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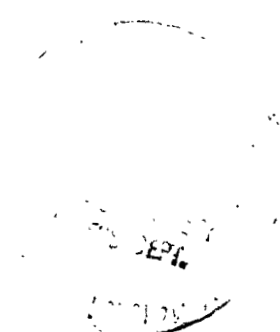
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Abstract

This paper presents a comparison of analysis and flight test data for a drone aircraft equipped with an active flutter suppression system. Emphasis is placed on the comparison of modal dampings and frequencies as a function of Mach number. Results are presented for both symmetric and antisymmetric motion with flutter suppression off. Only symmetric results are presented for flutter suppression on. Frequency response functions of the vehicle are presented from both flight test data and analysis. The analysis correlation is improved by using an empirical aerodynamic correction factor which is proportional to the ratio of experimental to analytical steady-state lift curve slope. In addition to presenting the mathematical models and a brief description of existing analytical techniques, an alternative analytical technique for obtaining closed-loop results is presented.

Introduction

With active control technology, supercritical airfoils, and new materials evolving into viable approaches for reducing aircraft mass, the confidence in accurately predicting the dynamic behavior of flexible aircraft incorporating these advancements is important. Because of the potentially destructive nature of flutter, the credibility of modern flutter prediction techniques are of increasing concern to the designer. To gain confidence in the analytical tools and techniques employed, comparisons of analytical results with those obtained from experiment are necessary.

Research efforts in the past at the NASA have been directed toward both developing analysis tools and correlation of wind-tunnel test results with analytical predictions. Wind-tunnel tests in the transonic region are traditionally limited by model size, dynamic pressure (for aeroelastic tests), and effects of interaction between the model and reflected shock waves. As a complement to wind-tunnel tests, flight tests are being conducted in a NASA program called Drones for Aerodynamic and Structural Testing (DAST).¹ In this program, unmaneuvered, remote-controlled drone aircraft are used as testbeds for high risk research such as flutter tests.

Some of the objectives of this flight program are to explore the accuracy and limitations of flutter predictions for a supercritical wing in the transonic region and to evaluate flutter suppression system (FSS) performance.

Reference 2 presents results of one of the first wind-tunnel tests of a flexible supercritical wing. The agreement between the measured flutter boundary and the calculated boundary was very good except in the critical transonic region. References 3 through 5 present results of studies of active control of wing/store flutter for a fighter configuration. Some of the comparisons of analysis and experiment that are given show reasonably good agreement. Reference 6 presents results for an active flutter margin augmentation system for a commercial transport with a conventional wing. Comparisons between analysis and flight data are given in terms of transfer functions and damping/frequencies and show good agreement. Wind-tunnel tests of a cantilever aeroelastic wing model of the present drone wing with an active flutter suppression system are reported in Ref. 7. However, in contrast to the drone aircraft wing section reported here, the wind-tunnel model did not have a supercritical shape. This study showed good agreement for flutter mode frequencies for both the FSS-off and FSS-on cases. Neither damping nor frequency response comparisons were made since these quantities were not experimentally measured. In summary, no flutter tests, either wind-tunnel or flight, of a supercritical wing configuration with active flutter suppression have been reported previously.

This paper is one of three companion papers that describe various aspects of the first few flights in the DAST program. Edwards⁸ presents details of the flutter test technique development and the implementation of the FSS on the vehicle. Bennett and Abel⁹ present the experimental frequency and damping estimates obtained using a post-flight parameter estimation technique. This paper is presented as an effort to correlate the results obtained by various analytical techniques with the experimental results given in Ref. 9. The study ranges from the comparison of open loop frequency responses to the comparison of flexible mode damping/frequency differences for both symmetric and antisymmetric motion. In addition, a description of the analytical techniques and mathematical models used in the study are given.

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Nomenclature

$C_{L\alpha}$	lift curve slope
c	reference chord
D	parameter scheduler
$D(s)$	denominator polynomial
G	shear modulus
G_a	antisymmetric filter
G_c	common filter
G_s	symmetric filter
g	acceleration due to gravity
$G(i\omega)$	vehicle frequency response
$H(i\omega)$	control law and actuator frequency response
i	$\sqrt{-1}$
k	reduced frequency
M	Mach number
$N(s)$	numerator polynomial
s	Laplace variable
t	time
V	free-stream velocity
w_g	gust velocity
$Y(s)$	transfer function input
\ddot{z}	vertical acceleration
β	aerodynamic lag
δ	control surface position
δ_c	actuator command
ζ	damping ratio
$\ddot{\theta}_x$	roll angular acceleration
ρ	free-stream density
ω	circular frequency

Matrices:

$[A]$	dynamics matrix
$[A_i]$	real aerodynamic coefficient matrix
$[B]$	control distribution matrix
$[C]$	state-coefficient output matrix
$[D]$	input-coefficient output matrix
$[D_s]$	structural damping matrix
$[K]$	generalized stiffness matrix
$[M]$	generalized mass matrix
$[Q]$	generalized aerodynamic force matrix
$\{q\}$	generalized coordinate vector
$\{u\}$	input vector
$\{X\}$	state vector
$\{Y\}$	output vector

Subscripts:

c	control system
FUS	fuselage
LW	left wing
m	motion
N	nominal
RW	right wing
v	vehicle
w_g	vertical gust
δ	control surface

Dots over symbols denote derivatives with respect to time.

Description of Drone Aircraft

The geometry of the drone aircraft is shown in Fig. 1. The wing has an aspect ratio of 6.8 with a supercritical airfoil designed for cruise at $M = 0.98$. The small outboard control surfaces on the wing have a dual function of both suppressing flutter through the control system and providing excitation for measurement purposes. Accelerometers located on the rear spar near the wing tip are used to sense motion for the FSS. Fuselage accelerometers are used to subtract out rigid body motion.

Bending stiffness of the wing is provided primarily by two steel spars located at 25 percent and 60 percent of the local wing chord. Torsional stiffness is provided primarily by fiberglass skins which are attached to the spars to form the wing box. The fibers in the skin are oriented parallel and normal to the wing elastic axis to create a torsional stiffness which is low enough to allow flutter to occur within the flight envelope. In addition, a ballast is added to the outboard section aft of the trailing edge to further reduce the flutter speed.

The drone has been modified to improve controllability and acquisition of research data. Flight tests involve air launches from beneath the wing of a B-52 and mid-air recovery by a helicopter. A description of the operational procedures is given in Refs. 1 and 8.

Analytical Techniques and Mathematical Models

Equations of Motion

The equations of motion are formulated through the conventional modal approach. The motion of the vehicle is described by a linear combination of its free vibration modes. The forcing functions to the system are both control and gust inputs. Therefore, the basic equations of motion can be written as

$$[M]\{\ddot{q}\} + [D_s]\{\dot{q}\} + [K]\{q\} + \frac{1}{2} \rho V^2 [Q_m(t)]\{q\} = -\frac{1}{2} \rho V^2 \left[\begin{matrix} Q_\delta(t) \\ Q_{w_g}(t) \end{matrix} \right] \left\{ \begin{matrix} \delta \\ w_g \end{matrix} \right\} \quad (1)$$

All coefficients with the exception of the aerodynamic terms in these equations are independent of time. The aerodynamic force matrices are complex valued and normally represented as tabular functions of reduced frequency, $k = c\omega/2V$. To obtain a set of constant coefficient differential equations, the unsteady aerodynamic force matrices are approximated by a rational polynomial in the Laplace variable s

$$[Q(s)] = [A_0] + [A_1] \left(\frac{cs}{2V} \right) + [A_2] \left(\frac{cs}{2V} \right)^2 + \sum_{j=1}^4 \frac{[A_{2+j}]s}{s + \frac{2V}{c} \beta_j} \quad (2)$$

After selecting a set of aerodynamic lags, β_j , the coefficient matrices $[A_i]$ are computed by a least squares curve fit in a fashion similar to that described in Ref. 10.

By equating derivatives to the powers of the Laplace operator s , the equations of motion for the basic vehicle can be written in standard state-space form as follows:

$$\begin{aligned} \dot{\{X_V\}} &= [A_V]\{X_V\} + [B_V]\{u_V\} \\ \{Y_V\} &= [C_V]\{X_V\} + [D_V]\{u_V\} \end{aligned} \quad (3)$$

where

$$\{X_V\} = \begin{Bmatrix} q \\ \dot{q} \\ X_{\beta_1} \\ \vdots \\ X_{\beta_n} \end{Bmatrix};$$

$$\{u_V\} = \begin{Bmatrix} c \\ \delta \\ \ddot{\delta} \\ w_g \\ \dot{w}_g \\ \ddot{w}_g \end{Bmatrix};$$

$\{Y_V\}$ = sensor outputs.

The matrix $[A_V]$ is a $6n \times 6n$ matrix where n is the number of modes.

Stability Analysis

A brief description of the analytical methods used to calculate stability with both FSS-off and FSS-on will be presented in this section. In particular, an alternate method for the FSS-on calculations which differs from that given in Ref. 10 is presented.

FSS-Off. The matrix $[A_V]$ in Eq. (3) is a function of Mach number, velocity, and density. For fixed values of Mach number, velocity, and density the eigenvalues of $[A_V]$ are the roots of the FSS-off flutter equation. Variation of the roots of the flutter equation as a function of altitude at constant Mach number can be computed. This approach yields a "matched point" solution without the necessity of cross-plotting several density calculations as with the conventional V-g flutter solution. The flutter boundary for the basic vehicle is determined by performing the calculations at several Mach numbers.

FSS-On. An alternate approach from that of Ref. 10 is employed to formulate the FSS-on equations. To use this approach, both the actuator and control law models are written in standard state-space form and then interconnected to the basic vehicle equations. This approach is used because it requires the solution of a much smaller order eigenvalue problem than the approach described in Ref. 10.

A single control input and single sensor output will be used to illustrate this technique, but the same technique can be extended to multiple controls and sensors. The control law and actuator dynamics are normally represented by transfer functions which can be combined and written in the form

$$\frac{S(s)}{Y(s)} = \frac{N(s)}{D(s)} \quad (4)$$

where

$N(s)$ = numerator polynomial in s

$D(s)$ = denominator polynomial in s

There are many ways of expressing this transfer function in state-space form. The technique used in this paper is described in Ref. 11. The development uses a Frobenius form realization to obtain the matrix coefficients. As shown in Ref. 11, the system defined by the transfer function of Eq. (4) can be written in state-space form as

$$\dot{\{X_C\}} = [A_C]\{X_C\} + \{B_C\}Y_C \quad (5)$$

$$\{u_C\} = [C_C]\{X_C\}$$

where

Y_C = transfer function input; and

$$\{u_C\} = \begin{Bmatrix} \delta \\ \ddot{\delta} \end{Bmatrix}$$

The matrix $[A_C]$ is an $nc \times nc$ matrix where nc is the order of $D(s)$ in Eq. (4). It is important to note that when using this approach the control displacement and its first and second derivatives must be made available as inputs to the basic vehicle system. This implies that the transfer function needs to be at least third order to obtain the higher derivatives and avoid any direct transfer term from Y_C .

The interconnection of systems is most often described by the "state augmentation technique." In the present application, consider the systems described by Eqs. (3) and (5). Defining the augmented state as

$$\{X\} = \begin{Bmatrix} X_V \\ X_C \end{Bmatrix} \quad (6)$$

and using the relations that $\{u_v\} = \{u_c\}$ and $Y_c = Y_v$, the closed loop or FSS-on dynamics are represented by

$$\begin{aligned} \dot{\{X\}} &= \begin{bmatrix} A_v & B_v C_c \\ B_c C_v & A_c \end{bmatrix} \{X\} \\ \dot{\{X\}} &= [A] \{X\} \end{aligned} \quad (7)$$

where $[A]$ is a $(6n+nc) \times (6n+nc)$ matrix. As in the FSS-off equations, the matrix $[A]$ in Eq. (7) is a function of Mach number, velocity, and density. For fixed values of Mach number, velocity, and density the eigenvalues of $[A]$ are the roots of the FSS-on flutter equation. Variations in Mach number and/or density can be used to predict modal dampings and frequencies and the FSS-on flutter boundary.

Structural Model

The structural model is a NASTRAN* finite element model of the wing, wing center section, fuselage, and empennage. The wing model was originally developed by Teledyne Ryan Aeronautical and then updated to a complete vehicle model by both Boeing-Wichita¹³ and NASA-Langley. A drawing of the updated NASTRAN model is presented in Fig. 2a.

The wing finite element idealization includes the leading and trailing edge structure which are modeled with elements providing stiffness only for translational degrees of freedom. The spars and ribs are modeled with rods and shear elements. The fiberglass wing skins are modeled using shear elements with rods added to represent the membrane stiffness. The wing center section is modeled with beam and plate elements lying in a horizontal plane. The wing is connected to the wing center section with single point connections at the front and rear spars.

Elastic axis representations employing beam elements are used in modeling the fuselage, vertical tail, and horizontal stabilizer. The connections between the wing center section and the fuselage are defined by constraint equations relating translations at the side of body to the motions of the elastic axis at the fuselage centerline. To simplify the structural representation, appropriate centerline constraints are used to define separately both symmetric and antisymmetric motion.

Ten symmetric and ten antisymmetric elastic modes were computed and used for analysis purposes. The symmetric modes cover a frequency range of 9.1 to 105.0 Hz. The antisymmetric modes cover a frequency range of 12.3 to 80.7 Hz. A ground vibration test was performed on the vehicle to measure both symmetric and antisymmetric modes and frequencies. A comparison between the measured and analytical frequencies is presented in Table 1. For both the symmetric and antisymmetric cases, the frequency of mode 1 which is primarily first wing bending is underpredicted, and the frequency of mode 3 which is primarily first wing torsion is overpredicted.

Unsteady Aerodynamic Model

The unsteady aerodynamic matrices in Eq. (1) are calculated using the doublet lattice program in the ISAC system.¹⁴ The frequency dependent unsteady aerodynamic pressure distributions are calculated by subdividing each lifting surface into an array of streamwise trapezoidal boxes. The wing and empennage sections are aerodynamically modeled as shown in Fig. 2b. The model includes 121 boxes for the wing, 14 boxes for the horizontal tail, and 8 boxes for the vertical tail. The aerodynamic effect of the fuselage is assumed to be negligible. Eight boxes are used to model the wing control surface shown by the shaded area. Unsteady aerodynamic forces are calculated for two rigid body modes, ten elastic modes, a control surface rotation, and a sinusoidal gust.

The aerodynamic lag terms β_j used in approximating the unsteady aerodynamics in the s-plane are arbitrarily selected from the range of reduced frequencies for which the aerodynamic matrices (Eq. (1)) are calculated. The selected values are 0.1, 0.2, 0.3, and 0.4.

Actuator Model

The control surface is driven by a rotary actuator located at the inboard edge of the control surface. The actuator is controlled by a servo-valve mounted in the wing center section. The servoactuators use actuator shaft position and differential load pressure feedback.

The original analytical model of the actuator, which is similar to that described in Ref. 15, was developed by Boeing-Wichita. The model predicted an actuator bandwidth near 100 Hz. However, when the system was assembled and bench-tested the bandwidth was approximately 70 Hz. The bench test indicated a hydraulic fluid mode near 110 Hz. This fluid mode is not accounted for in the analytical model. The presence of this mode contributed to the reduction in bandwidth. Furthermore, when the actuator was installed in the flexible wing, the final bandwidth reduced to approximately 55 Hz. At this point a significant effort was expended to increase the bandwidth and resulted in the implementation of several notch filters to stabilize high frequency (100-400 Hz) hydraulic modes.

Because of these complications, the original actuator model was not used in the analysis. Instead, the actuator model is derived from a measured frequency response to a one-degree-amplitude sinusoidal input with the actuator installed in the wing. The measured frequency response data was curve fit with a 6th order transfer function. This resulted in the following transfer function:

*NASTRAN: Registered trademark of the National Aeronautics and Space Administration

$$\frac{\delta(s)}{\delta_c(s)} = \frac{2.7 \times 10^{16}}{(s^2 + 76.8s + 8.7 \times 10^4)(s^2 + 589.4s + 1.9 \times 10^5)(s^2 + 502.7s + 1.6 \times 10^6)} \frac{\text{deg}}{\text{deg}} \quad (8)$$

Figure 3 presents a Bode plot showing the comparison between the frequency response using Eq. (8) and the measured data. It is interesting that later studies indicate little change in the response characteristics with the inclusion of airloads. However, a substantial decrease in bandwidth and corresponding increase in phase lag was observed when the response was measured for an 8° input amplitude.

FSS Control Law

The FSS was designed by Boeing-Wichita under contract to NASA-Langley. A description of the design of the initial control law is given in Ref. 13. However, between the second and third flights described in Ref. 8, Boeing-Wichita performed a redesign of the control law. A block diagram of the redesigned FSS is presented in Fig. 4. The redesigned control law is used in all analyses presented in this paper. The initial summation of sensors provides left and right wing vertical accelerations minus the rigid body accelerations due to vertical and roll accelerations. These signals are fed into the common filters G_c and then summed to form the symmetric filtered signal and differenced to form the antisymmetric filtered signal. These signals are filtered by the symmetric G_s and antisymmetric G_a filters, then summed and differenced, and multiplied by a gain of 0.5 to form the left and right control surface commands. The transfer functions for the common, symmetric, and antisymmetric filters are given by

$$G_c = \frac{2.25 \times 10^6 s(s^2 + 76s + 295^2)(s^2 + 120s + 306^2)}{(s + 2)(s + 295)^2(s^2 + 240s + 324^2)(s + 1500)^2} \\ \times \frac{(s^2 + 60s + 1037^2)(s^2 + 76s + 1269^2)}{(s^2 + 300s + 1037^2)(s^2 + 500s + 1269^2)} \\ G_s = \frac{-1.6(s^2 + 100s + 71^2)(s^2 + 100s + 168^2)}{(s^2 + 100s + 58^2)(s^2 + 100s + 112^2)} \\ G_a = \frac{-2.528(s^2 + 100s + 158^2)}{(s^2 + 80s + D)} \quad (9)$$

$$\text{where } D = 49000 - 804 \times \left(\frac{1}{2} \rho V^2\right)$$

Experimental Data and Discussion of Results

Experimental Data Acquisition

The control surfaces are used both for suppressing flutter and for exciting the wing responses. A fast swept sine wave, which varies from 10.0 to 40.0 Hz in 8 seconds, is used for exciting the wing. The two control surfaces are oscillated either in phase or out of phase for symmetric or antisymmetric motion, respectively. Measurements in the form of time histories of the responses, command signals, and control surface positions are made and relayed to ground facilities for recording. Frequency response functions of the flexible vehicle are obtained from the time histories through the use of Fast Fourier Transform (FFT) techniques. Both the modal characteristics and the frequency response functions are obtained from the data presented in Ref. 9.

Comparison of Analysis with Flight Test Results

It is well known that linear, thin-wing, potential flow theory starts to lose its validity when approaching the transonic region and is therefore not able to give satisfactory approximations to the aerodynamic forces. There are several empirical methods used to improve the representation of the aerodynamic forces in this region. Probably the most common method uses experimental steady-state aerodynamic force data. Using a relatively straightforward approach, a correction factor, based on the steady-state lift curve slope (C_{L_α}), is applied to all unsteady aerodynamic force data. Specifically, this correction factor is obtained by calculating, at each Mach number, the ratio of experimental-to-analytical C_{L_α} for the vehicle. Figure 5a presents a comparison of the experimental to analytical C_{L_α} for the drone aircraft (experimental values are taken from Ref. 16). The lift curve slope is underpredicted throughout the Mach number range. The ratio of experimental to analytical C_{L_α} varies from 1.03 ($M = 0.70$) to 1.18 ($M = 0.95$). This correction factor is then applied to the generalized aerodynamic force matrix by multiplying all elements by this factor, which has an effect equivalent to increasing density. Figure 5b demonstrates the use of the correction factor in the flutter analysis. The flutter Mach number is shifted more toward the flight test point resulting in a better agreement. Consequently, the correction factor will be used in all analytical results presented.

The predicted flutter boundary and the flight test points for the FSS-off case are presented in Fig. 6 for both symmetric and antisymmetric motion. The difference between the flight test and prediction at constant altitude is approximately 0.02 to 0.03 in Mach number. There could be many reasons for this difference, but a plausible explanation that was examined is the nonlinear behavior of the stiffness properties of the

fiberglass skins. More precisely, the shear modulus for the fiberglass skin was experimentally determined to vary substantially (nearly 50 percent) with shear stress as shown in Fig. 7. Since shear stress increases, and therefore the shear modulus decreases, as Mach number (dynamic pressure) increases, this effect could be important in flutter predictions. Therefore, flutter calculations (only for the symmetric case) were performed for a shear modulus variation of plus and minus 20 percent from the nominal value ($G_N = 0.58 \times 10^{10}$). The results of these calculations are also presented in Fig. 6. These parameter variations more than adequately bracket the data points obtained from the flight tests.

Frequency and damping characteristics were obtained from two flight tests. The symmetric and antisymmetric results shown in Figs. 8 and 9, respectively, were obtained from the first flight test at an altitude of 7.62 kilometers. The change in frequency of the bending (mode 1) and torsion (mode 2) modes with Mach number were well predicted for both symmetric and antisymmetric motion. However, the damping, over the same range of Mach numbers, was overpredicted by analysis for the flutter mode (mode 1) and underpredicted for the torsion mode (mode 2). From previous works^{3,6} where damping was experimentally and analytically available, the analysis usually overpredicted the damping of the flutter mode. This is consistent with the results presented here. These differences require further investigation and this discrepancy should be considered during any FSS design.

The frequency and damping of the dominant mode are presented for both the symmetric and antisymmetric cases in Figs. 10 and 11, respectively. The analysis and experiment correspond to the second flight test at an altitude of 4.57 kilometers. FSS-on results are presented only for the symmetric case. The change in frequency with Mach number is predicted well for both the FSS-off and FSS-on cases. However, analysis overpredicts the damping for both the FSS-off and FSS-on cases as in the 7.62 kilometer results. The experimental flutter speed is extrapolated to be approximately $M = 0.80$ for the FSS-off symmetric case (Fig. 10). An actual flutter point was encountered for the FSS-on case at $M = 0.82$. The analysis overpredicts the FSS-off flutter speed by 4 percent and overpredicts the FSS-on flutter speed by 2 percent.

An alternate method based on frequency response techniques was used to compare flight data to analysis. The frequency domain responses of the vehicle were obtained by FFT techniques from the time responses recorded during flight tests. To compare directly with the analysis open loop frequency response ($G(i\omega) \times H(i\omega)$), the experimentally obtained frequency response ($G(i\omega)$) was augmented by the control law and actuator dynamics ($H(i\omega)$). The results at $M = 0.70$ and 4.57 kilometers are plotted in the form of Nyquist diagrams presented in Figs. 12 and 13. Some smoothing was necessary to reduce the observed scatter apparent in the transformed (FFT) raw flight data. The lobes represent elastic resonance points in the vehicle dynamics and the fact that the lobes in the analysis do not coincide angularly with the ones from the flight data illustrate the phase differences between the two results. There appears to be a better agreement in frequencies for the symmetrical case than

the antisymmetrical case when comparing selected frequency points of the two results.

Concluding Remarks

A comparison of analysis and flight test data for a drone aircraft equipped with active flutter suppression has been presented. Although absolute values of the test results were not predicted, the analysis did predict the trends reasonably well. In particular, modal frequencies were predicted with good success. However, flutter mode damping was in most instances overpredicted. The flutter speed was overpredicted by 3 to 4 percent for the FSS-off case and overpredicted by 2 percent for the FSS-on case. This unconservatism is believed to be a result of the nonlinear torsional stiffness associated with the shear modulus of the fiberglass skin and the aerodynamics of the supercritical airfoil. These results suggest more care is needed in modeling structures in terms of nonlinear properties and understanding the aerodynamic mechanisms which contribute to the differences between analysis and test.

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Table 1 Analysis and GVT vibration frequencies				
Mode	Symmetric		Antisymmetric	
	Analysis	G.V.T.	Analysis	G.V.T.
1	9.1	9.6	12.3	13.5
2	16.5	16.2	21.7	19.3
3	29.6	29.1	30.0	27.0
4	34.0	31.2	34.2	31.0
5	39.8	40.6	36.0	--
6	48.6	48.5	48.3	48.4
7	65.2	--	53.7	--
8	71.1	--	54.0	--
9	78.9	--	77.4	--
10	105.0	--	80.7	--

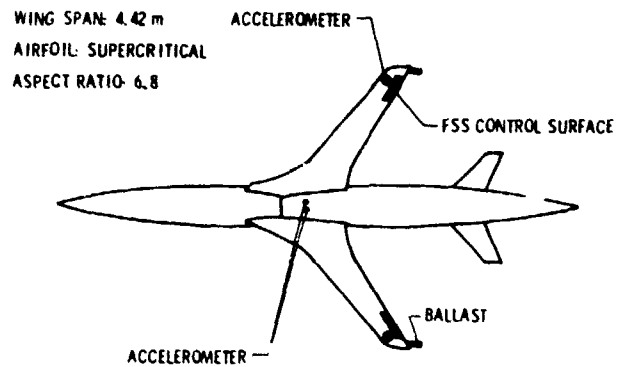


Fig. 1 Vehicle geometry.

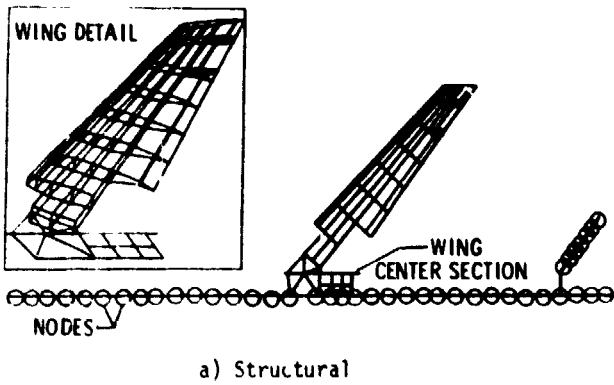


Fig. 2 Structural and aerodynamic models.

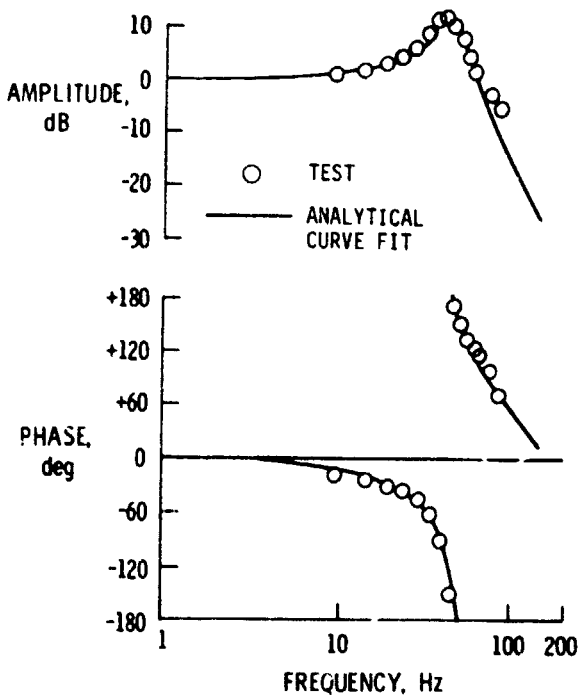


Fig. 3 Bode plot of actuator dynamics.

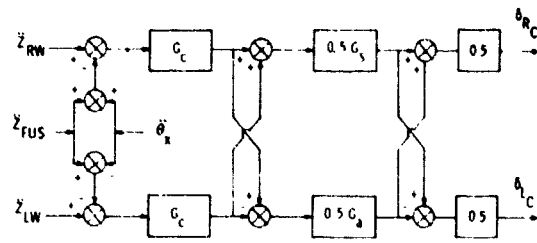
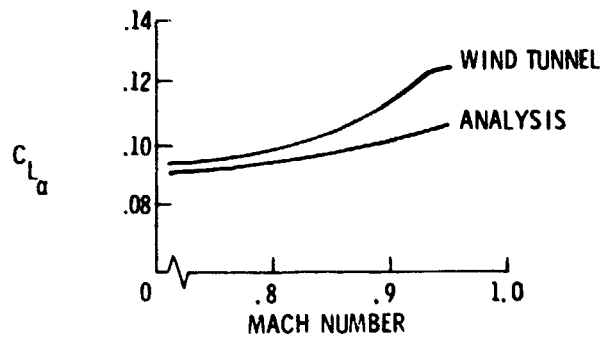


Fig. 4 Block diagram of the flutter suppression system.



a) Lift curve slope

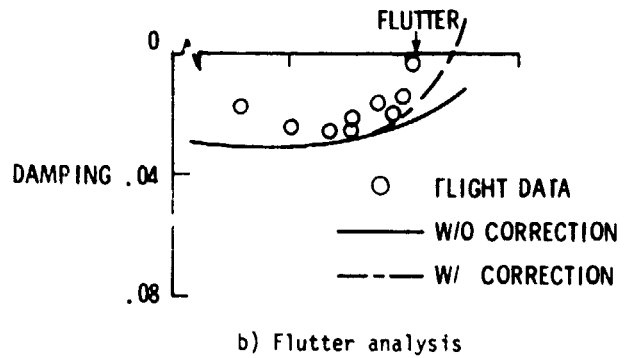


Fig. 5 Effect of aerodynamic correction factor.

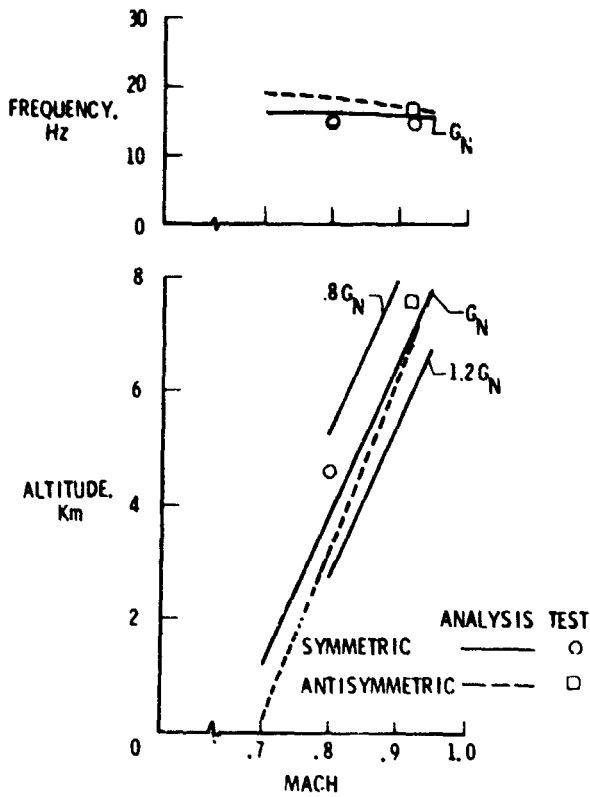


Fig. 6 Flutter boundary (FSS-off).

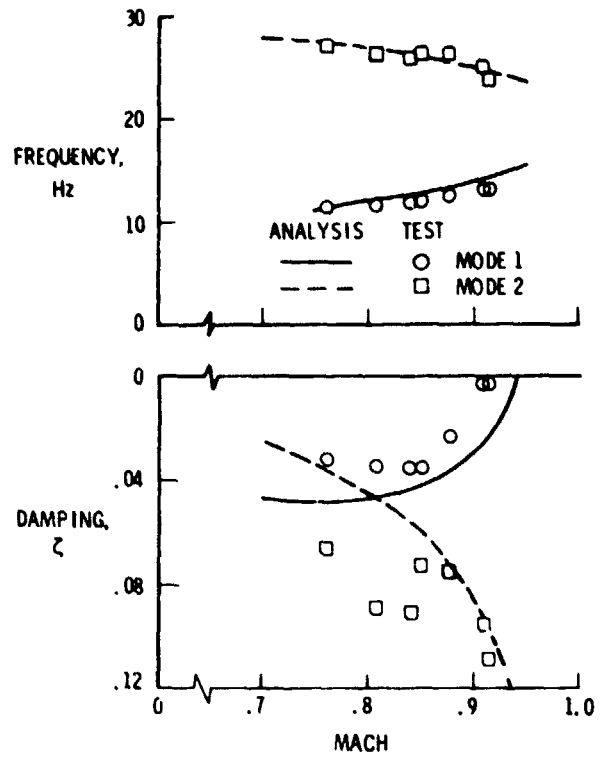


Fig. 8 Damping and frequency variations with Mach number at 7.62 km (symmetric; FSS-off).

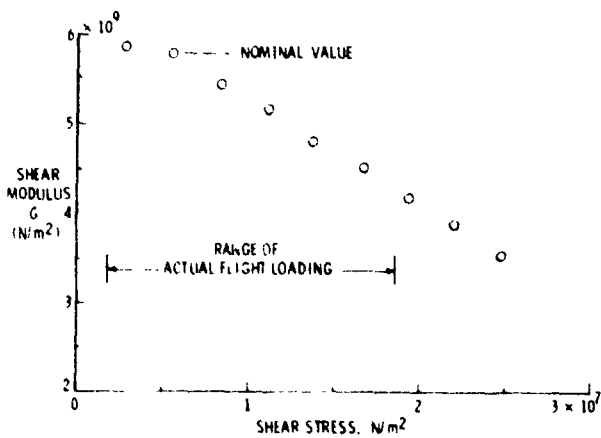


Fig. 7 Shear modulus variation with shear stress.

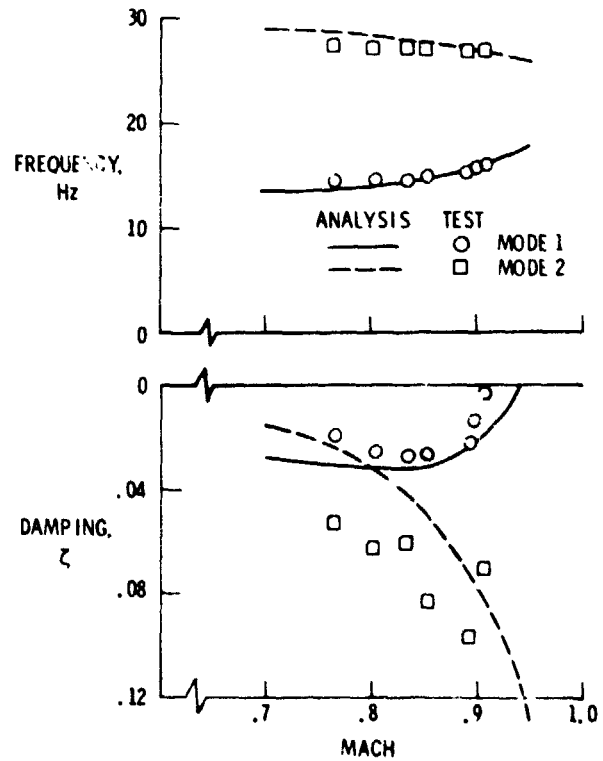


Fig. 9 Damping and frequency variations with Mach number at 7.62 km (antisymmetric; FSS-off).

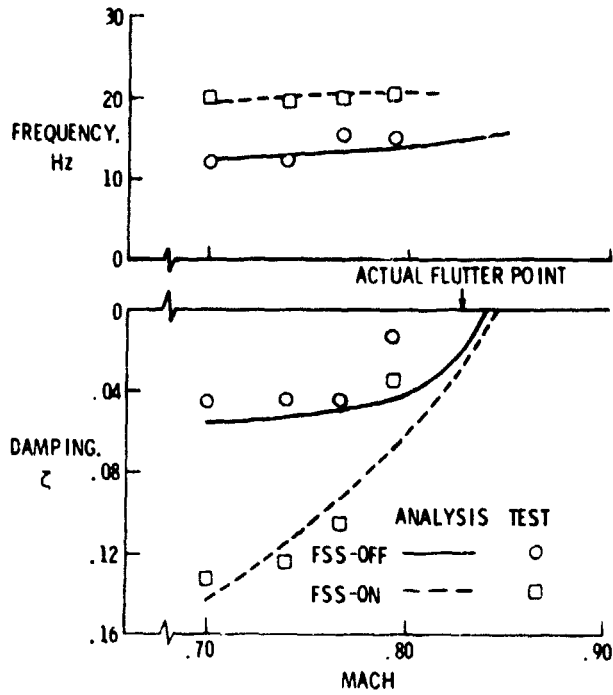


Fig. 10 Damping and frequency variations with Mach number at 4.57 km (symmetric).

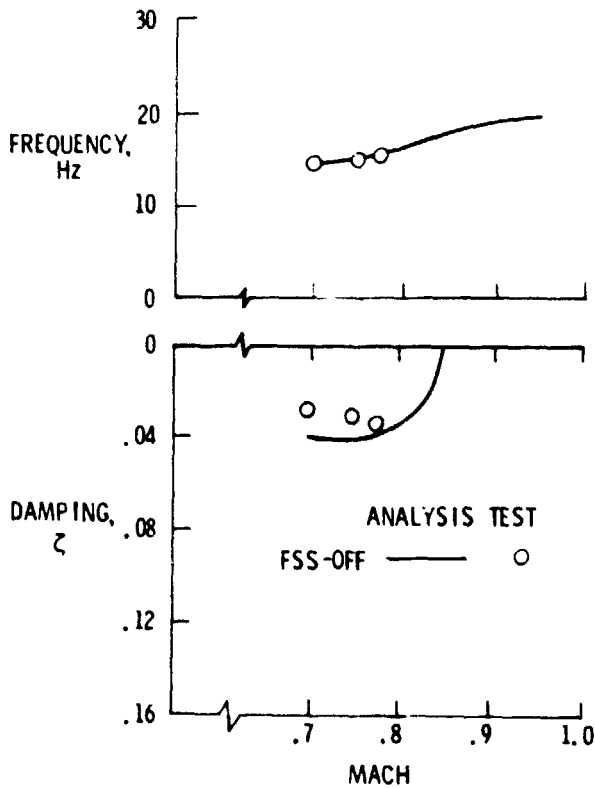


Fig. 11 Damping and frequency variations with Mach number at 4.57 km (antisymmetric).

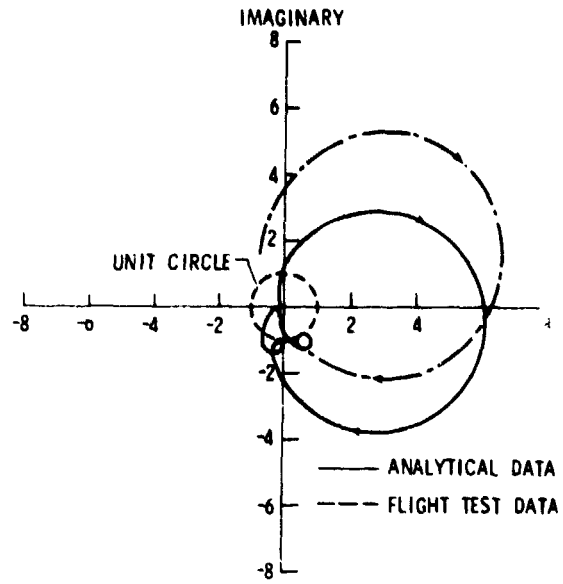


Fig. 12 Nyquist diagram of open loop system at $M = 0.70$ and 4.57 km (symmetric).

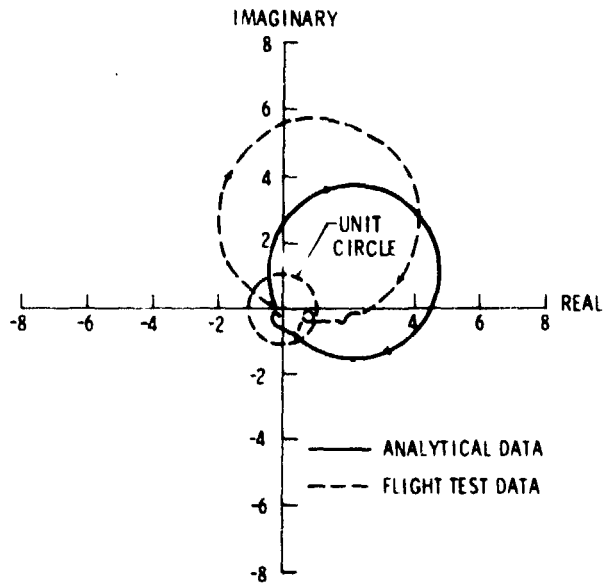


Fig. 13 Nyquist diagram of open loop system at $M = 0.70$ and 4.57 km (antisymmetric).