Computer Program for Pulsed Thermocouples With Corrections for Radiation Effects

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Summary

A pulsed thermocouple is used for measuring gas temperatures above the melting point of common thermocouples. This is done by allowing the thermocouple to heat until it approaches its melting point and then turning on the protective cooling gas. This method requires a computer to extrapolate the thermocouple data to the higher gas temperatures. In earlier work by this author the extrapolation was done by using a first-order exponential curve fit to predict the final thermocouple wire temperature. Since radiation effects were neglected, the gas temperature was not computed. Hand calculations had to be used to estimate the gas temperature. This report describes a method that includes the effect of radiation in the extrapolation. Computations of gas temperature are provided, along with the estimate of the final thermocouple wire temperature. Results from tests on high-temperature combustor research rigs are presented.

Introduction

An earlier investigation by the author (ref. 1) described the use of a pulsed thermocouple to measure gas temperatures above the melting point of common thermocouples. This method of measuring temperature is intended for the measurement of temperatures at the exit of experimental aircraft combustors at temperatures to \(2400\) K and pressures to \(4\) MPa (40 atm). The previous investigation described an approach that uses a thermocouple cooled by a small jet of inert gas. When a measurement is to be made, the cooling gas is turned off and the thermocouple output is sampled at a high rate and recorded. Just before the thermocouple reaches its melting point the cooling is reapplied to protect the thermocouple wire. The gas temperature can then be calculated by extrapolation from the initial heating curve. For the extrapolation to be valid, it must be based on a theoretical heating curve. The derivation of the theoretical equation is described here. All symbols are defined in appendix A.

The equation that describes the pulsed-thermocouple wire temperature can be derived from the basic heat transfer relations (ref. 3). Assume a bare wire thermocouple with infinitely long leads in a hot gas stream. This assumption causes the conduction effects to be neglected. Very little error is introduced if we neglect the transfer of heat to the junction by conduction along the wire for carefully designed probes. Thus in the absence of conduction, heat can be transferred to the wire by convection of the gas, by radiation from the gas, and by radiation from the duct walls. Also heat can be transferred away from the wire by radiation.

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The rate of heat storage in the wire will be equal to the rate of heat entering the wire minus the rate of the heat leaving the wire. The rate of heat storage \(q_s\) per
practice the $P_{sc}$ must be determined experimentally and generally falls in the range 0.8 to 1.0.

The rate of heat transfer by radiation $q_r$ is given by (ref. 3)

$$q_r = \sigma \varepsilon_w [(1 - \alpha) T_d^4 + \varepsilon_g T_g^4 - T_w^4] \pi D$$  \hspace{1cm} (4)

where $\sigma$ is the Stefan-Boltzmann constant, $\varepsilon_w$ is the emissivity of the wire, $\alpha$ is the effective absorptivity of the gas, $T_d$ is the duct temperature, and $\varepsilon_g$ is the emissivity of the gas. The first term in equation (4) represents the heat received by the wire from the hot walls of the duct. The second term represents radiant heat received from the gas. The third term represents radiant heat emitted from the wire.

Combining equations (1) to (4) gives

$$dt = \frac{-K_1 dT_w}{T_w^4 + K_2 T_w - K_3}$$  \hspace{1cm} (5)

where

$$K_1 = \frac{\rho CD}{4\sigma \varepsilon_w}$$  \hspace{1cm} (6)

$$K_2 = \frac{\text{Nu} K_g P_{sc}}{D \sigma \varepsilon_w}$$  \hspace{1cm} (7)

and

$$K_3 = K_2 T_g + [(1 - \alpha) T_d^4 + \varepsilon_g T_g^4]$$  \hspace{1cm} (8)

To solve equation (5), we integrate both sides of the equation. The integration is easier if we factor the denominator. The roots of a fourth-order equation can be found by algebraic methods (ref. 4). The roots of the equation are

$$T_w = \alpha_1 \pm i \beta, \alpha_2, \alpha_3$$  \hspace{1cm} (9)

where

$$\alpha_1 = \frac{1}{2} \sqrt{Y_1}$$  \hspace{1cm} (10)

$$\beta = \frac{1}{2} \left( Y_1 + 2K_2 \right)^{1/2}$$  \hspace{1cm} (11)

where $Y_1 = \frac{K_2}{\text{Nu} K_g P_{sc}}$. For an infinitely long wire in crossflow $P_{sc}$ is unity. The probe shape constant was introduced to take into account the fact that the presence of a probe to support the wire will cause a reduction in the effective Nusselt number of the thermocouple. In
\[ \alpha_2 = \frac{1}{2} \left[ \sqrt{Y_1} + \left( -Y_1 + \frac{2K_2}{\sqrt{Y_1}} \right) \right]^{1/2} \]

(12)

\[ \alpha_3 = -\frac{1}{2} \left[ \sqrt{Y_1} + \left( -Y_1 + \frac{2K_2}{\sqrt{Y_1}} \right) \right]^{1/2} \]

(13)

\[ Y_1 = \left[ \frac{K_3^2}{2} + \left( \frac{K_4^2}{4} + \frac{64K_3^3}{27} \right) \right]^{1/3} \]

\[ + \left[ \frac{K_2^2}{2} - \left( \frac{K_4^2}{4} + \frac{64K_3^3}{27} \right) \right]^{1/3} \]

(14)

Equation (5) can then be rewritten as

\[
\frac{d}{dt} = \frac{-K_1}{T_w - \alpha_1 - i\beta}(T_w - \alpha_1 + i\beta)(T_w - \alpha_2)(T_w - \alpha_3)
\]

(15)

or

\[
dt = \frac{H_1}{T_w - \alpha_2} + \frac{H_2}{T_w - \alpha_3}
\]

\[ + \left( \frac{H_3}{T_w - \alpha_1 - i\beta} + \frac{H_3^*}{T_w - \alpha_1 + i\beta} \right) dT_w \]

(16)

where

\[ H_1 = \frac{-K_1}{(\alpha_2 - \alpha_3)(\alpha_2 - \alpha_1)^2 + \beta^2} \]

(17)

\[ H_2 = \frac{-K_1}{(\alpha_3 - \alpha_2)(\alpha_3 - \alpha_1)^2 + \beta^2} \]

(18)

\[ H_3 = \frac{-K_1}{(\alpha_1 - \alpha_2 + i\beta)(\alpha_1 - \alpha_3 + i\beta)(2i\beta)} \]

(19)

If the denominator of \( H_3 \) is multiplied out, we get

\[ H_3 = \frac{-K_1}{E + iF} \]

(20)

\[ E = -2\beta^2(2\alpha_1 - \alpha_2 - \alpha_3) \]

(21)

and

\[ F = 2[\beta(\alpha_1 - \alpha_2)(\alpha_1 - \alpha_3) - \beta^3] \]

(22)

Thus \( H_3 \) can be rewritten as

\[ H_3 = \frac{-K_1E}{E^2 + F^2} + i \frac{K_1F}{E^2 + F^2} \]

\[ = H_{3A} + iH_{3B} \]

(23)

Equation (16) can then be integrated to get

\[ t = H_1\ln(\alpha_2 - T_w) + H_2\ln(T_w - \alpha_3) \]

\[ + H_{3A}\ln[(T_w - \alpha_1)^2 + \beta^2] \]

\[ + 2H_{3B}\tan^{-1}\left( \frac{\beta}{T_w - \alpha_1} \right) + H_4 \]

(24)

where \( H_4 \) is a constant of integration.

Equation (24) shows the theoretical relationship between the wire temperature \( T_w \) and the time \( t \). In general all the parameters in the equation are known except for the gas temperature \( T_g \), the probe shape constant \( P_{sc} \), and the integration constant \( H_4 \). After a measurement a set of wire temperature readings are known. The procedure used finds the values of \( T_g \), \( P_{sc} \), and \( H_4 \) that result in the best fit of the temperature data to the theoretical equation (eq. 24)). The next section describes the computer program written to fit equation (24) to the data.

Description of Computer Program

The FORTRAN IV computer program described in this report is designed to calculate gas temperature by using data taken from a separate pulsed-thermocouple controller. A listing of the program and its various subroutines is shown in appendix B. The program input requirement is a set of wire temperatures taken at regular time intervals, the Mach number, the total pressure, the wall temperature, and the probe shape constant. The computer program output is the extrapolated wire temperature and the computed gas temperature. In addition, if the probe shape constant has not been entered, the computer program will calculate and output \( P_{sc} \), the probe shape constant.

The program uses a curve-fitting procedure from reference 2 called the gradient-expansion method to
fit the theory to the input data. Two parameters, gas temperature TGAS and possibly probe shape constant PSC, are adjusted for best fit of the theory to the data. These parameters are adjusted until the sum of the squares of the differences between the measured wire temperature and the theoretical wire temperature is a minimum. The error, which is called CHISQR, is defined by

$$\text{CHISQR} = \sum_{i=1}^{n} \left( (T_{\text{data}})_i - (T_{\text{theory}})_i \right)^2$$

(25)

where \( T_{\text{data}} \) is the measured wire temperature, \( T_{\text{theory}} \) is the corresponding theoretical wire temperature, and \( n \) is the number of measured data points. Note that the theoretical wire temperatures must be evaluated point by point at the same values of the time parameter used for the measured data.

Both the gradient-expansion procedure and the evaluation of CHISQR require computation of theoretical wire temperature at every measurement time. In addition, the gradient-expansion method requires values for \( \frac{\partial T_w}{\partial T_g} \) and \( \frac{\partial T_w}{\partial P_{sc}} \) at every measurement time. These requirements create a difficulty because the analytical solution to the differential equation expresses time as a function of wire temperature in equation (24). The equation cannot easily be inverted to yield the needed wire temperature as a function of time and its derivatives. As a result a great amount of the computer time is devoted to numerically inverting the equation and evaluating the derivatives. Since theoretical wire temperature values at the measurement times are not available directly from equation (24), they are calculated by interpolating in a table of wire temperature–time pairs that do satisfy equation (24). This table must be regenerated whenever equation parameters are changed.

This procedure must be repeated once for every evaluation of wire temperature and twice for every evaluation of the derivatives. The derivatives are approximated by computing the differences in wire temperature that result for two values of the parameters TGAS and PSC: one value slightly above the present value and one value slightly below the present value.

The main computer program takes care of reading the input data, calling the curve-fitting routines, deciding when the curve fit is good enough, and writing the results. Initially input data of Mach number, pressure, and duct temperature are read in as well as 1000 readings of thermocouple wire temperature. The temperatures represented by these numbers are taken at equal time intervals before and during the temperature rise. The first 100 readings represent the thermocouple wire temperature while the cooling air is on. The rest of the 900 temperature readings are taken during the temperature rise of the thermocouple wire when the cooling air is turned off. If the cooling air is turned on again before the 900 readings are taken, the remaining readings are zero.

After the data are read in, a call to subroutine STCFIT determines the best estimate of the temperature ramp starting time. This is necessary because the theoretical curve is always forced to pass through this point.

With the starting time determined, the curve-fitting process begins. Repetitive calls to CURFIT and FDERIV result in adjustments to several parameters such that CHISQR is decreased. With every adjustment in the parameter values a call to CONGEN is needed to evaluate the constants in equation (24). The parameters adjusted include the gas temperature TGAS; the probe shape constant PSC; and FLAMDA, a parameter whose value controls the curve-fitting process. The probe shape constant is adjusted only if its value is not included in the input data. If the PSC is to be adjusted, the variable NTERMS is set equal to 2 by the computer program; otherwise NTERMS is set equal to 1 and only TGAS is adjusted. Thus the main program recalls FDERIV and CURFIT until the decrease in CHISQR is less than 1 percent. This value of 1 percent was chosen by trial-and-error methods to provide a wire temperature within 1 or 2 K of the ultimate wire temperature without using an unreasonable amount of computer time.

Subroutine CURFIT

Subroutine CURFIT makes a least-squares fit to a nonlinear function by using the gradient-expansion algorithm described in appendix C. The algorithm is really two curve-fitting techniques combined into one program. One of the techniques works well when the variables are far from the correct values, and the other works well when they are close to the final values. A parameter \( \lambda \) (called FLAMDA in the program) is used to change the curve-fitting routine gradually from one technique to the other.

The subroutine works by starting with FLAMDA = 0.001 (when FLAMDA is less than 1 the fitting technique that works close to the minimum is dominant—see appendix C). The error \( \chi^2 \) (appendix C) between the measured and theoretical data is called both CHISQI and CHISQR in the program. CHISQI is an initial value of \( \chi^2 \) calculated once when the subroutine is entered. The program makes changes in the wire temperature, the probe shape constant, and FLAMDA until a new value of \( \chi^2 \) (called CHISQR) starts to decrease, at which time FLAMDA is divided by 10 and the subroutine returns to the
calling program. It is the responsibility of the calling program to check CHISQR to see if the change in CHISQR since the last call to CURFIT is small enough to stop the program. If it is not, subroutine CURFIT should be called again without changing the value of the current FLAMDA.

Subroutine FDERIV

Subroutine FDERIV computes data needed by the curve-fitting routine CURFIT. The data needed are the derivatives of the wire temperature with respect to both gas temperature and the probe shape constant. Also needed are theoretical values of wire temperature evaluated at the measured time (the times corresponding to the measured wire temperatures). The derivatives are determined from (ref. 5)

\[
\frac{\partial T_w}{\partial T_g} = \frac{T_w(T_g + 1, \text{PSC}) - T_w(T_g - 1, \text{PSC})}{2} \quad (26)
\]

\[
\frac{\partial T_w}{\partial \text{PSC}} = \frac{T_w(T_g, \text{PSC} + 0.001) - T_w(T_g, \text{PSC} - 0.001)}{2(0.001)} \quad (27)
\]

If the probe shape constant is not to be calculated (NTERMS = 1), only equation (26) will be calculated. The theoretical values of wire temperature are generated from equation (24) with a call to subroutines TABL and INTRP.

The subroutine returns a 1000-by 3-element array. The derivative of the wire temperature with respect to the gas temperature at time \(I\) is returned in array DERIV(I,1). The derivative of the wire temperature with respect to the probe shape constant at time \(I\) is returned in array DERIV(I,2). The table of the computed wire temperatures at time \(I\) is returned in array DERIV(I,3).

Function XICALC

Function XICALC computes the sum of the squares of the differences between the measured wire temperature and the theoretical wire temperature (from the numerically inverted equation (24)). The sum of the squares of the differences will be

\[
XICALC = \sum_{I=\text{START}}^{\text{RANGE}} \left[ T_w(I) - T_{\text{theory}}(I) \right]^2 \quad (28)
\]

The program first calls subroutine CONGEN to generate new constants for equation (24) since the gas temperature and the probe shape constant may have changed. Subroutine TABL is then called to generate a table of theoretical temperatures and times. The interpolation necessary is done by this subroutine and not by subroutine INTRP because the output of this routine is a single number, the error XICALC, and not an entire table of numbers.

Subroutine TABL

The purpose of subroutine TABL is to generate values of theoretical wire temperatures and times for subroutine INTRP. Subroutine CURFIT, FDERIV, and function XICALC require a value of theoretical wire temperature at every measurement time. These wire temperatures must be obtained by inverting equation (24). However, because of the form of equation (24) a numerical inversion will have to be done. A call to subroutine TABL generates a table of temperature-time pairs that satisfy equation (24). Then a call to INTRP interpolates in this table to get temperatures at the measurement times.

To generate the interpolation table, a set of temperatures is needed to put into equation (24) to obtain computed times. The values of computed time that result from equation (24) should be as close as possible to the measured times for accurate interpolation by subroutine INTRP. The set of temperatures is determined one at a time, starting with a known point on the theoretical curve. Each succeeding temperature is computed from the previous one by using a linear approximation to the theoretical curve (fig. 2). The linear approximation will have a slope equal to the slope of the theoretical

![Figure 2](image-url)
curve at the previous temperature. Thus each succeeding temperature will be

$$T_{j+1} = T_j + \frac{t_{j+1} - t_j}{(dT/dT_w)_{T_w=T_j}}$$ (29)

where $j=1,2,3,\ldots, n$ measured data points. The times corresponding to the measured data points are $t_j$. The times $t_j'$ are computed by evaluating equation (24) with $T_w=T_j$. The derivative of equation (24) is

$$\frac{dt}{dT_w} = \frac{-H_1}{\alpha_2 - T_w} + \frac{H_2}{T_w - \alpha_3} + \frac{2[H_3A(T_w - \alpha_1) - H_3B\beta]}{(T_w - \alpha_1)^2 + \beta^2}$$ (30)

In the program $t_{j+1} - t_j'$ is defined as DELTIM and

$$\text{DELTMP} = \frac{\text{DELTIM}}{(dT/dT_w)_{T_w=T_j}}$$ (31)

The program starts by setting $T_j = T_1 = \text{TAVE}$, which is the temperature on the theoretical curve; and $t_j = t_1'$ is equal to MSTIME*START. The next temperature $T_{j+1}$ is evaluated by setting $t_{j+1} = t_2 = \text{MSTIME*(START+1)}$ in equation (29). What results is a table of theoretical time-temperature pairs that do satisfy equation (24), where the times are not exactly equal to the measurement times. The array of times is called TIMC, and the array of temperatures is called TC in the program. A linear interpolation will need to be done because temperatures at the exact measurement times are needed.

Subroutine INTRP

Subroutine INTRP is used to correct the table of theoretical temperatures (array TC) generated by subroutine TABL. Subroutine INTRP performs a linear interpolation between the calculated data points so that the calculated times (and corresponding temperatures) fall exactly on the measured time. The resulting interpolated values of temperature are stored in array TC.

Subroutine STCFIT

Subroutine STCFIT determines the starting point of the thermocouple temperature rise. The starting point is defined as the intersection of two straight lines. One line is the best fit through the data before the cooling is turned off. This line is called TAVE. The other line is the best fit through approximately the first 50 points of the temperature rise. Since a solenoid is used to turn the cooling air on and off, there will be some delay between when the power is removed and when the cooling air actually stops flowing. The solenoid power is turned off at data point 100, and the starting point search ranges between data points 100 and 130.

The starting point of the search process is shown in figure 3. A standard least-squares fit to a straight line of the data from point 100 to point 160 is performed. In general, point 100 is not the true starting point; so this line (line 1 in fig. 3) will not intersect the TAVE line at point 100. In fact, if the starting point of the data for the least-squares line is varied from 100 to 130, the intersection of the least-squares line (line 2) with TAVE will approach the true starting point and then back away. Therefore the intersection point will have a maximum as the starting point is varied. The output of this routine is this maximum value of the starting point. This represents the best approximation to the start of the ramp.

Subroutine CONGEN

Subroutine CONGEN computes the constants necessary to evaluate equation (24). Constants $K_1$, $K_2$, and $K_3$ are evaluated by using equations (6) to (8). The wire emissivity $\epsilon_w$ for clean platinum was found to be (ref. 6)

$$\epsilon_w = 0.085 + (0.76 \times 10^{-4})T_{wf}$$ (32)

where $T_{wf}$ is the final wire temperature in K. The other parameters used for platinum (type R) thermocouple wire are (ref. 7):
Wire density, kg/m³ .......................... 0.2078 × 10⁵
Stefan-Boltzmann constant, J/K⁴ sec m² .......................... 0.56697 × 10⁻⁷
Wire specific heat, J/kg K .................. 0.1427 × 10⁻³
Gas effective absorptivity ...................... 0
Gas effective emissivity ........................ 0
Wire diameter, m .......................... 0.8128 × (33)

\[ K_g = (0.3007 \times 10^{-3}) \times \text{TG}_\text{A}^{0.78} \text{J/(sec K m)} \]  (33)

\[ \text{Nu} = 188.41 \times (\sqrt{\text{WDIA} \times \text{MN} \times \text{P}}) \times \text{TG}_\text{A}^{-0.6} \\
\times \left[ 1 + 0.2 \times (\text{MN})^2 \right]^{-1/4} \]  (34)

where WDIA is the wire diameter, MN is the Mach number, P is the pressure in pascals, and TG$_A$ is in K. The gas effective absorptivity and emissivity are assumed to be zero. This corresponds to a transparent gas and the worst case for radiation effects.

The subroutine also computes $\alpha_1$, $\alpha_2$, $\alpha_3$, $\beta$, $H_1$, $H_2$, $H_3$A, and $H_3$B from equations (10) to (23). The value of $H_4$ is computed by putting the initial conditions into equation (24) and solving for $H_4$. The initial temperature is the average cooled temperature TAVE. The initial time is the measurement time interval MSTIME times START.

**Function EVALTM**

Function EVALTM evaluates equation (24) to obtain a calculated time for an input of wire temperature. The input wire temperature must be between the initial average cooled temperature TAVE and $\alpha_2$ in order to avoid taking the logarithm of a negative number. Values of $\alpha_1$, $\alpha_2$, $\alpha_3$, $\beta$, $H_1$, $H_2$, $H_3$A, and $H_4$ must have been previously calculated with a call to the CONGEN subroutine.

**Subroutine MATINV**

Subroutine MATINV does an inversion of a 1- or 2-degree matrix. For a 1-degree matrix only a simple reciprocal is needed. For a 2-degree matrix the adjoint matrix is calculated. Then each element is divided by the determinant to form the inverse matrix. The original matrix is then replaced by its inverse.

**Tests and Results**

A pulsed-thermocouple system was tested in a combustor rig at the Air Force Wright Aeronautical Laboratory (AFWAL) as part of a joint AF-NASA program on instrumentation. The system included a probe, a sample-and-hold voltmeter, a microcomputer-based controller, and a digital recorder, as shown in figure 4. Figure 5 shows the probe that was put into the combustor. The probe consisted of a water-cooled shell with a replaceable platinum (type R) thermocouple. Compressed-air cooling for the thermocouple was controlled by a fast-acting solenoid valve. The thermocouple voltage was converted to digital form by a sample-and-hold digital voltmeter. A microcomputer was used to control the voltmeter and turn the cooling air on and off. The time between data points (called MSTIME) was controlled at 0.0042 second. This value was chosen so that most of the ramp would be included in the 1000 data points. If a different probe with a different time constant were used, this MSTIME would have to be changed.

A full curve including the final wire temperature could be recorded for each pulse because the gas stream of the combustor configuration under test was not hot enough to require the cooling air to come on. The data were first processed by the computer program to compute the probe shape constant. The average computed probe shape constant for 20 pulses at fixed combustor conditions was 0.91, with a maximum deviation of 0.09. This deviation is the result of the fact that the burning process is not constant during the pulse and thus results in a temperature that can vary during the pulse by as much as 3.2 percent.

With the average probe shape constant of 0.91 the data were curve fit 60 percent of the way up the

![Figure 4. Block diagram of pulsed-thermocouple system.](image-url)
curve. It is estimated that at least 60 percent of the curve could be measured at the highest expected gas temperatures. A typical result is shown in figure 6. The solid line is the measured data (a total of 1000 data points). The triangles and squares represent the theoretical curve. The squares represent the portion of the curve that was used in the computation. The triangles represent the portion of the curve that was extrapolated by using the theoretical curve. The computed final wire temperature varied from 1525 K to 1581 K, with an average of 1561 K for the 20
readings. The actual final wire temperature varied from 1525 K to 1575 K because of fluctuations in the burning. A comparison between the final wire temperature computed using 60 percent of the ramp and the actual final wire temperature measured for the 20 readings showed a maximum deviation of 3 percent.

The average of the 20 computed gas temperatures was 1691 K, with a maximum deviation of 47 K, or 2.7 percent. The difference of 130 K between the computed wire temperature and the gas temperature is the radiation error. It is estimated that the radiation error can be computed to within about 20 percent, which for this case would be ±26 K.

Results for a pulsed-thermocouple probe different from the probe just described were obtained during a high-temperature combustor test at the Lewis Research Center as shown in figure 7. The probe shape constant for this geometry was determined at lower temperatures than shown in figure 7 to be 0.96. The gas temperature for the data shown in figure 7 was 2300 K, and the final computed wire temperature was 2190 K. The wire melts at 2215 K. The protective compressed air was set to turn on at about 2000 K in order to assure a long thermocouple life.

Concluding Remarks

The pulsed thermocouple was developed as an instrument to determine high gas temperatures. The pulsed feature is needed at temperatures above the melting point of common thermocouples or when streaking of a combustion process is occurring. The cooling gas was found to adequately protect the thermocouple during this high-temperature operation.

The computer program for computing gas temperature was designed to take the \( T^4 \) radiation error into account. The program requires as input the Mach number, the wall temperature, and the total pressure in addition to the thermocouple data. Tests at temperatures below the melting point of platinum thermocouples show that the pulsed-thermocouple system can compute the gas temperature to within about 4 percent with as little as 60 percent of the temperature step as input data.

Lewis Research Center
National Aeronautics and Space Administration
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# Appendix A

## Symbols

<table>
<thead>
<tr>
<th>Mathematical symbol</th>
<th>Computer symbol</th>
<th>Definition</th>
</tr>
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<tbody>
<tr>
<td>$a$</td>
<td>-----</td>
<td>parameter of function $\chi^2$</td>
</tr>
<tr>
<td>$C$</td>
<td>SPHT</td>
<td>specific heat of wire</td>
</tr>
<tr>
<td>$D$</td>
<td>WDIA</td>
<td>wire diameter, m</td>
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<tr>
<td>$H_1,H_2,H_3A$, $H_3B,H_4$</td>
<td>H1,H2,H3A, H3B H4</td>
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<td>-----</td>
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<td>MN</td>
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<td>PSC</td>
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<tr>
<td>$q_c$</td>
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<td>rate of heat transferred by convection into surface of wire, J/sec m</td>
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<td>rate of heat transferred by radiation, J/sec m</td>
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<td>$q_s$</td>
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</tr>
<tr>
<td>$T_w$</td>
<td>TWIRE</td>
<td>wire temperature</td>
</tr>
<tr>
<td>$T_{wf}$</td>
<td>TWF</td>
<td>final wire temperature</td>
</tr>
<tr>
<td>$t$</td>
<td>-----</td>
<td>time</td>
</tr>
<tr>
<td>$X,Y$</td>
<td>-----</td>
<td>general independent variables</td>
</tr>
<tr>
<td>$Y_1$</td>
<td>-----</td>
<td>intermediate constant</td>
</tr>
<tr>
<td>$\alpha$</td>
<td>ALPHAG</td>
<td>effective absorptivity of gas</td>
</tr>
<tr>
<td>$\alpha_1,\alpha_2$, $\alpha_3$</td>
<td>ALPHA1, ALPHA2, ALPHA3</td>
<td>intermediate constants</td>
</tr>
<tr>
<td>$\beta$</td>
<td>BETA</td>
<td>intermediate constant</td>
</tr>
<tr>
<td>$\epsilon_g$</td>
<td>EGAS</td>
<td>emissivity of gas</td>
</tr>
<tr>
<td>$\epsilon_w$</td>
<td>E1+E2*T</td>
<td>emissivity of wire</td>
</tr>
<tr>
<td>$\sigma$</td>
<td>SIGMA</td>
<td>Stefan-Boltzmann constant, J/K$^4$ sec m$^2$</td>
</tr>
<tr>
<td>$\chi^2$</td>
<td>CHISQR, CHISQ1</td>
<td>least-squares error</td>
</tr>
</tbody>
</table>
Appendix B
Computer Programs

ROUTINE FOR CURVE FITTING DATA FROM A PULSED THERMOCOUPLE.
THERMOCOUPLE DATA SHOULD BE CONSTANT FOR THE FIRST 100
DATA POINTS (THERMOCOUPLE COOLED). THE THERMOCOUPLE IS THEN
HEATED ON AN EXPONENTIAL HEATING CURVE (900 DATA POINTS).
THE PROGRAM NEGLECTS CONDUCTION ERRORS.

INPUTS REQUIRED ARE:
WIRE TEMPERATURE (1000 DATA POINTS) (K),
MACH NUMBER,
PRESSURE (Pa),
DUCT TEMPERATURE (K),
PROBE SHAPE CONSTANT,

ALSO THE FOLLOWING PARAMETERS MUST BE SET TO
THE PROPER VALUE DEPENDING ON THE TYPE OF
THERMOCOUPLE USED:
MSTIME = TIME BETWEEN MEASUREMENTS (THIS PROGRAM),
WDIA = WIRE DIAMETER, (SUBROUTINE CONGEN),
WDENS = WIRE DENSITY. (SUBROUTINE CONGEN),
SPHT = WIRE SPECIFIC HEAT, (SUBROUTINE CONGEN),
WIRE EMISSIVITY, (SUBROUTINE CONGEN),
EGAS = EMITTIVITY OF GAS, (SUBROUTINE CONGEN),
ALPHAG = ABSORPTIVITY OF GAS, (SUBROUTINE CONGEN),

REAL TWIRE(1000),TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMP

INTEGER START,RANGE

REAL TC(1000,2),TIMC(1000,2),DERIV(1000,3)

DATA MSTIME/0.42E-2/
MSTIME IS IN SECONDS.

INTEGER I,INTERMS
REAL CHISQO,FLAMDA,TGAS,TWF,PSC,CHISQR,X

COMMON /BLK1/TWIRE,TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMP

COMMON /BLK2/START,RANGE
COMMON /BLK3/TC,TMC,DERIV

C
C
C READ TEMPERATURE DATA (TWIRE)
C IN DEG. K. (1000 DATA POINTS).
C
DO 20 I=1,1000
    20 READ(1,80) TWIRE(I)
C
C READ INPUT DATA
C
WRITE(7,30)
    30 FORMAT(1X,'INPUT MACH NUMBER')
    READ(5,40) MN
WRITE(7,50)
    50 FORMAT(1X,'INPUT PRESSURE IN Pa.')
    READ(5,60) P
WRITE(7,70)
    70 FORMAT(1X,'INPUT DUCT TEMPERATURE IN DEG. K.')
    READ(5,80) TDUCT
WRITE(7,90)
    90 FORMAT(1X,'INPUT PROBE SHAPE CONSTANT.')
    READ(5,100) PSC
    100 FORMAT(F6.3)

C THE NEXT 3 STATEMENTS ARE NEEDED ONLY
C FOR THE EXAMPLE IN THIS REPORT.
C
WRITE(7,101)
    101 FORMAT(1X,'MACH NUMBER TEMPERATURE DEG. K.')
    READ(5,80) MTMP
C
C AVERAGE COOLED WIRE TEMPERATURE.
C
TAVE = 0.0
DO 110 I=1,99
    110 TAVE = TAVE+TWIRE(I)
    TAVE = TAVE/99.
C
C DETERMINE START OF RAMP.
C
115 CALL STCFIT
C
C TEMPERATURE OVER MELTING POINT?
C
118 DO 120 RANGE=100,1000
    IF (TWIRE(RANGE),LE,400.) GO TO 130
    120 CONTINUE
130   RANGE = RANGE-1
C   CURVE FIT.
C   CHISQR = 0,
C   FLAMDA = 0.001
C   TWF = TWIRE(RANGE)
C   TGAS = TWF
C
135   IF (PSC.NE.0.) GO TO 140
C   NTERMS = 2
C   PSC = 0.8
C
C   "FDERIV" COMPUTES THE DERIVATIVE OF TWIRE WITH RESPECT
C   TO TGAS & PSC. ALSO IT RETURNS VALUES OF CALCULATED
C   THEORETICAL WIRE TEMPERATURE AS A FUNCTION OF TIME.
C
140   CALL FDERIV(TGAS,PSC,NTERMS,TWF)
C
C   "CURFIT" MODIFIES TGAS AND PSC TO OBTAIN THE BEST
C   MATCH BETWEEN THE THEORETICAL CURVE AND THE ACTUAL DATA.
C
145   CALL CURFIT(NTERMS,PSC,TGAS,CHISQR,FLAMDA,TWF)
C
C   "CHISQR" IS THE ERROR BETWEEN THE THEORETICAL CURVE AND
C   THE ACTUAL MEASURED DATA. IF THERE IS LESS THAN A
C   ONE PERCENT CHANGE IN THE ERROR SINCE THE LAST
C   CALL TO CURFIT THEN THE PROGRAM IS FINISHED.
C
150   X = ABS((CHISQR-CHISQR)/CHISQR)
C   IF (X.LT.0.01) GO TO 150
C   CHISQR = CHISQR
C   TWF = ALPHA2
C   GO TO 140
155   WRITE(7,160) TGAS
160   FORMAT(1X,'GAS TEMPERATURE = ',F9.2,' K')
165   WRITE(7,170) ALPHA2
170   FORMAT(1X,'FINAL WIRE TEMPERATURE = ',F9.2,' K')
175   WRITE(7,180) PSC
180   FORMAT(1X,'PROBE SHAPE CONSTANT = ',F6.3)
185   STOP 123
END

C
C   SUBROUTINE CURFIT(NTERMS,PSC,CHISQR,FLAMDA,TWF)
C
PURPOSE

This subroutine makes a least squares curve fit to a non-linear function.

TIME = set of integers taken as independent variable.
TWIRE = array of wire temperature readings taken as independent variable.
START = integer value of time for start of data.
RANGE = integer value of time for end of data.
NTERMS = number of parameters (max. = 2).
TGAS = parameter 1: gas temperature.
PSC = parameter 2: probe shape constant.
A = array of parameters.
FLAMDA = proportion of gradient search included.
TWF = estimated final wire temperature.
CHISQR = chi square for fit.

SUBROUTINE CURFIT(NTERMS,PSC,TGAS,CHISQR,FLAMDA,TWF)

REAL TWIRE(1000),TAVE,ALPHA1,ALPHA2,ALPHA3,BETA, 
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMP

INTEGER START,RANGE

REAL TC(1000,2),TIMC(1000,2),DERIV(1000,3)

REAL BE(2),AL(2,2),PSC,TGAS,CHISQR1,CHISQR,FLAMDA,TWF, 
1 A(2),B(2),ARRAY(2,2)

INTEGER I,J,K,NTERMS,ERFLAG

COMMON /BLK1/TWIRE,TAVE,ALPHA1,ALPHA2,ALPHA3,BETA, 
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMP

COMMON /BLK2/START,RANGE

COMMON /BLK3/TC,TIMC,DERIV

AL = alpha matrix,
BE = beta matrix.

10 DO 20 J=1,NTERMS
BE(J) = 0.
DO 20 K=1,J
20 AL(J+K) = 0.

TRUNCATE TGAS SINCE SMALL CHANGES IN TGAS
CAUSE UNNECESSARY ITERATION

A(1) = AINT(TGAS)
A(2) = PSC

DO 60 I=START, RANGE
 DO 50 J=1, NTERMS
     BE(J) = BE(J) + (TWIRE(I) - DERIV(I, 3)) * DERIV(I, J)
 DO 40 K=1, J
40   AL(J,K) = AL(J,K) + DERIV(I,J) * DERIV(I,K)
50 CONTINUE
60 CONTINUE

DO 70 J=1, NTERMS
 DO 70 K=1, J
70 AL(K,J) = AL(J,K)

EVALUATE CHI SQUARE AT STARTING POINT

CHISQ1 = XICALC(A(1), A(2), ERFLAG, TWF)

DO 100 J=1, NTERMS
 DO 100 K=1, NTERMS

CALL MATINV(ARRAY, NTERMS)

B(2) = A(2)

DO 120 J=1, NTERMS
 B(J) = A(J)
 DO 120 K=1, NTERMS
120 B(J) = B(J) + BE(K) * ARRAY(J,K)

TRUNCATE B(1) & B(2) TO CONSIDER ONLY INTEGER VALUES
OF TEMPERATURE AND ONLY 2 SIGNIFICANT FIGURES FOR PSC.

B(1) = AINT(B(1))
B(2) = B(2) * 100.
B(2) = AINT(B(2))
B(2) = B(2) / 100.

CALL CHISQR FOR NEW PARAMETER VALUES.

CHISQR = XICALC(B(1), B(2), ERFLAG, TWF)

ERFLAG = 6 IF ALPHA2 IS TOO LOW.

IF (ERFLAG.EQ.6) GO TO 140
IF (CHISQ1 - CHISQR) 140, 150, 150
IF CHISQR INCREASED, INCREASE FLAMDA.

FLAMDA = 10.0 * FLAMDA
GO TO 80

IF CHISQR DECREASED, DECREASE FLAMDA & SET NEW
VALUES FOR TGAS & PSC.

TGAS = B(1)
PSC = B(2)
FLAMDA = FLAMDA / 10.

RETURN
END

SUBROUTINE FDERIV(TGAS, PSC, NTERMS, TWF)

PURPOSE
COMPUTE THE DERIVATIVE OF TWIRE WITH RESPECT TO BOTH
TGAS & PSC AND ALSO EVALUATE THE THEORETICAL EQUATION,

TGAS = ESTIMATED GAS TEMPERATURE (K),
PSC = PROBE SHAPE CONSTANT,
NTERMS = NUMBER OF TERMS (1 OR 2),
TWF = ESTIMATED FINAL WIRE TEMPERATURE (K).

COMMENTS
THIS PROGRAM COMPUTES THE DERIVATIVE OF TWIRE WITH
RESPECT TO TGAS AND STORES THE VALUES IN THE ARRAY
DERIV(I, 1) WHERE I = 1, 1000, THE DERIVATIVE OF TWIRE
WITH RESPECT TO PSC IS STORED IN DERIV(I, 2), THE
THEORETICAL EQUATION EVALUATED AT EACH MEASURED TIME
[MSTIME * (TIME INTEGER)] IS STORED IN DERIV(I, 3).

SUBROUTINE FDERIV(TGAS, PSC, NTERMS, TWF)

REAL TWIRE(1000), TAVE, ALPHA1, ALPHA2, ALPHA3, BETA,
1 H1, H2, H3A, H3B, H4, TDUCT, MSTIME, MN, P, MNN, MTMP

INTEGER START, RANGE
REAL TC(1000,2),TIMC(1000,2),DERIV(1000,3)
C
REAL DELTA(2),T,PC,TGAS,PSC,X,Y,Z,TWF
C
INTEGER F,FF,ERFLAG,I,L1
C
DATA DELTA/1.0,0.001/
C
COMMON /BLK1/TWIRE,TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
1   H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMP
C
COMMON /BLK2/START,RANGE
C
COMMON /BLK3/TC,TIMC,DERIV
C
C
COMPUTE DATA FOR DERIVATIVE OF TWIRE WITH
RESPECT TO TGAS.
C
COMPUTE DATA POINTS FOR T=TGAS+DELTA(1)
AND PC=PSC.
C
10  T = TGAS + DELTA(1)
PC = PSC
C
GENERATE CONSTANTS.
C
CALL CONGEN(T,PC,ERFLAG,TWF)
IF (ERFLAG.NE.0) STOP 997
C
GENERATE A TABLE OF COMPUTED TIMES AND TEMPERATURES.
TABL PUTS DATA INTO TC AND TIMC
C
CALL TABL(1)
C
COMPUTE DATA POINTS FOR T=TGAS-DELTA(1)
AND PC=PSC.
C
20  T = TGAS - DELTA(1)
PC = PSC
C
GENERATE CONSTANTS.
C
CALL CONGEN(T,PC,ERFLAG,TWF)
IF (ERFLAG.NE.0) STOP 996
C
C
GENERATE ANOTHER TABLE.
CALL TABL(2)
C
INTERPOLATE BETWEEN DATA POINTS SO THAT THE
CALCULATED TIMES (TIMC) CORRESPOND TO THE
MEASURED TIMES (I*MSTIME).
INTERPOLATE FOR T=TGAS+DELTA(1).
CALL INTRP(1,1)
INTERPOLATE FOR T=TGAS-DELTA(1).
CALL INTRP(2,1)
CALCULATE DERIVATIVE OF TWIRE WITH RESPECT TO Tgas AND STORE IT IN DERIV(I,1).
DERIV(START,1) = 0.
L1 = START + 1
DO 35 I=L1,RANGE
DERIV(I,1) = (TC(I,1)-TC(I,2))/(2.*DELTA(1))
CONTINUE
IF (NTERMS.EQ.1) GO TO 60

COMPUTE DATA FOR DERIVATIVES OF TWIRE WITH RESPECT TO PSC.

COMPUTE DATA POINTS FOR T=TGAS AND PC=PSC+DELTA(2).

T = TGAS
PC = PSC + DELTA(2)

GENERATE CONSTANTS.
CALL CONGEN(T,PC,ERFLAG,TWF)
IF (ERFLAG.NE.0) STOP 995

GENERATE A TABLE OF COMPUTED TIMES AND TEMPERATURES.
CALL TABL(1)

COMPUTE DATA POINTS FOR T=TGAS AND PC=PSC-DELTA(2),

T = TGAS
PC = PSC - DELTA(2)

GENERATE CONSTANTS.
CALL CONGEN(T,PC,ERFLAG,TWF)
IF (ERFLAG.NE.0) STOP 994

GENERATE ANOTHER TABLE.
CALL TABL(2)
INTERPOLATE BETWEEN DATA POINTS FOR PC=PSC+DELTA(2).
CALL INTRP(1,2)
INTERPOLATE FOR PC=PSC-DELTA(2)
CALL INTRP(2,2)
CALCULATE DERIVATIVE OF TWIRE WITH RESPECT
TO PSC AND STORE IN DERIV(I,2).
50 DERIV(START,2) = 0,
DO 55 I = L1,RANGE
   DERIV(I,2)=(TC(I,1)-TC(I,2))/(2.*DELTA(2))
   CONTINUE
55 GENERATE A TABLE OF ONLY THE FUNCTION (TWIRE VS. TIME),
GENERATE CONSTANTS FOR TGAS & PSC.
60 CALL CONGEN(TGAS,PSC,ERFLAG,TWF)
   IF (ERFLAG,NE,0) STOP 993
C COMPUTE A TABLE.
CALL TABL(1)
INTERPOLATE BETWEEN DATA POINTS
CALL INTRP(1,3)
STORE THE INTERPOLATED FUNCTION (TWIRE VS. TIME)
INTO DERIV(I,3),
70 DO 75 I=1,1000
   DERIV(I,3) = TC(I,1)
75 CONTINUE
RETURN
END

FUNCTION XICALC(TGAS,PSC,ERFLAG,TWF)
PURPOSE
TO COMPUTE CHI SQUARE FOR PRESENT PARAMETER VALUES,
C

TGAS = ESTIMATED GAS TEMPERATURE (K).
PSC = PROBE SHAPE CONSTANT.
TWF = ESTIMATED FINAL WIRE TEMPERATURE (K).

C

REAL TWIRE(1000), TAVE, ALPHA1, ALPHA2, ALPHA3, BETA,
H1, H2, H3A, H3B, H4, TDUCT, MTIME, MN, P, MNN, MTMP

C

INTEGER START, RANGE

C

REAL TC(1000, 2), TIMC(1000, 2), DERIV(1000, 3)

C

REAL X, Y, TCM, ER

C

INTEGER ERFLAG, J, I, L1, L2, K

C

COMMON /BLK1/TWIRE, TAVE, ALPHA1, ALPHA2, ALPHA3, BETA,
H1, H2, H3A, H3B, H4, TDUCT, MTIME, MN, P, MNN, MTMP

C

COMMON /BLK2/START, RANGE

C

COMMON /BLK3/TC, TIMC, DERIV

C

GENERATE CONSTANTS FOR THEORETICAL EQUATION.

C

10 CALL CONGEN(TGAS, PSC, ERFLAG, TWF)
IF (ERFLAG, NE, 0) RETURN

C

GENERATE TABLE OF THEORETICAL DATA POINTS OF
TEMPERATURE VS. TIME. THE VALUES OF TIME ARE ONLY
APPROXIMATELY EQUAL TO THE ACTUAL MEASURED TIMES.

C

CALL TABL(1)
J = START
L1 = START+1
L2 = RANGE+1
ER = 0.

C

INTERPOLATE SO THAT TIMES FOR THEORETICAL DATA
CORRESPOND TO DATA FOR ACTUAL MEASURED TIMES.

C

55 DO 90 I = L1, RANGE
X = MTIME*FLOAT(I)
DO 60 K = J, L2
IF (X, LT, TIMC(K, 1)) GO TO 70

60 CONTINUE
J = K-1

IF NO CHANGE IN TC(J,1) NO INTERPOLATION NECESSARY.

IF (TC(K,1),EQ,TC(J,1)) GO TO 85
TCM = TC(J,1)+(X-TIMC(J,1))*(TC(K,1)-TC(J,1))/1
(TIMC(K,1)-TIMC(J,1))

CALCULATE XI SQUARED. THIS IS INCLUDED INSIDE INTERPOLATION LOOP FOR CONVIENENCE.

ER = ER+(TWIRE(I)-TCM)**2
GO TO 90

TCM = TC(J,1)
GO TO 80

CONTINUE
XICALC=ER
RETURN
END

SUBROUTINE TAE(F)

PURPOSE

GENERATES A TABLE OF THEORETICAL TEMPERATURE VS. TIME DATA POINTS. THE COMPUTED TIME CORRESPONDS CLOSELY WITH THE MEASURED TIME BUT NOT EXACTLY.

F = SECOND INDEX ON TC AND TIMC, (F = 1 OR 2).

REAL TWIRE(1000), TAVE, ALPHA1, ALPHA2, ALPHA3, BETA, H1, H2, H3A, H3B, H4, TDUCT, MSTIME, MN, P, MNN, MTMP

INTEGER START, RANGE

REAL TC(1000,2), TIMC(1000,2), DERIV(1000,3)

REAL DELTIM, DELTMP, TCALC, ZY

INTEGER F, J, I, L1, L2

COMMON / BLK1/TWIRE, TAVE, ALPHA1, ALPHA2, ALPHA3, BETA, H1, H2, H3A, H3B, H4, TDUCT, MSTIME, MN, P, MNN, MTMP
COMMON /BLK2/START,RANGE

COMMON /BLK3/TC,TIMC,DERIV

100  TC(START,F) = TAVE
     TCALC = TAVE

MSTIME = ACTUAL TIME BETWEEN MEASURED DATA POINTS
         OF TWIRE.

DELTIM = MSTIME

FIX FIRST POINT.

TIMC(START,F) = EVALTM(TCALC)

INITIALIZE POINT COUNTER FOR ACTUAL MEASURED TIMES.

J = START
Z = ALPHA2-0.01

COMPUTE TABLE OF TEMPERATURES

110  L1 = START+1
     L2 = RANGE+1
     DO 140 I=L1,L2

COMPUTE TCALC AT POINT I - 1.

IF TCALC > Z WE ARE ON TOP FLAT PORTION OF CURVE
AND TEMPERATURE WILL NOT CHANGE ANY MORE.

IF (TCALC.GT.Z) GO TO 130

COMPUTE DERIVATIVE OF TIME WITH RESPECT TO
TEMPERATURE FOR THEORETICAL CURVE.

Y = -H1/(ALPHA2-TCALC) + H2/(TCALC-ALPHA3) +
   2.0*(H3A*(TCALC-ALPHA1)-H3B*BETA)/
   (TCALC-ALPHA1)**2+BETA**2)

DELTIM = CHANGE IN TIME FROM THE LAST DATA NECESSARY
TO MAKE THE CURRENT DATA POINT FALL APPROXIMATELY ON
THE ACTUAL MEASURED TIME.

DELTMP = THE CORRESPONDING CHANGE IN TEMPERATURE.
NOTE THAT THE THEORETICAL FUNCTION (EVALTM) GIVES
TIME BACK FOR AN INPUT OF TEMPERATURE.
DELTMP = DELTIM/Y
TCALC = TCALC+DELTMP
TC(I,F) = TCALC
TIMC(I,F) = EVALTM(TCALC)

C
C DETERMINE THE CHANGE IN TIME NECESSARY FOR THE NEXT
C CALCULATED DATA POINT TO FALL ON THE NEXT MEASURED
C DATA POINT.
C
120  J = J+1
DELTIM = MSTIME*FLOAT(J+1)-TIMC(I,F)
IF (DELTIM.LE.0.0) GO TO 120
GO TO 140

C
C
130  J = J+1
TC(I,F) = ALPHA2
TIMC(I,F) = MSTIME*FLOAT(J)
CONTINUE
RETURN
END

C
C SUBROUTINE INTRP(F,FF)
C
C PURPOSE
C TO INTERPOLATE BETWEEN DATA POINTS CALCULATED FROM
C THE THEORETICAL EQUATION, USED BY SUBROUTINE FDERIV.
C
C F = SPECIFIES THE SECOND INDEX (1 OR 2) FOR TC AND TIME.
C FF = SPECIFIES THE SECOND INDEX (1,2, OR 3) FOR DERIV.
C
C UPON RETURN TC(I,F) HAS THE INTERPOLATED VALUES.
C
C SUBROUTINE INTRP(F,FF)
C
C REAL TWIRE(1000),TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,F,MNN,MTMP
C INTEGER START,RANGE
C REAL TC(1000,2),TIMC(1000,2),DERIV(1000,3)
REAL X

INTEGER F, FF, J, I, K, L1

COMMON /BLK1/TWIRE, TAVE, ALPFA1, ALPFA2, ALPFA3, BETA,
1 H1, H2, H3A, H3B, H4, TDUCT, MTIME, MN, P, MNN, MNP, MTMP

COMMON /BLK2/START, RANGE

COMMON /BLK3/TC, TIMC, DERIV

10 L1 = START+1

DERIV(I, FF) HAS NOT BEEN USED YET AND
IS USED AS TEMPORARY STORAGE

160 DO 170 I=1, 1000
170 DERIV(I, FF) = TC(I, F)
175 J = START
180 DO 200 I=L1, RANGE
X = MTIME*FLOAT(I)
K = J-1
185 K = K+1
IF ( X .GE. TIMC(K, F) ) GO TO 180
J = K-1
IF ( DERIV(K, FF), EQ., DERIV(J, FF) ) GO TO 190

TC(I, F) = DERIV(J, FF) + (X-TIMC(J, F))*(
1 (DERIV(K, FF)-DERIV(J, FF))/
1 (TIMC(K, F)-TIMC(J, F))

GO TO 200
190 TC(I, F) = DERIV(J, FF)
200 CONTINUE
RETURN
END

SUBROUTINE STCFIT

PURPOSE
TO CALCULATE THE STARTING LOCATION OF THE RAMP. THIS
LOCATION IS CALLED "START".
INPUT IS TEMPERATURE DATA 'TWIRE'. THE DATA POINTS 
I=100 TO 160 WILL BE CURVE FIT TO A STRAIGHT LINE 
Y = M*X + B WHERE Y=TWIRE(I) AND X=I=TIME.

SUBROUTINE STCFIT

THIS ROUTINE CALCULATES THE STARTING LOCATION 
OF THE RAMP.

REAL TWIRE(1000),TAVE,ALPHA1,ALPHA2,ALPHA3,BETA, 
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMF 

INTEGER START,RANGE 

INTEGER I,J,K,IX 
REAL SUMX,SUMY,SUMXX,SUMYY,SUMXY,B,M,X 

COMMON /BLK1/TWIRE,TAVE,ALPHA1,ALPHA2,ALPHA3,BETA, 
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMF 

COMMON /BLK2/START,RANGE 

10 START = 100

USE K TO INCREMENT STARTING INDEX FOR TWIRE FROM 
100 TO 130.

DO 50 K=100,130

INITIALIZE VARIABLES

SUMX = 0. 
SUMY = 0. 
SUMXX = 0. 
SUMYY = 0. 
SUMXY = 0. 
J = 0

PERFORM STANDARD LEAST SQUARES CURVE FIT TO A 
STRAIGHT LINE.

20 DO 30 I=K,160 
   J = J+1 
   SUMXY = SUMXY+(FLOAT(I))*TWIRE(I) 
   SUMY = SUMY+TWIRE(I) 
   SUMX = SUMX+FLOAT(I) 
   SUMXX = SUMXX+(FLOAT(I))**2 
30 

CURVE FIT DATA TO Y = M*X + B
M = ((FLOAT(J))*SUMY-SUM*SUMY)/((FLOAT(J))*SUMXX-SUM*X*SUMX)  
B = (SUMY-M*SUMX)/(FLOAT(J))

STRAIGHT LINE CURVE FIT DONE.

DETERMINE IF LAST X WAS A MAXIMUM?  
IF NO CONTINUE.  
IF YES PROGRAM IS DONE AND "START" IS SET EQUAL  
TO X AS THE STARTING LOCATION OF THE RAMP.

SET Y = TAVE

40  X = (TAVE-B)/M  
IX = IFIX(X)  
IF (IX,GT,START) START = IX  
CONTINUE

RETURN
END

SUBROUTINE CONGEN(TGAS,PSC,ERFLAG,TWF)

PURPOSE  
THIS PROGRAM GENERATES THE CONSTANTS NEEDED FOR THE  
THEORETICAL TIME VS. TEMPERATURE EQUATION (EQU. 24).

TGAS = ESTIMATED GAS TEMPERATURE (K),  
PSC = PROBE SHAPE CONSTANT,  
ERFLAG = AN ERROR FLAG => SET=0 IF NO ERROR,  
TWF = ESTIMATED FINAL WIRE TEMPERATURE.

SUBROUTINE CONGEN(TGAS,PSC,ERFLAG,TWF)

REAL TWIRE(1000),TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
1 H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MIN,PN,MNN,MTMP

INTEGER START,RANGE

REAL NU,K1,K2,K3,PSC,WDIA,WDENS,SPHT,SIGMA,
1 E1,E2,EGAS,ALPHAG,X,AX,E,BZ,F,Y

DATA WDIA/0.8128E-3/, WDENS/0.20785E+5/, SPHT/0.1427E+3/,
1 SIGMA/0.56697E-7/, E1/0.85E-1/, E2/0.76E-4/,
1 EGAS/0.0/, ALPHAG/0.0/
TEMPERATURE = DEG, K.
WDIA = WIRE DIAMETER (METERS).
WDENS = WIRE DENSITY (Kg/m**3).
SPHT = WIRE SPECIFIC HEAT (J/(Kg*K)).
SIGMA = STEPHAN BOLTZMAN CONSTANT (J/(SEC.,K**4,m**2)).
EMISSIVITY OF WIRE = E1+E2*TWF. NO UNITS ON E1.
E2 HAS UNITS OF 1/DEG. K.
EGAS = EMISSIVITY OF GAS.
ALPHAG = ABSORPTIVITY OF GAS.
P = PRESSURE (PASCAL).

INTEGER ERFLAG

AZ,BZ,E,F ARE TEMPORARY VARIABLES.
AZ & E ARE AT THE SAME LOCATION TO SAVE SPACE.

EQUIVALENCE (AZ,E),(BZ,F)

COMMON /BLK1/TWIRE,TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
H1,H2,H3A,H3B,H4,TDUCT,MSTIME,MN,P,MNN,MTMP

COMMON /BLK2/START,RANGE

ERFLAG = 0
MNN = MN

THE FOLLOWING STATEMENT IS NEEDED ONLY
FOR THE EXAMPLE IN THIS REPORT.
THE MACH NUMBER WAS MEASURED DOWN STREAM OF THE PULSED
THERMOCOUPLE SITE WHERE THE GAS WAS COOLER. THIS NEXT
STATEMENT CONVERTS THE MACH NUMBER AT THE LOWER
TEMPERATURE TO THAT AT THE PULSED THERMOCOUPLE SITE.

MNN = MN*SQRT(TGAS/MTMP)

COMPUTE NUSSELT NUMBER.

NU = 188.41*(SQRT(MNN*P*WDIA))/
((TGAS**0.6)*((1.+2*MNN**2)**0.25))

K1 = WDIA*WDENS*SPHT/(4.*SIGMA*(E1+E2*TWF))

K2 = (TGAS**0.78)*NU*PSC*3.007E-4/
((WDIA*SIGMA*(E1+E2*TWF))

K3 = K2*TGAS+(1.-ALPHAG)*(TDUCT**4)+EGAS*(TGAS**4)

COMPUTE ALPHA1,ALPHA2,ALPHA3,BETA,H1,H2,H3A,H3B.
SCALE NUMBERS DOWN BY A FACTOR OF 10**-20 TO PREVENT OVERFLOW.

20  

\[ A_Z = (K2**2)*(1.0E-20)*(K2**2)/4, + \\
1  K3*(64.0E-20)*(K3**2)/27.0 \]

\[ B_Z = A_Z \\
Y = 1./3. \\
A_Z = (SQRZ(A_Z)*(1.0E+10) + (K2**2)/2.)***Y \\
B_Z = (K2**2)/2.0 - SQRZ(B_Z)*(1.0E+10) \\
X = ABS(B_Z) \]

ERROR IF B_Z IS POSITIVE

IF (X.EQ.B_Z) STOP 2  
B_Z = X**Y

EVALUATE ALPHA'S AND BETA.

\[ Y = A_Z-B_Z \\
A_LPHA1 = SQRZ(Y)/2. \\
BETA = Y+2.*K2/SQRZ(Y) \\
BETA = SQRZ(BETA)/2. \\
X = SQRZ(2.*K2/SQRZ(Y)-Y) \\
A_LPHA2 = -SQRZ(Y)/2.+X/2. \\
A_LPHA3 = -SQRZ(Y)/2.-X/2. \]

\[ H_1 = -K1/((A_LPHA2-A_LPHA1)**2+BETA**2)*((A_LPHA2-A_LPHA3)) \]
\[ H_2 = -K1/((A_LPHA3-A_LPHA1)**2+BETA**2)*((A_LPHA3-A_LPHA2)) \]
\[ E = -2.*(BETA**2)*((2.**A_LPHA1-A_LPHA2-A_LPHA3)) \]
\[ F = 2.*(BETA*(A_LPHA1-A_LPHA2)*((A_LPHA1-A_LPHA3)-BETA**3)) \]
\[ X = E**2+F**2 \]
\[ H_3A = -K1*E/X \]
\[ H_3B = K1*F/X \]

TIME = H1*ALOG(ALPHA2-T) + H2*ALOG(T-ALPHA3) + 
H3A*ALOG((T-ALPHA1)**2+BETA**2) + 
2.0*H3B*ATAN(BETA/(T-ALPHA1)) + H4

H4 = CONSTANT TO BE DETERMINED. 
TAVE = AVERAGE INITIAL WIRE TEMPERATURE. 
MSTIME IS TIME SCALE FACTOR. 
MSTIME*(TIME INTEGER) IS TIME SINCE START OF DATA. 
TIME = 0 AT DATA POINT I=0 
TIME = MSTIME AT DATA POINT I=1 etc. 

SET ERROR FLAG IF ALPHA2 IS LESS THAN TAVE SINCE IT WOULD REQUIRE TAKING THE LOG OF A NEGATIVE NUMBER.

30 IF (ALPHA2.GT.TAVE) GO TO 40
ERFLAG = 6
RETURN

C 40  X = TAVE-ALPHA1
C
C COMPUTE INTEGRATION CONSTANT H4 BY SETTING T=TAVE
C AT THE STARTING TIME AND SOLVING FOR H4.
C
H4 = MSTIME*FLOAT(START) - (H1*ALOG(ALPHA2-TAVE) +
  1 H2*ALOG(TAVE-ALPHA3) +
  1 H3A*ALOG((TAVE-ALPHA1)**2+BETA**2) +
  1 2.0*H3B*ATAN2(BETA,X))
C
RETURN
END

FUNCTION EVALTM(T)

PURPOSE
EVALUATE THE THEORETICAL EQUATION (TEXT EQU. 24)
FOR TIME AS A FUNCTION OF TEMPERATURE.

T = INPUT TEMPERATURE (K),

FUNCTION EVALTM(T)

REAL T,WIRE(1000),TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
  1 H1,H2,H3A,H3B,H4,TIUCT,MSTIME,MN,P,MNN,MTMP
C
REAL T,X
C
COMMON /BLK1/TWIRE,TAVE,ALPHA1,ALPHA2,ALPHA3,BETA,
  1 H1,H2,H3A,H3B,H4,TIUCT,MSTIME,MN,P,MNN,MTMP
C
10  X = T-ALPHA1
    EVALTM = H1*ALOG(ALPHA2-T)+
     1 H2*ALOG(T-ALPHA3)+
     1 H3A*ALOG((T-ALPHA1)**2+BETA**2)+
     1 2.0*H3B*ATAN2(BETA,X)+H4
RETURN
END
SUBROUTINE MATINV

PURPOSE:
   INVERT A 1 OR 2 DEGREE MATRIX.

SUBROUTINE MATINV (ARRAY,NORDER)

ARRAY = INPUT MATRIX WHICH IS REPLACED BY ITS INVERSE,
NORDER= DEGREE OF MATRIX.

REAL ARRAY(2,2),DET,X
INTEGER NORDER,I,J

10 IF (NORDER,EQ.1) GO TO 20
   IF (NORDER,EQ.2) GO TO 30
   STOP 800

CALCULATE INVERSE OF ONE DEGREE MATRIX.

20 ARRAY(1,1) = 1./ARRAY(1,1)
   RETURN

CALCULATE DETERMINANT FOR SECOND DEGREE MATRIX.

30 DET = ARRAY(1,1)*ARRAY(2,2)-ARRAY(1,2)*ARRAY(2,1)
   IF (DET,EQ.0) STOP 801

CALCULATE ADJOINT MATRIX

X = ARRAY(1,1)
   ARRAY(1,1) = ARRAY(2,2)
   ARRAY(2,2) = X
   ARRAY(1,2) = -ARRAY(1,2)
   ARRAY(2,1) = -ARRAY(2,1)

CALCULATE THE INVERSE OF SECOND DEGREE MATRIX.

DO 50 I=1,2
   DO 40 J=1,2
      ARRAY(I,J) = ARRAY(I,J)/DET
   CONTINUE
40 CONTINUE
50 CONTINUE

RETURN
END
Appendix C
Gradient-Expansion Method

This appendix describes the least-squares fit to a nonlinear function that uses the gradient-expansion algorithm taken from Bevington (ref. 2). The objective of the process is to search for the values of parameters in the theoretical equation that will minimize the sum of the squares of the difference between the data points and the theoretical nonlinear function. This sum to be minimized is defined as

\[ \chi^2 = \sum_{i=1}^{m} \left[ Y_i - Y(X_i) \right]^2 \]  

(C1)

where \( m \) is the number of data points, \( Y_i \) is the dependent variable, \( X_i \) is the independent variable, and \( Y(X) \) is the theoretical function with unknown parameters \( a_j \).

The quantity \( \chi^2 \) is regarded as a function of the parameters \( a_j \) of the fitting function \( Y(X) \). There are \( m \) data points \((X_i, Y_i)\). The idea is to choose the values of the \( n \) parameters \( a_j \) so that \( \chi^2 \) is a minimum.

The first approach is to take the gradient of \( \chi^2 \)

\[ \nabla \chi^2 = \sum_{j=1}^{n} \frac{\partial \chi^2}{\partial a_j} \cdot a_j \]  

(C2)

where the \( a_j \) are unit vectors. The gradient of \( \chi^2 \) gives the direction of the maximum rate of increase of \( \chi^2 \). We want to increment the parameters from some starting value \( x_0 \), so that \( \chi^2 \) decreases. Hence we write

\[ \delta a_j = -\left( \nabla \chi^2 \right)_j \Delta a_j \]

\[ = -\left( \frac{\partial \chi^2}{\partial a_j} \right) \Delta a_j \]  

(C3)

The \( \Delta a_j \) are size constants that must be supplied. The parameters \( a_j \) are incremented by \( \delta a_j \) and the process repeated. The minus sign insures that the increments are in a direction opposite to the gradient so that they are in the direction of most rapid decrease of \( \chi^2 \). However, the method tends not to work well near the actual minimum—it is better further away.

Another approach is to expand the fitting function \( Y(X) \) as a first-order Taylor series in the parameters

\[ Y(X) = Y_0(X) + \sum_{j=1}^{n} \frac{\partial Y_0(X)}{\partial a_j} \delta a_j \]  

(C4)

where \( Y_0(X) \) is the value of \( Y(X) \) at the starting point for the expansion. Then

\[ \chi^2 = \sum_{i=1}^{m} \left[ Y_i - Y_0(X_i) \right]^2 - \sum_{j=1}^{n} \frac{\partial Y_0(X_i)}{\partial a_j} \delta a_j^2 \]  

(C5)

We now want to minimize \( \chi^2 \) as a function of the increments \( \delta a_j \); so we take \( \frac{\partial \chi^2}{\partial \delta a_k} \) and set it equal to zero

\[ \sum_{i=1}^{m} 2 \left[ Y_i - Y_0(X_i) \right] \frac{\partial Y_0(X_i)}{\partial a_k} \]

\[ \quad = \sum_{j=1}^{n} \delta a_j \sum_{i=1}^{m} 2 \frac{\partial Y_0(X_i)}{\partial a_j} \frac{\partial Y_0(X_i)}{\partial a_k} \]  

(C6)

This gives a set of \( n \) linear equations for the \( n \) quantities \( \delta a_j \). Define

\[ \beta_k = \frac{\partial \chi^2}{\partial a_k} = \sum_{i=1}^{m} \left( Y_i - Y(X_i) \right) \frac{\partial Y_0(X_i)}{\partial a_k} \]  

(C7)

\[ \alpha_{jk} = \sum_{i=1}^{m} \frac{\partial Y_0(Y_i)}{\partial a_j} \frac{\partial Y_0(X_i)}{\partial a_k} \]  

(C8)

and

\[ \chi_0^2 = \sum_{i=1}^{m} \left[ Y_i - Y_0(X_i) \right]^2 \]  

(C9)

thus

\[ \beta_k = \sum_{j=1}^{n} \delta a_j \alpha_{jk} \quad k = 1, 2, \ldots, n \]  

(C10)
This can be put into the form of a matrix equation

\[ \beta = \delta a \cdot \alpha \]

or

\[ \beta \cdot \alpha^{-1} = \delta a \]

where \( \beta \) and \( \delta a \) are column matrices with \( n \) elements and \( \alpha \) is an \( n \)-by-\( n \) symmetric square matrix. This method tends to work well near the actual minimum but poorly far from the minimum.

By combining the two methods it is possible to obtain an algorithm that works well far from the minimum and also close to it. To combine the two methods, one writes (ref. 7)

\[ \beta = \alpha' \cdot \delta a \]

where

\[ \alpha_{jk} = \alpha_{kj} \quad \text{for} \quad j \neq k \]

and

\[ \alpha_{jj} = \alpha_{jj}(1 + \lambda) \quad \text{for} \quad \lambda \geq 0 \]

where \( \lambda \) is an arbitrary parameter that changes the method from the Taylor series to the gradient method. If \( \lambda \) is near zero, the method is the same as the Taylor series approach. If \( \lambda \) is large, the diagonal terms dominate and the equations are essentially

\[ \beta_j = \lambda \delta a_j \alpha_{jj} \]

or

\[ \delta a_j = \frac{1}{\lambda \alpha_{jj}} \beta_j = -\frac{1}{2 \lambda \alpha_{jj}} \delta x^2 \]

\[ \frac{\delta x^2}{2 \lambda \alpha_{jj}} \nabla x^2 \]

which result in the gradient method.

This technique can be used by starting with an arbitrary small value of \( \lambda \), such as 0.001. If the computed \( \delta a_j \) causes \( x^2 \) to increase instead of decrease, the initial guess at the \( a_j \) is not good enough, and \( x^2 \) is too far from the minimum for the second method to work. Then \( \lambda \) is increased by a factor of 10 and a new set of \( \delta a_j \) is found. Each time \( \lambda \) is increased the algorithm is more like just taking the gradient, which works well for \( a_j \) far from \( (a_j)_{\min} \). This continues until \( x^2 \) starts to decrease, at which time \( \lambda \) is divided by 10 at each iteration. By this time the minimum will have been found.
Appendix D
Typical Program Input and Results

This appendix provides an example of data used by the computer program. The following data were put into the computer program:

INPUT MACH NUMBER 0.0286
INPUT PRESSURE IN Pa. 99805.
INPUT DUCT TEMPERATURE IN DEG. K. 396.0
INPUT PROBE SHAPE CONSTANT 0.0
MACH NUMBER TEMPERATURE DEG. K. 415.8

The following data were put out by the computer program:

GAS TEMPERATURE = 1707.00 K
FINAL WIRE TEMPERATURE = 1565.79 K
PROBE SHAPE CONSTANT = 0.850

The following data were not put out by the computer program but may be useful:

CHISQR = 0.267E+05
TAVE = 677.0
ALPHA1 = 1183.2
ALPHA2 = 1565.8
ALPHA3 = -3932.2
BETA = 3218.3
H1 = -0.902
H2 = 0.259
H3A = 0.321
H3B = -0.260
H4 = 0.072

The 1000 data points of thermocouple wire temperature are shown in the following listing:
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<th>I</th>
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A pulsed thermocouple is used for measuring gas temperatures above the melting point of common thermocouples. This is done by allowing the thermocouple to heat until it approaches its melting point and then turning on the protective cooling gas. This method requires a computer to extrapolate the thermocouple data to the higher gas temperatures. In earlier work by this author the extrapolation was done by using a first-order exponential curve fit to predict the final thermocouple wire temperature. Since radiation effects were neglected, the gas temperature was not computed. Hand calculations had to be used to estimate the gas temperature. This report describes a method that includes the effect of radiation in the extrapolation. Computations of gas temperature are provided, along with the estimate of the final thermocouple wire temperature. Results from tests on high-temperature combustor research rigs are presented.