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**Abstract**

The $^{13}$C NMR chemical shifts of graphite intercalation compounds have been calculated. For acceptor types, the shifts come mainly from the paramagnetic (Ramsey) intra-atomic terms. They are related to the gross features of the two-dimensional band structures. The calculated anisotropy is about -140 ppm and is independent of the finer details such as charge transfer. For donor types, the carbon $2p\pi$ orbitals are spin-polarized because of mixing with metal conduction electrons, thus there is an additional dipolar contribution which may be correlated with the electronic specific heat. The general agreement with experimental data is satisfactory.
A wide variety of molecular and metallic species may be intercalated between the carbon layers of graphite. Because of the high electrical conductivities and the nearly two-dimensional characteristics, these graphite intercalation compounds (GIC) have been extensively studied.\textsuperscript{1,2} From the $^{13}$C nuclear magnetic resonance (NMR) spectra, the chemical shifts (CS) may be used as microscopic probes for the electronic structures. The method is also universally applicable to GIC of both the acceptor (with molecular species such as $\text{Br}_2$, $\text{HNO}_3$, $\text{AsF}_5$, $\text{SbCl}_5$, $\text{AlCl}_3$, etc.) and donor (with metallic species such as K, Cs, etc.) types. The pristine graphite itself has a layered structure with relatively weak interplanar interactions. Following the standard notations of the coordinate systems\textsuperscript{3}, the planes of the carbon atoms are parallel to the $xy$-planes and the $z$-axis is a threefold symmetry axis. The $^{13}$C CS-tensor $\delta$ is expected to be both diagonal and axially symmetric ($\delta$ is often denoted as the chemical shift shielding parameter $\sigma$ in the literature).\textsuperscript{4-7} The diagonal components $\delta_{zz}$, $\delta_{xx}$, and $\delta_{yy}$ will be abbreviated as $\delta_z$, $\delta_x$, and $\delta_y$. Since $\delta_x = \delta_y$, there are only two independent components, $\delta_z$ and $\delta_x$, which may be measured experimentally with the magnetic field $B$ parallel and perpendicular to the $z$-axis. They will be expressed in ppm ($10^{-6}$) with respect to tetramethylsilane (TMS) as the reference. It is convenient to define $\delta_i = (\delta_z + 2\delta_x)/3$ and $\delta_a = \delta_z - \delta_x$ as the isotropic and anisotropic parts of the CS. While structural changes do occur upon intercalation, the carbon layer structures remain intact for GIC, hence we also expect their CS tensor to have approximate axial symmetry.
There are experimental difficulties of $^{13}\text{C}$ NMR due to the low natural abundances (1%) of $^{13}\text{C}$ isotopes, unfavorable relaxation times, etc. Despite these difficulties, the general features of the experimental CS are now known.\textsuperscript{8-16} Because large single crystals of graphite are not readily available, many GIC samples are prepared by intercalation into highly oriented pyrolytic graphite (HOPG) which is polycrystalline with the $z$-axis of individual crystallites in good mutual alignment throughout the specimen ($\sim$2° spread), though the $x$- and $y$-axes, in the graphite planes, are disordered from crystallite to crystallite. For these GIC samples derived from HOPG, sharp resonance lines are usually observed when $B$ is perpendicular to the $z$-axis. Although structural changes do occur during intercalation\textsuperscript{2}, these experimental evidences ($S_x = S_y$) indicate that the axial symmetry of the CS tensor about the $z$-axis is indeed retained. Also, nearly universal values of $S_z \sim 40$ and $S_x \sim 180$ have been observed for all GIC of acceptor type (GICA); these shifts are nearly independent of the GICA composition. On the contrary, both $S_z$ and $S_x$ vary with the compositions of the GIC of donor types (GICD).

In the present work, we have calculated the CS for GIC and have considered the "band" and the "dipolar" contributions. We will show that the former term depends only on the gross features of the electronic $2p\sigma$ and $2p\pi$ band structures and is not very sensitive toward the finer details such as charge transfers or Fermi levels. We will then correlate the "dipolar" term with the electronic specific heat of GIC. It will then be shown that only "band" contribution may be present for GICA.
(since their electronic specific heats are very small) but both
terms may be important for GICD.

We will first consider the "band" contribution. The main
contribution to CS comes from the paramagnetic (Ramsey)
 intra-atomic term\textsuperscript{4-7} while the diamagnetic term (Lamb) may be
neglected because it is both small and nearly isotropic. For
molecules, the Ramsey term of CS is due to the deviations
from spherical symmetry of the electronic charge distribution\textsuperscript{4,5}
and may be calculated from the interaction between the occupied
and empty molecular orbitals. For solids, the molecular
orbitals are replaced by Bloch functions, thus we have\textsuperscript{17-19}

\begin{equation}
\delta_z = (2e^2 \hbar^2 / m^2 c^2) \langle r^{-3} \rangle_{2p} \sum_v \left[ \sum_c (E_c - E_v)^{-1} \langle \psi_v | \mathbf{L}_z | \psi_c \rangle \langle \psi_c | \mathbf{L}_z | \psi_v \rangle \right]
\end{equation}

where the factor of 2 is included for up and down spins, \( c \) is
light velocity, \( \hbar \) is Planck's constant divided by \( 2\pi \),
e and \( m \) are electronic charge and mass, \( \mathbf{L}_z \) is the \( z \)-component
of angular momentum operator in units of \( \hbar \), \( \langle r^{-3} \rangle_{2p} \) is the
expectation value of the inverse cube distance for the carbon
\( 2p \)-orbitals, \( \psi \) and \( E \) are the wave functions and their energies.
The subscripts $v$ and $c$ refer to valence (occupied) and conduction (unoccupied or excited) states. The functions $\psi_v$ and $\psi_c$ inside the square bracket have the same wave vector $k$. $\Sigma_c$ refers to summation over the discrete excited states. $\Sigma_v$ refers to integration or average over the $k$-space or Brillouin zone (BZ). Following Strong et al.\cite{17}, a reasonable choice of $\langle r^{-3} \rangle_{2p}$ is $1.8a_0^{-3} = 12.2 \text{Å}^{-3}$ for carbon 2p-orbitals where $a_0$ is the Bohr radius of the hydrogen atom, hence $(2e^2k^2/m^2c^2)\langle r^{-3} \rangle_{2p} = 5200 \text{ppm}$ when $E_c-E_v$ is expressed in units of eV. The expression for $\xi_x$ is rather similar with $L_z$ replaced by $L_x$.

Usually, informations are rather incomplete about the excited states, thus the energy denominators in (1) are often replaced by an average excitation energy $\Delta E$ so that the closure relation may be used.\cite{4,5} However, the choice of $\Delta E$ is somewhat arbitrary. For anisotropic solids such as graphite, there would be even more uncertainty in choosing two arbitrary values for $\xi_z$ and $\xi_x$.

Since CS would depend only on the general features of the energy bands, it may be suuffice to use the tight-binding approximation for two-dimensional graphite where the wave functions of the valence state and the first excited state have been given by Lomer.\cite{3} This limited information about the valence state $\psi_v$, $E_v$ and the first excited state (subscript $c=1$) $\psi_1$, $E_1$ (the lowest unoccupied state that satisfies the condition $\langle \psi_v | L | \psi_1 \rangle = 0$) permits us to give an over-estimate and under-estimate of CS:

(a) We may replace the energy denominator $E_c-E_v$ in (1) by the common value $\Delta E_1 = E_1 - E_v$. Using the closure relation,
we get:

$$S_z = 5200 \sum_v (\Delta E_1)^{-1} \langle \psi_v | L_z^2 | \psi_v \rangle$$  

(2)

We will refer to (2) as the average energy (AE) approximation. Clearly, AE is an over-estimate since $\Delta E_1$ is the lowest possible value of $E_c - E_v$ and all terms in the square bracket of (1) are positive. The equation is similar (with $L_z$ replaced by $L_x$) for $S_x$. However, a different value of $\Delta E_1$ must be used.

(b) We may truncate the series in (1) by including $\psi_1$ only under $\sum_c$:

$$S_z = 5200 \sum_v (\Delta E_1)^{-1} | \langle \psi_v | L_z | \psi_1 \rangle |^2$$  

(3)

We will refer to (3) as the first excited state (FE) approximation. Clearly, FE is an under-estimate since all terms in (1) are positive.

Because of the reflection symmetry in the xy-plane of graphite, the $\pi$-states are based entirely on the atomic 2$p_z$ (denoted as $p_0$, where the subscript 0 refers to the azimuthal quantum number) functions of the carbon atom, whereas the $\sigma$-states are based entirely on the atomic 2$s$, 2$p_+$ and 2$p_-$ (abbreviated as s, $p_+$, $p_-$) functions. Except for a small region near the center $\Gamma$ of the BZ, the states are arranged from lower to higher energies in the order of $\sigma$, $\pi$ (unprimed denotes valence or occupied orbitals), $\pi'$, $\sigma'$ (primed denotes conduction or unoccupied orbitals). The selection rules are $\sigma \pi'$, $\pi \sigma'$ for $S_x$ and $\sigma \sigma'$ for $S_z$, thus the denominator
\( \Delta E_1 \) is smaller in the former case. From either (2) or (3), it follows that \( S_x > S_z \) or \( S_a < 0 \) which is in accordance with the experimental results for GIC of acceptor type.

At the center \( \Gamma \) of the BZ, there is no mixing between the s and p-functions. The pure s-functions are omitted here because they do not contribute to CS. The occupied states are \( \sigma^\pm = (p_{A\pm} + p_{B\pm})/2^{\frac{1}{2}} \) (twofold degenerate) and \( \pi = (p_{A0} + p_{B0})/2^{\frac{1}{2}} \) where subscript A and B refer to the two atoms in the unit cell, \( t \) and 0 refer to the azimuthal quantum number \( m = \pm 1, 0 \) of the atomic 2p-functions. Similarly, the unoccupied states are \( \sigma'^\pm = (p_{A\pm} - p_{B\pm})/2^{\frac{1}{2}} \) (twofold degenerate) and \( \pi' = (p_{A0} - p_{B0})/2^{\frac{1}{2}} \). The atoms A and B have the same CS, hence it is sufficient to use \( L_{ZA} \) and \( L_{XA} \). For \( L_{ZA} \), the only nonzero matrix elements are the diagonal elements

\[ \langle p_{A\pm} \mid L_{ZA} \mid p_{A\pm} \rangle, \text{ thus the } \sigma^+ \sigma'^+ \text{ and } \sigma^- \sigma'^- \text{ excitations would contribute to } S_z. \]  

In the AE approximation, we have

\[ S_z = (2)(5200) \langle \tau^+_x \mid L_{ZA} \mid \tau^+_x \rangle <(E_{\pi'} - E_{\tau}) = 5200/(E_{\pi'} - E_{\tau}) \]  

(4)

On the other hand, the FE approximation is smaller by a factor of 2,

\[ S_z = (2)(5200) \langle \pi^+_x \mid L_{ZA} \mid \pi^+_x \rangle 2/(E_{\pi'} - E_{\tau}) = 2600/(E_{\pi'} - E_{\tau}) \]  

(5)

Similarly, the \( \sigma^+_x \pi^+ \) and \( \pi^+_x \pi^- \) excitations would contribute to \( S_x \). In the AE approximation, we have

\[ S_x = (5200) [2(E_{\pi'} - E_{\pi})^{-1} \langle \tau^+_x \mid L_{XA} \mid \pi^+_x \rangle + (E_{\tau'} - E_{\pi})^{-1} \langle \pi \mid L_{XA} \mid \tau^+_x \rangle] \]

\[ = 2600 \left[ (E_{\pi'} - E_{\pi})^{-1} + (E_{\tau'} - E_{\pi})^{-1} \right] \]  

(6)

Again, the FE approximation is smaller by a factor of 2,

\[ \delta_x = 5200 \left[ 2(E_{\pi'} - E_{\pi})^{-1} \langle \tau^+_x \mid L_{XA} \mid \pi^+_x \rangle^2 + 2(E_{\tau'} - E_{\pi})^{-1} \langle \pi \mid L_{XA} \mid \tau^+_x \rangle^2 \right] \]

\[ = 1300 \left[ (E_{\pi'} - E_{\pi})^{-1} + (E_{\tau'} - E_{\pi})^{-1} \right] \]  

(7)
Energy band calculations\textsuperscript{20, 21} give \( E_{\sigma'} - E_{\sigma} = 12 \text{ eV} \), \( E_{\pi'} - E_{\pi} = 15 \text{ eV} \). For center \( \Gamma \), we get \( \delta_z = 433, \delta_x = 347, \delta_a = \delta_z - \delta_x = 86 \) in the AE approximation and \( \delta_z = 216, \delta_x = 174, \delta_a = 42 \) in the FE approximation.

We will now calculate CS at one of the corners \( P \) of the BZ. We have the same \( \pi \) and \( \pi' \) wave functions as before, and \( E_{\pi} = E_{\pi'} \). There are now two \( \sigma \)-functions. One of them, denoted as \( \sigma_1' \), is twofold degenerate and is a mixture of s and p-functions. The other, denoted as \( \sigma_2' \), is a pure p-function. Similarly, there are the unoccupied \( \sigma_1' \) and \( \sigma_2' \) states. In the order of increasing energies, the \( \sigma \) and \( \sigma' \) functions and their energies are\textsuperscript{3}:

\[
E_1 = -\frac{1}{2} (E_s + E_p) - \frac{1}{2} \left[ (E_s - E_p)^2 + 2 \delta \right]^{\frac{1}{2}} \quad \sigma_1 = \alpha s + \beta p, \quad \sigma_1' = \beta s - \alpha p
\]

\[
E_2 = E_p + 3 \delta \quad \sigma_2 = (p_A - p_B)/2^{\frac{1}{2}}
\]

\[
E_1' = -\frac{1}{2} (E_s + E_p) + \frac{1}{2} \left[ (E_s - E_p)^2 + 2 \delta \right]^{\frac{1}{2}} \quad \sigma_1' = \beta s - \alpha p, \quad \sigma_1 = \alpha s + \beta p
\]

\[
E_2' = E_p - 3 \delta \quad \sigma_2' = (p_A - p_B)/2^{\frac{1}{2}}
\]

where \( E_s \) and \( E_p \) are one-electron Hamiltonian integrals for the 2s and 2p atomic functions, \( \delta \) is Hamiltonian integrals between neighboring atoms. They are denoted as \( E(s), E(p), H(s, p) \) by Lomer.\textsuperscript{3} These parameters are all negative. The mixing coefficients \( \alpha \) and \( \beta \) between the 2s and 2p functions are given by:

\[
\alpha = -3 \delta / C, \quad \beta = \left\{(E_s - E_p) + \left[(E_s - E_p)^2 + 2 \delta \right]^{\frac{1}{2}}\right\}/2C
\]
where the normalization constant \( C \) is chosen such that \( \alpha^2 + \beta^2 = 1 \).

The second functions of \( \sigma_1 \) and \( \sigma_1' \) do not contribute to \( CS \) of atom A since there is no angular momentum associated with the atomic \( s_\text{A} \) functions. The \( \sigma_1 \sigma_1' \) and \( \sigma_2 \sigma_2' \) excitations would contribute to \( \delta_\text{z} \). The \( \pi \pi', \pi' \pi, \pi' \sigma_1 \) and \( \pi' \sigma_2 \) excitations would contribute to \( \delta_\text{x} \). In the AE approximation, we get

\[
\delta_\text{z} = (5200) \left[ \left( E_1' - E_1 \right)^{-1} \langle \sigma_1 | L_z^2 | \sigma_1 \rangle \right] + \left( E_2' - E_2 \right)^{-1} \langle \sigma_2 | L_z^2 | \sigma_2 \rangle
\]

\[
= (2600) \left[ 2 \beta^2 (E_1' - E_1)^{-1} + (E_2' - E_2)^{-1} \right]
\]

(10)

\[
\delta_\text{x} = (5200) \left[ \left( E_\pi - E_\pi \right)^{-1} \langle \sigma_1 | L_x^2 | \sigma_1 \rangle \right] + \left( E_\pi - E_\pi \right)^{-1} \langle \sigma_2 | L_x^2 | \sigma_2 \rangle
\]

\[
= (2600) \left[ \beta^2 (E_\pi - E_\pi)^{-1} + \frac{1}{2} (E_\pi - E_\pi)^{-1} + (E_\pi' - E_\pi)^{-1} \right]
\]

(11)

In the FE approximation, we get

\[
\delta_\text{z} = (5200) \left[ \left( E_1' - E_1 \right)^{-1} \langle \sigma_1 | L_z^2 | \sigma_1 \rangle \right]^2 + \left( E_2' - E_2 \right)^{-1} \langle \sigma_2 | L_z^2 | \sigma_2 \rangle^2
\]

\[
= 5200 \alpha^2 \beta^2 (E_1' - E_1)^{-1} + 1300 (E_2' - E_2)^{-1}
\]

(12)

\[
\delta_\text{x} = (5200) \left[ \left( E_\pi' - E_\pi \right)^{-1} \langle \sigma_1 | L_x^2 | \pi' \rangle \right]^2 + \left( E_\pi' - E_\pi \right)^{-1} \langle \sigma_2 | L_x^2 | \pi' \rangle^2
\]

\[
= (1300) \left[ \beta^2 (E_\pi - E_\pi)^{-1} + \frac{1}{2} (E_\pi - E_\pi)^{-1} + \alpha^2 (E_\pi' - E_\pi)^{-1} + \frac{1}{2} (E_\pi' - E_\pi)^{-1} \right].
\]

(13)

For corner P, energy band calculations \( ^{20,21} \) gives \( E_1 = -22 \), \( E_2 = -20 \), \( E_\pi = E_\pi' = -9 \), \( E_1' = 5 \) and \( E_2' = 16 \) eV. In addition, it is necessary to evaluate the mixing coefficients \( \alpha \) and \( \beta \). An estimate of \( E_p = -6 \) may be obtained by averaging the \( \sigma \) and \( \sigma' \) energies at the center \( \Gamma \) of BZ. From the values of \( E_1 \) and \( E_1' \), we get \( E_s = -11 \) and \( H_{sp} = -4.4 \). Substitution into (8) gives
α=0.76 and β=0.65 which are consistent with the estimate 
α² = β² = 0.5 or α = β = 0.71 of Zunger. From (10) to (13), we get 
\( S_z = 189, S_x = 426, S_a = -237 \) for the AE approximation and 
\( S_z = 82, S_x = 177, S_a = -95 \) for the FE approximation.

For the symmetry point Q at the midpoint of the edge 
of BZ, a similar calculation gives 
\( S_z = 242, S_x = 446, S_a = -204 \) for the AE approximation and 
\( S_z = 94, S_x = 163, S_a = -69 \) for the FE approximation. On averaging over the entire BZ, we get the average values 
\( S_a = -170 \) and -70 for the AE and FE approximations respectively.

For GICA, the net effect of the charge transfer is the partial 
removal (usually several per cent) of the valence \( \pi \)-electrons 
near the corner P of the BZ. This charge transfer has no effect 
on \( S_z \) since contributions to \( S_z \) comes from the \( \sigma\sigma' \) excitations 
which are independent of the \( \pi \)-bands. The effect on \( S_x \) is also 
small. With the partial removal of the \( \pi \)-electrons, there is 
less \( \pi\sigma' \) excitation. However, the \( \pi \)-band is now partially 
vacant, hence some \( \sigma\pi \) excitations are also possible. There is 
considerable cancellation between these two contributions since 
the \( \pi \) and \( \pi' \) bands are close together near the corner P. Similarly, 
the charge transfer effects on both \( S_z \) and \( S_x \) may also be small 
for GICD.

In summary, the "band" contributions to \( S_x \) and \( S_z \) come 
from the entire BZ and are independent of charge transfer and other 
detailed features of the energy bands. We have found that \( S_a \) is 
between -170 and -70. These results are in general agreement 
with the experimental data on graphite and GICA (nearly universal 
values of \( S_x \sim 180 \) relative to TMS and \( S_a \sim -140 \)). We note that 
it is very difficult to determine \( S_z \) of pure graphite experimentally.
because of the weak signals; the available experimental data\textsuperscript{8-10,13,14} are not consistent with each other and further work is needed.

We will now consider the "dipolar" contributions to CS which may be important for GICD. Knight shift measurements\textsuperscript{23,24} indicate that there is considerable mixing between the metal conduction electrons and the carbo.: 2p\pi-electrons. The CS would come from the dipolar magnetic fields at the nuclear site due to the spin-polarized \pi-orbitals in the presence of the external magnetic field \( B \). For \( B \) parallel to the z-axis, we have

\[
\delta_z = \left( \frac{\eta \chi}{N} \right) \left< r^{-3} \right>_{2p} \left< 3\cos^2 \theta - 1 \right>
\]

where \( N \) is the Avagadro's number, \( \eta \) is the fraction of the polarized electrons that reside on carbon atoms, \( \chi \) is the Pauli spin susceptibility in cm\(^3\)/mole and \( \theta \) is the azimuthal angle for the \pi-orbital of the carbon atom. The first pair of angular brackets is radial averaging over the carbon atom \pi-orbital whereas the second pair is angular averaging. Since the angular dependence of the \pi-orbital wave function is given by \( \cos \theta \), we get

\[
\left< 3\cos^2 \theta - 1 \right> = \int_0^\pi \cos^2 \theta (3\cos^2 \theta - 1) \sin \theta \, d\theta / \int_0^\pi \cos^2 \theta \sin \theta \, d\theta = 4/5. \quad (14)
\]

Hence \( \delta_z = \left( \frac{4}{5} \right) \left( \frac{\eta \chi}{N} \right) \left< r^{-3} \right>_{2p} \). Because of the axial symmetry of \pi-orbitals about the z-axis, it follows that \( \delta_x = -\delta_z/2 \). Thus the "dipolar" contribution to CS anisotropy is

\[
\delta_a = \delta_z - \delta_x = \left( \frac{6}{5} \right) \left( \frac{\eta \chi}{N} \right) \left< r^{-3} \right>_{2p} \quad (15)
\]

where \( \left< r^{-3} \right>_{2p} = 12.2 \, \text{Å}^{-3} \). Because of the close proximity of \pi-orbital to the carbon nucleus, \( \delta_a \) is enhanced by approximately a factor of \( \nu \left< r^{-3} \right>_{2p} \sim 100 \) (where \( \nu \sim 9 \, \text{Å}^3 \) is the crystal volume per carbon atom) over the Pauli contribution to the dimensionless volume susceptibility. The experimental susceptibility is the sum of several contributions\textsuperscript{25}, thus the identification of a relatively
small Pauli component may be somewhat uncertain.

We note that the presence of the metallic electrons are also clearly demonstrated by the electronic specific heat\textsuperscript{11,26} which is proportional to temperature with the proportionality constant $\gamma$. This constant is nearly proportional to the metal concentration and may be related to the Pauli spin susceptibility $\chi$:\textsuperscript{25}

\begin{equation}
\chi = (3/\pi^2)(\mu/k_B)^2 \gamma
\end{equation}

where $\mu$ is the Bohr magneton, $k_B$ is the Boltzmann constant, $\gamma$ is in units of erg(mole)$^{-1}$(0K)$^{-2}$ and $\chi$ is in units of cm$^3$mole$^{-1}$. Thus the Pauli susceptibility $\chi$ may be eliminated between eq. (15) and (16). The result is $\delta_a = 333 \gamma$ when $\delta_a$ is expressed in ppm and $\gamma$ in units of (mJ)(mole)$^{-1}$(0K)$^{-2}$. By combining with the experimental value of the "band" contribution $\delta_a = -140$ which is independent of intercalant concentration, we get

\begin{equation}
\delta_a = -140 + 333 \gamma
\end{equation}

Using the experimental values\textsuperscript{11,26} of $\gamma$, we have calculated $\delta_a$ from (17) and have compared with some of the experimental values\textsuperscript{8-16} in the following table for the first and second stage GICD with the general composition $C_xK$ and $C_xCs$. From the Knight shift data on Cs, it is plausible that the polarized electrons are shared equally between the carbon and metal atoms, hence we have identified $\gamma$ with the mole fraction of carbon atoms in GIC, $\gamma = x/(x+1)$. The agreement is satisfactory. The variation in $\delta_i$ (isotropic part of CS with TMS as reference) is about a factor of 10 smaller than the anisotropic part $\delta_a$. This isotropic term is usually ascribed to polarizations of the carbon core s-electrons by the unpaired $\pi$-electron spins\textsuperscript{27}, and the factor of 10 reduction has been previously observed for similar atoms.\textsuperscript{28} The results for
GICA are also included for comparison. Very small values of $\gamma$ have been observed for these compounds, hence $\xi_a$ comes from the "band" contribution only.

\[
\begin{array}{cccccc}
\text{GICA} & \text{C}_8\text{K} & \text{C}_{24}\text{K} & \text{C}_8\text{Cs} & \text{C}_{36}\text{Cs} \\
\gamma \text{, mJ mol}^{-1}\text{K}^{-2} & 0 & 0.697 & 0.241 & 0.63 & 0.19 \\
\eta & \frac{8}{9} & \frac{24}{25} & \frac{8}{9} & 36/37 \\
\xi_a, \text{ calc.} & -140 & 66 & -63 & 46 & -79 \\
\xi_a, \text{ exptl.} & -140 & 57 & -63 & 35 & -93 \\
\xi_i, \text{ exptl.} & 133 & 113 & 121 & 107 & 123 \\
\end{array}
\]

Alternatively, the choice $\eta=1$ may be inferred from the x-ray photoelectron data\textsuperscript{29,30} which have indicated almost complete charge transfer of metallic electrons to carbon atoms. From (17), the calculated $\xi_a$ will be changed only slightly (about 20 ppm for $\text{C}_8\text{K}$ and $\text{C}_8\text{Cs}$, 3 ppm for higher stage compounds) and the agreement with experimental data is still satisfactory.

In conclusion, the CS in GIC may be expressed as the sum of "band" and "dipole" contributions. The "band" term is determined by the gross features of the two-dimensional band structure and is independent of the intercalant concentration. The "dipole" term is due to the mixing between metal and carbon electrons and is absent for GICA. It is related to the electronic specific heat for GICD. There is general agreement with the experimental data.

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REFERENCES

5. J. D. Memory, Quantum Theory of Magnetic Resonance Parameters, 
8. J. Conard, H. Estrade, P. Lauginie, H. Fuzellier, G. Furcin, 
10. J. Conard, P. Lauginie, H. Estrade-Szwarckopf, G. Hermann, 
11. U. Mizutani, M. Suganuma and T. Kondow, in Physics of 
    Intercalation Compounds, edited by L. Pietronero and E. 
12. H. A. Resing, D. C. Weber, M. Anderson, G. R. Miller, M. Moran, 
    C. F. Poranski Jr., and L. Mattix, Extended Abstracts (Amer. 
17. A. B. Strong, D. Ikenberry and D. M. Grant, J. Mag. Res. 2, 
    145 (1973); there appears to be a factor of 4 error in eq. (30).
25. Reference 1, p. 272.