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Computational Methods for Analyzing the Transmission Characteristics of a Beta Particle Magnetic Analysis System

Jag J. Singh

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Computational Methods for Analyzing the Transmission Characteristics of a Beta Particle Magnetic Analysis System

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NASA

National Aeronautics
and Space Administration

**Scientific and Technical
Information Branch**

1979

SUMMARY

Computational methods have been developed to study the trajectories of beta particles (positrons) through a magnetic analysis system as a function of the spatial distribution of the radionuclides in the beta source, size, and shape of the source collimator, and the strength of the analyzer magnetic field. On the basis of these methods, the particle flux, their energy spectrum, and source-to-target transit times have been calculated for Na^{22} positrons as a function of the analyzer magnetic field and the size and location of the target. These data will be useful in studies requiring parallel beams of positrons of uniform energy -- such as measurement of the moisture distribution in composite materials.

INTRODUCTION

Recent studies (refs. 1 to 4) have shown that positron lifetime in composite materials is a linear function of the moisture content of the material up to the saturation level. This linear dependence of positron lifetime on moisture content is now being utilized for determining the moisture distribution profile in the polymeric targets. Moisture distribution profile is an important factor in determining the environmental degradation of composites. It is believed (ref. 5) that nonuniform moisture distribution leads to higher internal fatigue of these materials. Ability to measure moisture profile nondestructively is of interest since continual absorption and desorption of moisture from composites in service usually leaves the moisture profile highly nonuniform inside the test specimen.

In the previously completed studies (ref. 4), moisture profiling has been accomplished by using a Na^{22} positron source in conjunction with thin steel beam energy degraders. This process is necessarily gross since the positron spectrum is continuous in energy. It would be preferable to use positron beams of well-defined energy to determine moisture content at the end of their range in the test specimen. One practical technique of doing so would require the use of a magnetic analyzer for selecting positrons of predetermined energy to probe the moisture content at different depths in the test specimen. This paper describes analytical methods for maximizing the number of energy-analyzed positrons arriving at the target, while minimizing the spread in their energy and arrival time. Use of magnetically analyzed positrons would be of practical interest only if the spread in their arrival time at the target is less than the time resolution of the lifetime measurement system (refs. 1 and 2) and their energy resolution is better than their range definition in the target. The computer programs needed to calculate various experimental parameters for optimizing the experimental methods have been developed by Gerald H. Mall and are included as an appendix to this paper.

SYMBOLS

| | |
|---------------|---|
| c | velocity of light |
| D | distance of center of the magnetic field circle from center of source |
| D_C | diameter of collimator |
| E_{β^+} | kinetic energy of positron |
| $F(Z,W)$ | Fermi function for positrons of energy W |
| L_C | length of collimator |
| l | particle's path length to target |
| m | slope of the tangent to particle's path at point of exit from magnetic field |
| m_0 | particle's rest mass |
| m^* | particle's relativistic mass |
| $N(W) dW$ | number of positrons in the energy range from W to $W + dW$ |
| P | transition matrix element for positron decay |
| R | radius of target |
| R_C | radius of collimator |
| R_F | radius of magnetic field region |
| R_p | radius of particle's circular path in the magnetic field |
| v | velocity of particle |
| W | total positron energy in units of m_0c^2 |
| W_0 | maximum value of W , corresponding to end-point beta-spectrum energy |
| x,y | coordinates of particle; the origin of the x-y coordinate system lies at the center of the source spot, with the x-axis passing through the center of the magnetic field region |
| Y_C | distance of center of magnetic field from target |
| θ | angle of particle emission from source, relative to x-axis |

| | |
|----------|---|
| τ | transit time of particle |
| τ_0 | universal time constant |
| ϕ | angle subtended by particle's circular path in magnetic field at circular path center |

PROBLEM STATEMENT

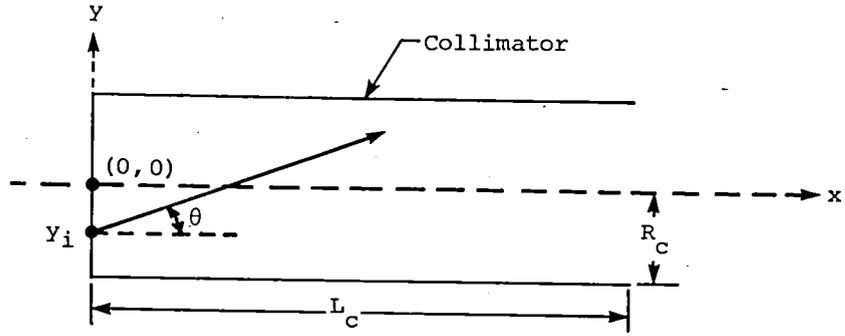
A well-collimated beam of Na^{22} positrons is allowed to enter a variable magnetic field. The purpose of this analysis is to follow the trajectories of positrons that escape the collimator and enter the magnetic field. The trajectory of each positron is followed until it arrives at the target located at different distances from the center of the magnetic field set to transmit preselected positron energies. Targets of sizes ranging in diameter from 1 cm to 2.54 cm, located at a number of distances from the center of the magnetic field, are selected for measuring the positron beam. Since the positron emission from the Na^{22} source nuclei is spherically symmetrical, the calculations are considered in a two-dimensional geometry, with the results generalized for a three-dimensional geometry. From the known energy spectrum of Na^{22} positrons (ref. 6), total transit times from the source to the target are calculated for each positron arriving at the target. Also, the energy spectrum of the positrons received at the target location for various settings of the magnetic field is determined to assess the definition of the range of the arriving positrons. Finally, the trajectories of all the positrons that are transmitted through preselected magnetic fields are followed for a distance of 20.0 cm from the center of the magnetic field. From these data, a distance at which the positron energy dispersion and the transit time spread are the smallest and the positron number density is the highest is determined for optimal location of the target.

THEORETICAL CALCULATIONS

For the purpose of developing an analytical basis for the computational methods, it is assumed that the positron source is deposited at the bottom of a long narrow cylinder. The open end of the cylinder serves to collimate the positron beam entering the magnetic field. The source profile is assumed to be one of the three practically attainable configurations illustrated in figure 1. The geometrical dimensions of the source collimator considered are shown in figure 2. Figure 3 shows a schematic diagram of the experimental system.

Particle's Trajectory Through Source Collimator

As depicted in the following sketch, the particle is emitted at point $(0, y_i)$ at an angle θ relative to the x-axis with a velocity v :



The path of the particle is described by

$$x = vt \cos \theta \quad (1a)$$

$$y = vt \sin \theta + y_i \quad (1b)$$

Solving equation (1a) for the product vt corresponding to $x = L_c$ yields

$$vt|_{x=L_c} = \frac{L_c}{\cos \theta} \quad (2)$$

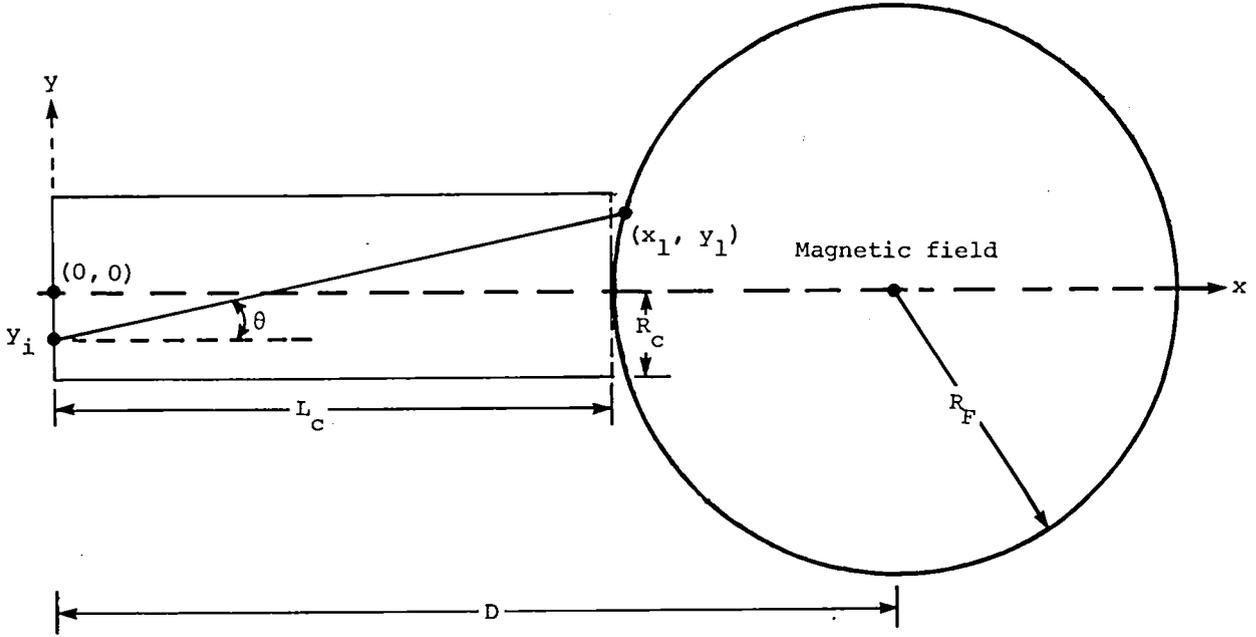
Substituting equation (2) into equation (1b) leads to

$$y_E = y|_{x=L_c} = L_c \tan \theta + y_i \quad (3)$$

If $y_E \geq R_c$ or if $y_E \leq -R_c$, the particle strikes the wall of the collimator and is not considered further. If $-R_c < y_E < R_c$, the intersection of the trajectory with the circular magnetic field is calculated.

First Intersection of Particle's Trajectory With Magnetic Field

In the following sketch, the particle's trajectory is represented by the line from $(0, y_i)$ to (x_1, y_1) , and the magnetic field by the circle:



The equations of the line and the circle are given by

$$y = x \tan \theta + y_i \quad (4a)$$

$$\left[x - (R_F + L_C) \right]^2 + y^2 = R_F^2 \quad (4b)$$

Substituting y from equation (4a) into equation (4b) yields

$$\left[x_1 - (R_F + L_C) \right]^2 + (x_1 \tan \theta + y_i)^2 = R_F^2 \quad (5)$$

Expanding equation (5) and making the substitution $R_F + L_C = D$ leads to

$$x_1^2(1 + \tan^2 \theta) + x_1(2y_i \tan \theta - 2D) + (D^2 + y_i^2 - R_F^2) = 0 \quad (6)$$

Equation (6) may be solved for x_1 :

$$x_1 = \frac{(2D - 2y_i \tan \theta) - \sqrt{(2y_i \tan \theta - 2D)^2 - 4(1 + \tan^2 \theta)(D^2 + y_i^2 - R_F^2)}}{2(1 + \tan^2 \theta)} \quad (7)$$

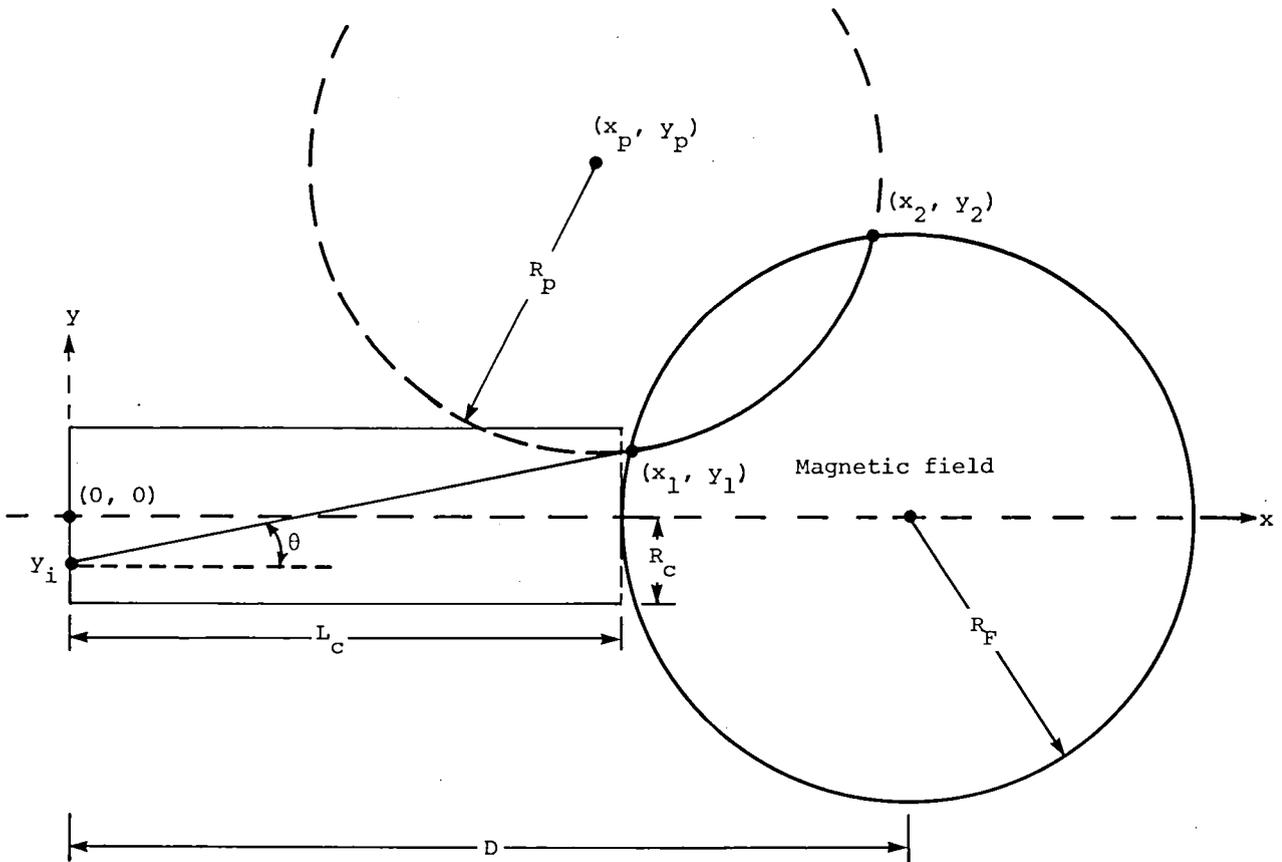
where the smaller root was selected by subtracting the square root. Expanding equation (7), collecting terms, and reducing the results leads to

$$x_1 = \frac{(D - y_i \tan \theta) - \sqrt{R_F^2 \sec^2 \theta - (y_i + D \tan \theta)^2}}{\sec^2 \theta} \quad (8)$$

Equation (8) gives the x-coordinate of the desired intersection point and the y-coordinate may be obtained by substituting equation (8) into equation (4a).

Second Intersection of Particle's Trajectory With Magnetic Field

Upon entering the magnetic field (assumed to be perpendicular to the plane of the paper), the particle experiences forces causing it to move in a circular path.



The forces experienced by the particle satisfy

$$qvH = m^*v^2/R_p \quad (9)$$

where q is the charge of the particle, m^* is the particle's mass at v , and H is the magnitude of the magnetic field. The radius of the particle's circular path is then given by

$$R_p = m^*v/qH \quad (10)$$

and the tangent to the path at the point (x_1, y_1) is the slope of the initial trajectory, $\tan \theta$. The equations of the two circles are given by

$$(x - D)^2 + y^2 = R_F^2 \quad (11a)$$

$$(x - x_p)^2 + (y - y_p)^2 = R_p^2 \quad (11b)$$

Solving equation (11a) for y and substituting into equation (11b) produces the following equation for x_2 :

$$(x_2 - x_p)^2 + \left[\pm \sqrt{R_F^2 - (x_2 - D)^2} - y_p \right]^2 = R_p^2 \quad (12)$$

Expanding equation (12) leads to

$$x_2^2 - 2x_2x_p + x_p^2 + R_F^2 - (x_2 - D)^2 \mp 2y_p \sqrt{R_F^2 - (x_2 - D)^2} + y_p^2 = R_p^2 \quad (13)$$

By making the substitutions $K = 2D - 2x_p$, $L = x_p^2 + y_p^2 + R_F^2 - R_p^2 - D^2$, equation (13) can be expressed as

$$Kx_2 + L = \pm 2y_p \sqrt{R_F^2 - (x_2 - D)^2} \quad (14)$$

Squaring both sides of equation (14) and collecting terms leads to the following expression for x_2 :

$$Sx_2^2 + Tx_2 + U = 0 \quad (15)$$

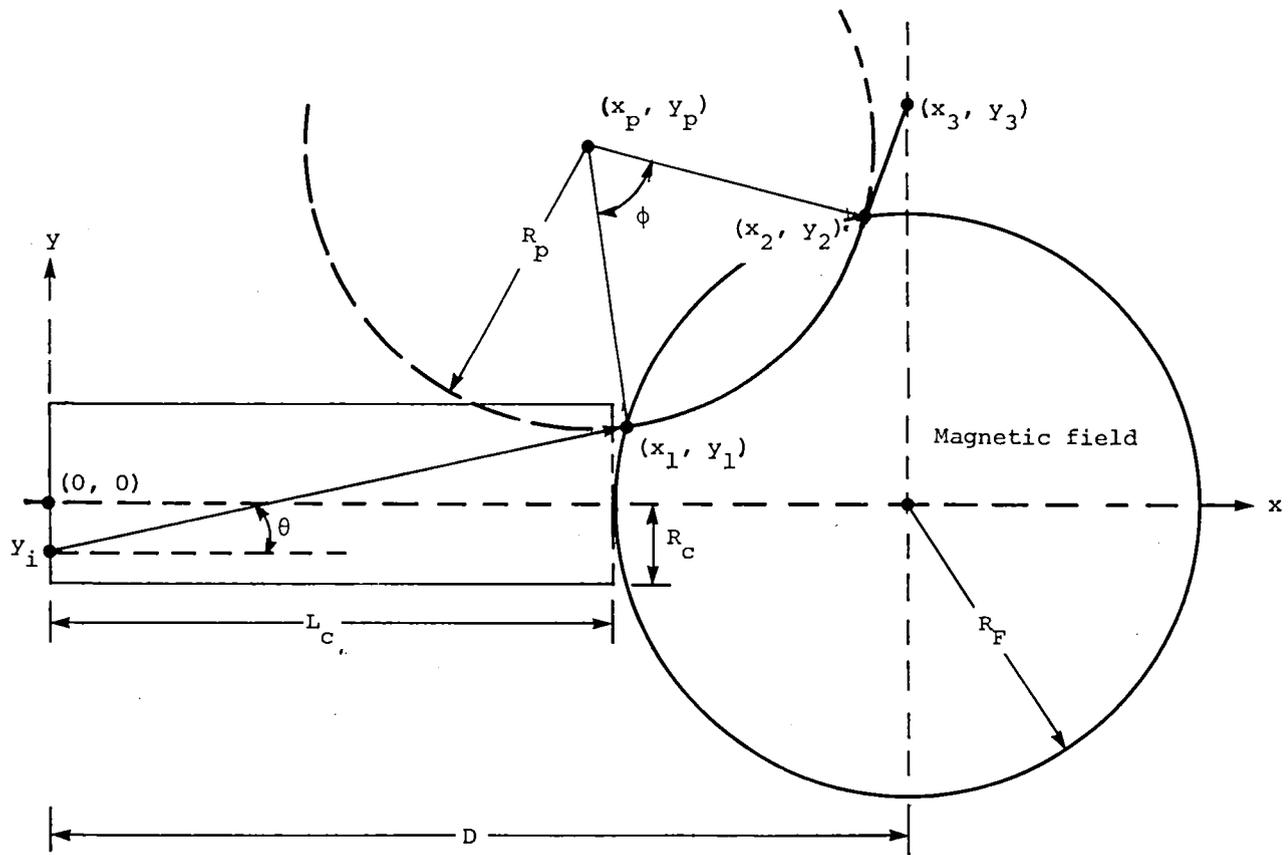
where $S = K^2 + 4y_p^2$, $T = 2KL - 8y_p^2D$, and $U = L^2 + 4y_p^2(D^2 - R_F^2)$. Then, x_2 is given the following equation where the positive sign of the square root has been selected:

$$x_2 = \frac{-T + \sqrt{T^2 - 4SU}}{2S} \tag{16}$$

The corresponding value for y_2 may be determined by substituting equation (16) into equation (11a).

Intersection of Particle's Trajectory With Line $x = L_C + R_F$

The following sketch illustrates the particle's trajectory after leaving the magnetic field:



Equation (11b) may be solved for y and differentiated with respect to x to find the slope of the tangent to the particle's circular path at (x_2, y_2) . The result is

$$m = \frac{x_2 - x_p}{\sqrt{R_p^2 - (x_2 - x_p)^2}} \quad (17)$$

The point of intersection of the particle's trajectory with the line $x = L_C + R_F$ is then given by

$$x_3 = L_C + R_F \quad (18a)$$

$$y_3 = y_2 - m(x_2 - L_C - R_F) \quad (18b)$$

Determination of Particle's Transit Time

The particle trajectory along the circular path traverses an angle ϕ , given by

$$\phi = \cos^{-1} \left[\frac{2R_p^2 - (x_2 - x_1)^2 - (y_2 - y_1)^2}{2R_p^2} \right] \quad (19)$$

The total path length is then given by the following expression:

$$\ell = R_p \phi + \sqrt{(x_1 - x_i)^2 + (y_1 - y_i)^2} + \sqrt{(x_3 - x_2)^2 + (y_3 - y_2)^2} \quad (20)$$

For a particle with kinetic energy T , the relativistic mass is given by

$$m^* = \frac{T + m_0 c^2}{c^2} \quad (21)$$

where m_0 is the rest mass and c is the velocity of light. The particle's velocity is given by the following equation:

$$v = \sqrt{1 - (m_0/m^*)^2} c \quad (22)$$

Equations (20) and (22) can be combined to yield the transit time:

$$\tau = \ell/v \quad (23)$$

Using equations developed in these sections the various parameters of interest have been calculated for a source, collimator, and magnetic field geometry of the type shown in figures 1 and 4, 2, and 3, respectively. The results are discussed in the following section.

DISCUSSION OF TYPICAL RESULTS

Trajectories of positrons for three different source configurations, transmitted through the magnetic analysis system set for 200 keV, are shown in figures 5 to 7. It is clear that none of the source configurations considered lead to actual positron focusing at distances up to 20 cm from the center of the magnetic field. The flux densities for the flat and convex sources were approximately equal and considerably higher than that for the concave source. Also the convex source provided slightly better energy and time resolution than the other configurations. It was therefore decided to use the convex source configuration (fig. 4) for detailed calculations of the transmitted positron beam parameters. Some typical results for the spectra of positrons from a convex source arriving at several target locations for two different size targets are shown in figures 8 to 10 and are summarized in tables I to III. These results include the effects of Na²² source deposit profile as well as the positron Coulomb correction factor (refs. 7 to 9). The number of positrons emitted at the source has been taken as 1000 in each case.

It is clear from the tables that for a 1-cm-diameter target located a distance from the magnetic field center $Y_C \geq 12$ cm, the dispersion in the arriving positron energy is less than 10% of the energy for which the magnetic field is set for all cases; that is, $\Delta E/E < 10\%$. Furthermore, the spread in transit time $\Delta\tau$ for positrons arriving at the 1-cm-diameter target is less than 100 ps for $Y_C \geq 12$ cm. This transit time spread is less than the time resolution of the lifetime measurement system. For a 2.54-cm-diameter target, $\Delta E/E \approx 10\%$ and $\Delta\tau < 100$ ps for $Y_C \geq 20$ cm for all magnetic field settings. Thus a target to magnetic field center distance of approximately 20 cm appears to be acceptable for both the 1-cm and the 2.54-cm targets. This conclusion is also corroborated by the data shown in figure 11. This figure shows the trajectories of all the positrons transmitted to the target through the magnetic analyzer when the latter is set for $E_{\beta+} = 300$ keV for 1-cm and 2.54-cm targets. (Similar results are obtained when the magnetic field is set for other positron energies.)

From the foregoing discussion, it is seen that for the 1-cm target, target to magnetic field center distances greater than 12 cm would be acceptable, though 12 cm would yield the highest number density at the target. For the 2.54-cm target, the minimum acceptable distance is 20 cm.

CONCLUSIONS

A computer program for optimizing a positron magnetic analysis system for studying moisture profile in composites has been developed. Three types of source configurations were considered. It appears that a convex source deposit

provides simultaneously the smallest energy and transit time spread and the largest positron flux density at the target locations up to 20 cm. For such a source, the energy dispersion in the positrons arriving at a 1-cm-diameter target located at or over 12 cm from the magnetic field center is less than 10% for all settings of the analyzing magnetic field. The transit time spread in the positrons striking the target under these conditions is less than 100 ps, which is less than the time resolution of the positron lifetime measurement system. An examination of the trajectories of the positrons from a convex source transmitted through the magnetic analyzer for several field settings confirms that a separation distance of approximately 12 cm also corresponds to the highest positron density in the transmitted beam. For a 2.54-cm-diameter target, the optimum distance would be 20 cm for the smallest energy and transit time spread and the largest positron number density at the target.

Langley Research Center
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Hampton, VA 23665
November 20, 1979

APPENDIX

COMPUTER PROGRAMS FOR CALCULATING POSITRON TRANSMISSION THROUGH A MAGNETIC ANALYSIS SYSTEM

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PROGRAM FOR CALCULATING ENERGY AND TRANSIT TIMES OF POSITRONS

Program Description

The computer program POSTRN is written in FORTRAN IV language for the Control Data CYBER 170 series digital computer system with network operating system (NOS) 1.2. The program including interactive graphics routines requires 60700 octal locations of core storage. A typical case in which 5000 positron trajectories are computed requires approximately 20 CPU seconds on the CYBER 173. Although the program has been designed to be used interactively, it can also be executed in a batch mode.

The program may be executed in three distinct ways. When the target location, target size, magnetic field setting, and desired number of trajectories are specified, the program uses a random number generator and two input probability distribution tables to determine the initial conditions of each positron. Programs have been developed to generate these tables giving probability versus energy and probability versus positron position within the collimator, and these programs are described subsequently in this appendix. When the program is executed in this manner, the output consists of several parameters describing those positrons which strike the target. Additionally, two disk files (TAPE1 and TAPE2) are generated which contain sufficient data to reconstruct the trajectory of each positron. TAPE1 contains data for positrons which miss the target and TAPE2 contains data for positrons which hit the target. The second mode of operation of the program requires as input one of the disk files TAPE1 or TAPE2. When this option is selected, program output consists of a plot of the trajectories of a specified number of positrons which have trajectory data on the disk file. The third mode of operation is similar to the second except that the user must provide trajectory data for each positron which is to be plotted. These data include the magnetic field setting and the energy, position within the collimator, and emission angle of each positron.

Description of FORTRAN Variables

The following list contains a description of the significant FORTRAN variables appearing in the program. The dimension for each array is beside the variable in parentheses. Each variable is also identified as I, input variable; P, program variable; or O, output variable. Variables which may be input or program variables, depending on the mode of operation, are identified by I/P.

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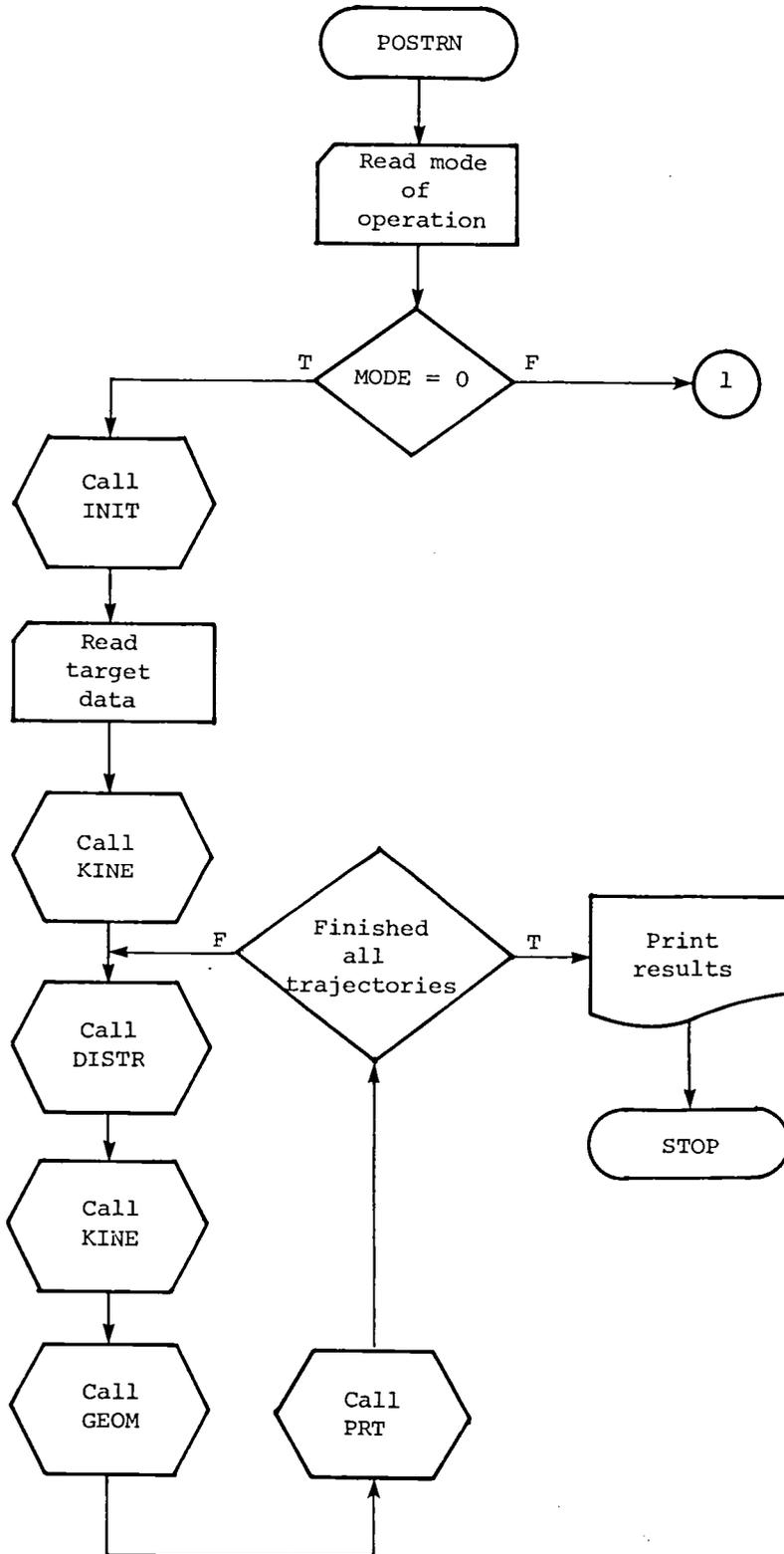
| <u>FORTRAN variable</u> | <u>Type</u> | <u>Description</u> |
|-------------------------|-------------|---|
| COLL | P | Collimator length |
| COLR | P | Collimator radius |
| EFLD | I | Energy for which the magnetic field is set |
| EMAX | O | Maximum energy of those positrons striking the target |
| EMIN | O | Minimum energy of those positrons striking the target |
| ENER | I/P | Positron energy |
| MODE | I | Mode of program operation |
| NHIT | O | Number of positrons which strike the target |
| NMISS | O | Number of positrons which miss the target |
| NTRY | I | Number of trajectories to be computed |
| RC | I | Target radius |
| RM | P | Magnetic field radius |
| THETA | I/P | Angle of emitted positron |
| TMAX | O | Maximum transit time of those positrons striking the target |
| TMIN | O | Minimum transit time of those positrons striking target |
| VEL | P | Positron velocity |
| X(1000) | P | x-coordinate of points to be plotted |
| Y(1000) | P | y-coordinate of points to be plotted |
| YC | I | Distance of target from center of magnetic field |
| YP | I/P | Position of emitted positron within the collimator |

Flow Charts and Listings of Program

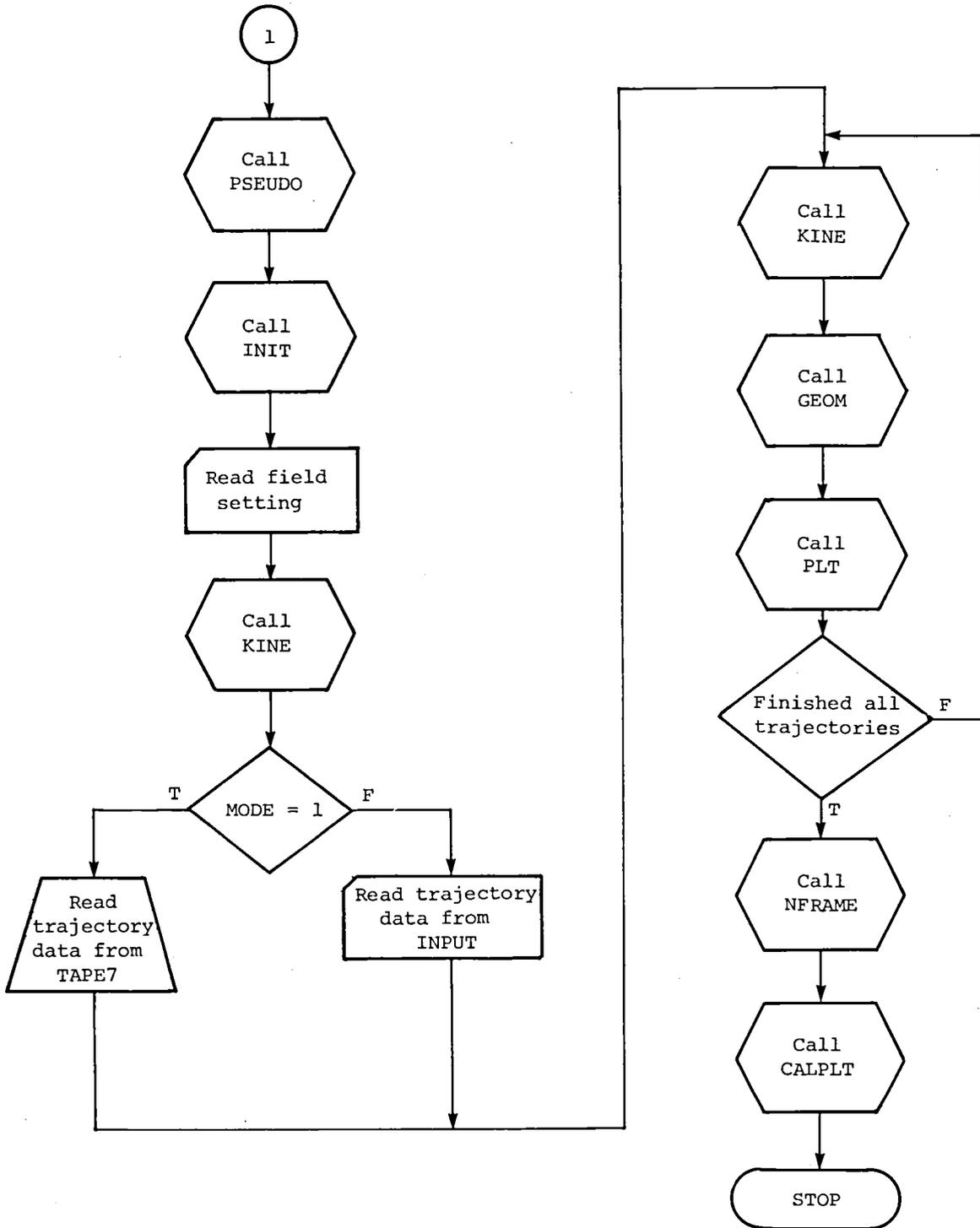
Program POSTRN.— POSTRN, the main program, performs all input and output operations, controls flow through the kinematics and geometry subroutines, and calls those subroutines which load the probability tables or generate graphic

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output depending on the specified mode of program operation. The flow chart and listing of POSTRN follow.



APPENDIX



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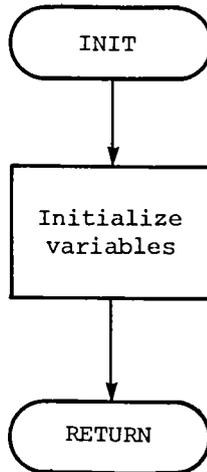
```

1      PROGRAM POSTRN(OUTPUT,INPUT,TAPE5=INPUT,TAPE6=OUTPUT,
      1 TAPE1,TAPE2,TAPE3,TAPE4,TAPE7)
      DIMENSION X(1000),Y(1000)
      COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,CLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
      2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
      DIMENSION EEE(100),YYY(100),TTT(100)
      COMMON/JTB/NFR,JREQ,IBAUD,JCCS,IJD,TFAC,IJTB(4)
10     NFR = 0
      JREQ = 0
      IBAUD = 120
      JCCS = 0
      TFAC = 0.016
15     READ(5,*) MODE
      IF(MODE.NE.0) GO TO 30
      CALL INIT
      READ(5,*) EFLD,YC,RC,NTRY
      IJK = 1
      ENER = EFLD
20     CALL KINE
      DO 20 I=1,NTRY
      CALL DISTR
      IJK = I + 1
      CALL KINE
25     CALL GEOM
      IF(X(1).LT.0.) GO TO 10
      CALL PRT
      GO TO 20
10 NMISS = NMISS + 1
30     CONTINUE
      WRITE(6,1) NHIT,NMISS,NTRY,EMIN,EMAX,TMIN,TMAX
      1 FORMAT(1H1,I10,* HITS*,I10,* MISSES*,I10,* TRYS*/
      1 1X,*MINIMUM ENERGY = *,E16.8,5X,*MAXIMUM ENERGY = *,E16.8/
35     2 1X,*MINIMUM TIME = *,E16.8,5X,*MAXIMUM TIME = *,E16.8)
      STOP
30     CONTINUE
      CALL PSEUDO
      CALL INIT
      READ(5,*) NTRY,EFLD
40     ENER = EFLD
      IJK = 1
      CALL KINE
      DO 40 I=1,NTRY
      IF(MODE.EQ.1) READ(7) EEE(I),YYY(I),TTT(I)
45     IF(MODE.EQ.2) READ(5,*) EEE(I),YYY(I),TTT(I)
30     CONTINUE
      DO 50 I=1,NTRY
      IJK = I + 1
      ENER = EEE(I)
50     YP = YYY(I)
      THETA = TTT(I)
      CALL KINE
      CALL GEOM
      IF(X(1).LT.0.) GO TO 50
55     CALL PLT
30     CONTINUE
      CALL NFRAME
      CALL CALPLT(0.,0.,999)
      STOP
60     END

```

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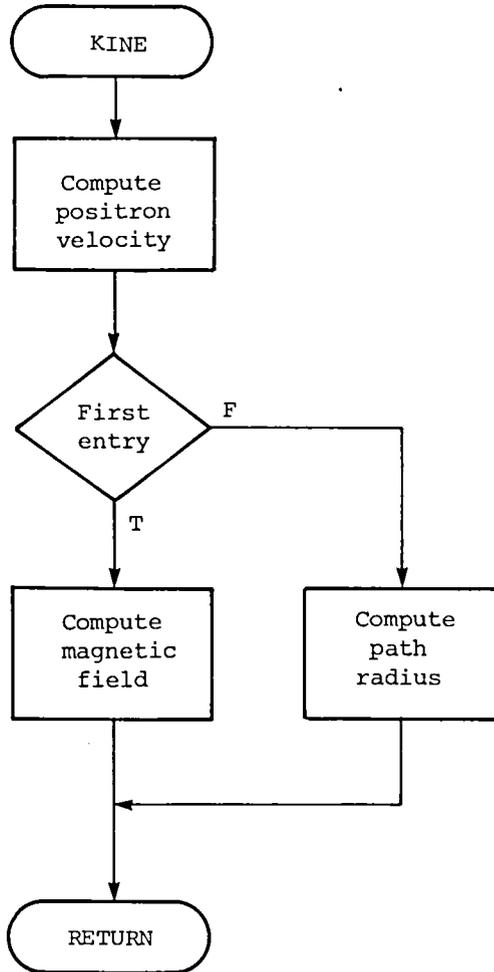
Subroutine INIT.— Subroutine INIT initializes the dimensions of the collimator and the magnetic field. Several variables used in the output summary are also initialized.



```
1      SUBROUTINE INIT
      DIMENSION X(1000),Y(1000)
      COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,COLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
      2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
      EMIN = 600.
      EMAX = 0.
      TMIN = 100.
      TMAX = 0.
10     NMISS = 0
      NHIT = 0
      RM = 3.48*2.54
      RP = RM
      COLL = 3.*2.54
15     COLR = 0.152*2.54/2.
      RETURN
      END
```

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Subroutine KINE.- Subroutine KINE computes the velocity of a relativistic positron of known energy. On the first call, the magnetic field strength required to focus positrons of the specified energy is determined. On subsequent calls, the radius of each positron's trajectory through the computed magnetic field is calculated.

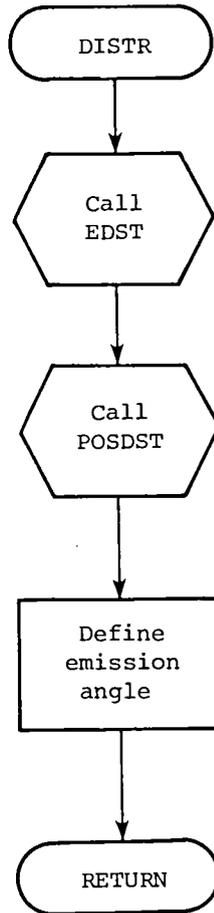


```

1      SUBROUTINE KINE
        DIMENSION X(1000),Y(1000)
        COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,CLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
        2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
        T = 1.6E-16*ENER
        XMAS = (T + 9.E+16*9.11E-31)/9.E+16
        TEMP = (9.11E-31/XMAS)**2
        TEMP = SQRT(1. - TEMP)
10     VEL = TEMP*3.E+08
        HR = XMAS*VEL/1.6E-19
        HR = 1.E+06*HR
        XMAS = 1000.*XMAS
        VEL = 100.*VEL
15     IF(IJK.EQ.1) H = HR/RP
        IF(IJK.NE.1) RP = HR/H
        RETURN
        END
  
```

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Subroutine DISTR.- Subroutine DISTR is the executive routine for assigning an energy and position within the collimator to each positron subject to user-defined probability distributions. An emission angle is also assigned subject to a uniform distribution $-\tan^{-1} (2 \cdot \text{COLR}/\text{COLL}) \leq \theta \leq \tan^{-1} (2 \cdot \text{COLR}/\text{COLL})$.

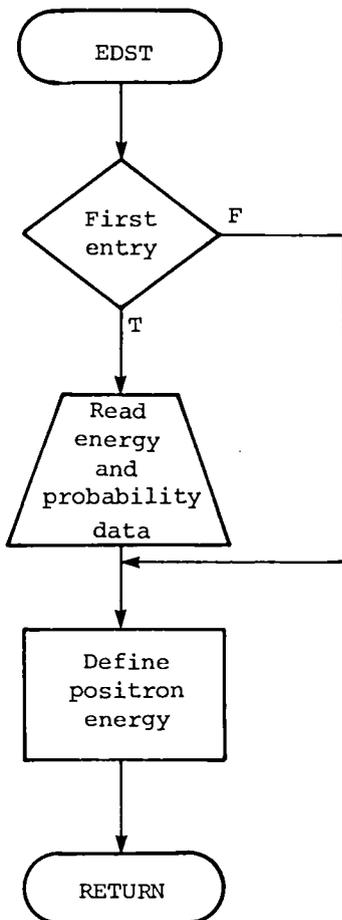


```

1      SUBROUTINE DISTR
        DIMENSION X(1000),Y(1000)
        COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,COLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
        2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
        CALL EDST
        CALL POSDST
        PI = 0.0506145
        TMX = -PI
10     THETA = TMX + 2.*PI*RANF(DUM)
        RETURN
        END
  
```

APPENDIX

Subroutine EDST.- Subroutine EDST assigns a random energy to each positron subject to a user-specified probability distribution. On the first entry the subroutine reads from TAPE3 an array of energies and a corresponding array of normalized cumulative probabilities. A program which generates these arrays is described subsequently in this appendix. A uniform random number between 0 and 1 is used to interpolate linearly and select an energy. On subsequent entries, the read operation is bypassed.



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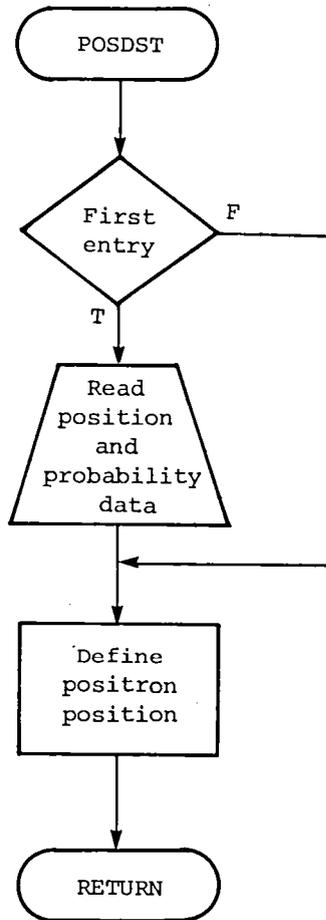
```

1      SUBROUTINE EDST
        DIMENSION X(1000),Y(1000)
        COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,COLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
        2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
        DIMENSION XX(300),YY(300)
        DATA NIN/0/
        IF(NIN.NE.0) GO TO 30
        K = 1
10     NIN = NIN + 1
        REWIND 3
        10 CONTINUE
        READ(3,1) XX(K),YY(K)
        1 FORMAT(2E16.8)
15     IF(EOF(3).NE.0) GO TO 20
        K = K + 1
        GO TO 10
20     K = K - 1
30     KM1 = K - 1
40     T = RANF(DUM)
        IF(T.LT.YY(1)) GO TO 40
        DO 50 I=1,KM1
        IF(T.GE.YY(I).AND.T.LE.YY(I+1)) GO TO 60
50     CONTINUE
        I = KM1
25     60 ENER = XX(I) + (XX(I+1)-XX(I))*(T - YY(I))/
        1 (YY(I+1) - YY(I))
        RETURN
        END

```

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Subroutine POSDST.- Subroutine POSDST assigns a random position within the collimator to each positron subject to a user-specified probability distribution. On the first entry the subroutine reads from TAPE4 an array of positions and a corresponding array of normalized cumulative probabilities. A program which generates these arrays is described subsequently in this appendix. A uniform random number between 0 and 1 is used to interpolate linearly and select a position. On subsequent entries, the read operation is bypassed.



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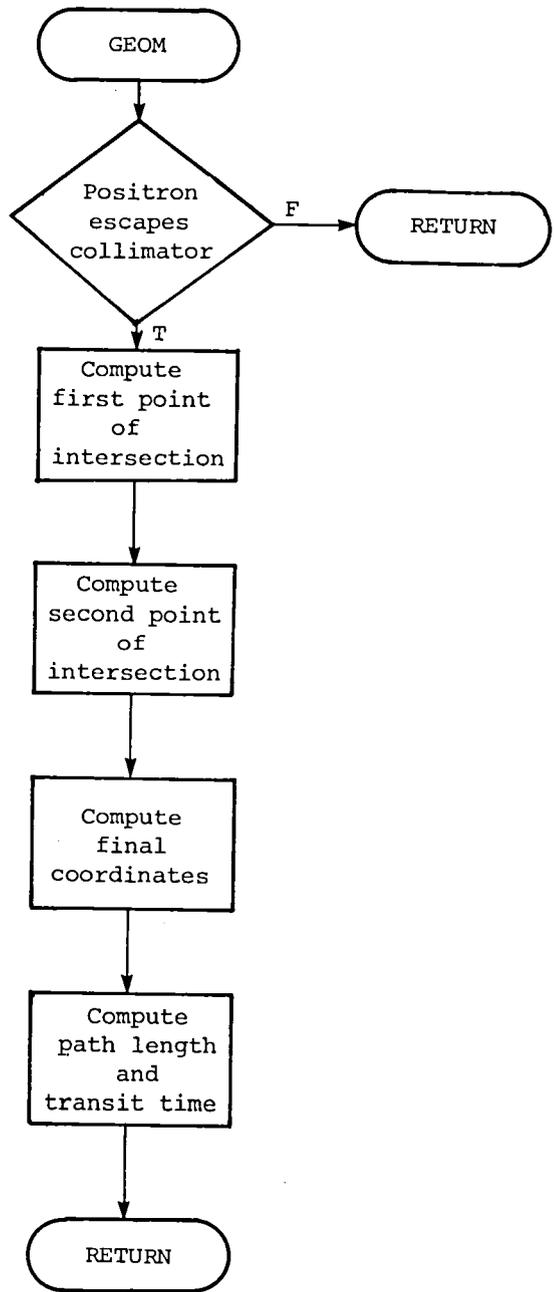
```

1      SUBROUTINE POSDST
      DIMENSION X(1000),Y(1000)
      COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,COLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
      2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
      DIMENSION XX(300),YY(300)
      DATA NIN/0/
      IF(NIN.NE.0) GO TO 30
      K = 1
10     NIN = NIN + 1
      REWIND 4
10     CONTINUE
      READ(4,1) XX(K),YY(K)
      1 FORMAT(2E16.8)
15     IF(EOF(4).NE.0) GO TO 20
      K = K + 1
      GO TO 10
20     K = K - 1
30     KM1 = K - 1
40     T = RANF(DUM)
      IF(T.LT.YY(1)) GO TO 40
      DO 50 I=1,KM1
      IF(T.GE.YY(I).AND.T.LE.YY(I+1)) GO TO 60
50     CONTINUE
      I = KM1
25     60 YP = XX(I) + (XX(I+1)-XX(I))*(T - YY(I))/
      1 (YY(I+1) - YY(I))
      RETURN
      END

```

APPENDIX

Subroutine GEOM.- Subroutine GEOM determines the path of each positron. For each particle which escapes the collimator, the first point of intersection with the magnetic field is calculated. The positron is then assumed to follow a circular path until it exits from the field at the second point of intersection of the path and magnetic field. From this point the particle travels in a straight line until it reaches a distance from the center of the field equal to the distance at which the target was placed. The coordinates of this final point are calculated as well as the total path length and transit time.



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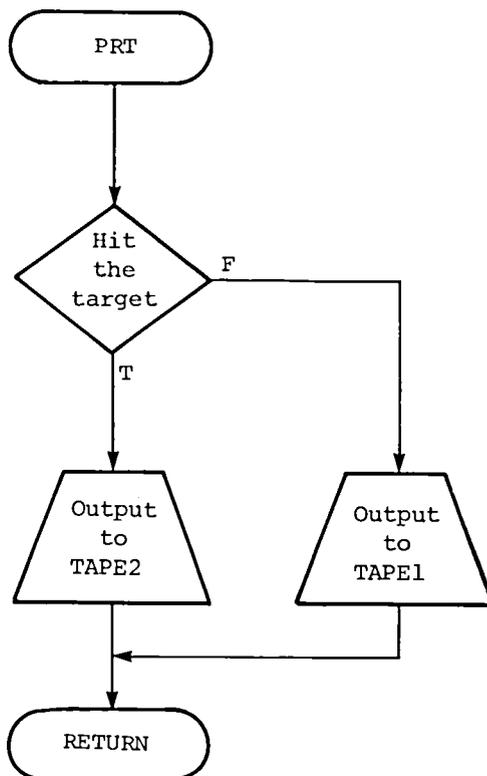
```

1      SUBROUTINE GEOM
      DIMENSION X(1000),Y(1000)
      COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,COLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
      2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
      IF(MODE.NE.0) THETA = THETA/57.29577951
      IF(MODE.NE.0) YC = 20.
      YE = COLL*TAN(THETA) + YP
10     X(1) = -1.
      IF(ABS(YE).GT.COLR) RETURN
      X(1) = 0.
      D = RM + COLL
      SECT = 1. + TAN(THETA)**2
      RAD = RM*RM*SECT - (YP + D*TAN(THETA))**2
15     X1 = (D - YP*TAN(THETA) - SQRT(RAD))/SECT
      Y1 = X1*TAN(THETA) + YP
      XCM = D
      YCM = 0.
      XCP = X1 - RP*SIN(THETA)
20     YCP = Y1 + RP*COS(THETA)
      ZK = 2.*(D - XCP)
      ZL = XCP**2 + YCP**2 + RM**2 - RP**2 - D**2
      ZS = ZK**2 + 4.*YCP**2
      ZT = 2.*ZK*ZL - 8.*D*YCP**2
25     ZU = ZL**2 + 4.*YCP**2*(D**2 - RM**2)
      X2 = (-ZT + SQRT(ZT**2 - 4.*ZS*ZU))/(2.*ZS)
      Y2 = SQRT(RM**2 - (X2 - D)**2)
      DEL = RP**2 - (X2 - XCP)**2
      IF(DEL.LT.0.00000001) GO TO 10
      SLOPE = (XCP - X2)/SQRT(DEL)
      IF(Y2.LT.YCP) SLOPE = -SLOPE
      B = Y2 - SLOPE*X2
      Y3 = YC
      X3 = (Y3 - B)/SLOPE
35     GO TO 20
10    X3 = X2
      Y3 = YC
20    CONTINUE
      DIST = (X2 - X1)**2 + (Y2 - Y1)**2
40     PHI = ACOS((2.*RP**2 - DIST)/(2.*RP**2))
      PL = PHI*RP
      PL = PL + SQRT(X1**2 + (Y1 - YP)**2)
      PL = PL + SQRT((X3 - X2)**2 + (Y3 - Y2)**2)
      TIME = PL/VEL
45     RETURN
      END

```

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Subroutine PRT.- Subroutine PRT generates two disk files which may be used as input to the program when operating in a different mode. If the final coordinates of the positron are within the target area, the energy, initial position within the collimator, initial path angle, transit time, and distance from the center of the target are output to TAPE2. For those particles which miss the target, the information is output to TAPE1.

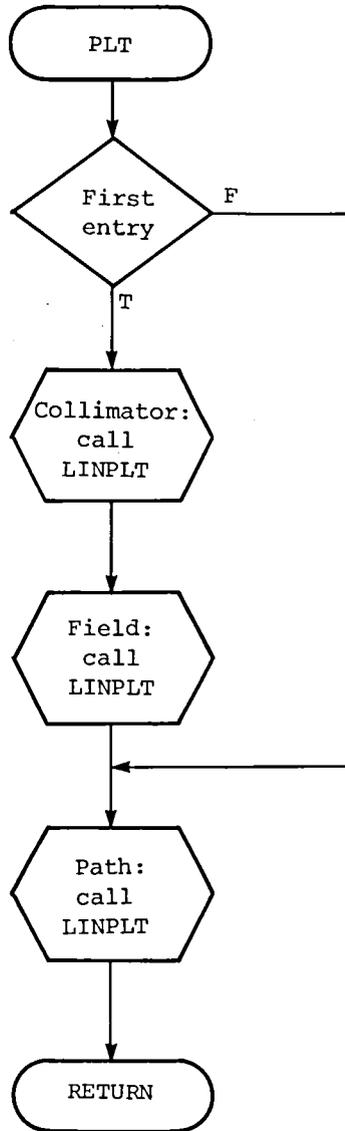


```

1      SUBROUTINE PRT
        DIMENSION X(1000),Y(1000)
        COMMON X,Y,ENER,XMAS,VEL,H,RM,RP,COLL,COLR,YP,THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
        2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
        DX = X3 - COLL - RM
        IF(ABS(DX).LT.RC) GO TO 10
C      A MISS
          THET = THETA*57.295777951
10      WRITE(1) ENER,YP,THET,TIME,DX
          NMISS = NMISS + 1
          GO TO 20
        10 CONTINUE
C      A HIT
15      THET = THETA*57.295777951
          NHIT = NHIT + 1
          IF(ENER.LT.EMIN) EMIN = ENER
          IF(ENER.GT.EMAX) EMAX = ENER
          IF(TIME.LT.TMIN) TMIN = TIME
20      IF(TIME.GT.TMAX) TMAX = TIME
          WRITE(2) ENER,YP,THET,TIME,DX
20      RETURN
        END
  
```

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Subroutine PLT.- Subroutine PLT produces graphic output. The plot includes the collimator, the magnetic field, and the trajectory of each positron.



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```

1      SUBROUTINE PLT
      DIMENSION X(1000),Y(1000)
      COMMON X,Y,ENER,XMAS,VEL,H, RM,RP, COLL, COLR, YP, THETA,
5      1 XP,X1,Y1,X2,Y2,XCM,YCM,XCP,YCP,X3,Y3,PL,TIME,IJK,
      2 EMIN,EMAX,NHIT,YC,RC,EFLD,TMIN,TMAX,NMISS,MODE
      DATA JIK/0/
      JIK = JIK + 1
      IF(JIK.GT.1) GO TO 30
10     XMIN = 0.
      YMIN = -10.
      XSF = 3.
      YSF = 3.
      X(1) = 0.
      Y(1) = -COLR
15     X(2) = 0.
      Y(2) = COLR
      X(3) = COLL
      Y(3) = COLR
      X(4) = COLL
20     Y(4) = -COLR
      X(5) = 0.
      Y(5) = -COLR
      X(6) = XMIN
      Y(6) = YMIN
25     X(7) = XSF
      Y(7) = YSF
      CALL LINPLT(X,Y,5,1,0,0,0,0)
      XX = COLL
      DX = 0.05
30     IK = 2
      X(1) = XX
      Y(1) = 0.
      XF = COLL + 2.*RM
10 CONTINUE
35     XX = XX + DX
      IF(XX.GT.XF) XX = XF
      YY = SQRT(RM**2 - (XX - XCM)**2)
      X(IK) = XX
      Y(IK) = YY
40     IK = IK + 1
      IF(XX.NE.XF) GO TO 10
      XX = XF
20 CONTINUE
45     XX = XX - DX
      IF(XX.LT.COLL) XX = COLL
      YY = -SQRT(RM**2 - (XX - XCM)**2)
      X(IK) = XX
      Y(IK) = YY
      IK = IK + 1
50     IF(XX.NE.COLL) GO TO 20
      IK = IK - 1
      X(IK+1) = XMIN
      Y(IK+1) = YMIN
      X(IK+2) = XSF
55     Y(IK+2) = YSF
      CALL LINPLT(X,Y,IK,1,0,0,0,0)

```

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```
30 CONTINUE
   X(1) = XP
   Y(1) = YP
60   X(2) = X1
   Y(2) = Y1
   IK = 3
   YY = Y1
40 CONTINUE
   YY = YY + DX
   IF(YY.GT.Y2) YY = Y2
   XX = SQRT(RP**2 - (YY - YCP)**2) + XCP
   X(IK) = XX
   Y(IK) = YY
70   IK = IK + 1
   IF(YY.NE.Y2) GO TO 40
   X(IK) = X3
   Y(IK) = Y3
75   X(IK+1) = XMIN
   Y(IK+1) = YMIN
   X(IK+2) = XSF
   Y(IK+2) = YSF
   CALL LINPLT(X,Y,IK,1,0,0,0,0)
   RETURN
80   END
```

Other subroutines.- The subroutines PSEUDO, NFRAME, CALPLT, and LINPLT are part of the NASA LaRC graphics output library.

Program Usage

Input

All program input is accomplished using FORTRAN list directed reads. The data may appear anywhere in the field and when more than one item is specified, the data items are separated by commas. The first item expected by the program is the variable MODE. Allowed values are 0, 1, and 2. When MODE = 0, the program randomly generates positrons, determines their path through the magnetic field, counts the number of particles which strike the target, and, for those which do strike the target, determines the minimum and maximum energies and transit times. When MODE = 1 or MODE = 2 the program produces graphic output. For MODE = 1, positron trajectories from an earlier MODE = 0 run are plotted. For MODE = 2, the user must specify the initial conditions of the positrons to be plotted.

Additional data are a function of the selected mode of operation. When MODE = 0, the program expects four additional values, separated by commas. The first value is the energy, in kilo-electron-volts, for which focusing is desired. The program uses this to determine the strength of the magnetic field required to focus any positron of this energy on the center of the target. The second value is the distance, in centimeters, from the center of the magnetic field at which the target is to be placed. The third value is the radius of

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the target, in centimeters. The final value is the number of positrons to be generated. The following is an example of acceptable data for MODE = 0:

300.,16.,0.5,5000

For MODE = 1, the user must transfer the data on TAPE1 (misses) or TAPE2 (hits) from a MODE = 0 run to TAPE7. The program then expects two additional data values. The first value is the number of paths to plot. The second value is the energy, in kilo-electron-volts, for which focusing is desired. An example follows:

4,300.

For MODE = 2, the program expects the same two data items as for MODE = 1. Additionally, for each path to be plotted the user must provide three items: the positron energy, in kilo-electron-volts; the initial y-coordinate of the positron within the collimator, in centimeters ($-\text{COLR} \leq \text{YP} \leq \text{COLR}$); and the initial path angle, in degrees ($-2.9^\circ \leq \theta \leq 2.9^\circ$). Sample MODE = 2 input follows:

3,300.

250.,0.,-1.

350.,0.,1.

300.,0.,0.

Output

Program output also varies with the mode of operation. For MODE = 0, printed output in the following form is generated:

| | | | |
|------------------|---------------|------------------|---------------|
| 196 HITS | 4804 MISSES | 5000 TRYS | |
| MINIMUM ENERGY = | .28077858E+03 | MAXIMUM ENERGY = | .32205860E+03 |
| MINIMUM TIME = | .12036213E-08 | MAXIMUM TIME = | .12665825E-08 |

The energies are in kilo-electron-volts and the times in seconds. For MODE \neq 0, no printed output is produced, but plots are drawn showing the collimator, the magnetic field, and the positron paths as illustrated in figure 12.

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PROGRAMS FOR GENERATING ENERGY AND POSITION DISTRIBUTIONS

Three programs have been written which generate the probability distribution tables required by program POSTRN. Listings of these programs, GENE, GENP, and CUMU, are included in this section of the appendix.

Program GENE tabulates the positron spectrum given by the following expression (ref. 10) at intervals of 2.5 keV:

$$N(W) dW = \frac{|P|^2}{\tau_0} F(Z,W) (W^2 - 1)^{1/2} (W_0 - W)^2 W dW$$

where $F(Z,W)$ for Na^{22} has been calculated using the numerical tables given in reference 11. Values of $F(Z,W)$ for 25 energy values are given in the GENE program listing. The resulting spectrum is illustrated in figure 13.

Program GENP tabulates the probability of a positron being emitted as a function of position within the collimator. It is assumed that the source is shaped like a segment of a sphere with a radius equal to the collimator radius and a height equal to 25% of the collimator diameter. Results are tabulated at 100 equally spaced intervals. The resulting distribution is illustrated in figure 14.

Program CUMU takes as input the output of either program GENE or program GENP. Output from CUMU consists of arrays of normalized cumulative probability distributions versus energy and position. Output from CUMU is supplied on TAPE3 and TAPE4 as input to program POSTRN.

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```

1      PROGRAM GENE(OUTPUT,TAPE6=OUTPUT,INPUT,TAPE5=INPUT)
      DIMENSION E(25),F(25),X(300),Y(300)
      COMMON/JTB/NFR,JREQ,IBAUD,JCCS,IJO,TFAC,IJTB(4)
5      NFR = 0
      JREQ = 0
      IBAUD = 120
      JCCS = 0
      TFAC = 0.016
10     E(1) = 10.1
      E(2) = 15.7
      E(3) = 22.5
      E(4) = 30.4
      E(5) = 39.4
      E(6) = 49.3
15     E(7) = 60.3
      E(8) = 72.2
      E(9) = 84.9
      E(10) = 98.4
      E(11) = 112.7
20     E(12) = 127.7
      E(13) = 143.4
      E(14) = 159.6
      E(15) = 176.5
      E(16) = 193.8
25     E(17) = 211.6
      E(18) = 248.6
      E(19) = 287.2
      E(20) = 327.1
      E(21) = 368.1
30     E(22) = 410.2
      E(23) = 453.1
      E(24) = 496.8
      E(25) = 541.1
      F(1) = 0.2525
35     F(2) = 0.3408
      F(3) = 0.4122
      F(4) = 0.4661
      F(5) = 0.5106
      F(6) = 0.5481
40     F(7) = 0.5760
      F(8) = 0.6017
      F(9) = 0.6222
      F(10) = 0.6414
      F(11) = 0.6551
45     F(12) = 0.6684
      F(13) = 0.6797
      F(14) = 0.6907
      F(15) = 0.6988
      F(16) = 0.7058
50     F(17) = 0.7130
      F(18) = 0.7231
      F(19) = 0.7319
      F(20) = 0.7391
      F(21) = 0.7449
55     F(22) = 0.7498
      F(23) = 0.7539
      F(24) = 0.7574
      F(25) = 0.7602
      W0 = 2.066

```

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```

60      ENER = 12.5
        K = 0
        XSUM = 0.
        YSUM = 0.
10     CONTINUE
65      T = ENER/1000.
        W = (T + 0.511)/0.511
        XN = W**2 - 1.
        XN = W*SQRT(XN)*(W0 - W)**2
70      DO 20 I=1,24
        IF(ENER.GE.E(I).AND.ENER.LE.E(I+1)) GO TO 30
20     CONTINUE
        I = 24
30     CONTINUE
75      S = F(I) + ((ENER - E(I))/(E(I+1) - E(I)))*(F(I+1) - F(I))
        XN = S*XN
        WRITE(6,1) ENER,XN
        K = K + 1
        X(K) = ENER
80      Y(K) = XN*91.538
        XSUM = XSUM + XN
        YSUM = YSUM + XN*ENER
1      FORMAT(2E16.8)
        ENER = ENER + 2.5
        IF(ENER.LE.540.) GO TO 10
85      XMIN = 0.
        XSF = 60.
        X(K+1) = XMIN
        X(K+2) = XSF
        CALL PSEUDO
90      Y(K+1) = 0.
        Y(K+2) = 10.
        CALL CALPLT(1.5,1.5,-3)
        CALL AXES(0.,0.,0.,10.,XMIN,XSF,1.,0.,
95      1 6HENERGY,0.2,-6)
        CALL AXES(0.,0.,90.,10.,Y(K+1),Y(K+2),1.,0.,
1      9HFREQUENCY,0.2,9)
        CALL LINPLT(X,Y,K,1,0,0,0,0)
        CALL NFRAME
        CALL CALPLT(0.,0.,999)
100     ENAVG = YSUM/XSUM
        WRITE(6,1) ENAVG
        STOP
        END

```

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```

1      PROGRAM GENP(OUTPUT,TAPE6=OUTPUT,INPUT,TAPE5=INPUT)
      DIMENSION X(300),Y(300)
      COMMON/JTB/NFR,JREQ,IBAUD,JCCS,IJO,TFAC,IJTB(4)
5      NFR = 0
      JREQ = 0
      IBAUD = 120
      JCCS = 0
      TFAC = 0.016
      K = 0
10     COLR = 0.152*2.54/2.
      A = COLR
      B = A/4.
      BB = (B**2 - A**2)/(2.*B)
      RSQ = A**2 + BB**2
15     DIST = -A
      DDIST = A/100.
10    CONTINUE
      FDIST = SQRT(RSQ - DIST**2) + BB
      WRITE(6,1) DIST,FDIST
20     K = K + 1
      X(K) = DIST
      Y(K) = 767.6*FDIST
      DIST = DIST + DDIST
      IF(DIST.LE.A) GO TO 10
25     1 FORMAT(2E16.8)
      XMIN = -0.2
      XSF = 0.04
      X(K+1) = XMIN
      X(K+2) = XSF
30     CALL PSEUDO
      Y(K+1) = 0.
      Y(K+2) = 10.
      CALL CALPLT(1.5,1.5,-3)
      CALL AXES(0.,0.,0.,10.,XMIN,XSF,1.,0.,
35     1 8HPOSITION,0.2,-8)
      CALL AXES(0.,0.,90.,10.,Y(K+1),Y(K+2),1.,0.,
1 9HFREQUENCY,0.2,9)
      CALL LINPLT(X,Y,K,1,0,0,0,0)
      CALL NFRAME
40     CALL CALPLT(0.,0.,999)
      STOP
      END

```

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```

1      PROGRAM CUMU(OUTPUT,TAPE6=OUTPUT,TAPE1)
      DIMENSION X(1000),Y(1000)
      K = 1
10     CONTINUE
5      READ(1,1) X(K),Y(K)
      1 FORMAT(2E16.8)
      IF(EOF(1).NE.0) GO TO 20
      K = K + 1
      GO TO 10
10     20 K = K - 1
      DO 30 I=2,K
      Y(I) = Y(I) + Y(I-1)
30     CONTINUE
      DO 40 I=1,K
15     Y(I) = Y(I)/Y(K)
40     CONTINUE
      DO 50 I=1,K
      WRITE(6,1) X(I),Y(I)
20     CONTINUE
      STOP
      END

```

REFERENCES

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10. Evans, Robley D.: The Atomic Nucleus. McGraw-Hill Book Co., Inc., c.1955, p. 558.
11. Tables for the Analysis of Beta Spectra. NBS Appl. Math. Ser. 13, U.S. Dep. Commer., June 2, 1952.

TABLE I.- SUMMARY OF CRITICAL PARAMETERS FOR POSITRONS ARRIVING AT
 TARGET FOR MAGNETIC FIELD SET FOR 200 keV

[Convex source configuration; 1000 positrons emitted at the source]

| Parameter | Target diameter | |
|--------------------------------------|-----------------|---------|
| | 1 cm | 2.54 cm |
| $Y_C = 9$ cm | | |
| E_{β^+} (min), keV | 175.83 | 154.91 |
| E_{β^+} (max), keV | 230.75 | 263.40 |
| τ^{β^+} (min), ps | 996.10 | 982.49 |
| τ (max), ps | 1085.64 | 1118.87 |
| No. of hits | 68 | 174 |
| $Y_C = 12$ cm | | |
| E_{β^+} (min), keV | 187.34 | 165.20 |
| E_{β^+} (max), keV | 218.70 | 247.30 |
| τ^{β^+} (min), ps | 1143.53 | 1123.25 |
| τ (max), ps | 1215.44 | 1241.45 |
| No. of hits | 44 | 140 |
| $Y_C = 16$ cm | | |
| E_{β^+} (min), keV | 187.94 | 174.86 |
| E_{β^+} (max), keV | 211.70 | 230.75 |
| τ^{β^+} (min), ps | 1350.55 | 1320.40 |
| τ (max), ps | 1411.32 | 1436.76 |
| No. of hits | 38 | 103 |
| $Y_C = 20$ cm | | |
| E_{β^+} (min), keV | 191.72 | 179.44 |
| E_{β^+} (max), keV | 208.06 | 224.71 |
| τ^{β^+} (min), ps | 1549.43 | 1518.31 |
| τ (max), ps | 1587.12 | 1610.38 |
| No. of hits | 28 | 77 |

TABLE II.- SUMMARY OF CRITICAL PARAMETERS FOR POSITRONS ARRIVING AT
 TARGET FOR MAGNETIC FIELD SET FOR 300 keV

[Convex source configuration; 1000 positrons emitted at the source]

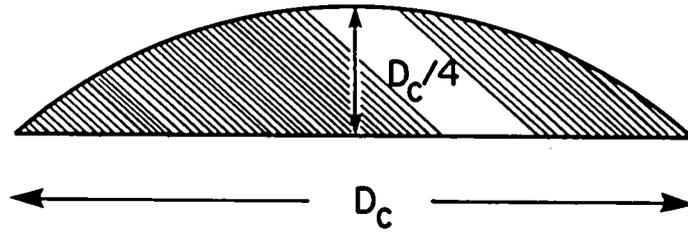
| Parameter | Target diameter | |
|--------------------------------------|-----------------|---------|
| | 1 cm | 2.54 cm |
| $Y_C = 9 \text{ cm}$ | | |
| E_{β^+} (min), keV | 272.96 | 235.58 |
| E_{β^+} (max), keV | 347.05 | 398.60 |
| τ^{β^+} (min), ps | 900.02 | 894.96 |
| τ (max), ps | 952.90 | 972.04 |
| No. of hits | 74 | 161 |
| $Y_C = 12 \text{ cm}$ | | |
| E_{β^+} (min), keV | 277.95 | 250.92 |
| E_{β^+} (max), keV | 328.33 | 365.64 |
| τ^{β^+} (min), ps | 1030.65 | 1025.05 |
| τ (max), ps | 1081.73 | 1101.24 |
| No. of hits | 58 | 127 |
| $Y_C = 16 \text{ cm}$ | | |
| E_{β^+} (min), keV | 285.76 | 262.36 |
| E_{β^+} (max), keV | 317.88 | 347.05 |
| τ^{β^+} (min), ps | 1210.39 | 1192.08 |
| τ (max), ps | 1253.52 | 1260.10 |
| No. of hits | 44 | 98 |
| $Y_C = 20 \text{ cm}$ | | |
| E_{β^+} (min), keV | 288.78 | 272.24 |
| E_{β^+} (max), keV | 315.12 | 332.66 |
| τ^{β^+} (min), ps | 1381.06 | 1367.96 |
| τ (max), ps | 1425.30 | 1433.93 |
| No. of hits | 31 | 84 |

TABLE III.- SUMMARY OF CRITICAL PARAMETERS FOR POSITRONS ARRIVING AT

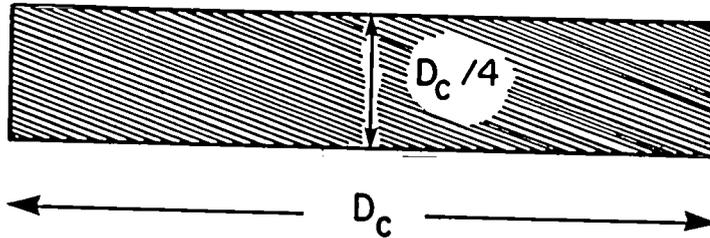
TARGET FOR MAGNETIC FIELD SET FOR 400 keV

[Convex source configuration; 1000 positrons emitted at the source]

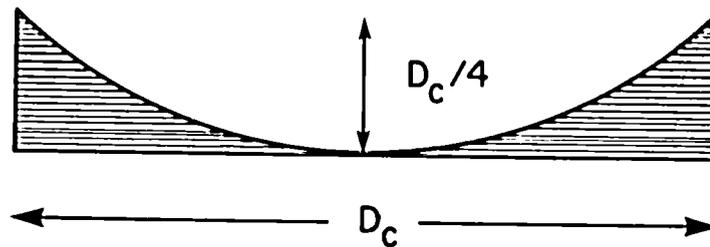
| Parameter | Target diameter | |
|--------------------------------------|-----------------|---------|
| | 1 cm | 2.54 cm |
| $Y_C = 9 \text{ cm}$ | | |
| E_{β^+} (min), keV | 352.47 | 314.06 |
| E_{β^+} (max), keV | 453.85 | 511.48 |
| τ_{β^+} (min), ps | 851.16 | 846.36 |
| τ (max), ps | 898.29 | 909.75 |
| No. of hits | 31 | 92 |
| $Y_C = 12 \text{ cm}$ | | |
| E_{β^+} (min), keV | 370.30 | 336.12 |
| E_{β^+} (max), keV | 431.36 | 472.16 |
| τ_{β^+} (min), ps | 970.47 | 970.47 |
| τ (max), ps | 1014.94 | 1023.05 |
| No. of hits | 23 | 64 |
| $Y_C = 16 \text{ cm}$ | | |
| E_{β^+} (min), keV | 377.53 | 352.47 |
| E_{β^+} (max), keV | 414.33 | 455.31 |
| τ_{β^+} (min), ps | 1133.09 | 1129.55 |
| τ (max), ps | 1175.51 | 1189.39 |
| No. of hits | 19 | 44 |
| $Y_C = 20 \text{ cm}$ | | |
| E_{β^+} (min), keV | 383.53 | 358.42 |
| E_{β^+} (max), keV | 414.33 | 434.74 |
| τ_{β^+} (min), ps | 1293.56 | 1288.62 |
| τ (max), ps | 1336.08 | 1342.42 |
| No. of hits | 17 | 30 |



(a) Convex configuration.



(b) Flat configuration.



(c) Concave configuration.

Figure 1.- Profiles of Na^{22} source deposited at bottom of collimator considered in present study.



Figure 2.- Schematic diagram of source collimator.

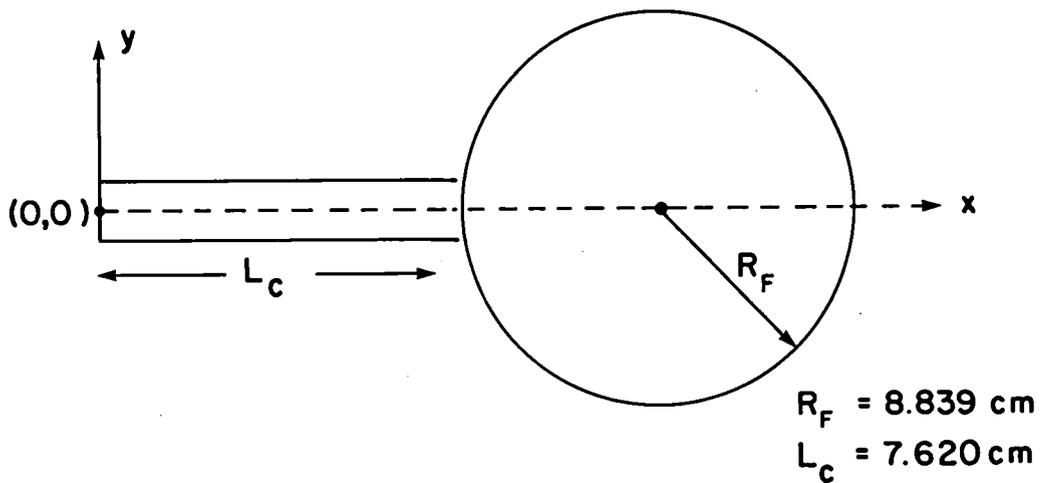


Figure 3.- Schematic drawing showing source collimator and magnetic field region.

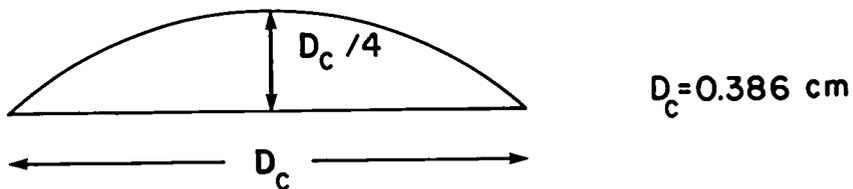
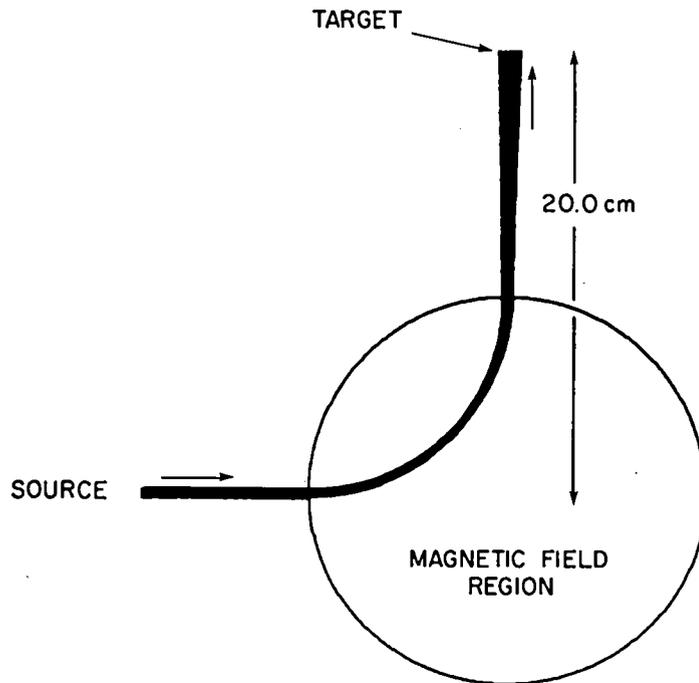
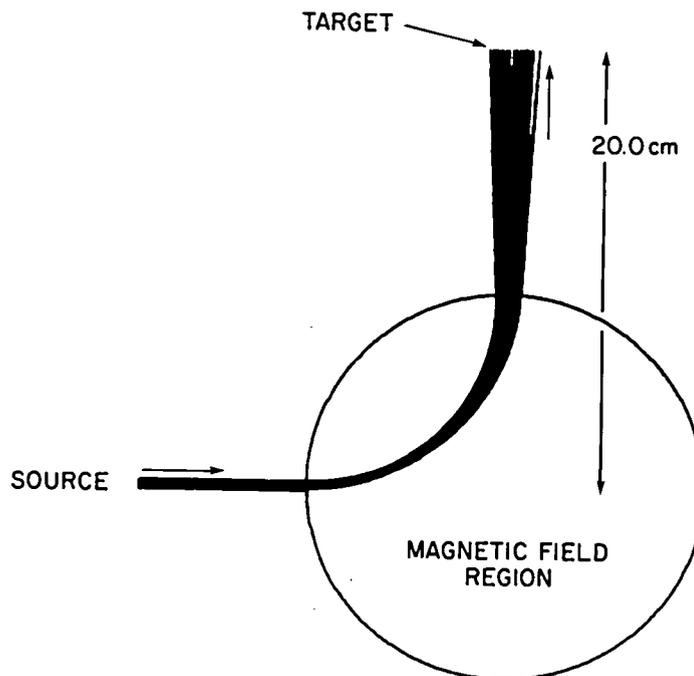


Figure 4.- Profile of Na^{22} source deposited at bottom of collimator selected for detailed study.

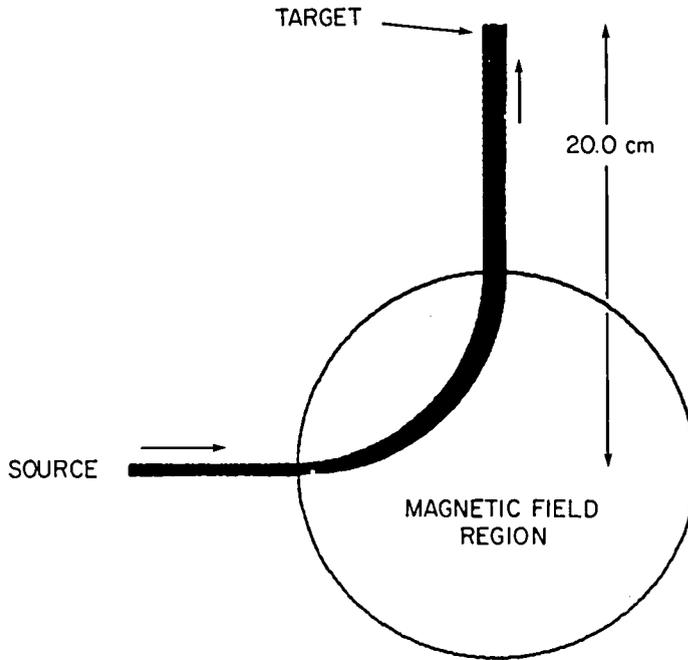


(a) 1-cm-diameter target.

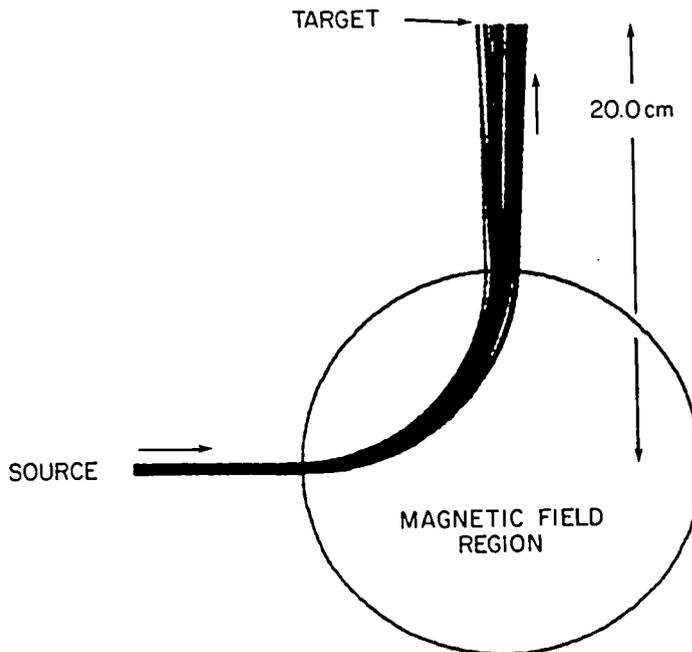


(b) 2.54-cm-diameter target.

Figure 5.- Trajectories of positrons striking target from convex source configuration. Magnetic field set for $E_{\beta^+} = 200$ keV.

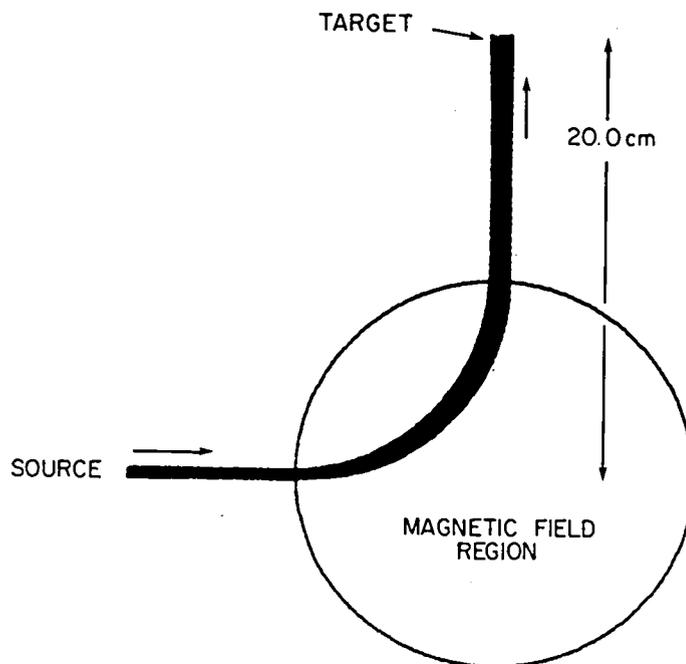


(a) 1-cm-diameter target.

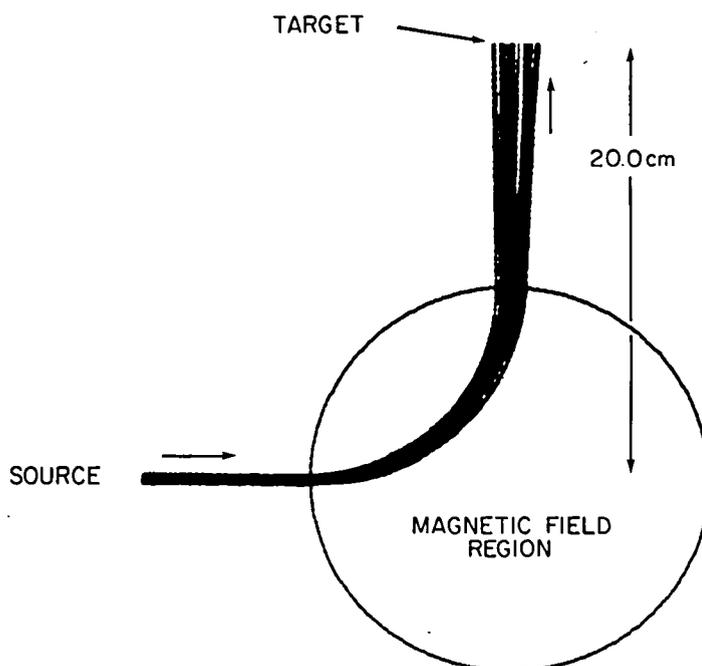


(b) 2.54-cm-diameter target.

Figure 6.- Trajectories of positrons striking target from flat source configuration. Magnetic field set for $E_{\beta^+} = 200$ keV.

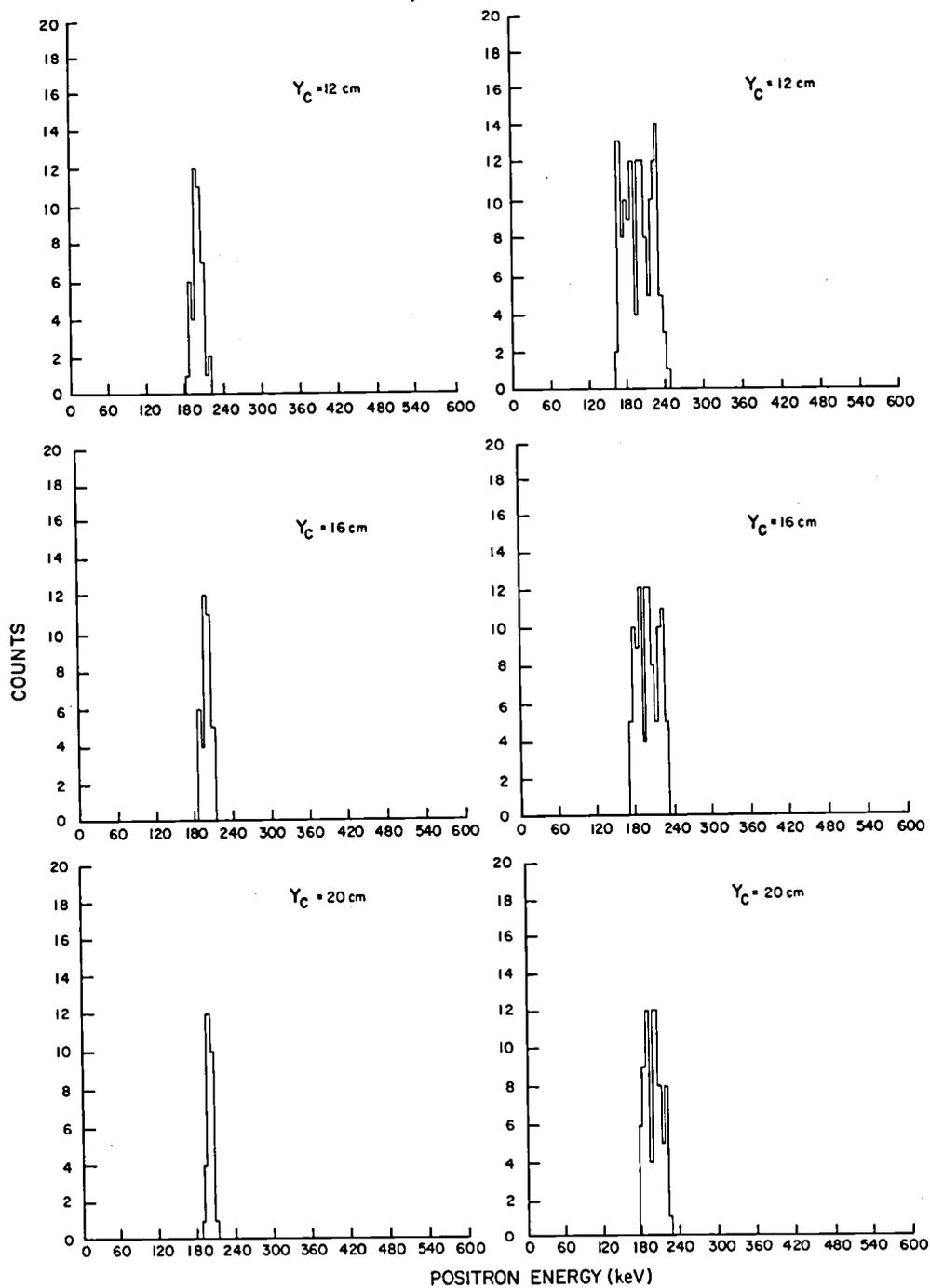


(a) 1-cm-diameter target.



(b) 2.54-cm-diameter target.

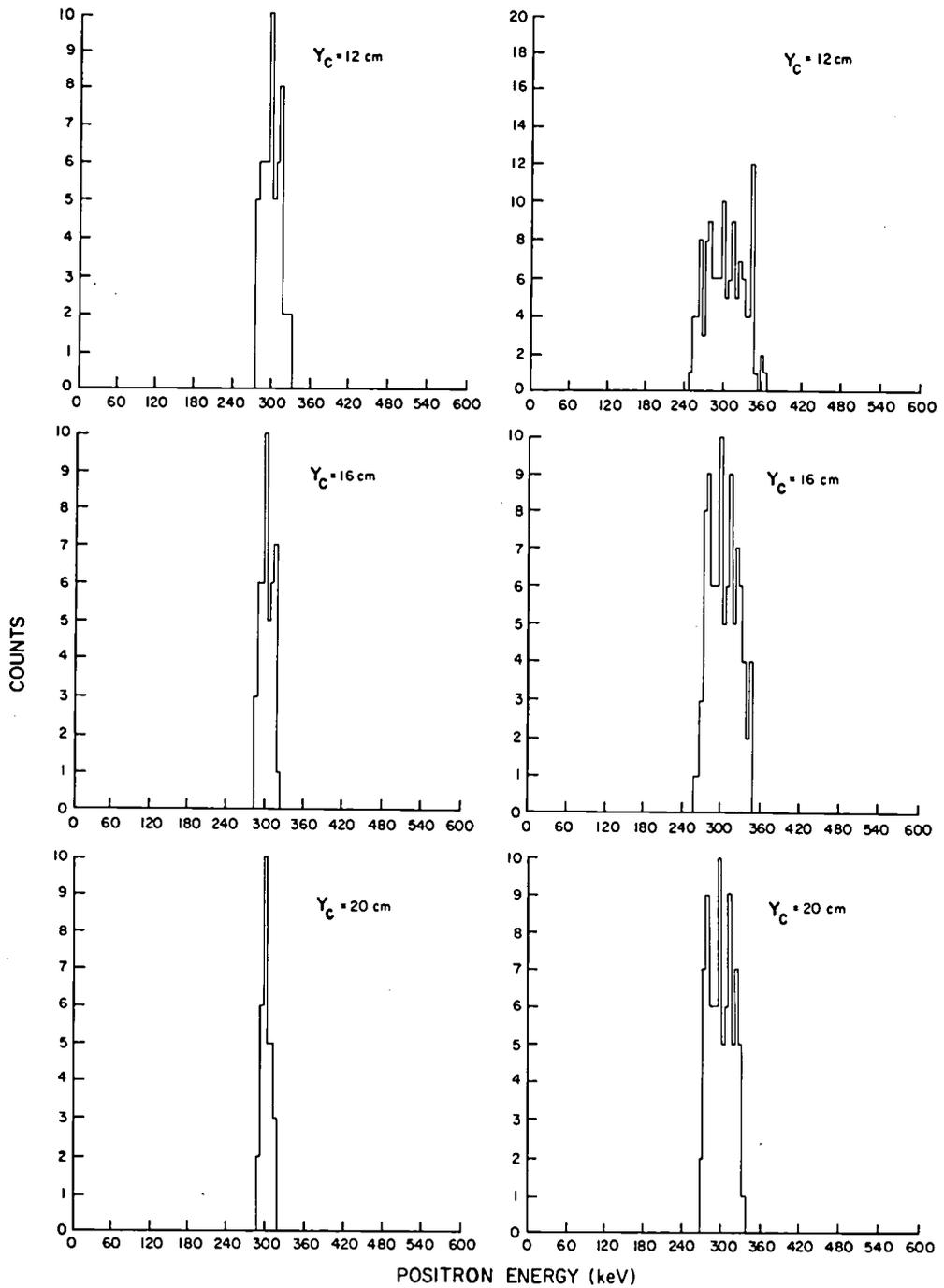
Figure 7.- Trajectories of positrons striking target from concave source configuration. Magnetic field set for $E_{\beta^+} = 200$ keV.



(a) $R = 0.50$ cm.

(b) $R = 1.27$ cm.

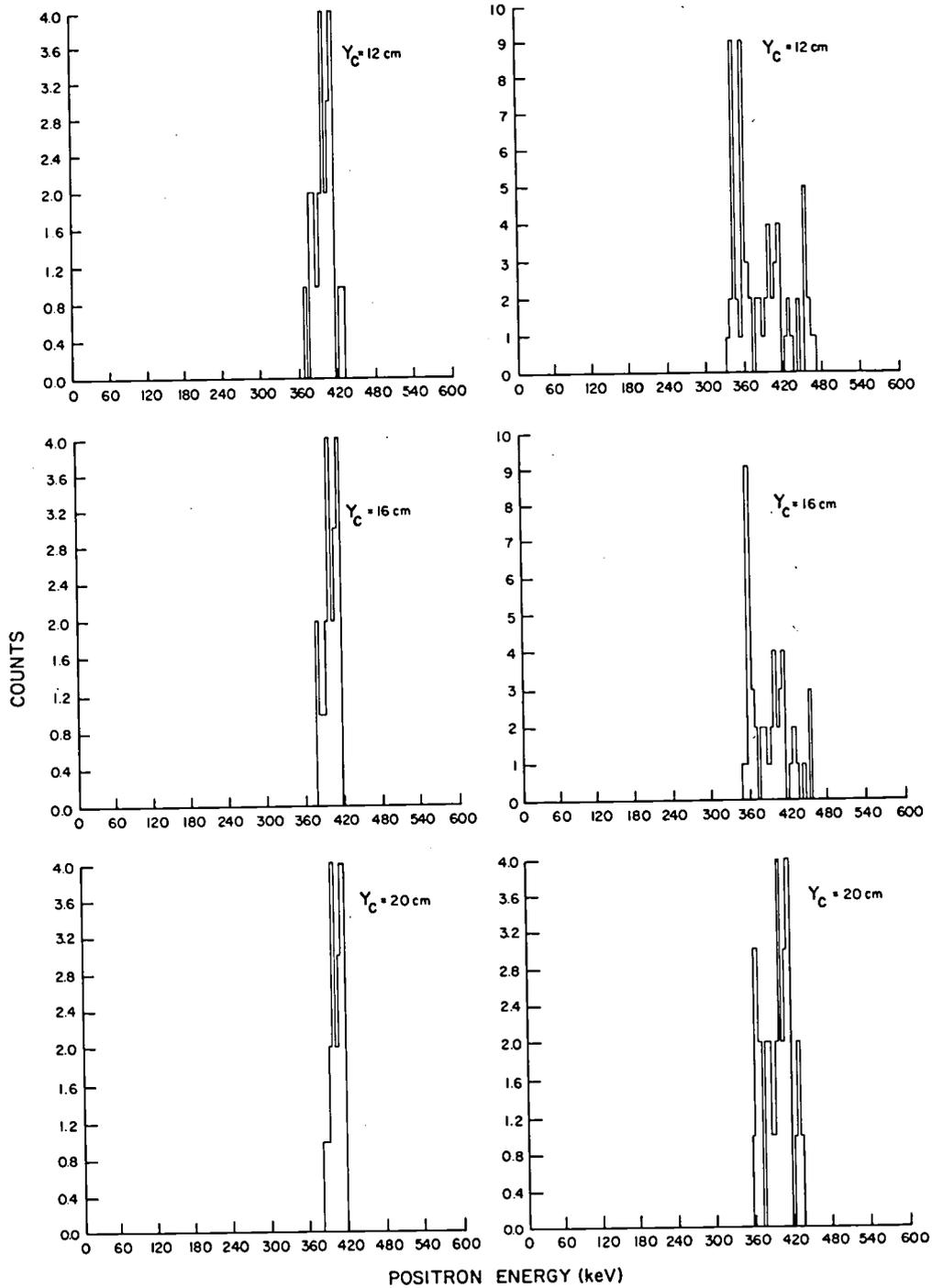
Figure 8.- Energy spectra of positrons arriving at target located at different distances from center of magnetic field set for $E_{\beta^+} = 200$ keV.



(a) $R = 0.50$ cm.

(b) $R = 1.27$ cm.

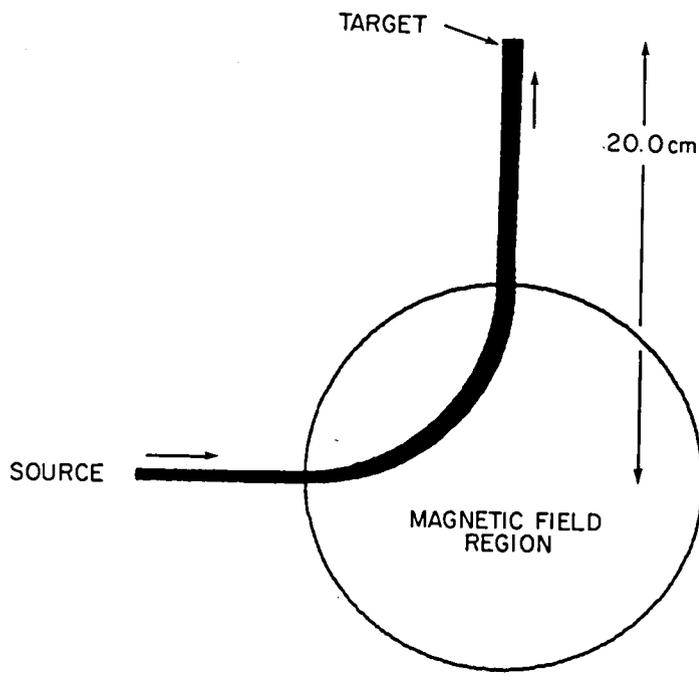
Figure 9.- Energy spectra of positrons arriving at target located at different distances from center of magnetic field set for $E_{\beta^+} = 300$ keV.



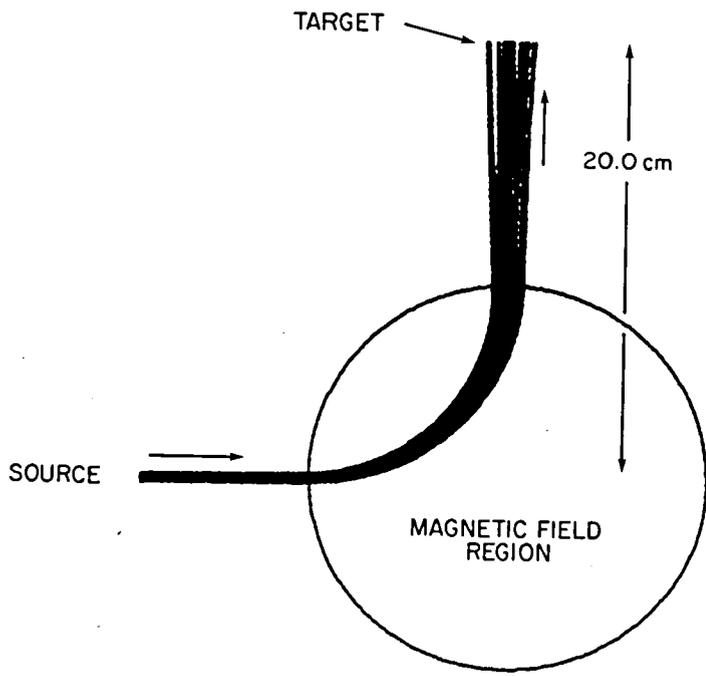
(a) $R = 0.50$ cm.

(b) $R = 1.27$ cm.

Figure 10.- Energy spectra of positrons arriving at target located at different distances from center of magnetic field set for $E_{\beta^+} = 400$ keV.



(a) 1-cm-diameter target.



(b) 2.54-cm-diameter target.

Figure 11.- Trajectories of positrons striking target located 20 cm from center of magnetic field set for $E_{\beta^+} = 300$ keV. Convex source configuration.

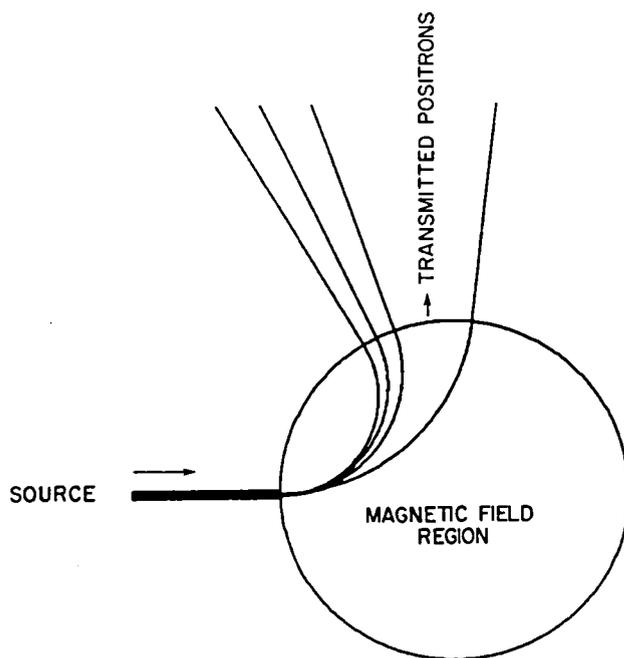


Figure 12.- Sample program output at 300 keV.

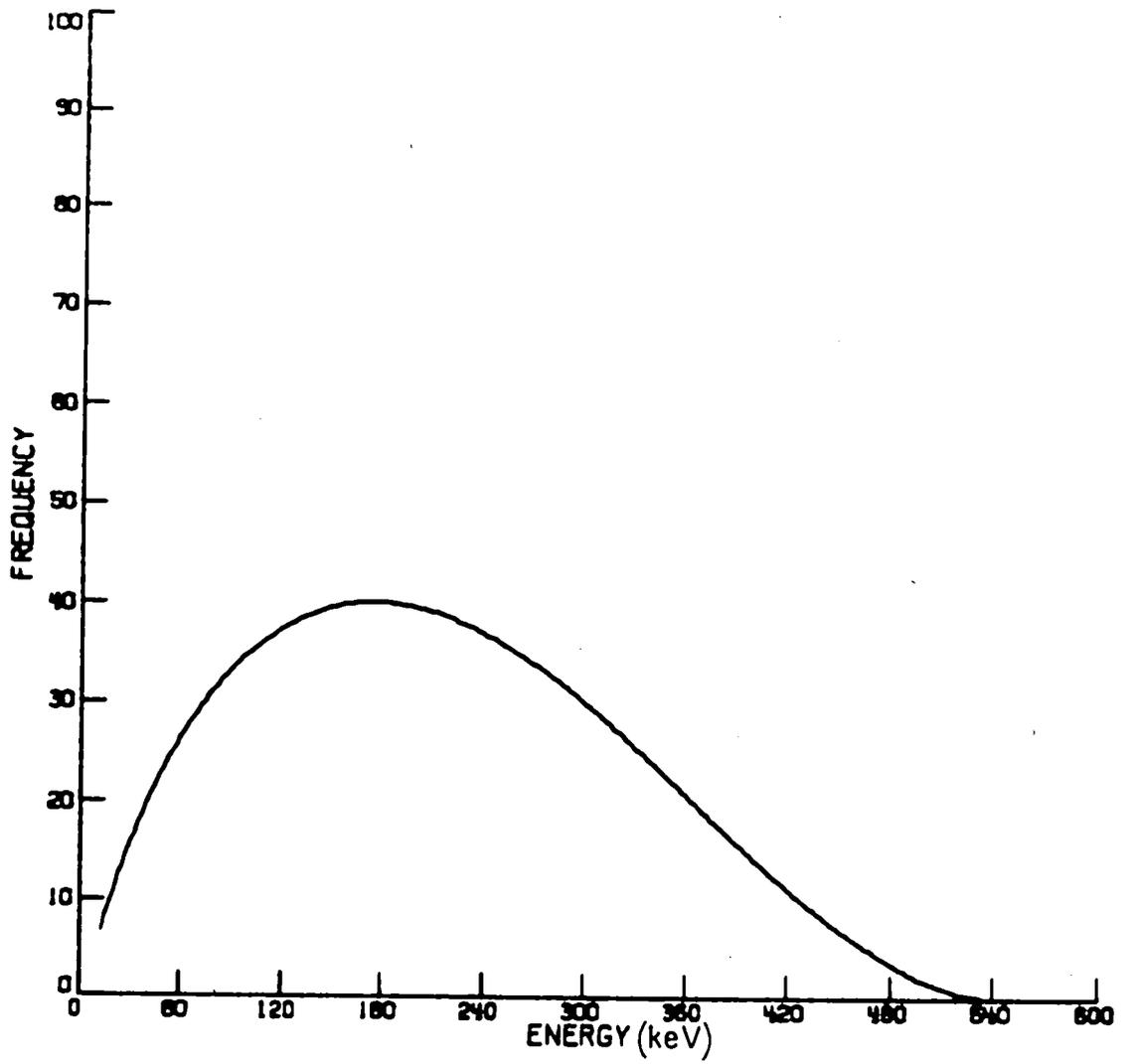


Figure 13.- Computed Na²² positron energy spectrum.

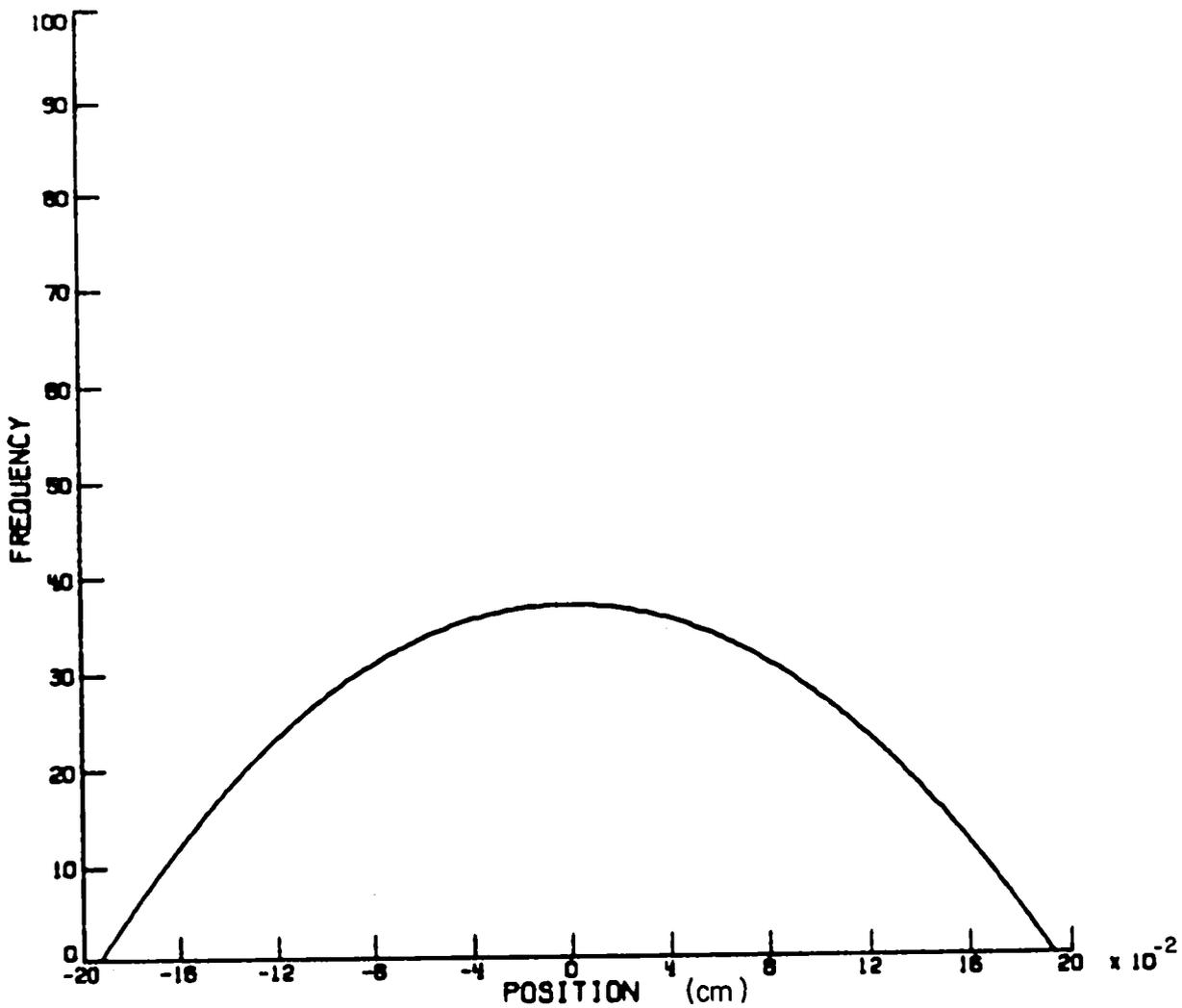


Figure 14.- Computed probability of a positron being emitted as a function of position within collimator.

| | | | |
|--|--|---|----------------------|
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