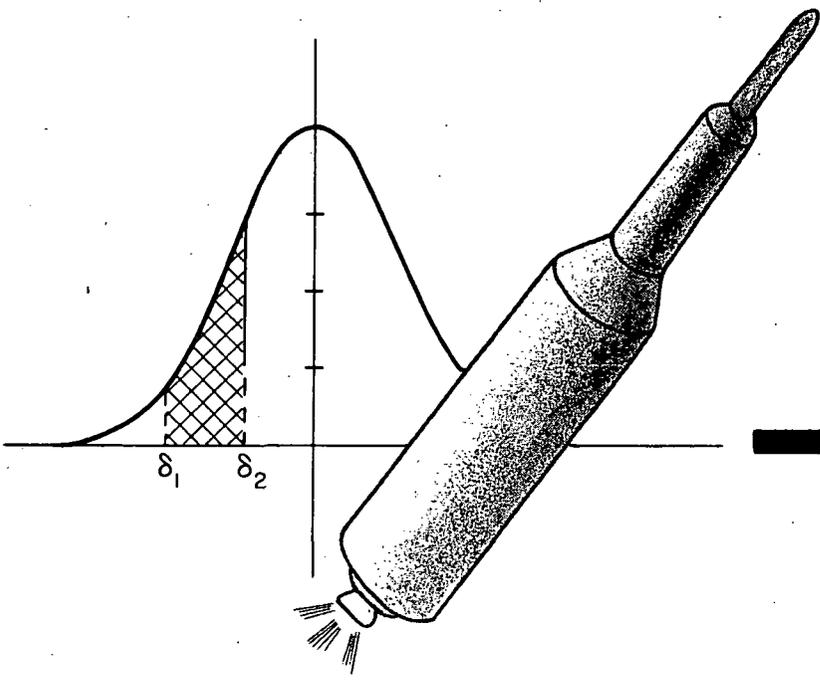


N73-13884  
CR-128655

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A USER'S GUIDE FOR A GENERALIZED  
INTERPLANETARY TRAJECTORY  
GENERATION PROGRAM  
June 17, 1972



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A USER'S GUIDE FOR A GENERALIZED  
INTERPLANETARY TRAJECTORY GENERATION PROGRAM

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Texas Center for Research

3100 Perry Lane

Austin, Texas 78731

A USER'S GUIDE FOR A GENERALIZED INTERPLANETARY  
TRAJECTORY COMPUTATION PROGRAM  
by Texas Center for Research

1.0 SUMMARY

The analysis, structure, and capability of a generalized precision interplanetary trajectory computation program are discussed, with emphasis being placed on the description of input and output. Sample cases showing input and output information are included.

2.0 INTRODUCTION

This report describes a computer program which was developed by the Texas Center for Research of Austin, Texas, for computing precision interplanetary trajectories. The program uses numerical optimization techniques and a variation of Newton's method to produce trajectory iterates which converge to an approximate n-body interplanetary trajectory. This approximate trajectory is then used with a perturbation procedure to produce the precision n-body interplanetary trajectory.

This procedure was first tested on earth-moon trajectories where it was used for a published paper (ref. 1) and several reports (refs. 2, 3).

The procedure developed in this study has been used to compute two-dimensional and three-dimensional earth-moon trajectories and two- and three-dimensional interplanetary n-body trajectories using circular planetary orbits, elliptic planetary orbits, and planetary orbits specified by the JPL Analytical Ephemeris. In addition, trajectories from earth to Venus to a heliocentric radius vector with a  $\Delta v$  at Venus have been computed. The program was then used in an iterative mode to minimize the  $\Delta v$  at Venus in order to simulate the free-fall fly-by. The modifications necessary to provide this capability will be made in the program and an addendum to the User's Manual will be provided.

This program, possibly with some minor modifications, will fill a variety of needs for precision interplanetary trajectories. This report will not cover all possible modifications of the program, but will indicate the types of modifications which could be made and will outline the potential uses of various modified versions of the program.

### 3.0 PROBLEM ANALYSIS

The trajectory computation problem for free-fall trajectories can be stated as follows:

Find the set of initial velocities  $v(t_i)$  such that a particle subject to the equations of motion

$$\begin{aligned}\dot{x} &= v \\ \dot{v} &= F(x,t)\end{aligned}\tag{1}$$

moves from a known initial point  $(x_i, t_i)$  to a known final point  $(x_f, t_f)$ . This is a classical two-point boundary value problem. It is not an optimal control problem as there are no controls in this formulation.

The details of the derivation will not be covered here. Only the important equations and how the technique works will be presented.

#### Primary Convergence Algorithm

Rewrite the equations as

$$\begin{aligned}\dot{x} &= u(t) \\ \dot{\lambda} &= F(x,t)\end{aligned}\tag{2}$$

with initial conditions  $\lambda(t_i) = u(t_i)$ . Clearly, we want to change the system in such a way that  $u(t)$  and  $\lambda(t)$  approach the same values, i.e.,  $u(t)$ . Note that if a  $u(t)$  history is guessed (a control program in the optimal control vocabulary) then a position history,  $x(t)$ , and a history of the variables  $\lambda(t)$  can be obtained.

The procedure used to drive  $u(t)$  and  $\lambda(t)$  to the same values requires that both of equations (2) above be linearized. The new values for  $u(t)$  and  $\lambda(t)$  are given by

$$\begin{aligned} u_{\text{new}}(t) &= u(t) + \delta u(t) \\ \lambda_{\text{new}}(t) &= \lambda(t) + \delta \lambda(t) \end{aligned} \quad (3)$$

Setting  $u_{\text{new}}(t)$  equal to  $\lambda_{\text{new}}(t)$ , one obtains

$$\delta u(t) = \delta \lambda(t) + \lambda(t) - u(t) \quad (4)$$

where  $u(t)$  and  $\lambda(t)$  are evaluated on the previous iterate.

In order to control stepsize and to keep consecutive iterates close enough to each other so that linearization is valid, the scalar variable  $P(0 \leq P \leq 1)$  is introduced into equation (4). Equation (4) is modified in the manner shown below

$$\delta u(t) = \delta \lambda(t) + P[\lambda(t) - u(t)] \quad (5)$$

Linearization of equations (2) leads to

$$\begin{aligned} \delta \dot{x} &= u(t) \\ \delta \dot{\lambda} &= F_x \delta x(t) \end{aligned} \quad (6)$$

where  $F_x$  is a 3 x 3 matrix of partial derivatives of the  $F_x$ 's with respect to the position coordinates. Substitution for  $\delta u$  in equation (6) from equation (5) leads to

$$\begin{aligned} \delta \dot{x} &= \delta \lambda(t) + P[\lambda(t) - u(t)] \\ \delta \dot{\lambda} &= F_x \delta x(t) \end{aligned} \quad (7)$$

It is now assumed that both  $\delta x(t)$  and  $\delta \lambda(t)$  are linear functions of the initial values of  $\delta x_i$  and  $\delta \lambda_i$ . Since  $\delta x_i = 0$ , it is possible to write

$$\begin{aligned}\delta x(t) &= A(t) \delta \lambda_i + M(t) \\ \delta \lambda(t) &= B(t) \delta \lambda_i + N(t)\end{aligned}\quad (8)$$

where  $A$  and  $B$  are  $3 \times 3$  time-dependent matrices and  $M(t)$  and  $N(t)$  are time-dependent 3-vectors. The conditions which define  $A$ ,  $B$ ,  $M$  and  $N$  will be determined below.

From the conditions  $\delta x(t_i) = 0$  and  $\delta \lambda(t_i) = \delta \lambda_i$ , the initial conditions on  $A$ ,  $B$ ,  $M$  and  $N$  can be defined; i.e.,

$$\begin{aligned}A(t_i) &= 0 \\ B(t_i) &= I \quad (\text{the } 3 \times 3 \text{ identity matrix}) \\ M(t_i) &= 0 \\ N(t_i) &= 0\end{aligned}\quad (9)$$

Taking derivatives of equations (8) and substitution of these results plus those of equations (8) into equations (7) leads to the relations

$$\begin{aligned}\dot{A} \delta \lambda_i + \dot{M} &= B \delta \lambda_i + N + P(\lambda - u) \\ \dot{B} \delta \lambda_i + \dot{N} &= F_x A \delta \lambda_i + F_x M\end{aligned}\quad (10)$$

Thus, the conditions defining  $A$ ,  $B$ ,  $M$  and  $N$  are

$$\begin{aligned}\dot{A} &= B & A(t_i) &= 0 \\ \dot{B} &= F_x A & B(t_i) &= I \\ \dot{M} &= N + P(\lambda - u) & M(t_i) &= 0 \\ \dot{N} &= F_x M & N(t_i) &= 0\end{aligned}\quad (11)$$

From these relations,  $A(t)$ ,  $B(t)$ ,  $M(t)$ , and  $N(t)$  can be determined. Then,

$$\delta \lambda_i = A^{-1}(t_f) [\delta x(t_f) - M(t_f)]\quad (12)$$

and

$$\delta\lambda(t) = B(t)A^{-1}(t_f)[\delta x(t_f) - M(t_f)] + N(t) \quad (13)$$

Also,

$$\delta u(t) = B(t)A^{-1}(t_f)[\delta x(t_f) - M(t_f)] + N(t) + P[\lambda(t) - u(t)] \quad (14)$$

The computational procedure is as follows:

1. Guess  $u(t)$
2. Integrate the equations below from  $t_i$  to  $t_f$  with the indicated initial conditions:

$$\begin{aligned} \dot{x} &= u(t) & x(t_i) &= x_i, \text{ given} \\ \dot{\lambda} &= F(x,t) & \lambda(t_i) &= u(t_i) \\ \dot{A} &= B & A(t_i) &= 0 \\ \dot{B} &= F_x A & B(t_i) &= I \\ \dot{M} &= N + P(\lambda - u) & M(t_i) &= 0 \\ \dot{N} &= F_x M & N(t_i) &= 0 \end{aligned} \quad (15)$$

3. Using the specified values for  $x_f$  and the integrated values of  $x(t_f)$ , calculate

$$\delta x_f = x_f - x(t_f) \quad (16)$$

4. Next, calculate

$$\delta\lambda_i = A^{-1}(t_f)[\delta x_f - M_f], \quad (17)$$

and

$$\delta u(t) = B(t)A^{-1}(t_f)[\delta x_f - M_f] + N(t) + P(\lambda(t) - u(t)) \quad (18)$$

5. Now set

$$\begin{aligned} \lambda_{i_{\text{new}}} &= \lambda_i + \delta\lambda_i \\ u_{\text{new}}(t) &= u(t) + \delta u(t) \end{aligned} \quad (19)$$

6. Return to step 2 and repeat the process until  $\lambda(t)$  is sufficiently close to  $u(t)$ .

Once the process described in steps 1 through 6 above converges, the values  $(x(t_i), \lambda_i)$  are used as initial conditions.  $(v(t_i) = \lambda_i)$  in the integration

$$\begin{aligned}\dot{x} &= v & x_i & \text{specified} \\ \dot{v} &= F(x, t) & v(t_i) &= \lambda_i\end{aligned}$$

This integration is employed in a standard perturbation procedure to refine the initial values  $v(t_i)$  (i.e., to refine  $\lambda_i$ ). The perturbation procedure normally converges in three or four iterates. Note that the standard perturbation algorithm is embedded in the computational algorithm being used. As  $\lambda(t) \rightarrow u(t)$ , the term  $P(\lambda - u)$  goes to zero and  $M$  and  $N$  become identically zero. When this occurs, equation (12) becomes the standard equation for  $\delta\lambda_i$  from the perturbation algorithm and equations (13) and (14) become the same.

#### 4.0 PROGRAM STRUCTURE

The n-body interplanetary trajectory program uses four basic groups of subroutines. The functions of the four groups of subroutines, and the individual subroutine names are listed in Table I.

The MAIN Program calls RDATA for input (see Table II for input structure) and then converts the planetocentric equatorial inputs to heliocentric ecliptic coordinates using subroutines COORDS and EPHEMP. Next MAIN calls subroutine DRIVER which controls the program from this point forward.

Subroutine DRIVER sets up the initial conditions for the first iteration and either generates a  $u(t)$  history from a heliocentric trajectory patched to hyperbolic orbits near the planets (subroutine VGUESS) or generates a  $u(t)$  history by integrating a guessed set of initial velocities for the transfer in the n-body solar system employing RK 3(4).

Once the initial  $u(t)$  history is known, DRIVER controls the iteration described under PROGRAM ANALYSIS until  $u(t)$  and  $\lambda(t)$  match sufficiently well to allow the shift to a standard perturbation procedure for final convergence. This iteration process employs RK 3(4) for integration and SPLINE for interpolation.

When the shift is made to the perturbation procedure, the RK 7(8) integrator is employed. Both RK 3(4) and RK 7(8) obtain derivatives from subroutine DERIV.

Slightly different derivatives are called for by each procedure, so DERIV contains a branch. Subroutine DERIV uses subroutine EPHEMP and SPLINE to aid in the generation of the required derivatives.

## 5.0 PROGRAM CAPABILITY

The program possesses four options which may be used as described below. The option chosen is indicated by the input parameter NØPT (NØPT=1,2,3, or 4). The options and their uses are described below.

### Option 1 (NØPT=1)

When operating in this mode, the program finds the n-body trajectory which goes from  $(t_i, x_i)$  to  $(t_f, x_f)$  where  $x_i$  is the specified initial position vector,  $t_i$  is the initial time,  $x_f$  is the specified final position vector, and  $t_f$  is the final time.

The initial and final position vectors are normally input in kilometers in planetocentric equatorial coordinates with  $x_i$  being measured in a system centered at planet NP(1) and  $x_f$  measured in a system centered at planet NP(2). The planetary indices used are Mercury=1, Venus=2, Earth=3, etc. If it is desired to input either  $x_i$  or  $x_f$  in heliocentric ecliptic coordinates in units of AU, the corresponding planetary index, NP(1) or NP(2), should be set to 10.

For this option, the program converts the initial and final position vectors into heliocentric ecliptic coordinates

(if necessary), and then obtains an initial velocity history from a Lambert solution to start the first order iteration process.

For  $N\emptyset PT=1$  runs, the required input data is:

TP(1)	$t_i$ ,	the initial time (J.D.)
TP(2)	$t_f$ ,	the final time (J.D.)
TP(3)		unspecified
XI(1)		the initial position vector in km. if
XI(2)	$x_i$ ,	$NP(1) \leq 9$ in AU if $NP(1) \geq 10$
XI(3)		
XF(1)		the final position vector in km if
XF(2)	$x_f$ ,	$NP(1) \leq 9$ in AU if $NP(1) \geq 10$
XF(3)		
VI(1)		unspecified
VI(2)		
VI(3)		
VF(1)		unspecified
VF(2)		
VF(3)		
NP(1)		initial planet index
NP(2)		final planet index
$N\emptyset PT$		integer, =1
NPRT		integer, print key
		NPRT=0, do not print last iterate
		NPRT=1, print every $NSPRT^{th}$ integra-
		tion steps on last iterate
		(converged iterate).
		NPRT=-1, print every $NSPRT^{th}$
		tion on every iterate.
NSPRT		integer, see NPRT

### Options 2 and 3 ( $N\emptyset PT=2/N\emptyset PT=3$ )

These options are available in order to allow a run which has been terminated prematurely (due to time limit, for example) to be restarted at the last iterate computed. If the program was still operating in the first order iteration mode when the termination occurred,  $N\emptyset PT=2$  should be used. If, however, the shift to the perturbation method has already occurred,  $N\emptyset PT=3$  should be used.

In either case, the initial position vector  $x_i$  (XI) and the initial velocity vector  $v_i$  (VI) will be in heliocentric ecliptic coordinates in units of AU and AU/year. These values should be used directly as input for the NØPT=2/NØPT=3 options. For these options  $t_i$  TP(1) and  $t_f$  TP(2) as in NØPT=1.

For NØPT=2/NØPT=3 runs, the required input data is:

TP(1)	$t_i$ ,	initial time
TP(2)	$t_f$ ,	final time
TP(3)		unspecified
XI(1)		initial position vector in AU
XI(2)	$x_i$ ,	
XI(3)		
XF(1)		final position vector in AU
XF(2)	$x_f$ ,	
XF(3)		
VI(1)		initial velocity vector in AU/yr.
VI(2)	$v_i$ ,	
VI(3)		
VF(1)		unspecified
VF(2)		
VF(3)		
NP(1)		unspecified (can be specified but will not be used)
NP(2)		
NØPT		integer, =2 or 3
NPRT		as in NØPT=1 option
NSPRT		

---

#### Option 4 (NØPT=4)

The purpose of this option is to enable the user to compute a trajectory which goes from  $(x_i, t_i)$  through a point  $(x_m, t_m)$  to a point  $(x_f, t_f)$  with a velocity impulse,  $\Delta V$ , at  $(x_m, t_m)$  and then, by allowing the point  $x_m$  to vary, make the  $\Delta V$  at  $(x_m, t_m)$  as small as possible. This option uses information obtained from two NØPT=1 runs as input data.

In order to obtain the information required for input for the NØPT=4 option, the user must set up an NØPT=1 run from  $(x_i, t_i)$  to  $(x_m, t_m)$  and get the converged initial position

and velocity components at  $t_i$  (in heliocentric ecliptic coordinates). Then he must set up an NØPT=1 run from  $(x_m, t_m)$  to  $(x_f, t_f)$  and get the converged final position and velocity components at  $t_f$  (in heliocentric ecliptic coordinates). These values are then used as XI, VI, XF, and VF (all 3-vectors) in the input for the NØPT=4 run. The times are input in the following manner:  $t_i$  TP(1),  $t_m$  TP(2),  $t_f$  TP(3). The quantities NP(1) and NP(2) do not need to be specified for an NØPT=4 run.

For NØPT=4 runs, the required input is:

TP(1)	$t_i$ ,	initial time (J.D.)
TP(2)	$t_m$ ,	fly-by time (J.D.)
TP(3)	$t_f$ ,	final time (J.D.)
XI(1)		initial position vector in AU (from
XI(2)	$x_i$ ,	NØPT=1 run)
XI(3)		
XF(1)		final position vector in AU (from
XF(2)	$x_f$ ,	NØPT=1 run)
XF(3)		
VI(1)		initial velocity vector in AU/yr. (from
VI(2)	$v_i$ ,	NØPT=1 run)
VI(3)		
VF(1)		final velocity vector in AU/yr. (from
VF(2)	$v_f$ ,	NØPT=1 run)
VF(3)		
NP(1)		unspecified
NP(2)		
NØPT		integer, =4
NPRT		as in NØPT=1
NSPRT		

TABLE IBasic Subroutine Groupings  
for the  
Interplanetary Trajectory Generation Program

1. Control Routines
  - MAIN
  - DRIVER
2. Numerical Integration Related Routines
  - RK34CØ
  - RK34
  - RK78CØ
  - RK78
  - DERIV
  - VGUESS
  - SPLINE
3. Celestial Mechanics Routines
  - CØNIC
  - ØETØRC
  - EPHEMP
  - TRFM
  - CØØRDS
  - EULER
  - DLMBRT
  - UVFNS

---

4. Vector Manipulation Routines
  - UNIT
  - DØT
  - VMAG
  - CRØSS

TABLE II

## Description of Input

<u>Variables</u>	<u>Description</u>
TP(1),TP(2), TP(3) Units (Julian days) Double Precision	Times in Julian days as required for the various problem options
XI(1),XI(2),XI(3) Units (km or AU) Double Precision	Initial position vector for NØPT=1, XI is expressed in planetocentric equatorial coordinates at planet NP(1) in kilometers for NØPT=2,3, or 4, XI is expressed in heliocentric ecliptic coordinates in A.U.
XF(1),XF(2),XF(3) Units (km or AU) Double Precision	Final position vector (see XI for details on units for various values of NØPT)
VI(1),VI(2),VI(3) Units (AU/yr) Double Precision	Initial velocity vector in heliocentric ecliptic coordinates in AU/yr. (Needed only for NØPT=2,3, or 4)
VF(1),VF(2),VF(3) Units (AU/yr) Double Precision	Final velocity vector in heliocentric ecliptic coordinates in AU/yr. (Needed only for NØPT=4)
NP(1),NP(2) Integers	Planet index denoting departure planet and target planet in system where Mercury=1, Venus=2, etc.
NØPT Integer	Option choice key, NØPT=1,2,3, or 4.
NPRT Integer	Print option key NPRT=-1 Print every NSPRT <sup>th</sup> integration step on every iterate NPRT=0 Print only summary information--do not print at intervals on any iterate NPRT=1 Print every NSPRT <sup>th</sup> integration step on the last (converged) iterate only
NSPRT Integer	The number of integration steps between prints if NPRT≠0. (see NPRT)

For more detail on input, see Section 5.0, PROGRAM CAPABILITY.

The data is input using the system routine RDATA. The call to RDATA is

```
CALL RDATA(TP,XI,XF,VI,VF,NP,NØPT,NPRT,NSPRT)
```

TABLE III

## Description of Output

<u>Output Quantity</u>	<u>Description</u>
ITERATION (Integer)	A counter which shows the current number of iterations used during any one mode of operation
FINAL TIME Units (years)	Elapsed time during the current iteration from the initial time
INTEGRATION STEPS (Integer)	Total number of steps taken by the integrator to complete the current iteration
XI Units (AU)	Initial position in heliocentric ecliptic coordinates (3 components)
XF Units (AU)	Final integrated position in heliocentric ecliptic coordinates (3 components)
VI Units (AU/yr)	Initial velocity (3 components)
VF Units (AU/yr)	Final integrated velocity (3 components)
DELXF Units (AU)	Terminal position miss defined as difference between input final position and integrated final position (3 components)
DUMAX Units (AU/yr)	Maximum change in the control velocity history for the next iteration of this mode 1 operation. This quantity is meaningless during modes 2 and 3 operation (3 components)
DELVI Units (AU/yr)	Change in the initial velocity for the next iteration (3 components)
MINIMUM APPROACH DISTANCES Units (AU)	The minimum radial distance obtained any-time during the current iteration from each of the solar system members listed
FINAL APPROACH DISTANCES Units (AU)	The final radial distance obtained during the current iteration from each of the solar system members listed

```

.....
X1 = -.982479456704+00 .....
X2 = .344286337007+00 .....
.....

```

FRACTION OF PLANETS GRAVITY USED = .010

```

ITERATION = 1      FINAL TIME =      .4790231348+00      INTEGRATION STEPS = 219
X1 = -.982479456704+00      -.182063401066+00      .941741311249-05
X2 = .344286337007+00      .634003659411+00      -.108192418573+01
X3 = .861511009512+00      -.562470706469+01      .109712549592+00
X4 = -.666332658713+01      .443790322097+01      -.958361371637-01
DELTA = .257090412503-07      -.132054308971-06      .259500398473+08
OJMAA = .488595990390-01      .407232377975-01      .546607941087-02
DELVI = .488595990390-01      .316451212251-01      -.548607941087-02

```

TABLE IV  
Sample Output

MINIMUM APPROACH DISTANCES

```

SUN      .691628671811+00
MERCUR   .359984324528+00
VENUS    .427313531536-04
EARTH    .200878015617-04
MARS     .794359201761+00
JUPITER  .459760460781+01

```

FINAL APPROACH DISTANCES

```

SUN      .721533629147+00
MERCUR   .118489543796+01
VENUS    .427313531536-04
EARTH    .924910014122+00
MARS     .219428371700+01
JUPITER  .5711313463052+01

```

FRACTION OF PLANETS GRAVITY USED = .010

```

ITERATION = 2      FINAL TIME =      .4790231348+00      INTEGRATION STEPS = 219
X1 = -.982479456704+00 .....
X2 = .344286337007+00 .....
.....

```

AF	•344266905746+00	•634002133016+00	-•168192166972-01
VI	•910373604551+00	-•559306194347+01	•104226470181+00
VF	-•658879353002+01	•448506633274+01	-•104642342968+00
DELTA	-•543029690999-06	•139433672918-05	-•285851609302-07
DUMAX	•575862282764-01	•452127513964-01	•650172116413-02
DELVI	•575862282764-01	•358478002525-01	-•658172118413-02

MINIMUM APPROACH DISTANCES

SUN	•691593971291+00
MERCUR	•359690333084+00
VENUS	•430164926567-04
EARTH	•228237195686-04
MARS	•794859201761+00
JUPITE	•459760480781+01

FINAL APPROACH DISTANCES

SUN	•721532558835+00
MERCUR	•114489432960+01
VENUS	•438164926567-04
EARTH	•924908537766+00
MARS	•219428246500+01
JUPITE	•571313546031+01

FRACTION OF PLANETS GRAVITY USED = .010

ITERATION#	3	FINAL TIME=	•4790231348+00	INTEGRATION	STEPS =	207
AI	-•982879456704+00		-•182063401026+00			•941741311249-05
AF	•344287135202+00		•6340000664213+00			-•108191856112-01
VI	•967956836827+00		-•555721414322+			•976447469973-01
VF	-•650915456123+01		•45348806150401			-•112689316691+00
DELTA	-•772485212640-06		•26631240576405			-•536511133812-07
DUMAX	•323319960530-02		•266572296601-02			•376956261976-03
DELVI	•301477449449-02		•181466442192-02			-•350419229487-03



TABLE V

## Sample Data Set for NØPT=4 Case

The following data set will be assumed to be punched on cards started in column 1 with no spaces unless indicated.

Card 1--TP

D1=2.4443299842D+06,2.4445049474D+06,2.4446221842D+06

Card 2--XI

D2=-9.824794567024D-01,-1.820634010750D-01,9.417413113088D-06

Card 3--XF

D3=2.610368798355D-01,-4.199897439846D-01,-3.679846437010D-02

Card 4--VI

D4=2.691275379500D+00,-4.440086578800D+00,-1.322957073700D-01

Card 5--VF

D5=8.7038421339D+00,4.30945631490D+00,-6.446741114700D-01

Card 6--NP

I6=10,10

Card 7--NØPT

I7=4

Card 8--NPRT

I8=1

Card 9--NSPRT

I9=10\$

Note that the dollar sign (\$) should follow immediately the last digit of data in each data set in RDATA input.