FLOW INDUCED FORCE OF LABYRINTH SEAL

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SUMMARY

This paper deals with theoretical analysis of flow induced instability force due to labyrinth seal. That is, approximate solution is given for the partial differential equation representing the flow in labyrinth seal and it is compared with the finite difference method which was proposed in the previous report in order to verify the accuracy of both methods. Then the effects of difference of inlet and outlet pressures of the seal, deflection of pressure and mass flow from the steady state, rotor diameter, seal clearance, seal interval and seal number on the flow induced force of the seal are investigated and it is known that some of these factors are very influential on the flow induced force.

INTRODUCTION

In the previous report (1), the fundamental equation of flow in the labyrinth seal was derived by considering the effect of the variation of gland cross section. The equation is numerically solved by using the finite difference method. Then the spring and damping coefficients of the labyrinth seal are calculated, and by using the result, the stability of a rotor system having labyrinth seal is discussed on this coefficient by using the energy concept. Furthermore, experiments are executed to observe the flow pattern in the gland and to study the characteristics of the flow induced forces in the labyrinth seal.

However, further theoretical investigations are required for the accuracy of numerical calculation and more detail description of labyrinth seal behaviour. Then in this paper the equation is solved by another method and the calculated results are compared with the solution of finite difference method. After that influences of deflection of steady state, pressure difference between inlet and outlet, labyrinth seal radius, seal clearance, seal pitch, number of seal chamber, seal strip height and divergence and convergence seal on the flow induced force, phase angle and leakage flow rate are investigated.

SYMBOLS

t time
w = Rsθ
Rs radius of labyrinth seal
p pressure in seal chamber
\( c \)  
peripheral velocity in seal chamber

\( \rho \)  
density of gas

\( q \)  
mass flow rate in axial direction

\( \varphi \)  
angle from x axis

\( f \)  
cross section of seal chamber

\( l \)  
length of strip pitch

\( h \)  
strip height

\( \delta \)  
radial clearance of seal

\( \tau' \)  
friction shear stress of stator surface

\( \tau'' \)  
friction shear stress of rotor surface

\( U' \)  
length of acting surface of shear stress (stator)

\( U'' \)  
length of acting surface of shear stress (rotor)

\( R \)  
gas constant

\( T \)  
absolute temperature of gas in seal

\( \gamma \)  
specific weight of gas

\( \psi \)  
flow coefficient

\( n \)  
ratio of specific heat

\( \omega \)  
rotating speed of rotor

\( \omega_n \)  
critical speed of rotor system

\( u = \omega R_s \)  
peripherical velocity of rotor

\( K \)  
coefficient of viscosity

\( \nu \)  
coefficient of kinematic viscosity

\( R_e \)  
Reynold's number

\( \lambda' \)  
friction coefficient (stator)

\( \lambda'' \)  
friction coefficient (rotor)

subscript *  
steady state
EQUATION OF MOTION

Derivation of Equation of Motion

Derivation of the equation which describes the labyrinth seal flow is followed the previous report (1), where the following assumptions are taken and schematic flow is shown in Fig.1;

(1) Fluid in labyrinth seal is assumed to be gas and its behavior is assumed to be ideal.

(2) Temperature of the fluid in the labyrinth seal is assumed to be constant.

(3) Cross section area of the seal gland is assumed to be constant in spite of the deflection of rotor and time derivative of cross section area is only considered.

(4) Change of flow state in the gland is assumed to be isentropic change.

Thus the fundamental equations with respect to continuity and momentum for the ith seal element are as follows;

\[ \frac{\partial (P_i C_i)}{\partial t} + f_i \frac{\partial (P_i C_i)}{\partial x} + (g_{i+1} - g_i) = 0 \]  \( \tag{1} \)

\[ f_i \frac{\partial (P_i C_i^2)}{\partial t} + \frac{\partial (P_i f_i C_i)}{\partial x} + (g_{i+1} C_i - g_i C_{i-1}) + \tau_{i+1} U_{i+1}' - \tau_i U_i'' = -\frac{\partial P_i}{\partial x} f_i \]  \( \tag{2} \)

As temperature of each gland is constant, the equation of state in the ith gland is given as;

\[ P_i = \gamma_i R T \]  \( \tag{3} \)

where

\[ \gamma_i = f_i g_i \]

Next, denoting the gas flow velocity through labyrinth nozzle (seal clearance) \( S \), the relation between axial flow rate and pressure in the gland is given by the thermodynamical energy equilibrium condition as,

\[ S = \sqrt{2g_i \frac{n}{n-1} P_i v_i \left[ 1 - \left( \frac{P_o}{P_i} \right)^{\frac{n-1}{n}} \right]} \]  \( \tag{4} \)

where \( n \) is isentropic exponent, \( v \) specific volume and subscripts I and 0 mean inlet and outlet of nozzle, respectively. Assuming that the flow is adiabatic change, i.e. \( v = (P_I / P_0)^{\frac{n}{n-1}} v_I \), the flow rate through the nozzle \( G \) [kgf/sec] is given as;

\[ G = \alpha F \sqrt{2g_i \frac{n}{n-1} P_i v_i \left( \frac{P_0}{P_i} \right)^{\frac{n}{n-1}} - \left( \frac{P_0}{P_i} \right)^{\frac{n+1}{n}}} \]  \( \tag{5} \)
where $F$ is nozzle cross area (area of seal clearance) and subscripts $(i, 0)$ correspond to the gland number of the labyrinth seal. Applying the above relation to flow in the labyrinth seal, we assume that the labyrinth seal is a series of the nozzle, pressure difference between each gland is small and enthalpy of each gland is nearly equal. Then leakage flow rate through labyrinth seal $G$ (kgf/sec) is given by using the relation

$$G_L = \alpha F \frac{\sqrt{g}}{R T} \sqrt{P_{i}^2 - P_{0}^2}$$

(6)

Applying Eq.(6) to the flow through the $i$th seal element, mass flow rate for unit area is given as;

$$\frac{\dot{q}_i}{S_i} = \frac{G_L}{F \cdot g} = \alpha \frac{\sqrt{P_{i}^2 - P_{0}^2}}{\sqrt{g} R T}$$

(7)

Rewriting Eq.(7), the following equation is obtained;

$$P_{i-1}^2 - P_{i}^2 = \frac{\dot{q}_i^2}{\mu^2 S_i}$$

(8)

where

$$\frac{1}{\mu^2} = \frac{g R T}{\alpha^2}$$

(9)

**Linearization of Fundamental Equation**

For the linearization of the equation, the perturbations of pressure, peripheral velocity and flow rate from those of the steady state and also mean value and perturbation of seal clearance are introduced as

$$P_i = P_{*i} (1 + \xi_i) \quad C_i = C_{*i} (1 + \psi_i)$$

$$\dot{q}_i = \dot{q}_{*i} (1 + \eta_i) \quad \delta_i = \delta_{*i} (1 + \psi_i)$$

(10)

where $P_{*i}, C_{*i}$ and $q_{*i}$ are pressure, peripheral velocity and axial flow rate of steady state in the $i$th gland and $\xi_i, \eta_i, \delta_i$ and $\psi_i$ are nondimensional perturbation terms of pressure, peripheral velocity and axial flow rate. Also $\delta_{*i}$ and $\psi_i$ are mean value and perturbation of seal clearance.

Assuming that the rotor is whirling along elliptical orbit, $\psi_i$ is represented as

$$\psi_i = -\frac{a_i}{\delta_{*i}} \cos w t \cos \varphi - \frac{b_i}{\delta_{*i}} \sin w t \sin \varphi$$

(11)
Then the cross section area in the $i^{th}$ gland is represented as

$$f_i = \left\{ h_i + \delta_{x_i}(1 + \psi_i) + h_i+1 + \delta_{x_{i+1}}(1 + \psi_{i+1}) \right\} \frac{l}{2} \quad (12)$$

As the change of state in the gland is isentropic, the following relations are obtained;

$$\frac{p_{x_i}}{p_{x_{i+1}}} = \frac{p_i}{p_{i+1}} = A \quad (13)$$

therefore

$$\frac{\partial p_i}{\partial t} = n \frac{p_i}{p_{x_i}} \frac{\partial f_i}{\partial t} \quad (14)$$

Substituting Eqs. (10) - (14) into Eqs. (1) and (2), and denoting the time derivative as $\frac{\partial}{\partial t}$ and spatial derivative as $\frac{\partial}{\partial x}$, the following equations are obtained for the $i^{th}$ seal element;

$$-\frac{\partial f_{x_i}}{\partial t} \frac{\partial u_{x_i}}{\partial t} + \frac{f_{x_i}}{n} \frac{\partial u_{x_i}}{\partial t} + \frac{f_{x_i} \psi_{x_i}}{n} \frac{\partial u_{x_i}}{\partial t} = \frac{gR T M^2}{8} \delta_{x_{i+1}} \left( P_{x_{i+1}} - P_{x_i} \right) \left[ a_i \cos \omega t \cos q + b_i \sin \omega t \sin q \right]$$

$$- \frac{gR T M^2}{8} \delta_{x_i} \left( P_{x_i} - P_{x_{i+1}} \right) \left[ a_i \cos \omega t \cos q + b_i \sin \omega t \sin q \right]$$

$$+ \frac{l}{2} \left[ -a_i \omega \sin \omega t \cos q + b_i \omega \cos \omega t \sin q - a_i \omega \sin \omega t \sin q + b_i \omega \cos \omega t \sin q \right] \quad (15)$$

$$\frac{\partial f_{x_i}}{\partial t} \frac{\partial u_{x_i}}{\partial t} + \frac{\partial f_{x_i} \psi_{x_i}}{\partial t} \frac{\partial u_{x_i}}{\partial t} - \frac{\partial f_{x_i} \psi_{x_i}}{\partial x} \frac{\partial u_{x_i}}{\partial x} = \frac{gR T M^2 \delta_{x_{i+1}}}{8} \left( P_{x_{i+1}} - P_{x_i} \right) \left[ a_i \cos \omega t \cos q + b_i \sin \omega t \sin q \right]$$

$$+ \frac{gR T M^2 \delta_{x_i}}{8} \left( P_{x_i} - P_{x_{i+1}} \right) \left[ a_i \cos \omega t \cos q + b_i \sin \omega t \sin q \right]$$

$$+ \frac{\partial f_{x_i} \psi_{x_i}}{\partial x} \frac{\partial u_{x_i}}{\partial x} - \frac{\partial f_{x_i} \psi_{x_i}}{\partial x} \frac{\partial u_{x_i}}{\partial x} = \frac{gR T M^2 \delta_{x_{i+1}}}{8} \left( P_{x_{i+1}} - P_{x_i} \right) \left[ a_i \cos \omega t \cos q + b_i \sin \omega t \sin q \right]$$

$$- \frac{\partial f_{x_i} \psi_{x_i}}{\partial x} \frac{\partial u_{x_i}}{\partial x} \left[ a_i \cos \omega t \cos q + b_i \sin \omega t \sin q \right]$$

$$+ \frac{l}{2} \left[ -a_i \omega \sin \omega t \cos q + b_i \omega \cos \omega t \sin q - a_i \omega \sin \omega t \sin q + b_i \omega \cos \omega t \sin q \right] \quad (16)$$
For the \( n \) stage labyrinth seal, Eqs. (15) and (16) are represented in matrix form as:
\[
\mathbf{T}\mathbf{w} + \mathbf{V}\mathbf{w}' + \mathbf{A}\mathbf{w} = \mathbf{G}_1[\mathbf{\alpha}_1\cos(\mathbf{\varphi} + \mathbf{\omega}t) + \mathbf{\alpha}_2\cos(\mathbf{\varphi} - \mathbf{\omega}t) - \mathbf{\beta}_1\cos(\mathbf{\varphi} + \mathbf{\omega}t) + \mathbf{\beta}_2\cos(\mathbf{\varphi} - \mathbf{\omega}t) + \mathbf{\beta}_3\cos(\mathbf{\varphi} + \mathbf{\omega}t) + \mathbf{\beta}_4\cos(\mathbf{\varphi} - \mathbf{\omega}t) + \mathbf{\beta}_5\cos(\mathbf{\varphi} + \mathbf{\omega}t) + \mathbf{\beta}_6\cos(\mathbf{\varphi} - \mathbf{\omega}t)]
\]
where \( \mathbf{P}_i = \mathbf{P}_0 \) (inlet pressure) and \( \mathbf{P}_n = \mathbf{P}_n \) (outlet pressure).

\section*{ANALYSIS}

\subsection*{Analysis of Steady State}

Setting the perturbation terms \( \xi_i \) and \( \eta_i \) zero in Eqs. (15) and (16) in order to obtain the pressure and peripheral velocity of each seal stage in steady state,
\[
\mathbf{q}_{i+1} = \mathbf{q}_i = \mathbf{q}_* \tag{18}
\]
\[
\mathbf{q}_{i+1}\mathbf{w}_i - \mathbf{q}_i\mathbf{w}_{i-1} + \tau'^{U_i} = 0 \tag{19}
\]

As the multi labyrinth stage is considered as a serial nozzle, the axial flow rate in the steady state is obtained as follows
\[
\mathbf{q}_i = \mathbf{\sigma}_i \sqrt{\frac{2\mathbf{h}}{\mathbf{n}_i}} \mathbf{P}_i^{'} \mathbf{C}_i \left[ (\frac{\mathbf{P}_i^{'}}{\mathbf{P}_i^{''}})^{\frac{2}{\mathbf{n}_i}} - (\frac{\mathbf{P}_i^{'}}{\mathbf{P}_i^{''}})^{\frac{n_i}{\mathbf{n}_i}} \right] \tag{20}
\]
where \( \mathbf{P}_i^{'} \) is the pressure of \((i-1)\)th chamber which is taken into account of draughty flow. The pressure is given by Ref. (3) as;
\[
\mathbf{P}_i^{' - } = \mathbf{P}_i + (\mathbf{P}_0 - \mathbf{P}_n)(1 - \mathbf{A}_i^2) \mathbf{\xi}_i \mathbf{\eta}_i \mathbf{\eta}_i^{'}, \quad \mathbf{\xi}_i = \frac{2\mathbf{A}_i^2}{\mathbf{A}_i + 1} \tag{21}
\]

where \( \mathbf{A}_i \) is determined by the shape of labyrinth seal;
\[
\mathbf{A}_i = \frac{\mathbf{\sigma}_i}{\mathbf{\xi}_i} \mathbf{C}_c \mathbf{\xi}_i \mathbf{\eta}_i \mathbf{\eta}_i^{'}, \quad \mathbf{\xi}_i = \frac{\mathbf{\sigma}_i}{\mathbf{\xi}_i} \mathbf{C}_c + \tan \mathbf{\eta}_i \tag{22}
\]

From these equations, state variables i.e. the pressure and flow rate satisfied the equilibrium condition, is obtained by iterating Eqs. (18) and (19). The peripheral velocity is obtained by using Eq. (19) in the same manner.
Analysis of Dynamic Force

Referring to the left term in Eq.(17), the solutions are assumed as follows

\[ P_i = P_{ai} + P_{ei} \left\{ A_{ai} \sin(\varphi + \omega t) + B_{ai} \cos(\varphi + \omega t) + A_{ei} \sin(\varphi - \omega t) + B_{ei} \cos(\varphi - \omega t) \right\} \quad (23) \]

\[ C_i = C_{ai} + C_{ei} \left\{ E_{ai} \sin(\varphi + \omega t) + F_{ai} \cos(\varphi + \omega t) + E_{ei} \sin(\varphi - \omega t) + F_{ei} \cos(\varphi + \omega t) \right\} \quad (24) \]

Representing the above coefficients in matrix form as

\[
\begin{bmatrix}
A_{ai} \\
B_{ai} \\
E_{ai} \\
F_{ai} \\
A_{ei} \\
B_{ei} \\
E_{ei} \\
F_{ei}
\end{bmatrix}
\]

the solution is written as;

\[ \mathbf{U} = \mathbf{\alpha}_1 \sin(\varphi + \omega t) + \mathbf{\beta}_1 \cos(\varphi + \omega t) + \mathbf{\alpha}_2 \sin(\varphi - \omega t) + \mathbf{\beta}_2 \cos(\varphi - \omega t) \quad (25) \]

Substituting Eq.(25) into Eq.(17), and separating it to the elements \( \cos(\varphi + \omega t) \), \( \sin(\varphi + \omega t) \), \( \cos(\varphi - \omega t) \) and \( \sin(\varphi - \omega t) \), the coefficients are represented as follows;

\[
\begin{align*}
\omega T \mathbf{\alpha}_1 + \mathbf{\nu} \mathbf{\alpha}_1 + \mathbf{A} \mathbf{\beta}_1 &= \mathbf{G}_1 \mathbf{\alpha}_1 + \mathbf{G}_2 \mathbf{\theta}_a + \mathbf{G}_3 \mathbf{b}_x + \mathbf{G}_4 \mathbf{b}_y \\
- \omega T \mathbf{\beta}_1 - \mathbf{\nu} \mathbf{\beta}_1 + \mathbf{A} \mathbf{\alpha}_1 &= \mathbf{G}_1 \mathbf{\alpha}_1 + \mathbf{G}_4 \mathbf{\theta}_a + \mathbf{G}_5 \mathbf{b}_x + \mathbf{G}_6 \mathbf{b}_y \\
- \omega T \mathbf{\alpha}_2 + \mathbf{\nu} \mathbf{\alpha}_2 + \mathbf{A} \mathbf{\beta}_2 &= \mathbf{H}_1 \mathbf{\alpha}_1 + \mathbf{H}_2 \mathbf{\theta}_a + \mathbf{H}_3 \mathbf{b}_x + \mathbf{H}_4 \mathbf{b}_y \\
\omega T \mathbf{\beta}_2 - \mathbf{\nu} \mathbf{\beta}_2 + \mathbf{A} \mathbf{\alpha}_2 &= \mathbf{H}_1 \mathbf{\alpha}_1 + \mathbf{H}_2 \mathbf{\theta}_a + \mathbf{H}_3 \mathbf{b}_x + \mathbf{H}_4 \mathbf{b}_y 
\end{align*}
\]

(26)  (27)  (28)  (29)

where \( \mathbf{G}_1 \sim \mathbf{G}_8 \) and \( \mathbf{H}_1 \sim \mathbf{H}_8 \) are coefficient matrices which is composed of the matrices \( \mathbf{A}, \mathbf{C}, \mathbf{F}, \mathbf{B} \) and \( \mathbf{L} \). Solving Eqs.(26),(27) and Eqs.(28), (29) simultaneously, \( \mathbf{\alpha}_1, \mathbf{\beta}_1, \mathbf{\alpha}_2 \) and \( \mathbf{\beta}_2 \) are obtained. Then velocity and pressure in the gland is calculated by substituting the coefficients into Eq.(23), and the flow induced force for whole seal is obtained by integrating the pressure for each gland.

Representing the \( x \) and \( y \) direction force by \( P_x \) and \( P_y \),

\[
\begin{align*}
P_x &= -R_S \sum_{i=1}^{n} \int_0^{2\pi} P_{ai} \xi_i l \cos \varphi \, d\varphi \\
P_y &= -R_S \sum_{i=1}^{n} \int_0^{2\pi} P_{ai} \xi_i l \sin \varphi \, d\varphi 
\end{align*}
\]

(30)

The flow induced force \( F \) and phase angle from \( x \) axis are described as

\[ F = \sqrt{P_x^2 + P_y^2} \]

\[ \phi = \tan^{-1} \frac{P_y}{P_x} \]

(31)

Spring coefficients and damping coefficients are obtained from Eqs.(30) and (31) in the following matrix form;
When Eqs. (30) and (31) are rewritten to Eq. (32) in the above, the following form is used as $t_L$,

$$t_L = t_{L1}a_x \cos \omega t + t_{L2}a_x \sin \omega t + t_{L3}b_x \cos \omega t + t_{L4}b_x \sin \omega t + t_{L5}a_x \cos \omega t$$

Using the above expression, the following relations are obtained between coefficients:

$$K_{xx} = K_{yy} \quad K_{0xx} = K_{0yy} \quad K_{xy} = -K_{yx} \quad K_{0xy} = K_{0yx}$$

$$C_{xx} = C_{yy} \quad C_{0xx} = C_{0yy} \quad C_{xy} = -C_{yx} \quad C_{0xy} = -C_{0yx}$$

NUMERICAL EXAMPLE AND DISCUSSIONS

The labyrinth seal having three teeth, which is shown in Fig. 2, is used as a numerical model. Details of the labyrinth seal is shown in Table 1 and for the parameter survey each element of seal dimension is independently changed in order to investigate the influence of its seal size on the flow induced force. In order to evaluate the seal force, circular orbit ($a_* = b_*$ in Eq.(17)) are used as a whirling of labyrinth seal, and its radius $a_* = 0.1 \text{mm}$. For the calculation the following equation is used as the friction coefficient between rotor and stator;

for stator surface:

$$R_{ei} = \frac{D(U - C \omega t)}{D \eta}$$

where $D$ is wetted perimeter length considered the rotating Reynold’s Number

$$R_{ei} < 1200, \quad \lambda_i = 64/R_{ei}, \quad R_{ei} \geq 1200, \quad \lambda_i = 0.3164 R_{ei}^{0.25}$$

for rotor surface:

$$R_{ei}^{0.25} = \frac{D(U - C \omega t)}{D \eta}, \quad R_{ei} < 1200, \quad \lambda_i'' = 64/R_{ei}, \quad R_{ei} \geq 1200, \quad \lambda_i'' = 0.3164 R_{ei}^{0.25}$$

Two methods are used for the numerical calculation, one is the method which is shown in previous section, the other is one which is the finite difference method that was proposed in Ref.(1). At first calculated results for these two methods are compared with and as the result it is known that these two results agree. Then parameter surveys of labyrinth seal are carried out to see the influence on the flow induced force which exits the instability.

In this paper critical speed of a rotor system is assumed to be 5000 rpm, and rotating speed of the rotor is represented in nondimensional expression as $\omega/\omega_0 (\omega (\text{rpm})$ rotating speed of the rotor). $F$ and PAI shown in figures mean the absolute value of the force induced by labyrinth seal and its phase angle from x axis.
a) Accuracy of Two Calculation Method

Fig. 3 shows the spring coefficients and damping coefficients which are obtained by two methods i.e. approximate method and finite difference method. In the figure, full line and broken line show the coefficients for inlet pressure \( P_0 = 196.13 \text{(kPa)} \) and \( 147.1 \text{(kPa)} \) respectively and outlet pressure \( P_3 = 98.07 \text{(kPa)} \) which are obtained by the approximate method and dotted line shows the coefficient for inlet pressure \( P_0 = 196.13 \text{(kPa)} \) and outlet pressure \( P_3 = 98.0 \text{(kPa)} \) which is obtained by the finite difference method. Also sign (-) on line means negative value.

From this figure full line and dotted line coincide very well so each calculation method has good accuracy.

b) Influence of Deflection of Steady

In this paper, pressure, peripheral velocity and flow rate in steady state are iteratively calculated, however, Kostyuk and other authors obtain them by another method and there are some discrepancies between both steady state values. Moreover, steady state values in field and the calculated values may have discrepancies. So it is important to investigate the influence of deflection from steady state on the accuracy of the induced force. Then Fig. 4 shows the flow induced force and its phase angle when pressure is a little deflected from steady state where the number on lines correspond to the number in Table 3.

The calculated data are shown in table 2. From the calculated result it is known that even if the flow rate is deflected 20% from the steady state, flow induced force and phase angle are not so influenced, and these curves become like one line, so this figure is neglected. On the other hand if the pressure is deflected 0.5% from steady state value, they are influenced very much. Also it is known that the induced force is not influenced by rotating speed of rotor but if it is separated to component, influence of rotating speed is remarkable as shown in Fig. 5. From the above discussions it is concluded that in the calculation of flow induced force, steady state values must be calculated with high accuracy.

c) Influence of Pressure Difference Between Inlet and Outlet

Fig. 6 shows the influence of pressure difference between inlet and outlet on the flow induced force when inlet (entrance) pressure is kept constant value (209 atm) and outlet (exit) pressure is changed. The figure shows that as the outlet pressure becomes large, i.e. the pressure difference between inlet and outlet becomes small, the flow induced force becomes small. As the flow rate decreases, the induced force proportionally decreases.

d) Influence of Labyrinth Seal Radius

Fig. 7 shows the influence of labyrinth seal radius on the flow induced force, where the seal radius is variable and other dimensions are fixed at constant values. Here flow rate means total flow rate as shown in the expression;

From the figure the flow induced force proportionally increases as the seal radius increases. Therefore, the flow induced force due to labyrinth seal can be easily evaluated by the similarity law, when the seal radius is known.

e) Influence of Seal Clearance

Fig. 8 shows the relation between seal clearance and the flow induced force and the leakage flow rate (leakage mass flow rate per unit length), where as the clearance increases, the induced force \( F \) decreases and the flow rate increases proportionally.
f) Influence of Seal Pitch

Fig.9 shows the influence of seal pitch on the flow induced force, when strip height is kept constant. From the figure it is known that as the seal pitch increases, the induced force increases, but the leakage flow rate $q$ does not so much changed. Thus the large pitch is not good from the view point of the flow induced force.

g) Influence of Number of Seal Chamber

Fig.10 shows the influence of chamber number on the induced force and the leakage flow rate, where inlet and outlet pressure is constant and each chamber dimension is also same. From the figure the induced force increases like hyperbolic as the chamber number increasing. On the other hand leakage flow rate decreased like exponentially. Therefore optimum number of chambers may be determined when inlet and outlet pressure are known.

h) Influence of Seal Strip Height

Fig.11(a)(b) shows the influence of seal strip height. In this case the induced force, phase angle and leakage flow rate do not influenced so much. But if this result is represented in the spring and damping coefficients, the influence due to the strip height is remarkable. Fig.11(c) shows the spring coefficients $k_{xx}$ and $k_{xy}$.

i) Influence of Divergence and Convergence Type Labyrinth Seal

Fig.12 shows a comparison of seal types, that is, comparison of straight through type, convergence type and divergence type. In this calculation there is little discrepancies of the flow rate $q$, so its figure is neglected. From Fig.12(a), it is known that the flow induced force for the convergence type is the smallest of all and divergence type is the largest inspite of the same leakage flow rate. From Fig.12(c), the flow induced coupling stiffness $k_{xz}$ is similar characteristics, but convergence type is excellent from the viewpoint of stability.

CONCLUSION

In this report the flow induced force due to labyrinth seal which is sometimes the cause of instability is studied for the special model. That is, the labyrinth seal of shroud of steam turbine is used as an example and its flow induced force and stiffness and damping coefficients are theoretically calculated by two method. Then influences of deflection of steady state, pressure difference between inlet and outlet, labyrinth seal radius, seal clearance, seal pitch, number of seal chamber and seal strip height on the flow induced force and leakage flow rate are investigated. As the results, it is known that the above factors are very influential on the flow induced force. However further theoretical and experimental investigation are required in order to obtain a more precise description of labyrinth seal characteristics.

REFERENCES

**Table 1. - Labyrinth Seal Dimensions for Numerical Example**

<table>
<thead>
<tr>
<th>Seal Radius (Rs)</th>
<th>1 (m)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Rotor Natural Frequency ($\omega_n$)</td>
<td>5000 (rpm)</td>
</tr>
<tr>
<td>Strip Height (h)</td>
<td>10 (mm)</td>
</tr>
<tr>
<td>Seal Clearance ($\delta$)</td>
<td>1.3 (mm)</td>
</tr>
<tr>
<td>Strip Pitch (l)</td>
<td>30 (mm)</td>
</tr>
<tr>
<td>Number of Chamber (N)</td>
<td>2</td>
</tr>
<tr>
<td>Seal Type</td>
<td>Straight-through type</td>
</tr>
</tbody>
</table>

**Table 2. - Gas Condition for Numerical Example**

<table>
<thead>
<tr>
<th>Operating Liquid</th>
<th>Air</th>
</tr>
</thead>
<tbody>
<tr>
<td>Liquid Temperature (T)</td>
<td>784.15 (K)</td>
</tr>
<tr>
<td>Entrance Pressure ($P_o$)</td>
<td>20482.0 (kPa)</td>
</tr>
<tr>
<td>Exit Pressure ($P_n$)</td>
<td>19678.4 (kPa)</td>
</tr>
</tbody>
</table>

**Table 3. - Deflection of Pressure and Flow Rate from Steady State**

<table>
<thead>
<tr>
<th>Steady State Value</th>
<th>Pressure of Each Stage (kPa)</th>
<th>Flow Rate (kg s/m²)</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>$P_0$</td>
<td>$P_1$</td>
</tr>
<tr>
<td>Steady State Value (1)</td>
<td>20482.0</td>
<td>20166.3</td>
</tr>
<tr>
<td>Deflection of Pressure $P_1$, $P_2$ (2)</td>
<td>20482.0</td>
<td>20267.2</td>
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<td>(3)</td>
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<td>Deflection of Flow Rate (2)</td>
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Figure 1. - Cross section of labyrinth seal and definition of the coordinate.

Figure 2. - Labyrinth seal model for numerical example.
Figure 3. - Accuracy of two calculating methods and influence of pressure difference between inlet and outlet.
Figure 4. - Influence of pressure deflection from steady state on the flow induced force.

Figure 5. - Influence of pressure deflection from steady state on the cross coupling stiffness.
Figure 6. - Influence of pressure difference between inlet and outlet.

Figure 7. - Influence of labyrinth seal radius.
Figure 8. - Influence of seal clearance.

Figure 9. - Influence of seal pitch length.
Figure 10. - Influence of number of seal chamber.

Figure 11. - Influence of seal strip height on the flow induced force and phase angle.
Figure 11. - Influence of seal strip height on flow induced cross-coupling.

Figure 12. - Influence of divergence and convergence type labyrinth seal.