potential drop of 1 keV situated below the rocket and above the effective scattering layer of the neutral atmosphere. In this T = 281 case, this field would be between 585 km and about 150 km.

For this T = 281 case we have again used the measured incident spectrum shown in Figure 10 but have included a vertical potential drop of 1 keV just below the rocket. The effect of this parallel field is to alter the pitch angle distribution of the incident particles in such a way as to reflect particles whose energy component parallel to the field is less than 1 keV and to reduce the energy of particles which traverse the field region by 1 keV. The Coulomb scattering, energy loss, and mirroring of these penetrating electrons is then calculated by the computer code and the resulting backscattered spectrum is then transformed back through the potential drop with the corresponding change in pitch angle distribution and increase in energy. The results of this calculation are shown by the red shaded area in Figure 10 for energies down to 1 keV. Note that the effect of the field is to increase the flux backscattered at the more nearly perpendicular pitch angles for electrons between 1 and ~10 keV in energy and also to cause the 100% backscatter below energies of a few keV. Qualitatively then, the calculations show that with the addition of a parallel downward directed electric field above the atmosphere to the general effects of coulomb scattering, the predicted backscattered flux is in good agreement with the measured backscattered flux for all energies.
Figure 11 shows the measured and calculated backscattered fluxes for the time period of very low intensity precipitation, $T = 479$ seconds. As before the measured incident spectrum is shown by the solid line and is used as the input for the calculation. The shaded area shows the predicted range of backscattered flux for the pitch angles measured by the downward viewing SPECS. The measured backscattered electrons are again shown by the "x's." In this calculation we have not included a parallel electric field. One can see that the agreement above 1 keV for coulomb scattering is quite good. Recall that this set of spectra was measured between the auroral arc regions. Although no parallel electric field is required to obtain agreement between calculated and observed backscattered spectra above 1 keV in this low intensity precipitation region, one can see that below about 500 eV the backscattered flux again approximately equals the incident flux indicating the possible need for a lower magnitude parallel field of about 500 eV to give agreement.

Figure 12 shows the set of spectra representing the second payload encounter with the intense precipitation region at $T = 670$ seconds. As above the shaded area represents the predicted backscattered fluxes with no parallel electric field and the "x's" represent the measured backscattered fluxes. As in the first case ($T = 281$) the agreement at higher energies (greater than about 1.5 keV) is good. However as in the previous encounter with the intense precipitation region, the backscattered fluxes below about 1.5 keV are comparable to the incident fluxes in contrast to the predictions
of the Coulomb scattering calculations which show the backscattered flux decreasing with decreasing energy below 1.5 keV. As before one can see that the addition of a parallel electric field of about 1 keV total potential drop below the rocket would cause an increase in the predicted backscattered fluxes at lower energies (less than about 6 or 8 keV) with a backscattered flux equal to the incident flux below about 1.5 keV in agreement with the measured fluxes.

We have found then that in regions of auroral precipitation the measured backscattered electron fluxes agree quite well with the backscattered fluxes predicted from Coulomb scattering alone for the higher energies greater than about 1.5 keV. However, for the lower energies, agreement between predicted and measured backscattered fluxes for Coulomb scattering alone is not sufficient but requires an effect such as a parallel downward directed electric field between the rocket and the atmosphere with a total potential drop of about 1 keV in regions of intense precipitation and about 500 eV in regions of less intense precipitation.

SUMMARY AND DISCUSSION

The features of the measured precipitated and backscattered auroral electron spectra can be summarized as follows:

(1) The spectra contain two main components—a "continuum" spectrum which remains fairly stable throughout the flight and a peak in
the 6 - 17 keV range which is present over the auroral arcs but which shifts to lower energy and intensity as the payloads move away from the arc. This peak, when present, carries the bulk of the total energy deposited.

(2) The incident electron fluxes at all energies appear to be isotropic in pitch angle to within a factor of 2 or 3 over the observed pitch angle range of $30^\circ \rightarrow 60^\circ$.

(3) The backscattered fluxes for energies below about 1 keV appear to be isotropic within a factor of 2 or 3 and equal in intensity to the precipitated fluxes. Above about 1 keV the fluxes show a larger variation with pitch angle as expected from the effects of atmospheric scattering and the magnitude of the backscatter ratio ($j(E)_{\text{up}}/j(E)_{\text{down}}$) at a given energy decreases with increasing energy.

(4) A theoretical backscatter calculation employing Coulomb scattering, continuous energy loss and magnetic mirroring and using the measured incident spectrum as input is able to reproduce the measured backscattered spectrum for energies greater than 1 keV in intense regions of precipitation and for energies of greater than 500 eV in regions of less intense precipitation. The measured backscattered spectra below about 1 keV, however, cannot be explained on the basis of Coulomb scattering alone.
In (1) above we mentioned the shift in energy of the peak in the incident electron spectrum as the payloads move in and out of the auroral arc region. If this shift in peak energy were displayed on an energy-time spectrogram it would resemble an "inverted V" structure similar to that discussed by Frank and Ackerson (1971). The widths of the structures observed here (50-100 km) are in reasonable agreement with the "typical" widths of 200 km discussed by Frank and Ackerson (1971). Exact agreement on these widths is not expected due to uncertainties resulting from high velocities of the auroral forms relative to the rocket payload velocities.

The most interesting feature of the electron spectra is the 100% backscatter ratio for energies less than 1 keV. As we mentioned in the previous section, there are at least three possible explanations for this large ratio — secondary electrons which are produced by the high energy primary electrons and which are scattered upward out of the atmosphere, a large variation in the ambient magnetic field caused by the nearby auroral electrojet which changes the electron pitch angles, or a downward-directed parallel electric field with a potential drop of 1 keV above the atmosphere which reflects lower energy electrons before they undergo scattering.

We discount the possibility that secondary electrons are responsible for the upward fluxes at energies between 40 - 100 eV. If these particles were secondaries, we would expect that the measured upward flux of these particles would be correlated with the measured incident primary energy flux of the particles greater than 400 eV. In Figure 13 we show the incident energy...
flux from Figure 1 along with the 40 eV precipitated and backscattered number flux. There is a good correlation between the 40 eV downgoing and upcoming flux, but there is no significant correlation between the precipitated primary energy flux and the 40 eV upcoming or backscattered flux. Thus, the experimental evidence argues against secondary electrons as a source of these low energy electrons.

Theoretical evidence also argues against secondaries as a source of the high intensity low energy upflux. Most of the secondaries that are produced are below 300 km and the photoelectron calculations of Nagy and Banks (1970) show that these electrons cannot escape the atmosphere because of the very high degree of scattering. However, these same calculations show that a 50 eV electron produced at an altitude greater than 300 km can escape significant scattering and travel freely upward. Cross-sections for secondary electron production by energetic primary electrons incident upon N₂ have been measured by Opal, et al, (1971). Taking the integrated N₂ density above 300 km for a medium density atmosphere as $2 \times 10^{15}$ atoms/cm² (Johnson, 1965), and the precipitated energy as 3 ergs/cm² - sec-str with an average energy of 10 keV we get an expected upward secondary flux at 40 eV from above 300 km of $2 \times 10^5$ electrons/cm² sec str eV. This is just about equal to the change seen in the 40 eV upward flux in the time period of $T = 340-440$ seconds as the primary energy flux drops from about 3 ergs/cm² - sec-str to .4 ergs/cm² sec str. Therefore, although secondary electron production above 300 km appears to contribute
to the upflux of low energy electrons, its contribution is only a minor part of the total observed upflux (.2 x 10^5 electrons/cm^2 sec str eV out of about 1.5 x 10^5 electrons/cm^2 sec str eV).

As in the case of the secondary electrons, we can also rule out the possibility of magnetic field perturbations as a cause of the large backscatter ratios. First invariant calculations show that an increase of a factor of more than three in the ambient magnetic field is required to cause a particle with a pitch angle of 30° at 585 km to mirror above the neutral atmosphere. Magnetic field changes of this magnitude are not observed in the vicinity of aurorae.

After discounting secondary electrons and magnetic field perturbations, the obvious explanation remaining is the presence of a parallel downward-directed electric field below the rocket of potential drop about 1 keV. A field of this sort would cause the incident auroral electrons with energies less than ~1 - 2 keV to mirror above the region of significant scattering in the atmosphere. The backscatter calculations of the previous section have shown that the addition of such a field would give good agreement between calculated and observed backscattered electrons.

There are, however, certain crucial limitations upon this electric field which must be met in this case. The first is that a total potential drop of ~1 keV must exist between the rocket altitude and the top of the scattering region (about 150 km for 1 keV electrons). The second is that
no large fraction of this potential drop can exist above 500 km, for there was not a systematic, large variation with altitude of the lowest energy flux in the altitude range 500 - 800 km. Thus for the purpose of discussion we postulate a parallel electric field with a total potential drop of about 1 keV located between 150 and 500 km altitude.

Recent experimental and theoretical investigations have provided a possible source for the postulated electric field. Zmuda et al., (1966), Cloutier et al (1970) and Choy et al (1971) have measured the presence of field-aligned currents composed of very low energy electrons flowing out of the auroral ionosphere. Kindel and Kennel (1971) have theoretically investigated the effects of these currents and have concluded that ion-acoustic and ion-cyclotron instabilities, which are driven by the field-aligned currents, can occur under the conditions found in the auroral ionosphere. The role of the instabilities is to furnish an impedance to the flow of field-aligned currents, and hence to cause a potential drop to develop.

There is, however, a certain difficulty with the application of this mechanism to our particular case. For the range of field-aligned current values expected just from flux balance in the aurora studied here \(10^8-10^9\) electrons/cm\(^2\)-sec, an examination of Figure 7 of Kindel and Kennel (1971) shows that the first instability expected would be the \(O^+\) or \(H^+\) ion-cyclotron mode, and then only at altitudes greater than 1000 km. We require that the electric field, and hence the instability region, exist below 500 km. A
possible explanation for this apparent inconsistency is as follows. The
critical current for instability is not only a function of altitude but
also of the electron to ion temperature ratio. As the electrons are heated
relative to the ions, ion Landau and cyclotron damping become less important
in stabilizing the waves. An examination of Figure 2 of Kindel and Kennel
shows that if the ratio of electron drift velocity to electron thermal
velocity is initially above a certain value (\(V_d/V_{th} \sim 0.03\) for \(T_e/T_i = 1\)),
then increasing the \(T_e/T_i\) ratio will cause waves to grow. Therefore, if
the electrons in the 150 - 500 km region could be heated sufficiently, then
the instabilities and hence the electric fields would exist and be confined
to this region. Possible mechanisms for accomplishing this heating include
an energy transfer from secondary electrons by means of electrostatic waves
(the "bump on tail" instability) or an upward transport of hot electrons
from below.

The above argument is necessarily qualitative, for data on crucial
parameters such as temperature and drift velocity of the electrons contrib-
buting to the field-aligned currents is lacking. It is offered here as a
plausible explanation for the required altitude structure of the postulated
parallel electric field. Whatever the precise cause, the existence of the
parallel downward-directed electric field appears necessary to adequately
explain the data shown above.
CONCLUSIONS

The physical processes which describe the interaction of auroral electrons with the atmosphere appear to be more complex than just the Coulomb scattering of the incident primary electrons with a subsequent loss of energy. The comparison of the measured backscattered electron spectra with spectra predicted using a theoretical scattering calculation has led to a discrepancy for energies below about 1 - 2 keV. It was found that the very high ratio (~ 100%) of backscattered to incident fluxes for these energies could be most reasonably explained by a parallel downward-directed electric field which prevents these lower energy electrons from entering the atmospheric scattering region. This parallel field with potential drop of about 1 keV is thought to have its origin in wave-particle interactions in the turbulent auroral ionosphere.
ACKNOWLEDGEMENTS

The authors are grateful to Dr. L. Westerlund, Mr. J. Sneddon, Mr. F. Abney, and Mr. R. Harrison for their work on the TWINS payload. We are also grateful to Dr. Martin Walt for helpful discussions during the course of the data analysis. This work was funded under contract NASW 2212 from the National Aeronautics and Space Administration.
REFERENCES


FIGURE CAPTIONS

FIGURE 1: Total directional incident and backscattered electron energy flux measured by TWINS 2A. The measurements were computed by integrating the measured electron spectra, appropriately weighted, over the energy range 400 eV to 20 keV.

FIGURE 2: Total directional incident and backscattered electron energy flux measured by TWINS 2B. The measurements were computed by integrating the measured electron spectra, appropriately weighted, over the energy range 400 eV to 20 keV.

FIGURE 3: Incident electron energy spectrum measured at T = 281 seconds. For reference with regard to the time of measurement for this figure and Figures 4 - 9 following, see Figure 2.

FIGURE 4: Backscattered electron energy spectrum measured at T = 281 seconds.

FIGURE 5: Incident and backscattered electron energy spectra measured at T = 414 seconds.

FIGURE 6: Incident and backscattered electron energy spectra measured at T = 479 seconds.

FIGURE 7: Incident and backscattered electron energy spectra measured at T = 630 seconds.

FIGURE 8: Incident and backscattered electron energy spectra measured at T = 670 seconds.
A series of three incident electron energy spectra measured in the period $T = 351$ to $T = 405$ seconds showing the shift in the energy of the peak as the payload exits the region over an auroral arc.

Comparison of measured and calculated backscattered electron spectra at $T = 281$ seconds, a region of intense precipitation.

Comparison of measured and calculated backscattered electron spectra at $T = 479$ seconds, the time of the minimum precipitation.

Comparison of measured and calculated backscattered electron spectra at $T = 670$ seconds, a period of intense precipitation.

A comparison of the total incident energy flux (from Figure 1) with the measured 40 eV incident and backscattered flux.
ALTITUDE (KM) 423 744 769 480
L 9.5 10.7 11.7 12.5

TWINS 2B

PRECIPITATED ENERGY FLUX
$30^\circ < \alpha < 60^\circ$

BACKSCATTERED ENERGY FLUX
$120^\circ < \alpha < 150^\circ$

erags/cm$^2$-sec-str

TIME AFTER LIFTOFF (SEC)

0 200 400 600 800 1000
T = 414 SEC
ALT = 755 KM

OBSERVED INCIDENT
PITCH ANGLES 31°-53°

OBSERVED BACKSCATTERED
PITCH ANGLES 128°-150°

FLUX (el/cm²-sec-str-eV)

ENERGY (eV)
ENERGY (eV)

10^6

10^5

10^4

10^3

10^2

10^1

FLUX (el/cm^2 - sec-str-eV)

T = 479 SEC
ALT = 785 KM

OBSERVED INCIDENT
PITCH ANGLES 29°-52°

OBSERVED BACKSCATTERED
PITCH ANGLES 128°-151°
--- OBSERVED INCIDENT PITCH ANGLES 31°-53°

--- OBSERVED BACKSCATTERED PITCH ANGLES 127°-150°

\( T = 630 \text{ SEC} \)

\( \text{ALT} = 745 \text{ KM} \)
INCIDENT - OBSERVED P.A. 32° - 54°
- BACKSCATTERED - OBSERVED P.A. 126° - 148°
- BACKSCATTERED - CALCULATED P.A. 126° - 149°, E_\parallel = 0

T = 670 sec
ALT = 705 Km
Measurements of Highly Collimated Short-Duration Burst of Auroral Electrons and Comparison with Existing Auroral Models

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The pitch-angle distributions, energy spectra, and temporal and spatial variations of auroral electrons with energies from 40 ev to 100 kev were measured with the Twins 2 rocket payloads at an altitude of 800 km. Enhancements of electrons with energies of ~5 kev, pitch angle $\alpha < 10^\circ$, and intensities of up to $5 \times 10^4$ particles cm$^{-2}$ sec$^{-1}$ ster$^{-1}$ were observed occasionally. The enhancements were sometimes by a factor of 10, even while electrons with energies of ~40 ev and 100 kev at $30^\circ \leq \alpha \leq 60^\circ$. The bursts were thought to be temporal in origin, and they lasted for the minimum resolvable time of ~0.1 sec up to times of several seconds. The altitude of the source of the bursts of field-aligned electrons is believed to be less than 1 Rs, and their occurrence is not clearly related either to the ever-present auroral continuum nor to the electrons with a few kiloelectron volts that cause auroral luminosity.

We examine several existing theories of auroral particle acceleration and/or precipitation and show them each to be inadequate (as presently formulated) to individually explain all the observations. A potpourri model of adaptations of all these theories is shown to be adequate, but its reality is not otherwise established. Field-line merging and diffusion by wave-particle interactions are considered to be commonly occurring phenomena.

Direct measurements of fluxes of auroral electrons with peak intensity at small pitch angles [Hoffman and Evans, 1968] marked an important experimental breakthrough in auroral physics. Although such a phenomenon had been predicted [see O'Brien, 1967, Figure 6 (12); Speiser, 1965], no experimental evidence supported it for higher-energy electrons (>40 kev) measured at low altitudes [O'Brien, 1964]; indeed, the phenomenon thought had been unlikely if the electrons came from a source high above the atmosphere [Chamberlain, 1961]. Accordingly, with reports of their findings, Hoffman and Evans [1968] formulated a model for the origin of the auroral electrons. To do so, they had to adopt the 'very tenuous assumption' that the variations they observed were of temporal rather than of spatial origin. By contrast, Eather and Akasofu [1969] argue that the effects were spatial and not temporal and hence that the theory was not proved.

We briefly report here preliminary results of a rocket experiment explicitly planned to distinguish decisively between spatial and temporal effects in auroral-electron fluxes. We find that, up to altitudes of ~800 km and separation distances of up to ~1 km, enhanced fluxes of ~5-kev electrons with small pitch angles were observed simultaneously, on occasion for a time as short as <0.1 sec. We report a summary of the experimental facts and analyze the implications in detail. The effect is believed to be of a temporal origin, rather than of a spatial one. We therefore call the enhancements 'bursts.'

Thus our findings may be regarded as validating the very tenuous assumption of Hoffman and Evans [1968]. However, we argue against their subsequent theoretical conclusions as an explanation for the bursts observed here.

<table>
<thead>
<tr>
<th>Name</th>
<th>SPEC 1</th>
<th>SPEC 2</th>
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<th>Low DED 2</th>
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<th>Magnetometer</th>
</tr>
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<td></td>
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<td>8.2</td>
<td>3.7</td>
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</table>

* Orientation refers
EXPERIMENTAL DETAILS

The effect was found in analysis of data from the Twins 2 payloads, carried by a Javelin rocket into an aurora above Fort Churchill, Canada, at 0459 UT on March 2, 1968. The auroral event consisted of a multiple arc structure of IBC II intensity moving slowly northward, requiring approximately 1 hour to cross from the southern to the northern horizon. The arcs were relatively poorly defined, with the region between the arcs filled with a stable IBC I glow. It should be emphasized that this was not the classic midnight breakup aurora, although there was significant movement of ray structure along the arcs. The Churchill magnetometer showed a small negative bay of 120°, and the riometer absorption was 1 db. The flight azimuth was 67°, and thus the payloads crossed the visible arcs at an angle of approximately 30°.

The top and bottom Twins packages (2A and 2B, respectively) carried very similar detector complements to measure auroral electrons and protons over a wide range of energies and pitch angles. Here we are concerned only with the auroral electrons measured by detectors listed in Table 1. The two payloads were spin stabilized at about 9 Hz and separated by a spring at an altitude of about 250 km. They thus achieved a relative separation velocity of (4 ± 1) m sec⁻¹ along their trajectory. Their relative and absolute velocity orthogonal to the geomagnetic field vector B are the relevant factors here, and their values are (1 ± 0.2)m sec⁻¹ and (500 ± 100)m sec⁻¹, respectively.

The concept behind this approach was our endeavor to have the two payloads pass, one after the other, through the same region of space, thereby giving a unique opportunity to distinguish spatial from temporal auroral variations, a problem generally acknowledged but often ignored.

The first Twins was launched successfully a year prior to Twins 2, and preliminary analysis of data from one of its payloads was reported [Chappell, 1968; Westerlund, 1968, 1969]. Several modifications were made to Twins 2 as a result of study of Twins 1 data, and hence we concentrate on Twins 2.

Analysis of payload orientation and the range of pitch angles studied by the switching proton-electron Channeltron spectrometer (SP1ECS) instruments [O'Brien et al., 1967] was carried out in much the manner described by Westerlund [1968]. On Twins 2 we had the added advantage that the separation-velocity transducer operated. The rocket axis was at an angle of (13° ± 2°) to B, so the upward-viewing SP1ECS scanned the pitch angle (θ) from ~30° to 60°, having an opening angle of 2° in θ.

Each differential energy detector (DED) is structurally identical. Particles are collimated by a series of slits and then pass between a pair of electrostatic deflection plates. Particles

<table>
<thead>
<tr>
<th>Table 1. Electron Detectors on a Twin</th>
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<tr>
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</tr>
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</tr>
<tr>
<td>Low DED 2</td>
</tr>
<tr>
<td>8001 Geiger</td>
</tr>
</tbody>
</table>

* Orientation refers to the rocket spin axis. θ = 0° is 'up.'
of the proper energy and charge sign (for example, electrons with \( E = 5 \pm 1 \text{ keV} \)) are deflected such that they are incident on the aperture of the Bendix Channeltron [Evans, 1965]. The center energy of the instrument's passband could be adjusted by appropriate selection of the deflection-plate voltage. Each Twins payload carried three such sensors, two oriented parallel to the spin axis (\( \theta = 0^\circ \)) and the third oriented antiparallel to the spin axis (\( \theta = 180^\circ \)). (Refer to Table 1 and Appendix 1.)

Nominally, since each DED was mounted parallel to the spin axis, it should have viewed the same pitch angle, namely \( \alpha = 13^\circ \) for the two up-viewing DED's and \( 167^\circ \) for the down-viewing DED. In fact, the field of view of the DED is not cylindrically symmetrical (see Figure 1), and indeed it is also not parallel to the instrument mounting bracket and thus not parallel to the rocket axis. The center of the field of view is only \( 3^\circ \) from the rocket axis, and prior to launch we judged this to be trivial and quite tolerable. In fact, the precision of this serendipity is what led to the discovery reported here.

Counts from all sensors were fed to individual 12-bit accumulators, with the exception that channels 3 and 4 of each SPECS were commutated. Magnetometer outputs were digitized and transmitted as a 6-bit word. Each pulse-code modulation system had a frame length of 216 bits, and a bit rate of about 20,000 bits per sec, and thus the sampling rate of each sensor was about 100 times per sec. The deflection voltage changed on each SPECS each 0.6 sec, giving a complete spectrum every 3.6 sec and the background (plate voltage zero) every 40 sec.

Fig. 1. Angular response of the DED sensor.

**Experimental Results**

On various occasions during the flight, the output of the low DED's (5 keV) on each payload showed a marked modulation at twice the spin period of nominal 9 Hz. The modulation, seen simultaneously with both payloads, would sometimes be in phase (Figure 2) and sometimes not in phase (Figure 3). Owing to various dynamical effects at separation, the final spin rates of the two Twins payloads differed by approximately 0.02 Hz. In all instances, however, the modulation seen on each Twin was quite clearly in phase with the maximums and minimums of the 45° magnetometer on that Twin. This magnetometer lies in the plane of the \( \theta \) or long axis of the DED field of view (Figure 1), and hence, when the magnetometer is at a maximum or minimum, the plane defined by the long angular DED acceptance opening, the \( \theta \) plane, very closely contains the geomagnetic \( B \) vector. At these times, the detector instantaneously samples pitch angles ranging from \( -1^\circ \) to \( +1^\circ \) or from \( +7^\circ \) to \( +25^\circ \). The combination of the above quoted conditions, namely the presence of the modulation, the phasing of the modulation with respect to the magnetometers, the in-phase and out-of-phase nature of the modulations matching with the in-phase and out-of-phase nature of the payload spins, and finally the characteristics and orientation of the DED angular response function all lead to the conclusion that, on the occasion during a spin when the DED count was enhanced, it was receiving electrons from a smaller or larger pitch angle than \( 13^\circ \) where the intensity was greater. The modulation of the DED data at exactly twice the spin rate is therefore due to the fan-shaped angular response of the DED sampling a strongly anisotropic pitch-angle distribution, and the envelopes of the peaks of the modulation show the temporal history of these anisotropy events.

Because the most common pitch-angle distributions of magnetospheric electrons are peaked at \( \alpha \sim 90^\circ \) [see O'Brien, 1964], one might at first think, therefore, that the occasional enhancements of the DED sensors were due to its acceptance of electrons with pitch angles larger than \( \alpha \sim 13^\circ \), say at \( \alpha \sim 25^\circ \) (see Figure 1). In fact, we now show it to be due instead to the acceptance of electrons with pitch angles smaller than \( \alpha \sim 13^\circ \) and, indeed,
with $\alpha \lesssim 10^\circ$. We show this causation simply by demonstrating that the fluxes of electrons of comparable energy at $\alpha \sim 30^\circ$ as viewed by SPECS are too small by factors of about 10 to give the observed enhancements of the DED rates.

From preflight calibrations with a 3.3 meV source, it may be noted that we found (at a deflection voltage of $-3500$ volts, relevant here) that the sum of channels 1 and 2 in SPECS was about 36 per sec, whereas that of the low DED's was about 35 per sec. Thus the relation $(\text{ch.1 + ch.2}) \text{ counts} = \text{low DED counts}$ seems applicable here.

We also examined flight data to show that this relation continued to exist. We chose not only isotropic fluxes (Figure 4), but instead, we plotted a random set of 1-sec sums of low DED against 49-frame (0.5 sec) sums of $(\text{ch.1 + ch.2}) \text{ counts}$ at $-3500$ volts at the same times (Figure 5). It can be seen from Figure 5 that for many of the plotted points (say about)

![Diagram showing electron flux modulation](image)

**Fig. 2.** Short-lived enhancement of the flux of 5-keV electrons observed simultaneously in two payloads about 1 km laterally apart. The modulation at twice the spin rate is due to the fan-shaped response of Figure 1 coupled with a strongly anisotropic flux. (See text.)
Fig. 3. Similar to Figure 2 but for a time when the magnetometers of the two payloads are out of phase. The modulation is shown to be due to the cause mentioned in the legend of Figure 2 rather than to a periodic modulation of an electron flux.

Fig. 4. Essentially the same as Figure 3, with relative geometrical alinement of the two payloads taken relative to the geomagnetic field. The magnetometer output from channel 1 at -3500 volts and channel 2 at -1500 volts were recorded for two complete spins in frame 1 and 2. The counts in channels 1 and 2 have been normalized to those that followed a spin.

half) the preflight examined the greatest number of counts, and we find a strong modulation effect the whole number of points in the range $R \approx (1.0 : 1.4$).

This value of $R$ is, to both Twins 2A and 2B spectra where the ratio of channel 2 to that of channel 1 for a count between $\sim 1:8$ and $\sim 2:1$ over only a 2-kev range of the modulation.

Consider the event...
Essentially isotropic flux used to derive relative geometric factors for SPECS and low DED after launch. The graph shows the magnetometer output and the sum of channels 1 and 2 at -3500 volts deflection. Since there are two complete spins in -23 frames, counts in frame 1 have been added to those in frame 24 and to those that follow to improve statistical accuracy.

half) the preflight relation is valid. We examined the greatest deviations from the unit slope line in the sense of anomalously high DED counts, and we find some, but not all, having strong modulation effects. We attribute the general scatter to this effect and to differences in pitch-angle and spectral distributions. Of the whole number of points, about two-thirds are in the range $R \approx (1.0 \pm 0.4^\circ)$, where

$$R = \frac{\text{sum of channels 1 and 2 rates}}{\text{low DED rates}}$$

This value of $R$ is, for this study, applicable to both Twins 2A and 2B over a wide range of spectra where the ratio of the rate of channel 1 to that of channel 2 varies, for example, between $\sim 1:8$ and $\sim 4:1$, i.e., a 30-fold change over only a 2-kev range. We now examine several of the modulation events in detail.

Consider the event at $T + 624$ sec (Figures 2 and 6 and Table 2). From the combined counts of 1557 counts in channels 1 and 2 in 49 frames, one would predict a low DED rate of 32 counts per frame, or using the above range in $R$, a count of $(32 \pm 8^\circ)$. In fact (Figure 2), the peak rate was 150, and the minimum was around 30. Thus one can indeed use the SPECS data to predict low DED rates when it is viewing larger pitch angles, but the flux at lower pitch angles is abnormally high. The above example for Twins 2A is given by 2B also. From SPECS sampled at the -3500 volts level 1 sec later, one would predict a rate of $(12.7 \pm 4^\circ)$ for the low DED, whereas in fact the maximum was 202, and the minimum was around 30.

Consider again the event at $T + 616$ sec (Figure 8). From the SPECS, one would predict a DED rate of $(27 \pm 9^\circ)$, whereas in fact we found maximums of 242 and minimums around 20.

In every instance of significant spin-induced modulation we examined, we find the same conclusion, namely that electron fluxes at $\alpha \sim 30^\circ$ are inadequate by factors of about 10 to explain the low DED count rates.

Of course the off-axis and fan-shaped angular response of the DED (Figure 1), which accounts for the modulation at twice the spin rate, means
that at one instant the detector is sampling pitch angles of $1^\circ$ to $19^\circ$, and half a revolution later the detector is sampling pitch angles of $7^\circ$ to $25^\circ$. At the minimums in the modulations, the pitch angles being sampled are from $\sim 12^\circ$ to $15^\circ$ (see Figure 1). Hence we can make the important conclusion that the DED enhancement is due to electrons with $\alpha \lesssim 10^\circ$. This conclusion is not only important in itself, but, since only portions of the DED field of view detects these electrons, the enhancements represent fluxes even more intense than the ratio of count rates would imply.

Therefore we find that the flux of precipitated electrons at an altitude of $\sim 750$ km is enhanced by a factor of $\sim 10$ at pitch angles $\alpha \lesssim 10^\circ$, whereas electrons of comparable energy ($\sim 4$ to $7$ kev) at pitch angles of $\sim 30^\circ$ to $60^\circ$ are not appreciably affected.

**Effects at Other Electron Energies**

We now briefly examine the behavior of other electrons at the times of such a dramatic surge. We have yet to detect any significant effect, although further studies are being pursued. To illustrate the problem, the outputs of all sensors on Twins 2 for a few frames before and after the large pulse of Figure 2 are listed in Tables 2 and 3.

Generally we do find that, during such enhancements, the pitch-angle distribution is erratic (Figure 6) or more clearly peaked at $\alpha \sim 60^\circ$ rather than at $\alpha \sim 30^\circ$ (Figure 7). We thus have a distribution with one peak at $\alpha > 60^\circ$ and another intense peak at $\alpha > 10^\circ$, the two distributions apparently behaving independently, as is evident from Table 2 and the following related discussion. (Note that, since the existence of the modulation is dependent on the existence of pitch-angle anisotropies, we do not find any clear-cut examples of modulations in the DED data when the SPECS indicates relative isotropy over $30^\circ < \alpha < 60^\circ$.) However, we are not yet able to test the important question of whether enhance-

**TABLE**

<table>
<thead>
<tr>
<th>Frame Number</th>
<th>Magnetometer</th>
</tr>
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<tbody>
<tr>
<td>1011</td>
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<td>1012</td>
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</tr>
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<td>1021</td>
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<tr>
<td>1022</td>
<td>21</td>
</tr>
</tbody>
</table>

*Note that the 2B measurements at $\alpha < 10^\circ$ and isotropic fluxes.*

As we examine Table 2, several facts. First, the effect in the high DEDs went up to 202 counts per frame; in this instance the enhancements were such an enhancement was not treated here.

We would not expect to detect an enhancement of electrons for several frames in the enhanced precipitation.
ing pursued. To its of all sensors before and after listed in Tables hiring such en.

distribution is marly peaked at 30° (Figure 7). with one peak at > 10°, in the same way over 30° < a < 60° never arise over otherwise isotropic fluxes.) As we examine Tables 2 and 3, we note several facts. First, there is no perceptible effect in the high DED when the low DED went up to 202 counts per frame. The fields of view of both detectors are identical, and thus in this instance the a < 10° enhancement must have been confined to energies E < 7 kev. However, in other portions of the flight there was such an enhancement in high DED on Twins 2B, but, owing to the failure of the similar device on 2A after separation, we cannot prove these effects to be temporal, and they are not treated here.

We would not expect the down-viewing DED to detect an enhanced flux of backscattered electrons for several reasons. The solid angle of the enhanced precipitation is small, and the backscattered flux would be dispersed in pitch angle and degraded in energy, with resultant dispersal of arrival times back at the payload. Thus we see no increase in the down-viewing DED count rates. (See Table 3.)

* Note that the 2B sensors were sampled 4 msec (about half a frame) before those of 2A. See Figure 2.

As one views Table 3, it might be thought significant that there was apparently an increased upflux seen in channels 1, 2, and 5 of the down-viewing SPECS in the frame after the 'pulse.' Indeed, each count rate was more than twice its 49-frame average of 3.7, 3.3, and 5.8, respectively. The number of times the counts of 11, 7, and 16, respectively, were equaled or exceeded in the other 48 frames were 0, 6; and 1 times, respectively. Thus the probability of all three occurring by chance in this one frame after the pulse is of the order of 1 part in 10 thousand. Actually, it is even less than that, because the other high count rates

TABLE 2. Twins 2A Sample Responses around a 'BURST' at T + 624

<table>
<thead>
<tr>
<th>Frame Number</th>
<th>Magneto-meter</th>
<th>Low DED 1</th>
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<td></td>
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<td></td>
<td></td>
<td>2A 2B**</td>
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- Note that the 2B sensors were sampled 4 msec (about half a frame) before those of 2A. See Figure 2.

TABLE 3. Twins 2B Sample Responses around a 'BURST' at T + 624

<table>
<thead>
<tr>
<th>Frame Number</th>
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<tr>
<td>123</td>
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- Note that the 2B sensors were sampled 4 msec (about half a frame) before those of 2A. See Figure 2.
occurred, as one would expect, at smaller pitch angles (\( \alpha \sim 120^\circ \)), whereas the quoted numbers were for the maximum pitch angle sampled, namely \( \alpha = 150^\circ \).

Therefore, from these data alone, one cannot rule out the chance that there was an enhancement of up-moving 100-ev electrons close to \( B \) accompanying the pulse of 5-kev precipitating electrons. Since the rocket was at an altitude of over 700 km at the time, such electrons would have taken \( \sim 100 \) msec to travel up from a \( \sim 200\)-km source altitude. In fact (Table 3), they were present less than 10 msec after the pulse, with no time dispersion at all. We therefore conclude that this upsurge, although statistically significant at the 99.9% level, cannot be explained as backscattered electrons whose intensity increased because of the pulse of 5-kev electrons precipitated.

We have paid particular attention to this pulse in Twins 2B, because it is the only isolated pulse seen in the data to date. Yet it is undoubtedly a real effect (see Appendix 1). A comparably large pulse of 200 counts occurred in Twins 2B some 8 sec earlier accompanied by nearby pulses of 58 and 65 counts, and, of course, with pulses in Twins 2A (Figure 8). It is interesting to note that in this instance the upflux of electrons seen by channels 1, 2, and 5 at the same voltage level in the frames around the modulation peaks was very close to the 49-frame averages, and thus our concept that the apparent enhancement of Table 3 was due to chance is reinforced. This relationship is also observed for other bursts we have studied.

Since pulses have been found at all voltage settings of SPECS (see Appendix 1), we have been able to examine the related behavior of electrons over the complete energy range. In no instance were we able to detect consistently significant effects for different pulses, other than the pitch-angle distribution discussed above.

**Temporal Variations**

We have also examined the temporal variation in the development and progress of the modulation effects, and again no definite pattern has appeared. For example, on a given Twin there may be one large isolated modulation (Figure 2), a few isolated modulations (Figure 8), or a long train of modulations lasting several seconds (Figure 3). Whereas the low DED on Twins 2A appears a little more sensitive to the modulation than that on Twins 2B (which has a slightly higher-energy passband as in Table 1), it does appear that, whenever modulations are seen by 2B, they are present in 2A also. Separation of the two payloads ranges up to \( \sim 1 \) km in the plane orthogonal to \( B \), and the preliminary survey does not reveal any marked deviation from the above at any separation. However, a lengthy cross-correlation study is underway to examine those findings precisely, and at this stage we use only clear-cut isolated bursts, such as those of Figures 2 and 8, to justify the statement that the effect is seen simultaneously at separations (perpendicular to \( B \)) of up to \( \sim 1 \) km. If the effect had been a purely spatial one, i.e., if the two payloads had passed successively through a single stationary spatial structure, Twins 2A would have observed it about 1 sec before Twins 2B.

Of course, the possibility remains that the simultaneous enhancements were due to spatial structures moving relative to the trajectory of the payloads. For the leading edge of such structures to arrive at the two payloads separated by 1 km within the minimum resolvable time of 0.05 sec, the velocity of the structures relative to the payloads would have to be at least 20 km/sec. (Recall that the absolute velocity of the payloads orthogonal to \( B \) was 0.5 km/sec.) The effects could also be produced by spatial structures nearly aligned with and moving normal to the trajectory of the payloads. Although we concede that such situations could exist, it must be understood that there could have been no stationary spatial structures, for no time delays on the order of 1 sec were seen. Therefore, if one assumes that the effects were due to spatial structures, one must assume an extremely fortuitous combination of structure orientation and velocity, a situation that seems highly unlikely. We therefore argue that the enhancements are truly temporal in origin.

In summary, we have proven that, on occasions sometimes as short as \( \leq 0.1 \) sec but sometimes for several seconds, there is an \( \sim 10\)-fold increase in the flux of precipitated 5-kev electrons in a very narrow range of \( \alpha \leq 10^\circ \). The effect is seen simultaneously at altitudes of up to \( \sim 800 \) km by two payloads up to 1 km laterally distant from one another. To first order, the remaining auroral electrons display no significant change in pitch angle.

**SUMMARY OF EXPERIMENTAL RESULTS**

To summarize the enhancements of the electrons with pitch angle found to occur sporadically to 800 km. Since the...
...appears a little more sensitive to the
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constraints on any postulated mechanism of auroral acceleration, and now discuss the relevance of this finding to existing auroral theories. We refer to this phenomenon as a burst and its characteristics are summarized in Table 4.

Confinement to Small Pitch Angles

If we consider either a short-duration or a long-duration burst (see Figures 2 and 3), we find a remarkable pitch-angle distribution of 5-kev electrons, shown here as Figure 9. The extent of the confinement to very small values of $\alpha$, i.e., near-parallel to the geomagnetic field $B$ is quite extreme. For example, the count rate of the low DED sensor falls by a factor of 2 when it samples 5-kev electrons with $12^\circ \leq \alpha \leq 15^\circ$ instead of $1^\circ \leq \alpha \leq 10^\circ$, and we have concluded that the peak flux is at $\alpha \leq 10^\circ$. In other bursts, the modulation factor (which is a factor of ~2 in Figure 8) is as large as a factor of 10, and thus the confinement to near $B$ is even more extreme (Figure 10).

As seen in Figures 9 and 10, therefore, the flux of electrons with energy $E \sim 5$ kev has a pitch-angle distribution with a broad maximum at $\alpha \geq 60^\circ$ and a very narrow maximum at $\alpha \leq 10^\circ$. By contrast (Figure 9), the flux of very-low-energy electrons (40-50 ev) has a maximum only at $\alpha \leq 30^\circ$. This feature is important and is discussed later.

We should add that three other groups [Riedler, 1966; Lampton et al., 1967; Hoffman and Evans, 1968] also reported fluxes peaked at small pitch-angles. However, because those studies had neither the capability of separating spatial from temporal effects nor of studying electrons with energy down to ~40 ev, only with the Twins 2 data can we examine the effect more accurately and indeed show some previous interpretations to be in error (see below).

Temporal Extent of the Bursts

Although the bursts may persist for up to ~10 sec, one occurred with a duration of $\leq 0.1$ sec (see Figures 2 and 3). If one assumes that the electrons responsible had energies over much of each low DED passband (namely 3.7-6.7 kev and 4.0-7.6 kev), as is indicated by the fact that both the slightly different sensors detected the bursts, this short-duration pulse is extremely interesting. The velocities of the electrons at the lower and higher edges of the energy passband are, respectively,

$$V_1 \sim 3.5 \times 10^5 \text{ cm/sec}$$

and

$$V_2 \sim 5 \times 10^5 \text{ cm/sec}$$

Hence, for the electrons to arrive at the rocket payload simultaneously within $<0.1$ sec, if they were initially 'produced' at the same time and place, that location must be less than about 5000 km above the rocket.

The rocket at the time of this short burst was near its peak altitude of 800 km. Thus, given the above assumptions, the 'source' would have to have been at an altitude of $<6000$ km. To test this further, we sought such short-lived bursts in the lower-altitude data. Unfortunately, although the strong spin-modulation due to the fan-shaped off-center detector was still present on occasions, indicating peak fluxes at $\alpha \leq 10^\circ$, at no time was the modulation as brief as in the burst of $T + 624$ sec (Figure 10). Thus this one example remains. Since the envelope of the burst is not precisely defined, we will generally argue for a source altitude of $\leq 1 R_s$, given the above assumptions. We discuss this further below and compare it with estimates made by others.

With the source altitude of $\leq 1 R_s$, the count rate of the high DED sensor of ~12-kev electrons would be expected to have detected the burst some 50 msecs or 5 telemetry frames earlier, given the usual assumptions of simultaneous acceleration at the same source. In fact, there was no detectable effect at all.

In a general consideration of the temporal nature of bursts, which is of course the en-
AURORAL ELECTRON BURSTS

C. B. G. Wilson and A. M. McDiarmid

AURORAL ELECTRON BURSTS in error (see

in error). Bursts that persist for up to a duration of $\leq 0.1$ sec assume that the energies over some band (namely, $\leq 10^5$ eV) is indicated by different sensors. It has been stated that auroral pulse velocities of the edges of the bursts are short for about $0.1$ sec, if they last at all, and that the average duration is about $\leq 0.1$ sec. In fact, the rocket would move about $6000$ km. Thus, even for short bursts of duration $0.1$ sec, it is clear that the rocket would have moved about $6000$ km.

Fortunately, the time and space scales imply that the count

rate per frame is generally related to the count

rate per frame. If as in the cases studied in this one

study, $\alpha = 10^\circ$, where $\alpha$ is the pitch angle of the incident particles, we are confronted at once with the fact that there is a slight but definite difference in the responses of the similar sensors on the two payloads. These differences were of two types.

First, the low DED on Twin 2A was more prone to show strong spin modulation than the low DED on Twin 2B (for example, see Figures 2 and 8). Second, whereas the Twin 2A sensor showed nearly equal modulation at twice the spin rate, the modulation on Twin 2B was definitely unequal, with a more pronounced maximum when the magnetometer count rate (and hence pitch angle) was a minimum.

Because we have already seen how closely collimated the bursts are (Figures 9 and 10 and discussion), we believe that these differences arise from slight ($\sim 1^\circ$ to $2^\circ$) differences in the angular responses and mechanical mounting of the two sensors. It may be noted, for example, that in-flight calibration of the magnetometers implied that the Twin 2A rocket axis was about...
Fig. 10. Pitch-angle distribution during an auroral burst of 5-kev electrons. As the rocket spun, the off-axis wedged-shaped response of the low DED caused it to sample electrons over the ranges shown, namely (a) 1°-19°, (b) 7°-25°, and (c) 12°-15°. The SPECS instrument gave the value shown (d) at α = 30°. The hypothetical pitch-angle distribution is derived by combining all the measurements.

2° closer to B than was that of Twin 2B. A similar difference, in the same sense, was found on the Twin 1 flight [Westerlund, 1968]. Since there was negligible coning, these angles must actually be the same, and thus we can assert a possible sensor alignment error of ~2° in the correct sense for the observed differences in the responses. Since the bursts studied were of low-energy electrons, there may also have been an increased sensitivity of the Twin 2A sensor because its energy passband was of slightly lower energy.

Given the above conditions, we have found no consistent pattern of growth, duration, and decay of the bursts. Whatever the cause, the burst must be capable of lasting as short a time as ≤0.1 sec or as long a time as ~10 sec. The burst was quite a common feature of the Twins 2 flight, and, just from visible scanning of the data, it is prominent in at least 10% of the data.

**Electron Spectra**

The electron spectra averaged over pitch angles 30° ≤ α ≤ 60° ranged during the flight over the limits shown in Figure 11. This figure is derived by considering the maximum and minimum 49-frame (=0.5 sec) sums of counts from each channel. The individual spectrum during a 3.8-sec scan can be extremely complex and can have two clear peaks, one at say E ~ 50 ev and the other at E ~ 7 kev, or it may just have the 'continuum' spectrum discussed by Westerlund [1969] with a clear maximum in the range from 40 to 50 ev and merely a slight inflection in the vicinity of a few kilo-electron volts.

We have not been able to discern any consistent pattern of association of one particular spectral form with bursts, except for the fact that bursts of ~5-kev or ~12-kev electrons at α ≤ 10° are found only in association with electrons with 30° ≤ α ≤ 60°, which indeed have spectra at lower energies. In contrast, when the burst pitch-angles is of the order of 7°-25°, there must be an additional 'auroral' spectrum of a few kilo-electron volts, and there were occasional electrons at α ≤ 10° giving a secondary effect at ~5 kev.

We conclude that it is not possible to attribute an auroral peak to electrons with energies lower than 5 kev. There is possibly a contribution from bursts at these energies, but this contribution is small, and we have no evidence that it was present, or that it had much effect on the rapidly changing SPECS data.

**VARIATION**

Figure 9 shows another variation in pitch-angle and low-energy bursts, the minimum at α = 30° for low-energy bursts and ~30°. There is no explanation for this particular variation.

According to the model of Chamb, whereby the auroral trough is able to pass through a line of ions, and the ionospheric drop in density is the initial excitation, the electron energies may have a distribution of energies with an isotropic pitch angle. Finally, all the particles were at an altitude of 1000 km. This figure shows the likely periods of rapid increase (9), and...
have spectra with high intensities at these energies. In other words, bursts are not evident when the electron spectrum at larger pitch angles is of the continuum type. However, it must be emphasized that there were many 'auroral' spectra, i.e., with a strong peak at a few kiloelectron volts at $30^\circ \leq \alpha \leq 60^\circ$, when there were occasionally bursts of $\sim 12$-kev electrons at $\alpha \leq 10^\circ$ without any accompanying effect at $\sim 5$ kev.

We conclude, therefore, that the presence of an auroral peak in the spectrum at several kiloelectron volt energy and $30^\circ \leq \alpha \leq 60^\circ$ is probably a necessary but not sufficient condition for bursts to be found. We do not believe that this conclusion should be given undue weight because, when the continuum spectrum was present, the count rate of the low DED was low, and unambiguous detection of a burst in the rapidly varying small counts of the sensor was necessarily difficult.

**VARIATION OF SPECTRA WITH PITCH ANGLE**

Figure 9 clearly shows the typically different pitch-angle dependence of auroral ($\sim 5$ kev) and low-energy ($\sim 50$ ev) electrons. During bursts, the auroral electrons may have maximums at $\alpha \leq 10^\circ$ and at $\alpha \geq 60^\circ$, whereas the low-energy electrons have a maximum at $\alpha \leq 30^\circ$. There are considerable variations in detail, but here we wish to discuss the implications of this particular spectral pitch-angle variation.

According to Chamberlain [1969], knowledge like the above could be used to 'test any given model of electric acceleration.' We now proceed to such tests.

Chamberlain [1959] considers a mechanism whereby a weak parallel electrostatic field is able to precipitate particles initially trapped on a line of force. In this calculation, the potential drop was assumed to be small relative to the initial kinetic energies of the particles. According to the results of Chamberlain's calculation, the lowest energy particles, those with energies comparable to the potential drop, would have a pitch-angle distribution more or less isotropic with a tendency to be peaked at small pitch angles, whereas the higher-energy particles would be peaked at larger pitch angles. This finding is in agreement with our data for periods when no bursts are present (see Figure 9), and therefore a small potential drop ($\sim 40$ volts) could explain the situation in which the low-energy electrons ($\sim 50$ ev) are peaked at small pitch angles and the higher-energy electrons ($\sim 5$ kev) are peaked at larger pitch angles. However, if one wishes to extend this mechanism to explain the field-aligned bursts of high-energy electrons, one must assume a potential drop on the order of 5 kev or greater to account for the extreme collimation of the 5-kev electrons. However, such a large potential drop would also drastically reduce the flux of lower-energy electrons and at the same time increase the flux of higher-energy electrons ($>10$ kev) at large pitch angles. In fact, neither of these effects are observed in our data. In summary, we can explain the pitch-angle distributions of the electrons in the absence of bursts by a small potential drop ($\sim 40$ volt) along the field lines, but we are unable to explain the bursts by appealing to a much stronger transient potential drop because we do not observe corresponding effects on electrons at other energies.

It is interesting to examine the bursts in the light of Perkins' [1968] study of plasma-wave instabilities above an aurora. Perkins predicted that, if the 'monoenergetic' auroral electron flux were to grow above $\sim 7 \times 10^6$ particles cm$^{-2}$ sec$^{-1}$ ster$^{-1}$ at $\alpha \sim 90^\circ$, electrostatic plasma
waves could 'grow to an amplitude large enough to cause stochastic acceleration of a few electrons to energies of \( \sim 40 - 100 \) kev on a time scale of \( \sim 10^4 \) sec.' Perkins intended his studies to be applicable to the breakup phase of an aurora, which was not true for the Twins 2 flight, but it still appears to be useful to examine briefly the possible role of this mechanism.

First, since we did not examine electrons at \( \alpha = 90^\circ \), we are not in a position to determine whether the critical flux was reached. From study of the data, we believe, however, that the flux was generally less than \( 10^7 \) particles \( \text{cm}^{-2} \text{sec}^{-1} \) in the range from \( 1 \) to \( 10 \) kev. However, let us suppose that it did reach the necessary critical level, whether or not it was the specified figure. Perkins [1968] estimates a growth rate on a time scale of the order of 10 m/sec, and that certainly satisfied our minimum observed time scale of \( \sim 100 \) m/sec. Let us also assume that somehow the process is capable of enduring for up to \( \sim 10 \) sec with no perceptible damping or effect on the \( \sim 12 \)-kev electrons or any electrons in the \( 30^\circ \leq \alpha \leq 60^\circ \) range.

However, there are two unassailable criticisms of the applicability here of the process envisaged by Perkins [1968]. First, he predicts that the 'very energetic downgoing electrons should be produced roughly isotropically but with a weak concentration of pitch angles towards \( 90^\circ \)' Our Figures 9 and 10 refute this prediction. Second, Perkins predicts that the 'upwards-going very energetic electrons (should have) the pitch-angle distribution peak at \( 180^\circ \)'. Our data shown above refutes this completely.

In fact, the data are much more in accord with what one might call an 'upside-down' Perkins process with peak flux at \( \alpha \leq 10^\circ \) rather than at \( \alpha \sim 180^\circ \). However, such an upside-down process is, according to Perkins, not physically realizable when one considers the direction of propagation for which amplification of the wave packets will occur. In addition, of course, the occurrence of bursts over the altitude range of \( \sim 200 - 800 \) km, with no clear-cut differences, poses a problem for any process so dependent on local ionospheric properties.

Another theoretical model is that of Speiser [1965], wherein a small transverse \( B \) in the magnetotail causes ejection of particles energized in their drift across the postulated electrostatic potential. However, as pointed out by many others [Soumet et al., 1967; Hoffman and Evans, 1968] the magnetotail magnetic field is measured to be quite different from the assumed by Speiser [1963]. Second, the extreme collimation we observe in the bursts makes it seem very unlikely that they could have traveled such great distances [see Speiser, 1967], as does their occasional short duration (see above).

### Field-Line Merging

It is apparent from the above discussions that the theories or models treated [see O'Brien, 1967] cannot account for our observations. Nor can the observations be ignored because we have shown that they are consistent with, and more detailed than, other independent experiments [Hoffman and Evans, 1968; Riedler, 1966; Lampton et al., 1967] and the previous Twins 1 study (Westerlund, 1969). Thus, for example, Westerlund [1969] found the preference of the \( \sim 50 \)-ev electrons to have one maximum at \( \alpha \leq 30^\circ \), but he was not in a position to clearly detect the bursts of more energetic electrons. On the other hand, Hoffman and Evans [1968] found the field-aligned bursts, but they were unable to measure pitch-angle distributions of lower-energy electrons or, for that matter, to prove that the bursts were temporal. Thus we need another explanation of the observations. One proposed source of energization of auroral particles is the dissipation of magnetic energy involved when geomagnetic-field lines merge [see Axford, 1969; Piddington, 1968]. However, analysis of the conditions for merging of geomagnetic with interplanetary-field lines is extremely complex and, to our knowledge, imperfectly understood. We now will set up post hoc conditions for the source mechanisms, and then briefly examine whether these can be applicable to magnetic-field merging.

The source(s) of electrons must be able to explain the following types of fluxes:

1. A continuum flux extending more-or-less monotonically from a peak at \( \sim 40 \) ev up to energies of \( \geq 20 \) kev, which is generally present at night over Fort Churchill, with considerable fluctuations in intensities at given energies. This continuum has a pitch-angle distribution that is peaked at \( \alpha \leq 30^\circ \) (ev) and at \( \alpha \geq 60^\circ \) (10 kev). The continuum changes over the altitude and not over the invariant surface of 1962 of \( 70^\circ \leq \Lambda \leq 75^\circ \).

2. An auroral spectrum of electrons primarily with \( \alpha = 90^\circ \) is greatly enhanced as one approaches the continuum. Thus the observed continuum has a high flux of \( 40 - 100 \)-kev electrons, predominantly lower than \( 60^\circ \) and another, a burst.

Now, in agreement with O'Bric and Riedler, who used Twins 1 data to show that the continuum is consistent with either of the thermal or hot continua found in the magnetosheath, we propose the auroral spectrum to be composed of (a) isotropic, (b) anisotropic with one component, and (c) anisotropic with two components at \( 0^\circ \) and \( 60^\circ \) and another, a burst.

We believe the auroral energies concentrated in the bursts, to be due to pitch-angle diffusion, as proposed by Kennel [1969]. However, it is clear that the auroral burst has an exponential spectrum accompanying the diffusion or from 'fast ion' unknown to us and we do not have any evidence that these are trapped in the same continuum as the aurora.
Auroral Electron Bursts

3 at ~40 ev up to fluxes: s must be able to Ic-field merging. 1. We now will set the source mech-

2. An auroral spectrum in which the flux of electrons primarily with energies of ~5 kev is greatly enhanced and apparently superim-

3. Energy spectra such as those found by Frank [1968] and Bame et al. [1967] in the magnetosheath and by Montgomery et al. [1968] in the plasma sheet are similar to those when geomagneti-

4. We therefore tentatively conclude that when the Twins 2 observed anything except the field-

5. Now we examine the cause of the bursts. Axford [1969] has noted that, if an electron is able to make many traversals of the entire field line during the rapid contraction of the field lines after reconnection, they will take on a pitch-angle distribution peaked along the magnetic-field direction. But our observation of bursts of only ~0.1-see duration implied source altitudes of ~1 Rs. Furthermore, since we stated above, the flux of electrons at 30° ≤ a ≤ 60° remains apparently unaffected during a burst, whereas any such field-line shrinking would greatly affect both the pitch-

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and, of course, cannot say that one or another of the postulated source mechanisms never occurs.)

**Discussion**

We examined the various postulated acceleration/precipitation processes known to us and found each individually to be inadequate. We now briefly examine the adequacy of a mixture of several processes.

The following processes may occur. First, thermalized solar wind can, on occasion, gain access to the plasma sheet, thus giving it the continuum electron population. This process must be quite common, because, as we have discussed elsewhere [O'Brien, 1970], with the electron content of the order of $10^{4}$ particles in a tube of force through 1 cm² at 800 km above Fort Churchill, the observed precipitation $10^{3}$ to $10^{4}$ electrons/cm² sec would completely drain such a tube in some 10⁴ to 10⁵ sec. The rocket flight time was of the order of 10⁵ sec, and appreciable precipitation of at least $10^{4}$ particles/cm² sec was observed throughout. The visible auroral display also continued for several hours. It is apparent then that auroral-particle depletion of the ambient plasma in the outer regions of the magnetosphere poses a complex problem of replenishment.

Since the backscattered (i.e., upmoving) flux of electrons above 40 ev was measured to be of the order of 10% of the precipitated flux, any significant replenishment by backscattered electrons must be by those with energy $E \lesssim 40$ ev. It would thus be relatively slow, with a 10-ev electron taking tens of seconds to reach the equator. Nevertheless, such a return current could suffice, although an initial large surge might fade away through this depletion.

It is often suggested [Axford, 1969; Peddington, 1968] that convection of 'new' tubes of force across the polar cap will bring new supplies of plasma to the aurora. However, because one frequently can observe long multiple east-west arcs with no significant latitudinal variation of intensity and persisting for hours [Chamberlain, 1961], it is difficult to see how this resolves the depletion problem we discussed. Indeed, even if one turns to an external source of electrons, such as the solar wind, one can only just satisfy the particle demand. We have shown elsewhere [O'Brien, 1964] that the energy dissipated in the auroral world wide is, on the average, about 1% of the energy brought by the solar wind to the magnetosheath. By a comparable calculation, one can show that the worldwide deposition of auroral electrons into the atmosphere is of the order of 10% of the incident flux of solar-wind electrons. (A simple way to relate the energy and flux percentages is to compare the average solar-wind proton energy of $\sim 1$ keV with the average auroral-electron energy of $\sim 100$ ev.)

In the context of the above discussion, it appears to us to be most likely that the major source of replenishment of electrons in auroral tubes of force is the flux of backscattered solar-wind electrons with energies of the order of 1 to 10 ev. The role of convection would be then in smearing this upflux over many 'tubes of force.' If this is the situation, it implies that a particularly bright auroral form at a given location should be of short duration, although it may well be repeated in some tens of seconds.

Let us suppose then that, for some 1-10% of the time, the plasma sheet is populated by solar-wind electrons through direct connection with interplanetary-field lines. The rest of the time the field lines are closed. Of course, they might always be closed with particle population furnished by an $E \times B$ drift.

We further suppose that a Taylor-Hones acceleration mechanism can occur, but it occurs in a microscopic extremely variable manner. Thus the magnetospheric structure averaged over time is like the Taylor-Hones equipotential pattern, but at any instant positive or negative potential differences of kilovolts between neighboring tubes of force may exist. This would explain the extreme variability of the flux of $1$ to $20$ keV, which had no discernable latitudinal dependence.

We then need moderate diffusion processes [Kennel, 1969] to precipitate the continuum and then change to strong diffusion when the more energetic electrons are present. Alternatively, of course, as Kennel [1969] comments 'the high-energy tail (caused by strong diffusion) can be more extensive than in weak diffusion.'

To produce the $\alpha \lesssim 30°$, we present a potential drop ($\sim 200$ ev) through the magnetic field line up to $L \sim 5$, and then $E \times B$ must extend to values of the depletion rate of the order of a few meters per second.

But we had better undertake to explain the bursts. After all, we had the liberty we took when we turned to the current system. In the context of the above discussion, it was the shortest-duration limitation of excess electron current $R_x$. We have chosen this origin because it permits two payload $\sim 1$ sec before auroral bursts. Because of the high fluxes, it is difficult to explain the simultaneous replenishment by both payloads. Each merely happens to come through two small ($\sim 50$ meters diameter) electron plasma clouds which do not appear to have the region of order 10% of the payload. In spite of these events, although small, the current theory can explain above events, because the region of order 10% of the payload does not appear to have the region of order 10% of the payload.

However, we can indicate that our concept may be valid. To measure the time-averaged flux at a rocket of this size, we had to deduce from $R_x$ [1967a, 1967b]. We had all indicated a spatial origin; no spatial origin; these events support our findings.

Accordingly, we can 'explain' the bursts by merging and current systems occurring at various altitudes.
Auroral Electron Bursts

To produce the low-energy electron peak at \( \alpha \leq 30^\circ \), we presumably should add a small potential drop (\( \sim 40 \) volts) along the geomagnetic-field line up toward the plasma sheet. It must extend to very high altitudes to satisfy the depopulation rate. It could thus be of the order of a few microvolts per meter or less.

But we had been unable to set up a model explaining the bursts, no matter how much liberty we took with magnetospheric topology and current systems. A major stumbling block was the shortest-duration burst, with its implied limitation of the source altitude to \( \leq 1 R_s \). We have claimed it to be temporal in origin because it was seen simultaneously by two payloads \( \sim 1 \) km apart, whereas for several seconds before and after there were no such bursts. Because of the great variability of the fluxes, it is difficult to assess the probability that the simultaneous detection of this burst by both payloads was due to chance and that each merely happened to pass simultaneously through two small spatially localized regions (\( \sim 50 \) meters diameter) where a highly collimated electron flux was present. This event does not appear to be very likely, both because the region would have to be only of the order of a few cyclotron diameters and because we have simultaneously seen several other events of only \( \sim 0.1 \)- to \( 0.2 \)-sec duration. In spite of these factors, the possibility is real, although small, and, if an otherwise acceptable theory can explain all the observations discussed above except for the altitude limitation, it presumably should be accepted temporarily.

However, we have further relevant data that indicate that our altitude restriction to \( \leq 1 R_s \) may be valid. When attempts were made to measure the time delays between the 'arrival' at a rocket of electrons of several energies and thus to deduce a source altitude [Evans, 1968, 1967a, 1967b; Lampton, 1967; Winiecki, 1967], all indicated a source of \( \leq 1 R_s \). Those experiments were based on the assumption that variations observed were of temporal rather than spatial origin; nevertheless, they do tend to support our finding.

Accordingly, although we might prefer to 'explain' the bursts as being due to field-line merging and contraction, the necessity for this occurring at various latitudes between \( 71^\circ \) and \( 73^\circ \) at altitudes of \( \leq 1 R_s \) seems to us to place untenable requirements on magnetospheric topology.

**Practical Implications of These Results**

Some additional interesting and nontrivial implications of this study are now mentioned briefly. We have proven basically that the flux of auroral electrons at \( \alpha \sim 0^\circ \) can be enhanced with no appreciable effect on the flux at larger pitch angles. With a smaller \( \alpha \), an electron of a given energy will, of course, penetrate deeper into the atmosphere. Yet many studies equate deeper penetration with a higher energy and neglect this pitch-angle dependence.

To cite one example, we discuss Eather's [1969] study of O I (6300-A) pulsations. From the interrelation of pulse rate and quenching times, Eather concludes that the faster pulsations (\( \sim 1-5 \) sec) are associated with higher-energy electrons. In fact, he states that this should be an integral feature of any theory of pulsations. This conclusion is based on a peak O I emission at \( \sim 97 \) km versus \( 118 \) km. Of course, such an extended atmospheric penetration is readily achieved with a particle of the same energy but smaller pitch angle (such as reported here), and thus we presently see no need to make such energy-period dependence a prerequisite of satisfying auroral theories. Intuitively one might feel the need for such an interdependence, but the data we present here show that past studies have not delineated it. We have also shown [O'Brien, 1971] the extreme uncertainties involved in deducing characteristics of auroral electrons from ground-based studies.

**Concluding Comments**

Far from clarifying the cause and source of auroral electrons, the additional unique capabilities of the Twin 2 experiments have served to show that no single theory known to us satisfies all the observations. However, we find that a potpourri of several theories can explain the data, if we considerably modify each model.

Thus, if we invert the sense of the plasma instabilities visualized by Perkins [1968], we might explain only the field-aligned bursts. If we allow the magnetic-field lines through the aurora generally to close in the plasma sheet...
but frequently to meet interplanetary-field lines in the magnetosheath [see Dungey, 1968], we can explain the continuum and its replenishment, given weak diffusion by wave-particle interactions [Kennel, 1969] to precipitate it.

To provide the auroral-electron peak with a few kiloelectron volts, we need either strong diffusion, acceleration in a very disordered Taylor-Hones model field, or both of these.

Finally, to explain the pitch-angle distribution of low-energy electrons (40-100 ev), we need a parallel electrostatic field at high altitudes with a potential drop of the order of 40 volts.

Such a post hoc mixture of the various models is aesthetically far from attractive and scientifically very unsatisfactory. It is possible that it can be replaced by a new unified theory. We wish to emphasize, however, that every individual finding of our experiment is substantiated by other experiments, and we have brought the whole together with the Twins 2. Therefore, any postulated source mechanism must be capable of simultaneously producing electron fluxes of the nature outlined here. We regard it as most likely that there are two or more mechanisms involved, although whether there are the five or six mechanisms discussed above is not clear, nor is the extent to which these processes are coupled or uncoupled clear. These uncertainties remain a matter for further study.

**APPENDIX.**

**Search for possible spurious effects.** It is clear from the exact relationship of the modulation of counting rates and the magnetometer orientation that the large modulations of Figures 2, 3, and 8 must arise from a peculiar property of each payload. The fact that, when the sporadic effect is seen in one payload, it is seen in the other implies that the auroral radiation must be different at those times. For reasons given in the text, we judged that, when the modulation is present, it is caused by the off-axis fan-shaped response of the DED sensor sweeping in and out of an abnormally large flux of electrons with pitch angles $\alpha \geq 10^\circ$. Here we examine the following other possible explanations:

1. 'Stray' electrical or magnetic fields in the payloads.
2. Malfunction of the DED's.
3. Saturation effects in any sensor.
4. Spurious counts from any of sources, e.g., corona, pulse-code-modulation.
5. Interaction of one payload with the other.
6. Modulation of the ambient auroral electrons by the transmitted and modulated $K_{p}$ power ($\approx$ 5 watts).
7. Peculiar shielding of the sensors by 'stray' material or by a particular payload configuration.
8. 'Background' radiation or protons.
9. Acceptance of lower- or higher-energy electrons.
10. Periodicities in the auroral fluxes at small pitch angles.

We find that none of these causes can possibly give rise to the observed effects. If the enhancements were due to any of the possible effects numbered 1 through 6, a spurious effect would be required to be consistent in synchronization with the payload spin rate. We cannot conceive of such a fortuitous circumstance. We can further assert that no stray material was intermittently shielding the detectors (possible cause 8), since the absence of coning showed that the payloads maintained their structural integrity after separation. Contamination from intense fluxes of particles outside the energy passbands of the DED's can be discounted because the other particle detectors would have detected such fluxes. Finally, if the effects were due to natural periodicities in the incident electron flux at small pitch angles, the periodicities would have to lie in phase with the individual payload spin rates, a highly unlikely possibility. We therefore conclude that our original interpretation is correct.

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AURORAL ELECTRON BURSTS

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Auroral activity is caused by changes in the magnetosphere (Alfvén, 1950). A direct measurement through auroral fluxes of low energy (10 kev) particles established the auroral form (Albert, 1959).

Ground-based observations of auroral bays are important in understanding the auroral processes. The auroral particle fluxes observed by rocket-based experiments are significant in studying the high-energy particles.