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THE INTRINSIC VALUES AND COLOR EXCESSES OF (B-V) FOR 115 F-K SUPERGIANTS

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ABSTRACT

Color excesses in B-V are determined indirectly from a study of Strömgren's b-y color for a sample of F0 - K5 supergiants. The resulting E(B-V)'s are estimated to have an expected precision of ± 0.05 . With the calculated color excesses and the observed values of B-V given in various catalogs, the run of B-V with spectral type is obtained. This B-V/(spectral type) relationship is compared with those found previously by other investigators.

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I. INTRODUCTION

The supergiant FO-K5 subgroup is of much importance for three basic reasons. First, they are useful in the determination of galactic structure and kinematics, either as a self-enclosed unit (Bidelman 1958), or as a sub-unit to a larger sample (Humphreys 1970). Their worth in such studies is somewhat limited by the difficulty of assigning to individual stars accurate absolute magnitudes and reddening corrections. In addition, these stars are not of zero age and some have moved appreciably away from their place of origin. However, as evolutionary models for stars become more precise it is possible that they will be fine markers for the determination of a galactic diffusion constant. A second asset of the group, which results from their extreme brightness, is their detectability to great distances, even in extra-galactic objects. Thus, they can be used to determine both the galactic contribution and the projected-space variations of the internal reddening in extra-galactic objects, and to evaluate distance moduli (Dachs 1972). Finally, they are a class of stars whose atmospheres are inherently interesting:

- a/ They are the stable equivalents of the cepheid variables, which they resemble closely; in fact,

the correspondence is so tight that some appear to reside within the instability strip, but they themselves are stable (Schmidt 1972a).

- b/ The combination of physical conditions in their atmospheres is unique and an understanding of it is most intriguing - high radiation pressure, extended atmospheres, very low gravity, micro- and macro-turbulent motions, mass loss by stellar winds, radiative diffusion, effectiveness of convection, etc.

Recently, the attacks on the understanding of the physical make-up of supergiant atmospheres are of ever increasing sophistication and completeness (see, for example, Bell and Rogers 1969; Auman 1969; Parsons 1969, 1971; Böhm-Vitense 1972; Schmidt 1972b,c).

To implement fully projects in any of the above three categories it is essential to have accurate observational assessments of intrinsic colors. The UBV system's B-V color is one of prime importance. It is measurable to a faint limit, either photoelectrically or photographically; is used extensively, and have a large body of data already carefully cataloged (Blanco et al. 1968). Theoretically B-V is a color reasonably easily simulated from atmospheric model results, and thus it is a useful fiducial mark to tie theory to observation.

There have been numerous efforts to ascertain $(B-V)_0$ as a function of spectral type (see, for example, Kraft and Hiltner 1961; Fernie 1963; Schmidt-Kaler 1965; Johnson 1966; FitzGerald 1970). The results, while in general accord, are divergent up to 0.1^m or more at a single spectral type. Most of the troubles doubtless arise from the insufficiency of U-B, B-V and spectral type data as a complete set, an aspect of the problem succinctly analyzed by FitzGerald (1970).

In this paper we approach the finding of $(B-V)_0$ by an indirect route. Using a G-band measure and a blue-minus-yellow color gives precise color excesses which are transformable to $E(B-V)$'s. With these excesses and values of B-V listed in catalogs, a first value for $(B-V)_0$ for particular stars is found. We determine an additional estimate for $(B-V)_0$ using the linear correlation between the intrinsic value of our blue-minus-yellow color and B-V. A $(B-V)_0$ for a given star is the simple average of the two estimates. The $(B-V)_0$'s so determined are averaged for stars of the same spectral type to give the run of $(B-V)_0$ with spectral type. The findings are compared to those found by investigators.

II. THE DATA

Kelsall (1971) observed F-K supergiants on a composite photometric system made up of the four-color Strömberg system (1966) and the ABC system of Crawford (1961). From these results we utilize the yellow magnitude, y , the blue-minus-

yellow color, $b-y$, and the G-band index,

$G = B-A = 2.5 \text{Log} (I_{4297}/I_{4377})$. The UBV data consists of the V magnitude and the B-V color. These are from the catalog of Blanco et al. (1968), and for four stars from the listing of Humphreys (1970). In the cases where there are multiple listings in the Blanco et al. catalog, a simple unweighted average is taken.

The yellow magnitude of the Strömgren system is linearly related to the V of the UBV system, as is seen in Fig. 1. The RMS deviation about the straight line $y = V$ (which is as adequate as any fit), is $0^{\text{m}}.04$, and with average deviation of $+0^{\text{m}}.004$ which is below significance. In fact, the fit $y = V$ is at the expected noise level of $\sim 0^{\text{m}}.02$, if only seven of the one hundred and six stars are removed from the sample. Recognizing that some supergiants exhibit small fluctuations in V, and that the 1-to-1 fit is near perfect, we accept the yellow magnitude y as identical to V, and list it as such here. The comparison between the colors $b-y$ and B-V is illustrated in Fig. 2. There is an obvious linear relationship. The differential reddening effect on the two colors is evidenced by the successive sliding of the earlier spectral types under the later spectral types.

III. THE DETERMINATION OF THE $E(B-V)$ 'S.

The dependence of the G-band index on $b-y$ is clearly linear, as is shown in Fig. 3. The diffusion of the points toward the right is the result of reddening. The slope of the reddening line is 0.15, with an estimated error of ± 0.02 . It appears that a linear blue-most envelope is viable. The simple linear approximation is supported by: (1) the data

plotted in Fig. 3; (2) the tightness of the $G/(b-y)$ loops for cepheids, all of whose semi-major axes have approximately identical inclinations in the $G, b-y$ plane (Williams 1966; Kelsall 1971); (3) the determination of G on $b-y$ found by convoluting atmospheric simulations with the filter functions appropriate for these colors (Bell and Parsons 1972). From the data shown in Fig. 3 we find reasonable the relation --

$$G_{\circ} = 0.532(b-y)_{\circ} - 0.091. \quad (1)$$

The zero point is a revision of that given earlier (Kelsall 1971) which appears, from various evidences, to give a too conservative estimate for the blueward position (Kelsall 1972). An excess in $b-y$ is obtained by simply extrapolating the observed position of the star back along the reddening trajectory until it intersects the thermal locus --

$$E(b-y) = (b-y)_{\text{obs}} - 2.618(G_{\text{obs}} - 0.15(b-y)_{\text{obs}} + 0.091). \quad (2)$$

Clearly from the coefficient of the second term on the rhs of Eq. (2) we see that small errors in the observed quantities and the zero point contribute significantly to errors in the excess. We should expect the excess for a particular star to be accurate to the order of three or four hundredths of a magnitude. Of course, this is all within the simplex frame work of a line thermal locus. Doubtless the locus is a band of finite width, but the implementation of such a refinement in the analysis is presently beyond us.

The transformation of the $b-y$ excesses to excesses in $B-V$ is not as trivial as expected. Epstein (1969) in a fine study of the RR-Lyrae stars finds the ratio of the color

excesses, $r_E = E(B-V)/(b-y)$, to be 1.82; while Strömberg (1966) and Crawford et al. (1971), working mainly with B-F stars, use an r_E of 1.43. Crawford et al. do point out that this ratio does not appear good for their supergiant results, though, of course, these results are for supergiants of much earlier types than those considered here. Kelsall (1972) finds for the cepheids that a fine compromise for r_E is 1.22. This ratio brings his excesses in $b-y$ onto the excellent base of the Kraft-Sandage-Tammann system of B-V colors for cepheids (Sandage and Tammann 1971). To decide on an r_E suitable for the F-K supergiants we have five pieces of information. First is the result for the cepheids stated above, a result which should be applicable to the range F3 to G2 for the supergiants. A second point is that the convolution of the filter functions, simulated spectra in the range early F's to early G's, and van de Hulst's curve #15 for the interstellar reddening leads to a predicted r_E of 1.27 (Bell and Kelsall 1973). Third is a comparison of our $E(b-y)$'s against the theoretical/observational evaluations for $E(B-V)$ made recently by Schmidt (1972b). Unfortunately, the overlap in the projects is only ten stars, but from these we find an r_E of 1.18, a value germane to stars of spectral types F0 to G2. We should note that Schmidt found it difficult to reconcile our results for the supergiants and cepheids with his predictions. However, he compared an $E(B-V)_{\text{Kelsall}}$, not our measured quantity $E(b-y)$, to an $E(B-V)_{\text{Schmidt}}$. The transformation $E(b-y) \rightarrow E(B-V)$ is not stated in his article.

Probably an ad hoc transformation was used ($R_E \sim 1.4$) with no effort made to ascertain a more appropriate regression line. However, a puzzle remains in attempting to coalesce Schmidt's and our work for the regression lines linking our excesses are not identical for the cepheids ($r_E \sim 1.0$) and the supergiants ($r_E \sim 1.2$), but this might arise simply from there being too few stars in common in both groups (see also Kelsall 1972). The fourth datum is the comparison of our $E(b-y)$'s with the extensive observational compilation of $E(B-V)$'s by Buscombe (1970). The scatter is quite large, but the evaluation of r_E is in the range of 1.3, whether all the stars in common (92) are used or the comparison is made for limited spectral type groupings. Except for Buscombe's results which cover the full range F0-K5, the other clues to the proper transformation of the excesses are restricted to the F through the early G stars. To include the late G and K stars we use as our fifth indicator the dependence of B-V on b-y at a particular spectral type. Here, of course, the assumptions are that the spectral type designations are accurate, and that the differential positioning of the stars in the (B-V), (b-y) plane is purely a result of reddening. The G5 stars give an $r_E = 1.24$, and the K3 stars require an $r_E = 1.21$. It should be noted that groupings of stars of earlier spectral types produce r_E 's more in the range of 1.3. All five evidences presented above appear to indicate that the compromise used for the cepheids (Kelsall 1972) is also eminently suitable for the supergiants. Thus, we accept the relation --

$$E(B-V) = 1.22E(b-y). \quad (3)$$

We can use the linear dependence of B-V on b-y to give us a second, though not totally independent, estimate for E(B-V). From the reddening-corrected (b-y)'s for the stars least reddened in b-y and the unreddened values of B-V using Eq. (3), we find that --

$$(B-V)_0 = 1.594(b-y)_0 + 0.053. \quad (4)$$

A relation virtually identical to that found by Kelsall (1971) for the cepheids, when that relation is modified to our new zero point in b-y --

$$(B-V)_0 = 1.606(b-y)_0 + 0.071. \quad (5)$$

The relation (4) is quite insensitive to the imposed reddening corrections, for if no corrections are made in either color we find for the blueward-most stars --

$$(B-V) = 1.596(b-y) + 0.033. \quad (6)$$

Once the intrinsic b-y is fixed we determine our second, smoothing estimate for E(B-V) from --

$$E(B-V) = (B-V)_{\text{obs}} - (1.594(b-y)_0 + 0.053). \quad (7)$$

The errors here should be comparable to those for the transformed excesses in b-y, which we estimate to be of the order of $\pm 0.^m05$.

We adopt as a final E(B-V) for any star the arithmetic average of Eq. (3) and Eq. (7). These accepted values for the excesses are compared in Fig. 4 to those given by Buscombe (1970) and Schmidt (1972b). The agreement in trend is satisfactory in both cases, as is demonstrated by the distribution of the points about the line of slope = 1. Relative to Buscombe's results we systematically estimate lower

reddening values. There is no clear cut dependence on spectral type, but such correlations are hard to discern because of the large scatter. Comparison with Schmidt is limited by the few stars in common, but we note the obvious discrepant values for the two G supergiants α and β Aqr, while the F stars are in total agreement except for zero point. This discord in the excesses for these two stars is in common with the estimates for the excesses made by Parsons and Bell (1972) using spectrum synthesis techniques. However, Schmidt does consider his method to be difficult to apply to stars with $(B-V)_0$'s greater than ~ 0.8 , a value close to that which we predict for these two stars (α Aqr, 0.90; β Aqr. 0.72).

IV. THE RESULTS

In Table I we present the results for our 115 stars. A brief description of the columns is as follows:

- Col. 1 - name of the star;
- Col. 2 - HD or BD catalog number;
- Col. 3 - MKK classification;
- Cols. 4-5 - new galactic coordinates;
- Col. 6 - yellow magnitude from Kelsall (1971);
- Col. 7 - B-V as listed in Blanco et al. (1968), or Humphreys (1970);
- Col. 8 - our adopted value for $E(B-V)$, enclosed in parentheses if it is purely a direct transformation of $E(b-y)$ via Eq. (3) (all subsequent entries are also enclosed in parentheses);

- Col. 9 - visual magnitude corrected for reddening,
 $V_c = V - 3.3E(B-V)$, following Martin (1971);
- Col. 10 - intrinsic value for B-V;
- Col. 11 - calculated distance from the sun based on
 Blaauw's (1963) calibration of absolute
 magnitudes.

In Fig. 5 we show our approximation for the run of $(B-V)_0$ on spectral type. At each spectral type the values for $(B-V)_0$ are the average of the data given in Table I. Most data points are given a weight of one, except those in parentheses which are given half weight. The results for the luminosity class II stars are excluded totally; our data is too scant on this group for any conclusions, except to say that they do appear different from the luminosity class I stars. The large circles in Fig. 5 are these weighted average points, which we call the "raw" run. The series of connected straight lines in the figure are our eye-estimate for the run, which we designate as the "smooth" run. If the values of the smooth run in B-V are plotted against their b-y equivalents, it is found that the points in the $(B-V), (b-y)$ plane are scattered about the intrinsic relationship given by Eq. (4). Assuming equal errors in the two colors we can adjust the runs on spectral type so each spectral type point falls on this intrinsic line of Eq. (4). We call such an adjustment the photometrically "perfect" run. In Table II we compare our three runs to those previously

determined by other investigators. Our findings are in accord in rough aspect with the others, with the largest divergences occurring in the early K results.

V. CONCLUSIONS

Adequate transformations can be found to convert the results in the Stromgren-Crawford system onto the UBV system; in particular, the color excesses for B-V. The usefulness of the inclusion of Crawford's G-band index in conjunction with a blue-minus-yellow color is the ability to determine precise color excesses for F-K supergiants. As B-V and b-y are linearly related it then would appear reasonable for UBV observationalists to consider carefully the option of expanding their system to the UBV:AB system. In addition the index G is a measure of a physically important parameter.

The ultimate aim of finding an improved run of $(B-V)_0$ on spectral type is accomplished. The run, while considered precise, does not support any other single previous determination. Our run must be considered as yet another approximation to a difficult, but important problem.

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VII. FIGURE CAPTIONS

- Fig. 1 The V magnitude from Blanco et al. (1968) versus the y magnitude from Kelsall (1971).
- Fig. 2 Comparison of the two blue-minus-yellow colors B-V and b-y.
- Fig. 3 The G-band index versus b-y.
- Fig. 4 Relationship between the B-V color excesses of Kelsall to those of Buscombe (1970) and Schmidt (1972b).
- Fig. 5 The run of $(B-V)_0$ with spectral type.

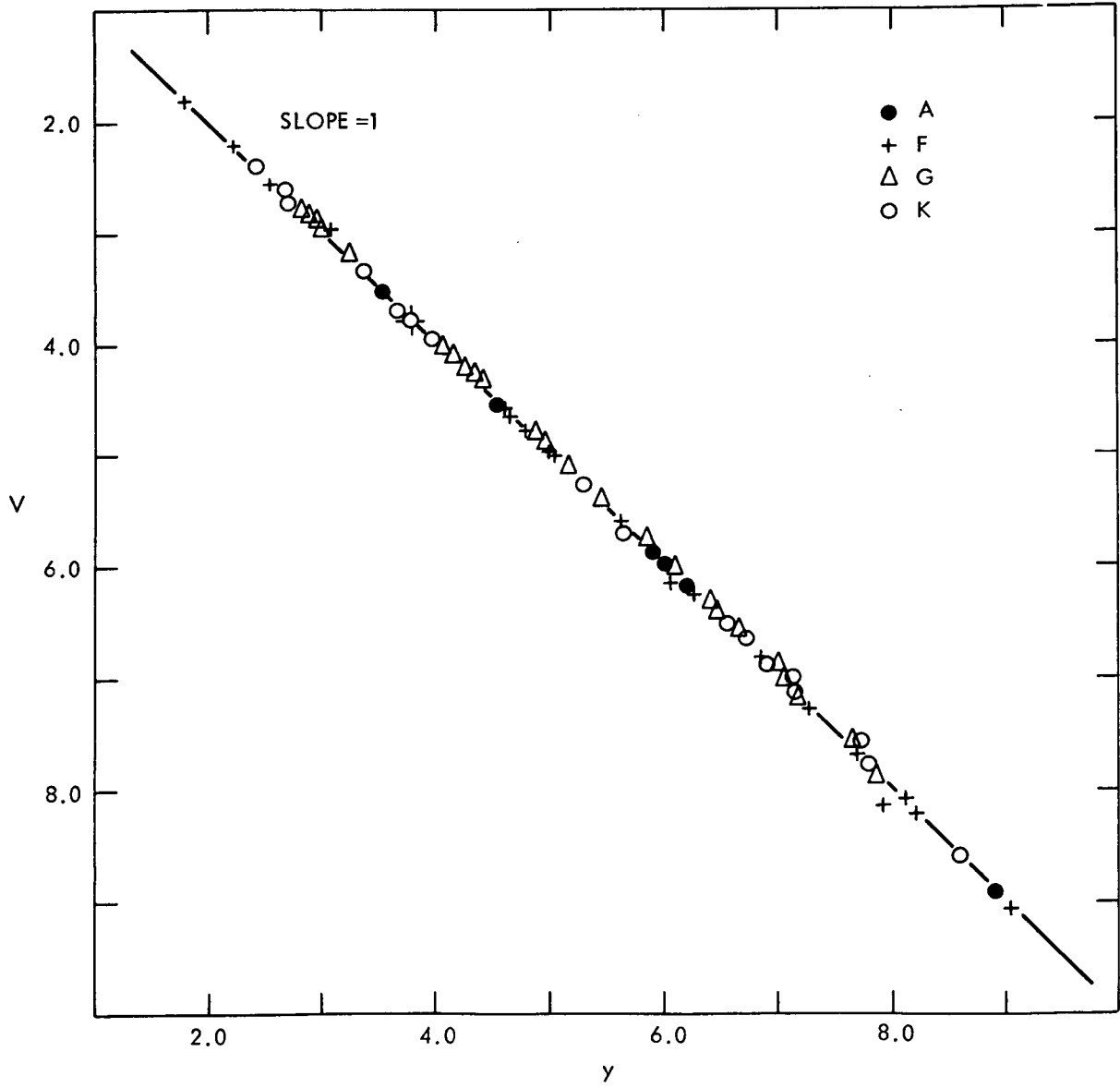


Fig. 1

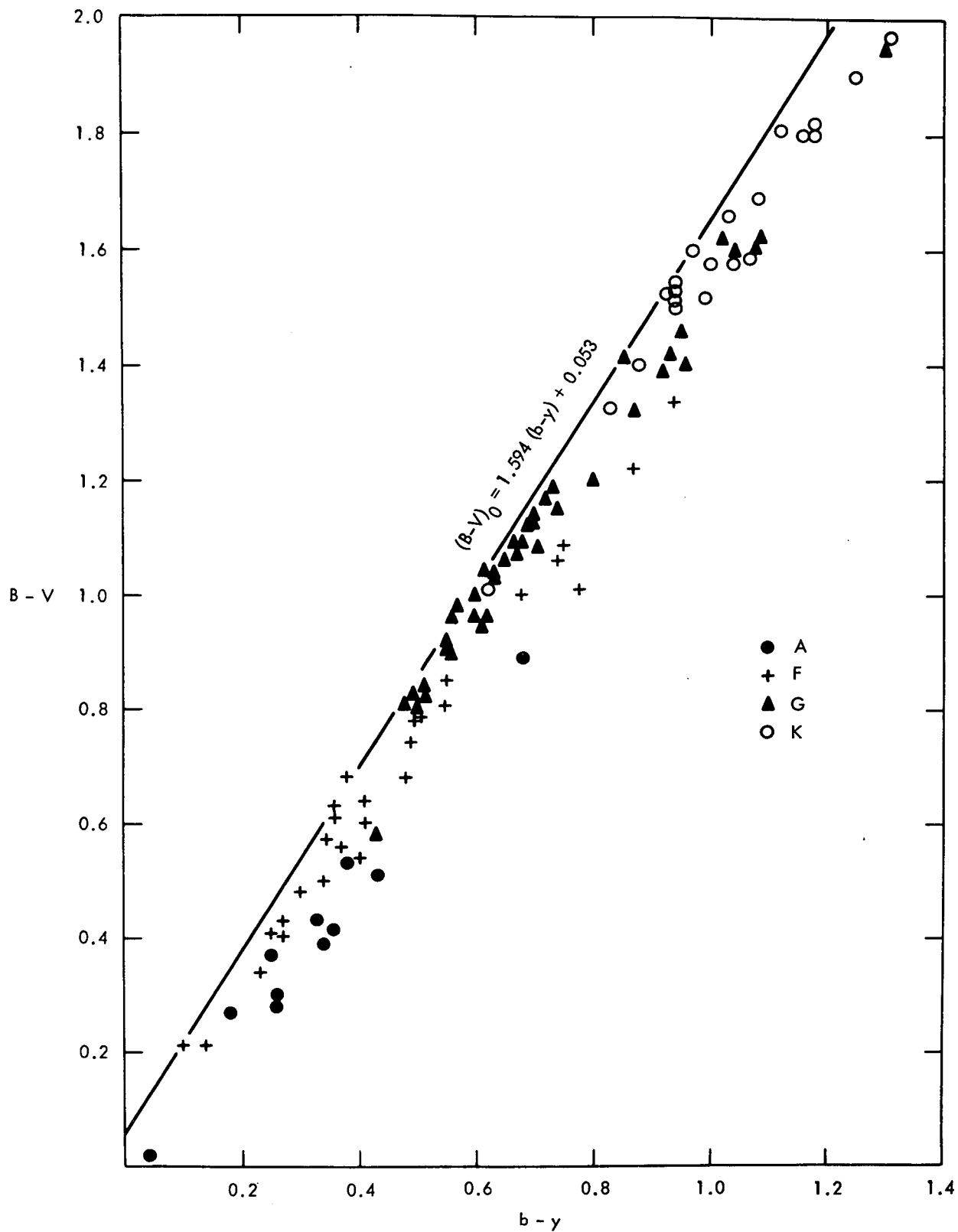


Fig. 2

Fig. 3

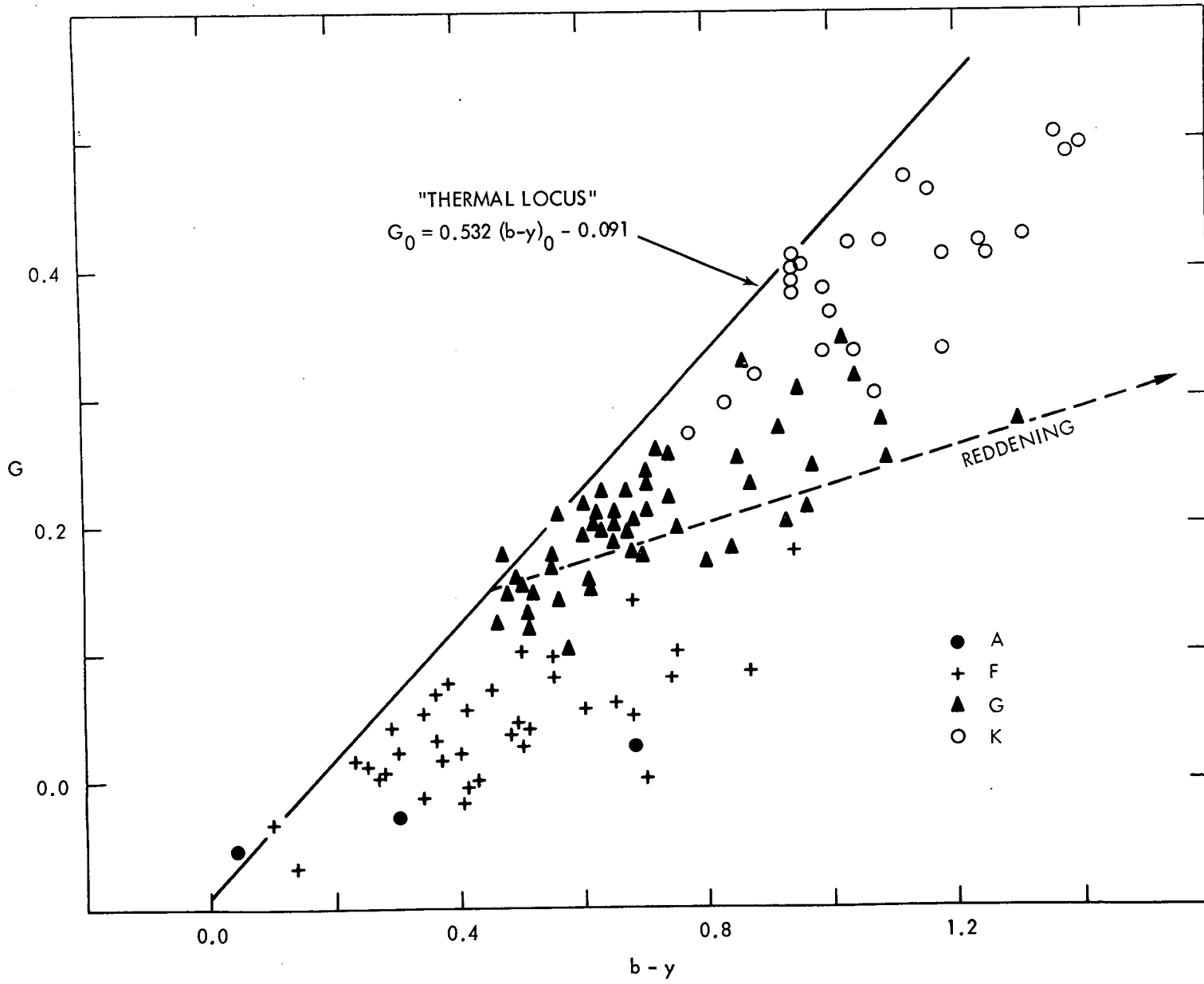
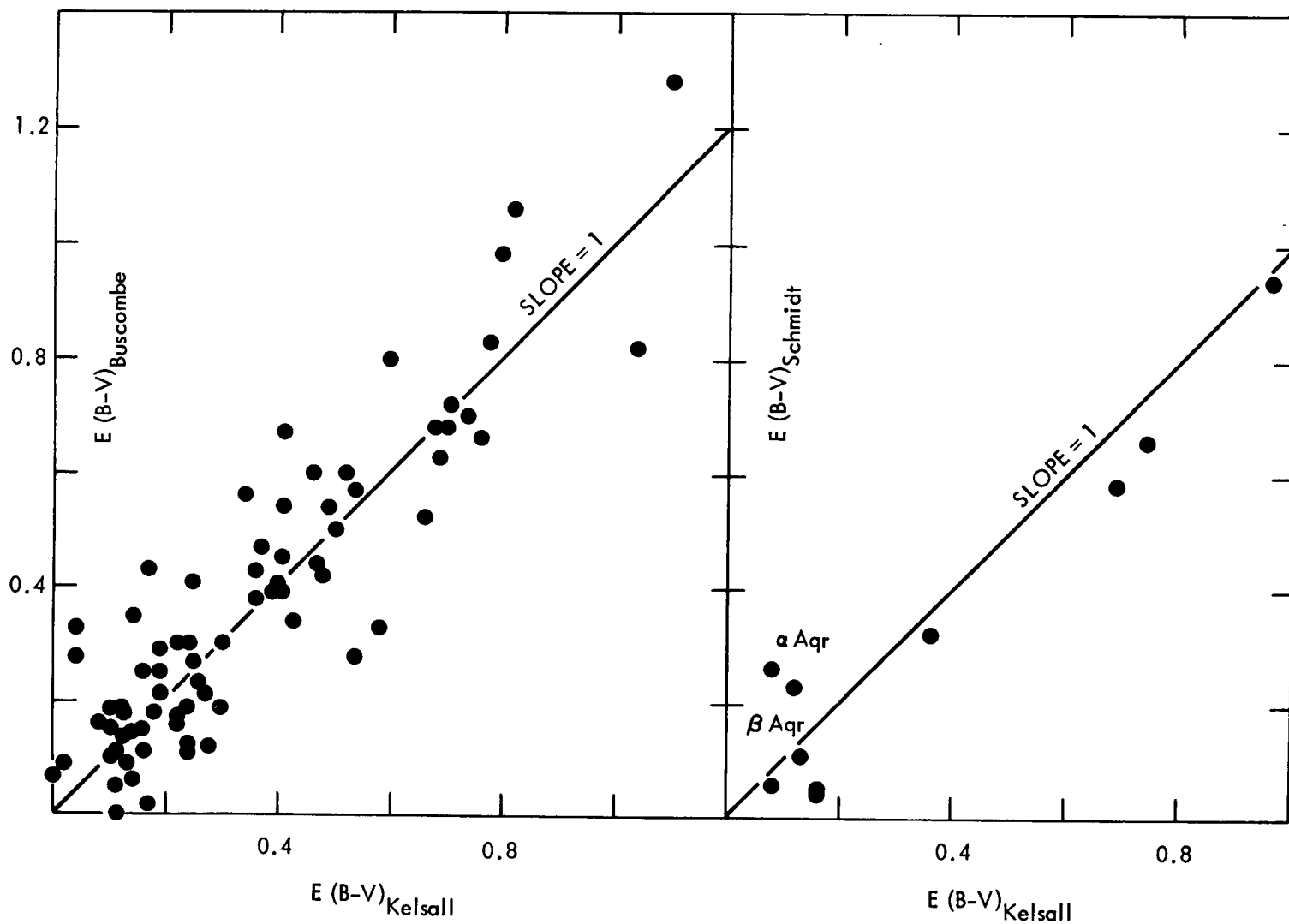


Fig. 4



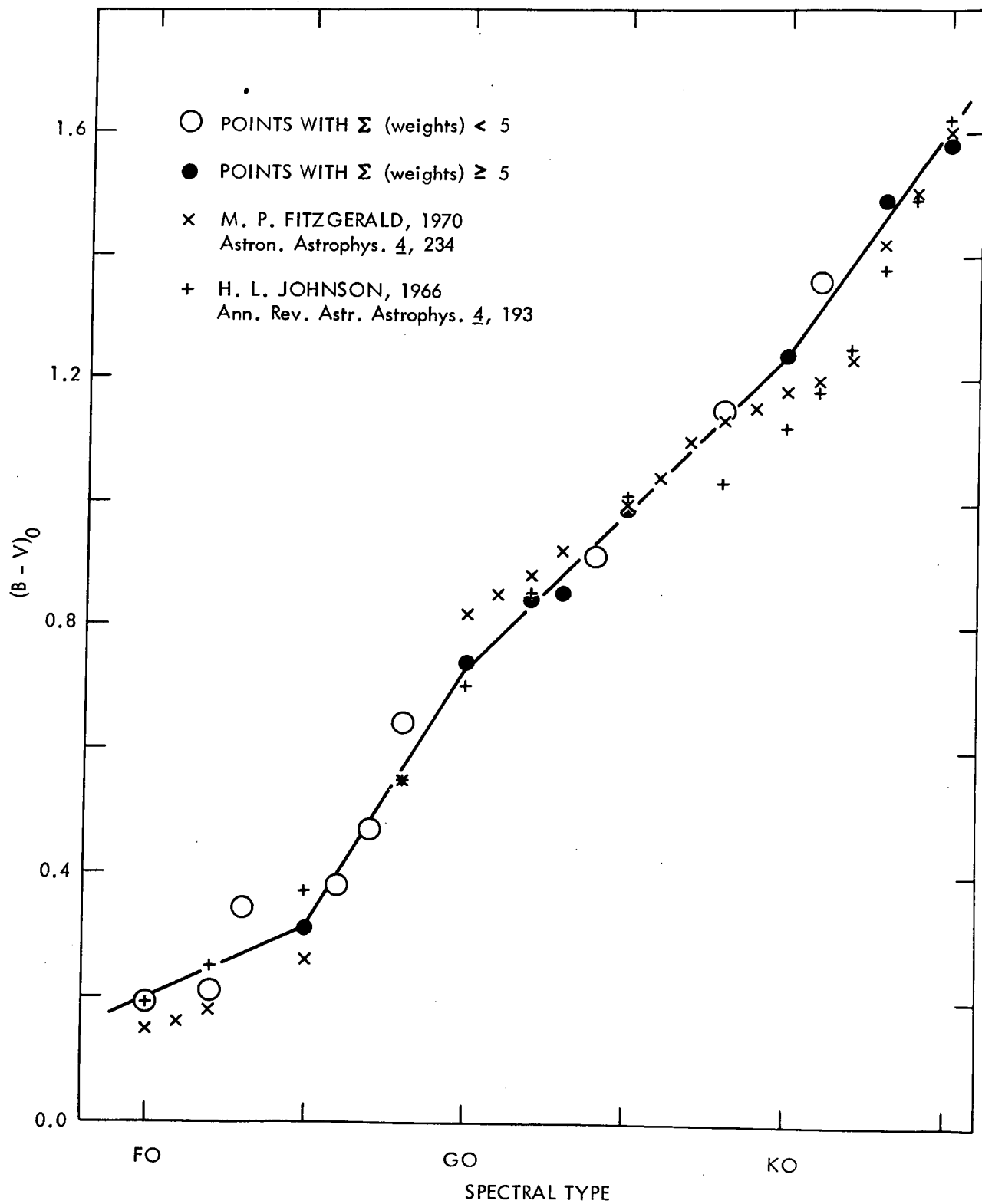


Fig. 5

TABLE I. Data for the 115 supergiants.

| Name | HD/BD | MK Class | l | b | V | B-V | E(B-V) | V_c | $(B-V)_o$ | r(kpc) |
|--------|---------------------|-------------|-------|-------|------|------|--------|--------|-----------|--------|
| 22 And | 571 | F2 II | 115.5 | -16.2 | 5.04 | 0.40 | 0.15 | 4.54 | 0.25 | 0.26 |
| | 4362 | G0 Ib | 122.3 | -3.3 | 6.41 | 1.09 | 0.25 | 5.58 | 0.84 | 1.04 |
| | 6474 | G0 Ia | 124.6 | 1.0 | 7.64 | 1.62 | 0.80 | 5.00 | 0.82 | 3.98 |
| φ Cas | 7927 | F0 Ia | 126.7 | -4.4 | 5.00 | 0.68 | 0.41 | 3.64 | 0.27 | 2.68 |
| | 8906 | F3 Ib | 127.6 | -2.5 | 7.13 | 0.74 | 0.41 | 5.77 | 0.33 | 1.19 |
| | 8992 | F6 Ib | 127.9 | -3.7 | 7.78 | | (0.54) | (6.00) | (0.29) | (1.32) |
| | 9250 | G0 Ib | 127.5 | 1.1 | 7.18 | 1.42 | 0.71 | 4.90 | 0.73 | 0.76 |
| | 9366 | K3 Ib | 128.7 | -7.8 | 6.92 | | (0.48) | (5.34) | (1.41) | (0.89) |
| | 10494 | F5 Ia | 129.1 | -0.4 | 7.28 | 1.22 | 0.96 | 4.17 | 0.28 | 2.98 |
| | 11092 | K5 Iab-Ib | 129.2 | 2.7 | 6.57 | 2.08 | 0.50 | 4.98 | 1.60 | 1.50 |
| | 11800 | K5 Ib | 131.1 | -1.6 | 7.79 | 2.04 | 0.37 | 6.57 | 1.67 | 1.56 |
| | 12014 | K0 Ib | 131.6 | -2.5 | 7.72 | 1.97 | 0.58 | 5.80 | 1.39 | 1.10 |
| | 59 ^o 389 | F0 Ib | 131.7 | -1.4 | 9.05 | 1.01 | 1.04: | 5.61: | -0.05: | 1.15: |
| 14662 | F7 Ib | 135.9 | -5.2 | 6.28 | 0.85 | 0.36 | 5.09 | 0.49 | 0.87 | |
| 14 Per | 16901 | G0 Ib | 143.3 | -14.1 | 5.45 | 0.90 | 0.22 | 4.72 | 0.68 | 0.70 |
| | 17378 | A5 Ia | 138.5 | -2.2 | 6.25 | 0.89 | 0.78 | 3.67 | 0.11 | 1.88 |
| η Per | 17506 | K3 Ib | 139.1 | -3.2 | 3.77 | 1.69 | 0.19 | 3.14 | 1.50 | 0.32 |
| | 17958 | K3 Ib | 136.0 | 4.7 | 6.22 | 2.06 | 0.46 | 4.70 | 1.60 | 0.66 |
| | 17971 | F5 Ia | 137.8 | 1.2 | 7.74 | 1.06 | 0.74 | 5.29 | 0.32 | 4.99 |
| | 18391 | G0 Ia | 139.5 | -1.0 | 6.94 | 1.95 | 1.10 | 3.31 | 0.85 | 1.83 |
| | 20123 | G5 II | 144.9 | -5.7 | 5.05 | 1.15 | 0.27 | 4.16 | 0.88 | 0.20 |
| α Per | 20902 | F5 Ib | 146.6 | -5.9 | 1.81 | 0.48 | 0.16 | 1.28 | 0.32 | 0.15 |
| υ Per | 23230 | F5 II | 153.8 | -9.6 | 3.80 | 0.43 | 0.16 | 3.27 | 0.27 | 0.13 |
| | 25030 | K1 Ib | 149.8 | -0.5 | 8.60 | 1.50 | 0.04 | 8.47 | 1.46 | 3.75 |
| | 25056 | G0 Ib | 148.8 | 0.8 | 7.05 | 1.20 | 0.54 | 5.27 | 0.66 | 0.90 |
| | 25291 | F0 II | 145.5 | 5.0 | 5.05 | 0.50 | 0.34 | 3.93 | 0.16 | 0.19 |
| | μ Per | 26630 | G0 Ib | 154.0 | -1.8 | 4.15 | 0.96 | 0.25 | 3.32 | 0.71 |
| | 31118 | K5 Ib | 162.4 | 0.0 | 7.09 | 1.81 | 0.11 | 6.72 | 1.70 | 1.67 |
| ι Aur | 31398 | K3 II | 170.6 | -6.2 | 2.68 | 1.53 | 0.00 | 2.68 | 1.53 | 0.10 |
| β Cam | 31910 | G0 Ib | 149.6 | 11.4 | 4.05 | 0.90 | 0.12 | 3.65 | 0.78 | 0.43 |
| ε Aur | 31964 | F0 Iap | 162.8 | 1.2 | 3.05 | 0.54 | 0.30 | 2.06 | 0.24 | 1.29 |
| | 33299 | K1 Ib | 174.2 | -5.3 | 6.72 | 1.58 | 0.40 | 5.40 | 1.18 | 0.91 |
| α Lep | 36673 | F0 Ib | 220.9 | -25.1 | 2.59 | 0.21 | 0.00 | 2.59 | 0.21 | 0.29 |
| | 36891 | G3 Ib | 169.5 | 4.4 | 6.11 | 1.03 | 0.16 | 5.58 | 0.87 | 1.04 |
| | 37819 | F5 Ib | 179.6 | -0.5 | 8.12 | 0.56 | 0.30 | 7.13 | 0.26 | 2.22 |
| | 38247 | G8 Iab | 188.7 | -5.4 | 6.64 | 1.62 | 0.36 | 5.45 | 1.26 | 2.04 |
| | 38808 | G3 Ib-II | 184.5 | -1.7 | 7.55 | | (0.15) | (7.06) | (0.90) | (1.35) |

TABLE I. Data for the 115 supergiants (continued)

| Name | HD/ BD | MK Class. | l | b | V | B-V | E(B-V) | V _c | (B-V) _o | r |
|---------|--------|--------------|-------|------|------|------|--------|----------------|--------------------|--------|
| | 39416 | G3 Ib-II | 184.2 | -0.5 | 7.52 | | (0.23) | (6.76) | (0.80) | (1.18) |
| | 39949 | G2 Ib | 182.7 | 1.3 | 7.25 | 1.08 | 0.36 | 6.06 | 0.72 | 1.29 |
| | 43282 | G5 Ib-II | 192.0 | 1.2 | 7.74 | | (0.35) | (6.58) | (0.95) | (1.08) |
| | 44033 | K3 Ib | 196.4 | 0.0 | 5.65 | 1.60 | 0.00 | 5.65 | 1.60 | 1.02 |
| | 45829 | K0 Iab | 203.5 | -1.1 | 6.66 | 1.58 | 0.26 | 5.80 | 1.32 | 2.40 |
| 25 Gem | 47731 | G5 Ib | 186.5 | 10.4 | 6.44 | 1.09 | 0.13 | 6.01 | 0.96 | 1.26 |
| ε Gem | 48329 | G8 Ib | 189.6 | 9.6 | 3.00 | 1.41 | 0.14 | 2.54 | 1.27 | 0.26 |
| | 48616 | F5 Ib | 209.4 | -0.1 | 6.85 | 0.79 | 0.48 | 5.26 | 0.31 | 0.94 |
| | 58526 | G3 Ib | 222.0 | 5.0 | 5.99 | 0.92 | 0.11 | 5.62 | 0.81 | 1.06 |
| | 59067 | G8 Ib+B | 227.4 | 2.7 | 5.86 | 0.58 | 0.28 | 4.93 | 0.30 | |
| ζ Mon | 67594 | G2 Ib | 223.1 | 16.5 | 4.37 | 0.96 | 0.02 | 4.30 | 0.94 | 0.58 |
| | 74395 | G2 Ib | 233.3 | 21.0 | 4.64 | 0.84 | 0.17 | 4.08 | 0.67 | 0.52 |
| | 77912 | G8 Ib-II | 184.3 | 42.1 | 4.56 | 1.04 | 0.07 | 4.33 | 0.97 | 0.34 |
| ε Leo | 84441 | G0 II | 206.8 | 48.2 | 2.97 | 0.80 | 0.08 | 2.70 | 0.72 | 0.09 |
| | 148743 | A7 Ib | 8.0 | 26.7 | 6.52 | 0.37 | 0.22 | 5.79 | 0.15 | 1.31 |
| β Dra | 159181 | G2 II | 79.6 | 33.3 | 2.81 | 0.96 | 0.11 | 2.44 | 0.85 | 0.08 |
| | 161796 | F3 Ib | 77.4 | 31.0 | 7.11 | | (0.07) | (6.88) | (0.42) | (1.98) |
| 89 Her | 163506 | F2 Ia | 51.4 | 23.2 | 5.45 | 0.34 | 0.02 | 5.38 | 0.32 | 5.70 |
| 108 Her | 168913 | F9 Ib | 57.4 | 19.2 | 5.62 | 0.21 | 0.15 | 5.12 | 0.06 | 0.88 |
| 45 Dra | 171635 | F7 Ib | 86.2 | 25.0 | 4.78 | 0.61 | 0.10 | 4.45 | 0.51 | 0.65 |
| | 172365 | F9 Ib | 36.3 | 5.1 | 6.36 | 0.79 | 0.24 | 5.57 | 0.55 | 1.08 |
| | 173638 | F2 Ib-II | 23.4 | -3.6 | 5.72 | | (0.46) | (4.20) | (0.09) | (1.38) |
| | 174104 | G0 Ib | 58.6 | 13.4 | 8.37 | | (0.10) | (8.04) | (0.65) | (3.22) |
| | 179784 | G5 Ib | 48.8 | 2.1 | 6.68 | 1.39 | 0.40 | 5.36 | 0.99 | 0.94 |
| | 180028 | F6 Ib | 41.0 | -2.4 | 6.93 | 0.81 | 0.39 | 5.64 | 0.42 | 1.12 |
| | 180583 | F6 Ib-II | 60.6 | 7.5 | 6.06 | 0.63 | 0.28 | 5.13 | 0.35 | 1.18 |
| | 182296 | G3 Ib | 44.3 | -3.1 | 7.02 | 1.32 | 0.47 | 5.47 | 0.85 | 0.99 |
| ū Aql | 182835 | F2 Ib | 37.3 | -7.6 | 4.66 | 0.60 | 0.43 | 3.24 | 0.17 | 0.37 |
| | 183864 | G2 Ib | 59.6 | 3.2 | 7.33 | | (0.56) | (5.48) | (0.66) | (0.99) |
| α Sge | 185758 | G0 II | 54.5 | -2.1 | 4.40 | 0.81 | 0.08 | 4.13 | 0.73 | 0.18 |
| γ Aql | 186791 | K3 II | 48.7 | -7.1 | 2.71 | 1.53 | 0.08 | 2.44 | 1.45 | 0.09 |
| | 187203 | G0 Ib | 49.1 | -7.5 | 6.44 | 0.96 | 0.27 | 5.55 | 0.69 | 1.02 |
| | 187299 | G5 Iab-Ib | 61.5 | -0.3 | 7.14 | 1.60 | 0.66 | 4.96 | 0.94 | 1.24 |
| | 226223 | F6 Ib | 73.7 | 6.0 | 9.16 | | (0.43) | (7.74) | (0.14) | (2.94) |
| | 190113 | G5 Ib | 72.1 | 2.7 | 7.84 | 1.46 | 0.34 | 6.71 | 1.12 | 1.75 |
| | 331777 | F8 Ia | 69.1 | 0.5 | 7.92 | 1.54 | 0.82: | 5.21: | 0.72: | 4.39: |
| | 190323 | G0 Ia-Iab | 54.7 | -8.6 | 6.85 | | (0.22) | (6.12) | (0.65) | (4.85) |
| | 190446 | F6 Ib | 76.2 | 4.9 | 8.22 | 0.57 | 0.14 | 7.75 | 0.43 | 2.95 |

TABLE I. Data for the 115 supergiants (continued)

| Name | HD/BD | Class ^{MK} | l | b | V | B-V | E(B-V) | V _c | (B-V) _o | r |
|----------------|----------------------|---------------------|-------|-------|------|------|--------|----------------|--------------------|--------|
| | 191010 | G3 Ib | 64.2 | -3.5 | 8.17 | | (0.12) | (7.77) | (0.89) | (2.84) |
| | 37 ^o 3827 | F3 Ib | 75.6 | 2.3 | 8.13 | | (0.63) | (6.05) | (0.27) | (1.35) |
| 22 Vul | 192713 | G2 Ib | 63.5 | -6.3 | 5.16 | 1.04 | 0.18 | 4.56 | 0.86 | 0.65 |
| α^1 Cap | 192876 | G3 Ib | 32.3 | -24.2 | 4.25 | 1.07 | 0.24 | 3.45 | 0.83 | 0.39 |
| α^2 Cyg | 192909 | K3 Ib-II | 83.7 | 7.0 | 3.96 | 1.52 | 0.31 | 2.93 | 1.21 | 0.17 |
| 35 Cyg | 193370 | F5 Ib | 73.4 | 0.5 | 5.16 | 0.64 | 0.24 | 4.37 | 0.40 | 0.62 |
| | 193469 | K5 Ib | 76.8 | 1.7 | 6.38 | 1.90 | 0.54 | 4.59 | 1.36 | 0.63 |
| γ Cyg | 194093 | F8 Ib | 78.2 | 1.9 | 2.23 | 0.67 | 0.13 | 1.80 | 0.54 | 0.19 |
| 41 Cyg | 195295 | F5 II | 70.9 | -5.0 | 4.01 | 0.41 | 0.10 | 3.68 | 0.31 | 0.16 |
| 44 Cyg | 195593 | F5 Iab | 76.4 | -1.4 | 6.21 | 1.00 | 0.76 | 3.70 | 0.24 | 1.05 |
| 47 Cyg | 196093/4 | K2 Ib+B | 75.4 | -2.9 | 4.59 | 1.59 | 0.58 | 2.67 | 1.01 | |
| θ Del | 196725 | K3 Ib | 58.0 | -16.6 | 5.66 | 1.53 | 0.10 | 5.33 | 1.43 | 0.88 |
| | 200102 | G1 Ib | 86.1 | -0.7 | 6.63 | 1.06 | 0.22 | 5.90 | 0.84 | 1.20 |
| | 200805 | F5 Ib | 86.8 | -1.2 | 8.31 | | (0.47) | (6.76) | (0.23) | (1.87) |
| ξ Cyg | 200905 | K5 Ib | 86.0 | -2.1 | 3.68 | 1.66 | 0.11 | 3.31 | 1.55 | 0.35 |
| ζ Cyg | 202109 | G8 II | 76.8 | -12.4 | 3.22 | 1.00 | 0.06 | 3.02 | 0.94 | 0.12 |
| | 202314 | G2 Ib | 76.7 | -12.9 | 6.16 | 1.12 | 0.12 | 5.76 | 1.00 | 1.13 |
| | 204022 | G0 Ib | 92.9 | 0.1 | 7.43 | | (0.58) | (5.52) | (0.84) | (1.01) |
| β Aqr | 204867 | G0 Ib | 48.0 | -37.9 | 2.91 | 0.84 | 0.12 | 2.51 | 0.72 | 0.25 |
| | 205349 | K1 Ib | 90.8 | -4.3 | 6.23 | 1.82 | 0.41 | 4.87 | 1.41 | 0.71 |
| | 206312 | K1 II | 93.9 | -2.6 | 7.13 | | (0.15) | (6.64) | (1.08) | (0.62) |
| ϵ Peg | 206778 | K2 Ib | 65.6 | -31.5 | 2.40 | 1.54 | 0.04 | 2.27 | 1.50 | 0.22 |
| 9 Peg | 206859 | G5 Ib | 72.0 | -26.5 | 4.34 | 1.17 | 0.12 | 3.94 | 1.05 | 0.49 |
| 12 Peg | 207089 | K0 Ib | 76.6 | -22.8 | 5.27 | 1.40 | 0.19 | 4.64 | 1.21 | 0.64 |
| | 207489 | F5 Ib | 88.3 | -11.3 | 7.23 | | (0.24) | (6.44) | (0.45) | (1.61) |
| | 207647 | G4 Ib | 95.4 | -3.2 | 7.02 | | (0.36) | (5.83) | (0.78) | (1.16) |
| | 208606 | G8 Ib | 103.5 | 5.5 | 6.12 | 1.60 | 0.49 | 4.50 | 1.11 | 0.63 |
| α Aqr | 209750 | G2 Ib | 59.9 | -42.0 | 2.94 | 0.98 | 0.08 | 2.67 | 0.90 | 0.27 |
| ζ Cep | 210745 | K1 Ib | 103.1 | 1.7 | 3.37 | 1.58 | 0.17 | 2.81 | 1.41 | 0.28 |
| | 216206 | G4 Ib | 104.1 | -7.7 | 6.25 | 1.14 | 0.16 | 5.72 | 0.98 | 1.11 |
| | 216946 | K5 Ib | 104.6 | -7.0 | 4.98 | 1.80 | 0.19 | 4.35 | 1.61 | 0.56 |
| | 217476 | G0 Ia | 108.2 | -2.7 | 5.00 | 1.40 | 0.70 | 2.69 | 0.70 | 1.37 |

TABLE I. Data for the 115 supergiants.

| Name | HD/BD | MK Class | l | b | V | B-V | E(B-V) | v_c | (B-V) _o | r |
|---------|----------------------|-------------|-------|-------|-------|------|--------|--------|--------------------|--------|
| 56 Peg | 218356 | K0 Ibp | 95.1 | -31.7 | 4.76 | 1.33 | 0.19 | 4.13 | 1.14 | 0.51 |
| | 219135 | G0 Ib | 109.6 | -3.8 | 7.62 | | (0.30) | (6.63) | (0.73) | (1.68) |
| | 60 ^o 2532 | F7 Ib | 112.8 | 0.4 | 8.30 | 1.09 | 0.69 | 6.02 | 0.40 | 1.33 |
| | 221861 | K0 Iab | 116.9 | 9.7 | 5.84 | 1.80 | 0.68 | 3.59 | 1.12 | 0.87 |
| 104 Aqr | 222574 | G0 II | 59.4 | -71.4 | 4.80 | 0.82 | 0.05 | 4.63 | 0.77 | 0.22 |
| ψ And | 223047 | G5 Ib | 111.3 | -15.0 | 4.96 | 1.12 | 0.24 | 4.17 | 0.88 | 0.54 |
| ρ Cas | 224014 | F8 Iap | 115.3 | -4.5 | 4.58 | 1.30 | 0.52: | 2.86: | 0.74: | 1.49: |
| | 224165 | G8 Ib | 113.2 | -14.5 | 6.00. | 1.19 | 0.14 | 5.53 | 1.05 | 1.01 |

TABLE II. The dependence of $(B-V)_0$ on spectral type for supergiants.

| Spectral Type | --- "Raw" | Kelsall "Smooth" | --- "Perfect" | FitzGerald (1970) | Johnson (1966) | Schmidt-Kaler (1965) | Fernie (1963) | Kraft & Hiltner (1961) |
|---------------|--------------|---------------------|------------------|----------------------|-------------------|-------------------------|------------------|---------------------------|
| FO | 0.19 | 0.20 | 0.19 | 0.15 | 0.19 | 0.17 | 0.25 | |
| 1 | | 0.22 | 0.22 | 0.16 | | | 0.28 | |
| 2 | 0.21 | 0.25 | 0.25 | 0.18 | 0.25 | 0.25 | 0.31 | |
| 3 | 0.34 | 0.27 | 0.28 | | | | 0.34 | |
| 4 | | 0.29 | 0.30 | | | | 0.37 | |
| 5 | 0.31 | 0.31 | 0.33 | 0.26 | 0.37 | 0.36 | 0.40 | 0.36 |
| 6 | 0.38 | 0.40 | 0.41 | | | | 0.46 | 0.42 |
| 7 | 0.47 | 0.48 | 0.48 | | | 0.48 | 0.52 | 0.48 |
| 8 | 0.64 | 0.57 | 0.56 | 0.55 | 0.55 | 0.58 | 0.61 | 0.56 |
| 9 | | 0.65 | 0.64 | | | | 0.70 | 0.68 |
| GO | 0.74 | 0.73 | 0.72 | 0.82 | 0.70 | 0.67 | 0.77 | 0.72 |
| 1 | | 0.78 | 0.77 | 0.85 | | | 0.86 | |
| 2 | 0.84 | 0.83 | 0.83 | 0.88 | 0.85 | 0.84 | 0.93 | 0.87 |
| 3 | 0.85 | 0.88 | 0.88 | 0.92 | | | 0.96 | 0.94 |
| 4 | 0.91 | 0.93 | 0.92 | | | | 1.00 | 0.98 |
| 5 | 0.99 | 0.98 | 0.97 | 1.00 | 1.01 | 1.08 | 1.03 | 1.00 |
| 6 | | 1.03 | 1.01 | 1.04 | | | | |
| 7 | | 1.08 | 1.05 | 1.10 | | | | |
| 8 | 1.15 | 1.14 | 1.10 | 1.14 | 1.03 | | 1.11 | 1.06 |
| 9 | | 1.19 | 1.14 | 1.16 | | | | |
| KO | 1.24 | 1.24 | 1.19 | 1.18 | 1.12 | | 1.16 | 1.11 |
| 1 | 1.36 | 1.32 | 1.27 | 1.20 | 1.18 | 1.45 | 1.24 | 1.17 |
| 2 | | 1.39 | 1.36 | 1.23 | 1.25 | 1.50 | 1.34 | 1.24 |
| 3 | 1.49 | 1.46 | 1.44 | 1.42 | 1.38 | 1.55 | 1.46 | 1.40 |
| 4 | | 1.53 | 1.53 | 1.50 | 1.49 | | 1.56 | |
| 5 | 1.58 | 1.61 | 1.62 | 1.60 | 1.62 | 1.60 | 1.66 | 1.62 |