MAXIMUM LIKELIHOOD ESTIMATION WITH EMphasis ON AIRCRAFT FLIGHT DATA

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ABSTRACT

Accurate modeling of flexible space structures is an important field that is currently under investigation. Parameter estimation, using methods such as maximum likelihood, is one of the ways that the model can be improved. The maximum likelihood estimator has been used to extract stability and control derivatives from flight data for many years. Most of the literature on aircraft estimation concentrates on new developments and applications, assuming familiarity with basic estimation concepts. This paper presents some of these basic concepts.

The paper briefly discusses the maximum likelihood estimator and the aircraft equations of motion that the estimator uses. The basic concepts of minimization and estimation are examined for a simple computed aircraft example. The cost functions that are to be minimized during estimation are defined and discussed. Graphic representations of the cost functions are given to help illustrate the minimization process. Finally, the basic concepts are generalized, and estimation from flight data is discussed. Specific examples of estimation of structural dynamics are included. Some of the major conclusions for the computed example are also developed for the analysis of flight data.

INTRODUCTION

Accurate modeling of flexible space structures is an important area that is currently under investigation. The mathematical modeling of these structures can be improved using parameter estimation. Such techniques have been successfully used to estimate aircraft stability and control derivatives and refine aircraft mathematical models. Some of the experience gained in the aircraft problem can be applied directly to analysis of flexible space structures.

The maximum likelihood estimator has been used to obtain stability and control estimates from flight data for nearly 20 years. The results of many applications have been reported worldwide. Reference 1 contains a representative list of some of these reports. Several good texts (including Refs. 2 and 3) contain thorough treatments of the theory of maximum likelihood estimation. Experience reports (Refs. 1, 4, and 5) pointing out practical considerations for

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applying the maximum likelihood estimator have also been published. Stability and control derivatives estimated from flight data are currently required for correlation studies with predictive techniques, handling qualities documentation, design compliance, aircraft simulation enhancement and refinement, and control system design. Correlation, simulation, and control system design applications (including the space shuttle) are discussed in Ref. 6. Current studies have concentrated on estimation model structure determination (Refs. 7 and 8), equation error with state reconstruction (Refs. 9 to 11), and maximum likelihood estimation in the frequency domain (Refs. 12 and 13).

Most of the reports in the estimation area concentrate on new developments and applications, assuming familiarity with the basic concepts of maximum likelihood estimation. In this paper some of these basic concepts are reviewed, concentrating on simple, idealized models. These simple models provide insights applicable to a wide variety of real problems.

This paper also presents some of the basics of maximum likelihood estimation as applied to the aircraft problem. It briefly discusses the maximum likelihood estimator and the aircraft equations of motion that the estimator uses. The basic aspects of minimization and estimation are then examined in detail for a simple computed aircraft example. Finally, the discussion is expanded to the general aircraft estimation problem including specific examples of estimation of structural dynamics.

SYMBOLS

\( A, B, C, D, F, G \) system matrices
\( a_n \) normal acceleration positive upward, g
\( a_x \) longitudinal acceleration, g
\( a_y \) lateral acceleration, g
\( a_z \) normal acceleration positive upward, g
\( b \) reference span, ft
\( C_L \) coefficient of rolling moment
\( C_n \) coefficient of yawing moment
\( C_X \) coefficient of axial force
\( C_Y \) coefficient of side force
\( C_Z \) coefficient of normal force
\( f(\cdot), g(\cdot) \) general functions
measurement noise covariance matrix

acceleration due to gravity, ft/sec²

approximation to the information matrix

moment of inertia about subscripted axis, slug-ft²

general index

cost function

sidewash factor

rolling moment divided by $I_x$, deg/sec²

rolling moment, ft-lb

rolling moment due to yaw jet, ft-lb

pitching moment divided by $I_y$, deg/sec²

mass, slug

number of time points or cases or yawing moment divided by $I_z$, deg/sec²

state noise vector or number of unknowns

estimated roll rate due to turbulence, deg/sec

roll rate, deg/sec

pitch rate, deg/sec

dynamic pressure, lb/ft²

innovation covariance matrix

yaw rate, deg/sec

reference area, ft²

time increment, sec

time, sec

control input vector

forward velocity, ft/sec
$x$ state vector

$x_{ay}, y_{ay}, z_{ay}$ distance between lateral accelerometer and the center of gravity along the appropriate axis, ft

$z$ observation vector

$\tilde{z}_{\xi}$ predicted Kalman-filtered estimate

$\alpha$ angle of attack, deg

$\beta$ angle of sideslip, deg

$\hat{\beta}_g$ estimated angle of sideslip due to turbulence, deg

$\Delta$ time sample interval, sec.

$\delta$ control deflection, deg

$\delta_a$ aileron deflection, deg

$\delta_e$ elevon deflection, deg

$\delta_r$ rudder deflection, deg

$\eta$ measurement noise vector

$\theta$ pitch angle, deg

$\mu$ mean

$\xi$ vector of unknowns

$\sigma$ standard deviation

$\tau$ time, sec

$\phi$ transition matrix or bank angle, deg

$\psi$ integral of transition matrix, or heading angle, deg

$\omega$ frequency, rad/sec

Subscripts:

$p, q, r, a, \dot{\alpha}, \dot{\beta}, \xi, \delta, \delta_a, \delta_r, \delta_e$ partial derivative with respect to subscripted quantity

$0$ bias or at time zero

$m$ measured quantity
OTHER NOMENCLATURE:

~ predicted estimate
* estimate
t transpose
' indicates moment in ft-lb

MAXIMUM LIKELIHOOD ESTIMATION

The concept of maximum likelihood is discussed in this section. First the general heuristic problem is discussed, and then the specific equations for obtaining maximum likelihood estimates for the aircraft problem are given. In the following sections, both the concepts and the computations involved in a simple but realistic example are discussed in detail.

The aircraft parameter estimation problem can be defined quite simply in general terms. The system investigated is assumed to be modeled by a set of dynamic equations containing unknown parameters. To determine the values of the unknown parameters, the system is excited by a suitable input, and the input and actual system response are measured. The values of the unknown parameters are then inferred based on the requirement that the model response to the given input match the actual system response. When formulated in this manner, the problem of identifying the unknown parameters can be easily solved by many methods; however, complicating factors arise when application to a real system is considered.

The first complication results from the impossibility of obtaining perfect measurements of the response of any real system. The inevitable sensor errors are usually included as additive measurement noise in the dynamic model. Once this noise is introduced, the theoretical nature of the problem changes drastically. It is no longer possible to exactly identify the values of the unknown parameters; instead, the values must be estimated by some statistical criterion. The theory of estimation in the presence of measurement noise is relatively straightforward for a system with discrete time observations, requiring only basic probability.

The second complication of real systems is the presence of state noise. State noise is random excitation of the system from unmeasured sources, the standard example for the aircraft stability and control problem being atmospheric turbulence. If state noise is present and measurement noise is neglected, the analysis results in the regression algorithm.

When both state and measurement noise are considered, the problem is more complex than in the cases that have only state noise or only measurement noise. Reference 14 develops a mixed continuous/discrete maximum likelihood formulation that allows for both state and measurement noise. This formulation has a continuous system model with discrete sampled observations.
The final problem for real systems is modeling. It has been assumed throughout the above discussion that for some value (called the "correct" value) of the unknown parameter vector, the system is correctly described by the dynamic model. Physical systems are seldom described exactly by simple dynamic models, so the question of modeling error arises. No comprehensive theory of modeling error is available. The most common approach is to ignore it: Any modeling error is simply treated as state noise or measurement noise, or both, in spite of the fact that the modeling error may be deterministic rather than random. The assumed noise statistics can then be adjusted to include the contribution of the modeling error. This procedure is not rigorously justifiable, but, combined with a carefully chosen model, it is probably the best approach available.

With the above discussion in mind, it is possible to make a more precise, mathematically probabilistic statement of the parameter estimation problem. The first step is to define the general system model (aircraft equations of motion). This model can be written in the continuous/discrete form as

\[ x(t_0) = x_0 \]  \hspace{1cm} (1)

\[ \dot{x}(t) = f[x(t), u(t), \xi] + F(\xi)n(t) \]  \hspace{1cm} (2)

\[ z(t_i) = g[x(t_i), u(t_i), \xi] + G(\xi)\eta_i \]  \hspace{1cm} (3)

where \( x \) is the state vector, \( z \) is the observation vector, \( f \) and \( g \) are system state and observation functions, \( u \) is the known control input vector, \( \xi \) is the unknown parameter vector, \( n \) is the state noise vector, and \( \eta \) is the measurement noise vector. The state noise vector is assumed to be zero-mean white Gaussian and stationary, and the measurement noise vector is assumed to be a sequence of independent Gaussian random variables with zero mean and identity covariance. For each possible estimate of the unknown parameters, a probability that the aircraft response time histories attain values near the observed values can then be defined. The maximum likelihood estimates are defined as those that maximize this probability. Maximum likelihood estimation has many desirable statistical characteristics; for example, it yields asymptotically unbiased, consistent, and efficient estimates (Ref. 15).

If there is no state noise and the matrix \( G \) is known, then the maximum likelihood estimator minimizes the cost function

\[ J(\xi) = \frac{1}{2} \sum_{i=1}^{N} [z(t_i) - \tilde{z}(t_i)]^* (GG^*)^{-1} [z(t_i) - \tilde{z}(t_i)] \]  \hspace{1cm} (4)

where \( GG^* \) is the measurement noise covariance matrix, and \( \tilde{z}(t_i) \) is the computed response estimate of \( z \) at \( t_i \) for a given value of the unknown parameter vector \( \xi \). The cost function is a function of the difference between the measured and computed time histories.
If Eqs. (2) and (3) are linearized (as is the case for the stability and control derivatives in the aircraft problem),

\begin{align*}
x(t_0) &= x_0 \\
\dot{x}(t) &= Ax(t) + Bu(t) + \xi_n(t) \\
z(t_i) &= Cx(t_i) + Du(t_i) + G\eta_i
\end{align*}

For the no-state-noise case, the \( \tilde{z}_\xi(t_i) \) term of Eq. (4) can be approximated by

\begin{align*}
\tilde{x}_\xi(t_0) &= x_0(\xi) \\
\tilde{x}_\xi(t_{i+1}) &= \psi \tilde{x}_\xi(t_i) + \psi [u(t_i) + u(t_{i+1})]/2 \\
\tilde{z}_\xi(t_i) &= C\tilde{x}_\xi(t_i) + Du(t_i)
\end{align*}

where

\[ \psi = \exp \left[ A(t_{i+1} - t_i) \right] \]

\[ \psi = \int_{t_i}^{t_{i+1}} \exp \left( AT \right) \, dt \]

When state noise is important, the nonlinear form of Eqs. (1) to (3) is intractable. For the linear model defined by Eqs. (5) to (7), the cost function that accounts for state noise is

\[ J(\xi) = \frac{1}{2} \sum_{i=1}^{N} [z(t_i) - \tilde{z}_\xi(t_i)]^* R^{-1} [z(t_i) - \tilde{z}_\xi(t_i)] + \frac{1}{2} N \ln |R| \]

where \( R \) is the innovation covariance matrix. The \( \tilde{z}_\xi(t_i) \) term in Eq. (11) is the Kalman-filtered estimate of \( z \), which, if the state noise covariance is zero, reduces to the form of Eq. (4). If there is no state noise, the second term of Eq. (11) is of no consequence (unless one wishes to include elements of the \( G \) matrix as unknowns), and \( R \) can be replaced by \( GG^* \) which makes Eq. (11) the same as Eq. (4).

To minimize the cost function \( J(\xi) \), we can apply the Newton-Raphson algorithm which chooses successive estimates of the vector of unknown coefficients, \( \hat{\xi} \). Let \( L \) be the iteration number. The \( L+1 \) estimate of \( \hat{\xi} \) is then obtained from the \( L \) estimate as follows:

\[ \hat{\xi}_{L+1} = \hat{\xi}_L - [V_{\xi}^* J(\hat{\xi}_L)]^{-1} [V_{\xi} J(\hat{\xi}_L)] \]
The first and second gradients are defined as

$$V_{\xi J}(\xi) = - \sum_{i=1}^{N} [z(t_i) - \tilde{z}(t_i)]^* (GG^*)^{-1} [V_{\xi \tilde{z}}(t_i)]$$

$$V_{\xi J}(\xi) = \sum_{i=1}^{N} [V_{\xi \tilde{z}}(t_i)]^* (GG^*)^{-1} [V_{\xi \tilde{z}}(t_i)]$$

$$- \sum_{i=1}^{N} [z(t_i) - \tilde{z}(t_i)]^* (GG^*)^{-1} [V_{\xi \tilde{z}}(t_i)]$$

The Gauss-Newton approximation to the second gradient is

$$V_{\xi J}(\xi) \approx \sum_{i=1}^{N} [V_{\xi \tilde{z}}(t_i)]^* (GG^*)^{-1} [V_{\xi \tilde{z}}(t_i)]$$

Figure 1 illustrates the maximum likelihood estimation concept. The measured response of the aircraft is compared with the estimated response, and the difference between these responses is called the response error. The cost functions of Eqs. (4) and (11) include this response error. The Gauss-Newton computational algorithm is used to find the coefficient values that maximize the cost function. Each iteration of this algorithm provides a new estimate of the unknown coefficients on the basis of the response error. These new estimates of the coefficients are then used to update the mathematical model of the aircraft, providing a new estimated response and, therefore, a new response error. The updating of the mathematical model continues iteratively until a convergence criterion is satisfied. The estimates resulting from this procedure are the maximum likelihood estimates.

The Gauss-Newton approximation, which is sometimes referred to as modified Newton-Raphson, is computationally much easier than the Newton-Raphson approximation because the second gradient of the innovation never needs to be calculated. In addition, it can have the advantage of speeding the convergence of the algorithm, as is discussed in the SIMPLE AIRCRAFT EXAMPLE section.

The maximum likelihood estimator also provides a measure of the reliability of each estimate based on the information obtained from each dynamic maneuver. This measure of the reliability, analogous to the standard deviation, is called the Cramér-Rao bound (Ref. 16) or the uncertainty level. The Cramér-Rao bound is computed by current programs should generally be used as a measure of relative accuracy rather than absolute accuracy. This bound is obtained from the approximation of the information matrix, $H$. This matrix equals the approximation to the second gradient given by Eq. (14b). The bound for each unknown is the square root of the corresponding diagonal element of $H$. That is, for the $i$th unknown, the Cramér-Rao bound is $\sqrt{H(i,i)}$. 

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The Maine-Iliff formulation (Ref. 14) and minimization algorithms discussed above are implemented with the Iliff-Maine code (MMLE3 maximum likelihood estimation program). The program and computational algorithms are described fully in Ref. 17. All the computations shown and described in the remainder of the paper use the algorithms exactly as described in Ref. 17.

**AIRCRAFT EQUATIONS OF MOTION**

For the discussion that follows in later sections of this paper, some knowledge of the aircraft equations of motion is assumed. To clarify some of the discussion, the aircraft equations are discussed briefly in this section.

First, the axis system on which the aircraft equations of motion are based is discussed. Figure 2(a) shows the aircraft reference body-axis system and the conventional control surfaces. The origin of the body-axis system is at the center of gravity. The sign convention for this axis system is defined by the right-hand rule with the x-axis defined as positive forward on the aircraft. The longitudinal acceleration $a_x$ and nondimensional axial force coefficient $C_x$ are defined along this axis, and the roll rate $p$ and rolling moment $L'$ are defined about this axis. The y-axis is defined as positive out the right wing. The lateral acceleration $a_y$ and nondimensional side force coefficient $C_y$ are defined along this axis, and the pitch rate $q$ and pitching moment $M'$ are defined about this axis. The z-axis is defined as positive out the bottom of the aircraft. The normal acceleration $a_z$ and nondimensional normal force coefficient $C_z$ are defined along this axis, and the yaw rate $r$ and yawing moment $N'$ are defined about this axis. The normal acceleration is sometimes defined as positive upward but is then referred to as $a_N$. The three moments $(L', M', \text{ and } N')$ are usually divided by the corresponding moments of inertia $I_x$, $I_y$, and $I_z$, and are then referred to without the prime as $L$, $M$, and $N$. These quantities are nondimensionalized $(C_L$, $C_M$, and $C_N$, respectively) for use in the equations of motion soon to be discussed. The primary control about the roll axis (x-axis) is the aileron $(\delta_a)$, about the pitch axis (y-axis) is the elevator $(\delta_e)$, and about the yaw axis (z-axis) is the rudder $(\delta_r)$. Some aircraft have other controls, but in this paper these will only be defined where they are discussed (the reaction control jets on the space shuttle, for example).

The Euler angles $\phi$, $\theta$, and $\psi$ define the aircraft attitude with respect to the earth. These angles define the rotations which transform earth-fixed axes to the aircraft reference body-axis system of Fig. 2(a). The order of rotation must be about the z-axis ($\psi$), then the y-axis ($\theta$), and finally the x-axis ($\phi$) for the aircraft equations of motion that will be written subsequently.

For stability and control analysis, the velocity of the aircraft with respect to the air (not with respect to the earth) is of primary interest. Figure 2(b) shows the relationship between the aircraft axis system and the flow angles. The flow angle in the x-z plane is the angle of attack $(\alpha)$, and the flow angle in the x-y plane is the angle of sideslip $(\beta)$. A more rigorous and
detailed definition is required for the derivation of the equations of motion, but the above definitions are sufficient to define the following equation of motion.

Generalized nonlinear equations of motion are given in detail in Ref. 17, which fully describes the Iliff-Mine code (MMLE3 program). All computations and aircraft examples in this paper use the linearized form for the lateral-directional equations. These equations are given below and referred to in the remainder of the paper.

\[ \dot{\beta} = \frac{\text{sgn}}{\text{sgn}V} (C_Y + \dot{\beta}_0) + \frac{q}{V} \cos \theta \sin \phi + p \sin \alpha - r \cos \alpha \]  

\[ \dot{p} = p + r \cos \phi \tan \theta + q \sin \phi \tan \theta + \dot{\phi}_0 \]  

where

\[ C_Y = C_{Y\beta} \beta + C_{Yp} \frac{pb}{V} + C_{Yr} \frac{rb}{V} + C_{Y\delta} \delta + C_Y \]  

\[ C_I = C_{I\beta} \beta + C_{Ip} \frac{pb}{V} + C_{Ir} \frac{rb}{V} + C_{I\delta} \delta + C_I \]  

\[ C_n = C_{n\beta} \beta + C_{np} \frac{pb}{V} + C_{nr} \frac{rb}{V} + C_{n\delta} \delta + C_n \]  

where the \( \delta \) term is summed over all controls.

The observation equations are

\[ \beta_m = \kappa \beta - \frac{z_B}{V} p + \frac{x_B}{V} r \]  

\[ P_m = P \]  

\[ r_m = r \]  

\[ \phi_m = \phi \]  

\[ a_m = \frac{g}{mg} C_Y - \frac{z_{av}}{g} p + \frac{x_{av}}{g} r - \frac{y_{av}}{g} (p^2 + r^2) \]  

\[ I_m = p + P_0 \]  

\[ r_m = r + r_0 \]
The state, control, and observation vectors for the lateral-directional mode can then be defined as

\[ x = (\beta \ p \ r \ \phi) \]  
\[ u = (\delta_a \ \delta_r) \]  
\[ z = (\beta_a \ P_m \ r_m \ \phi_m \ \delta_m \ P_m \ r_m) \]  

\[ (29) \]  
\[ (30) \]  
\[ (31) \]

**SIMPLE AIRCRAFT EXAMPLE**

The basic concepts involved in a parameter estimation problem can be illustrated by using a simple example representative of a realistic aircraft problem. The example chosen here is representative of an aircraft that exhibits pure rolling motion from an aileron input. This example, although simplified, typifies the motion exhibited by many aircraft in particular flight regimes, such as the F-14 aircraft flying at high dynamic pressure, the F-111 aircraft at moderate speeds with the wing in the forward position, and the T-37 aircraft at low speed.

Derivation of an equation describing this motion is straightforward. Figure 2(c) shows a sketch of an aircraft with the x-axis perpendicular to the plane of the figure (positive forward on the aircraft). The rolling moment \( L' \), roll rate \( p \), and aileron deflection \( \delta_a \) are positive as shown. For this example, the only state is \( p \) and the only control is \( \delta_a \). The result of summing moments is

\[ I_x p = L'(p, \delta_a) \]  
\[ (32) \]

The first-order Taylor expansion then becomes

\[ \dot{p} = I_{p} \dot{p} + L_{p} \delta_a \]  
\[ (33) \]

where

\[ L' = I_x L \]

Since the aileron is the only control, it is notationally simpler to use \( \delta \) instead of \( \delta_a \) for the discussion of this example. Equation (33) can then be written as

\[ \dot{p} = I_{p} \dot{p} + L_{p} \delta \]  
\[ (34) \]

An alternate approach that results in the same equation is to combine Eq. (16) with Eq. (20), substituting for \( C_z \), and then eliminate the terms that are zero for our example. This yields

\[ \dot{p} I_x = -gsb \ C_{b} \ \frac{p b}{2} + C_{t} \ \delta \]  
\[ (35) \]
where \( p \) is the roll rate and \( \delta \) is the aileron deflection. Rearranging terms, the equation can be put into the dimensional derivative form of Eq. (34).

Equation (34) is a simple aircraft equation where the forcing function is provided by the aileron and the damping by the damping-in-roll term, \( L_p \). In subsequent sections we examine in detail the parameter estimation problem where Eq. (34) describes the system. For this single-degree-of-freedom problem, the maximum likelihood estimator is used to estimate either \( L_p \) or \( L_\delta \) or both for a given computed time history.

We will assume that the system has measurement noise, but no state noise as in Eqs. (1), (2), and (3). Equation (4) then gives the cost function for maximum likelihood estimation. The weighting \( G^2 \) is unimportant for this problem, so let it equal 1. For our example, Eqs. (2) and (3) become \( x_1 = p_1 \) and \( z_1 = z_1 \). Therefore, Eq. (4) becomes

\[ J(L_p, L_\delta) = \frac{1}{2} \sum_{i=1}^{n} [p_i - \tilde{p}_i(L_p, L_\delta)]^2 \]  

where \( p_i \) is the value of the measured response \( p \) at time \( t_i \) and \( \tilde{p}_i(L_p, L_\delta) \) is the computed time history of \( \tilde{p} \) at time \( t_i \) for \( L_p = \hat{L}_p \) and \( L_\delta = \hat{L}_\delta \). Throughout the rest of the paper, where computed data (not flight data) are used, the measured time history refers to \( p_i \), and the computed time history refers to \( \tilde{p}_i(L_p, L_\delta) \). The computed time history is a function of the current estimates of \( L_p \) and \( L_\delta \), but the measured time history is not.

The most straightforward method of obtaining \( p_i \) is with Eqs. (3) and (8). In terms of the notation stated above,

\[ \tilde{p}_{i+1} = \psi \tilde{p}_i + \psi (\delta_{i+1} + \delta_{i+1}/2) \]  

where

\[ \psi = \exp(L_p \Delta) \]

\[ \phi = \frac{L_\delta [1 - \exp(L_p \Delta)]}{L_p} \]

and \( \Delta \) is the length of the sample interval \( (t_{i+1} - t_i) \). Simplifying the notation

\[ \delta_{i+1/2} = (\delta_{i+1} + \delta_{i+1}/2) \]  

then

\[ \tilde{p}_{i+1} = \psi \tilde{p}_i + \psi \delta_{i+1/2} \]
The maximum likelihood estimate is obtained by minimizing Eq. (36). The Gauss-Newton method described earlier is used for this minimization. Equation (12) is used to determine successive values of the estimates of the unknowns during the minimization.

For this simple problem, \( \hat{\xi} = [\hat{L}_p, \hat{L}_g] \) and successive estimates of \( \hat{L}_p \) and \( \hat{L}_g \) are determined by updating Eq. (12). The first and second gradients of Eq. (12) are defined by Eqs. (13) and (14). The complete set of equations is given in Ref. 17.

The entire procedure can now be written for obtaining the maximum likelihood estimates for this simple example. To start the algorithm, an initial estimate of \( L_p \) and \( L_g \) is needed. This is the value of \( \xi_0 \). With Eq. (12), \( \xi_1 \) and subsequently \( \xi_L \) are defined by using the first and second gradients of \( J(L_p, L_g) \) from Eq. (36). The gradients for this particular example from Eq. (13) and (14b) are

\[
V_{\xi L}(\xi_L) = -\sum_{i=1}^{N} (p_i - \tilde{p}_i) V_{\xi} \tilde{p}_i
\]

\[
V_{\xi L}^2(\xi_L) = \sum_{i=1}^{N} (V_{\xi p_i})^2 (V_{\xi \tilde{p}_i})
\]

With the specific equations defined in this section for this simple example, we can now proceed in the next section to the computational details of a specific example.

**Computational Details of Minimization**

In the previous section we specified the equations for a simple example and described the procedure for obtaining estimates of the unknowns from a dynamic maneuver. In this section we give the computational details for obtaining the estimates. Some of the basic concepts of parameter estimation are best shown with computed data where the correct answers are known. Therefore, in this section we study two examples involving computed time histories. The first example is based on data that have no measurement noise, which results in estimates that are the same as the correct value. The second example contains significant measurement noise; consequently, the estimates are not the same as the correct values. Throughout the rest of the paper, where computed data are used, the term "no-noise case" is used for the case with no noise added and "noisy case" for the case where noise has been added.

Since we are studying a simple computed example, it is desirable to keep it simple enough to complete some or all of the calculations on a home computer or, with some labor, on a calculator. With this in mind, the number of data points needs to be kept small. For this computed example, 10 points (time samples) are
used. The simulated data, which we refer to as the measured data, are based on Eq. (34). We use the same correct values of \( L_p \) and \( L_6 \) (-0.2500 and 10.0, respectively) for both examples. In addition, the same input (\( \delta \)) is used for both examples. The sample interval (\( \Delta \)) is 0.2 sec, and the initial conditions are zero. Tables of all the significant intermediate values are given with each example. These values are given to four significant digits, although to obtain exact same values with a computer or calculator requires the use of 13 significant digits, as in the computation of these tables. If the four-digit numbers are used in the computation, the answers will be a few tenths of a percent off, but will still serve to illustrate the minimization accuracy. In both examples, the initial values of \( L_p \) and \( L_6 \) (or \( \xi_0 \)) are -0.5 and 15.0, respectively.

### Example With No Measurement Noise

The measurement time history for no measurement noise (no-noise case) is shown in Fig. 3. The aileron input starts at zero, goes to a fixed value, and then returns to zero. The resulting roll-rate time history is also shown. The values of the measured roll rate to 13 significant digits are given in Table 1 along with the aileron input.

Table 2 shows the values for \( \hat{L}_p \), \( \hat{L}_6 \), and \( \hat{J} \) for each iteration, along with the values of \( \psi \) and \( \xi \) needed for calculations of \( \bar{P}_1 \). In three iterations the algorithm converges to the correct values to four significant digits for both \( L_p \) and \( L_6 \). \( \hat{L}_6 \) overshoots slightly on the first iteration and then comes quickly to the correct answer. \( \hat{L}_p \) overshoots slightly on the second iteration.

Figure 4 shows the match between the measured data and the computed data for each of the first three iterations. The match is very good after two iterations. The match is nearly exact after three iterations.

Although the algorithm has converged to four-digit accuracy in \( L_p \) and \( L_6 \), the value of the cost function, \( \bar{J} \), continues to decrease rapidly between iterations 3 and 4. This is a consequence of using the maximum likelihood estimator on data with no measurement noise. Theoretically, using infinite accuracy the value of \( J \) at the minimum should be zero. However, with finite accuracy the value of \( \bar{J} \) becomes small but never quite zero. This value is a function of the number of significant digits that are being used. For the 13-digit accuracy used here, the cost eventually decreases to approximately \( 0.3 \times 10^{-28} \).

### Example With Measurement Noise

The data used in this example (noisy case) are the same as those used in the previous section, except that pseudo-Gaussian noise has been added to the roll rate. The time history is shown in Fig. 5. The signal-to-noise ratio is quite low in this example, as is readily apparent by comparing Figs. 3 and 5. The exact values of the time history to 13-digit accuracy are shown in Table 3. The values of \( L_p \), \( L_6 \), \( \psi \), \( \xi \), and \( J \) are shown for each iteration in Table 4.
The algorithm converges in four iterations. The behavior of the coefficients as they approach convergence is much like the no-noise case. The most notable results of this case are the converged values of \( L_p \) and \( L_6 \), which are somewhat different from the correct values. The match between the measured and computed time histories is shown in Fig. 6 for each iteration. No change in the match is apparent for the last two iterations. The match is very good considering the amount of measurement noise.

In Fig. 7, the computed time history for the correct values of \( L_p \) and \( L_6 \) is compared to that for the noisy-case estimates of \( T \), \( \nu \), and \( L_6 \). Because the algorithm converged to values somewhat different from the correct values, the two computed time histories are similar but not identical.

The accuracy of the converged elements can be assessed by looking at the Cramér-Rao inequality (Refs. 16 and 17) discussed earlier. The Cramér-Rao bound can be obtained from the following approximation to the information matrix.

\[
H = 2(J_{\text{minimum}})(V_e^2)^{-1}/(N-1)
\]

The Cramér-Rao bounds for \( L_p \) and \( L_6 \) are the square roots of the diagonal elements of the \( H \) matrix, or \( \sqrt{H(1,1)} \) and \( \sqrt{H(2,2)} \), respectively. The Cramér-Rao bounds are 0.1593 and 1.116 for \( \hat{L}_p \) and \( \hat{L}_6 \), respectively. The errors in \( L_p \) and \( L_6 \) are less than the bounds.

**Cost Functions**

In the previous section we obtained the maximum likelihood estimates for computed time histories by minimizing the values of the cost function. To fully understand what occurs in this minimization, we must study in more detail the form of the cost functions and some of their more important characteristics. In this section, the cost function for the no-noise case is discussed briefly. The cost function of the noisy case is then discussed in more detail. The same two time histories studied in the previous section are examined here. The noisy case is more interesting because it has a meaningful Cramér-Rao bound and is more representative of aircraft flight data.

First we will look at the one-dimensional case where \( L_6 \) is fixed at the correct value, because it is easier to grasp some of the characteristics of the cost function in one dimension. Then we will look at the two-dimensional case, where both \( L_p \) and \( L_6 \) are varying. It is important to remember that everything shown in this paper on cost functions is based on computed time histories that are defined by Eq. (36). For every time history we might choose (computed or flight data), a complete cost function is defined. For the case of \( n \) variables, the cost function defines a hypersurface of \( n + 1 \) dimensions. It might occur to us that we could just construct this surface and look for the minimum, avoiding the need to bother with the minimization algorithm. This is not a reasonable
approach because, in general, the number of variables is greater than two. Therefore, the cost function can be described mathematically but not pictured graphically.

One-Dimensional Case

To illustrate the many interesting aspects of cost functions, it is easiest to first look at cost functions having one variable. In an earlier section, the cost function of \( L_p \) and \( L_0 \) was minimized. That cost function is most interesting in the \( L_p \) direction. Therefore, the one-variable cost function studied here is \( J(L_p) \). All subsequent discussions are for \( J(L_p) \) with \( L_0 \) equal to the correct value of 10. Figure 8 shows the cost function plotted as a function of \( L_p \) for the case where there is no measurement noise (no-noise case). As expected for this case, the minimum cost is zero and occurs at the correct value of \( L_p = -0.2500 \). It is apparent that the cost increases much more slowly for a more negative \( L_p \) than for a positive \( L_p \). In fact, the slope of the curve tends to become less negative where \( L_p \) is more negative than -1.0. Physically this makes sense since the more negative values of \( L_p \) represent cases of high damping, and the positive \( L_p \) represents an unstable system. Therefore, the \( p_i \) for positive \( L_p \) becomes increasingly different from the measured time history for small positive increments in \( L_p \). For very large damping (very negative \( L_p \) the system would show essentially no response. Therefore, large increases in damping result in relatively small changes in the value of \( J(L_p) \).

In Fig. 9, the cost function based on the time history with measurement noise (noisy case) is plotted as a function of \( L_p \). The correct value of \( L_p (-0.2500) \) and the value of \( L_p (-0.3218) \) at the minimum of the cost (3.335) are both indicated on the figure. The general shape of the cost function in Fig. 9 is similar to that shown in Fig. 8. Figure 10 shows the comparison between the cost functions based on the time histories with and without measurement noise. The comments relating to the cost function of the no-noise case also apply to the cost function based on the noisy case. Figure 10 shows clearly that the two cost functions are shifted by the difference \( \lambda \) in the value of \( L_p \) at the minimum and increased by the difference in the minimum cost. One would expect only a small difference in the value of the cost when far from the minimum. This is because the "estimated" time history is so far from the measured time history that it becomes irrelevant as to whether the measured time history has noise added. Therefore, for large values of cost, the difference in the two cost functions should be small in comparison to the total cost.

Figure 11 shows the gradient of \( J(L_p) \) plotted as a function of \( L_p \) for the noisy case. This is the function for which we were trying to find the zero (or equivalently, the minimum of the cost function) using the Gauss-Newton method of a previous section. The gradient is zero at \( L_p = -0.3218 \), which corresponds to the value of the minimum of \( J(L_p) \).

The difference between the Newton-Raphson method (Eq. (14a)) and the Gauss-Newton method (Eq. (14b)) of minimization has been mentioned previously.
For this simple one-dimensional case, we can easily compute the second gradient both with the second term of Eq. (14a) (Newton-Raphson), and without the second term (Gauss-Newton, Eq. (14b)). Figure 12 shows a comparison between the Newton-Raphson and the Gauss-Newton approximation second gradients. The Gauss-Newton second gradient (dashed line) always remains positive because it is the sum of quadratic terms (squared for the one-dimensional example). The Newton-Raphson second gradient can be positive or negative, depending upon the value of the second partial with respect to $L_p$. Other than the difference in sign for the more negative $L_p$, the two curves have similar shapes.

As stated earlier, the Gauss-Newton method can be shown to be superior to Newton-Raphson in certain cases. We can demonstrate obvious cases of this with our example. An easy way to select a spot where problems with the Newton-Raphson method will occur is to look for places where the second gradient (slope of the gradient) is near zero or negative. Figure 11 has such a region near $L_p = -1.0$. If we choose a point where the gradient slope is exactly zero, we are forced to divide by zero in Eq. (12) with the Newton-Raphson method. This point is at $L_p = -1.13$ in Fig. 12. If the value of the slope of the gradient is negative, then the Newton-Raphson method will go to very negative values of $L_p$.

For very negative values of $L_p$, the cost becomes asymptotically constant and the gradient becomes nearly zero. In that region, the Newton-Raphson algorithm would diverge towards negative infinity. If the slope of the gradient is positive but small, we still have a problem with the Newton-Raphson method. Figure 13 shows the first iteration starting from $L_p = -0.95$ for both Gauss-Newton and Newton-Raphson. The Newton-Raphson method selects a point where the tangent of the gradient at $L_p = -0.95$ intersects the zero line. This results in the selection of an $L_p$ of approximately 2.6 in the first iteration. From that value it requires many iterations to return to the actual minimum. On the other hand, the Gauss-Newton method selects a value for $L_p$ of approximately -0.09 and converges to the minimum to four-digit accuracy in two more iterations. With more complex examples a comparison of the convergence properties of the two algorithms becomes more difficult to visualize, but the problems are generalizations of the situation we have observed with the one-dimensional example.

The usefulness of the Cramér-Rao bound was discussed in the Example With Measurement Noise section. At this point it is useful to digress briefly to discuss some of the ramifications of the Cramér-Rao bound for the one-dimensional case. The Cramér-Rao bound only has meaning for the noisy case. In the noisy example, the estimate of $L_p$ is -0.3218 and the Cramér-Rao bound is 0.0579. The calculation of the Cramér-Rao bound was defined in the previous section for both one-dimensional and two-dimensional examples. The Cramér-Rao bound is an estimate of the standard deviation of the estimate. One would expect the scatter in the estimates of $L_p$ to be of about the same magnitude as the estimate of the standard deviation. For the one-dimensional case discussed here, the range ($L_p \pm 0.3218$) plus or minus the Cramér-Rao bound (0.0579) nearly includes the correct value of $L_p$ (-0.2500). If noisy cases are generated for many time histories (adding different measurement noise to each time history), then the sample mean and sample standard deviation of the estimates for these cases can be calculated. Table 5 gives the sample mean, sample standard deviation, and the
standard deviation of the sample mean (standard deviation divided by the square root of the number of cases) for 5, 10, and 20 cases. The sample mean, as expected, gets closer to the correct value of -0.2500 as the number of cases increases. This is also reflected by the decreasing values in column 4 of Table 5, which are estimates of the error in the sample mean. Column 3 of Table 5 shows the sample standard deviations, which indicate the approximate accuracy of the individual estimates. This standard deviation, which stays more or less constant, is approximately equal to the Cramér-Rao bound for the noisy case being studied here. In fact, the Cramér-Rao bounds for each of the 20 noisy cases used here (not shown in the table) do not change much from the values found for the noisy case being studied. Both of these results are in good agreement with the theoretical characteristics (Ref. 16) of the Cramér-Rao bounds and maximum likelihood estimators in general.

The examples shown here indicate the value of obtaining more sample time histories (maneuvers). More samples improve confidence in the estimate of the unknowns. The same result holds true in analyzing actual flight time histories (maneuvers); thus it is always advisable to obtain several maneuvers at a given flight condition to improve the best estimate of each derivative.

The size of the Cramér-Rao bounds and of the error between the correct value and the estimated value of \( L_p \) is determined to a large extent by the length of the time history and the amount of noise added to the correct time history. For the example being studied here, it is apparent from Fig. 5 that the amount of noise being added to the time history is large. The effect of the power of the measurement noise (\( G^* \), Eqs. (3) and (4)) on the estimate of \( L_p \) (that is, \( \hat{L}_p \)) for the time history is given in Table 6. The estimate of \( L_p \) is much improved by decreasing the measurement noise power. A reduction in the value of \( G \) to one-tenth of the value in the noisy example being studied yields an acceptable estimate of \( L_p \). For light data, the measurement noise is reduced by improving the accuracy of the output of the measurement sensors.

Two-Dimensional Case

In this section the cost function (which is dependent on both \( L_p \) and \( L_\delta \)) is studied. The no-noise case is examined first, followed by the noisy case.

No-noise case. Even though the cost function is a function of only two unknowns, it is much more difficult to visualize than the one-unknown case. The cost function over a reasonable range of \( L_p \) and \( L_\delta \) is shown in Fig. 14. The cost increases very rapidly in the region of positive \( L_p \) and large values of \( L_\delta \). The reason is just an extension of the argument for positive \( L_p \) given in the previous section. The shape of the surface can be depicted in greater detail if we examine only the values of the cost function less than 200 for \( L_p \) less than 1.0. Figure 15 shows a view of this restricted surface from the upper end of the surface. The minimum must lie in the curving valley that gets broader as we go to the far side of the surface. Now that we have a picture of the surface, we can look at the isoclines of constant cost on the \( L_p \)-versus-\( L_\delta \) plane. These isoclines are shown in Fig. 16. The minimum of the cost function
is inside the closed isocline. The steepness of the cost function in the positive-$L_p$ direction is once again apparent. Inside the closed isocline the shape is more nearly elliptical, indicating that the cost is nearly quadratic here, so fairly rapid convergence in this region would be expected. The $L_p$ axis becomes an asymptote in cost as $L_6$ approaches zero. The cost is constant for $L_6 = 0$ because no response would result from any aileron input. The estimated response is zero for all values of $L_p$, resulting in constant cost.

Figure 16 shows the region of the minimum value of the cost function, which, as seen in the earlier example (Table 2), occurs at the correct values for $L_p$ and $L_6$ of -0.2500 and 10, respectively. This is also evident by looking at the cost function surface shown in Fig. 17. The surface has its minimum at the correct value. As expected, the value of the cost function at the minimum is zero.

Noisy case. As shown before in the one-dimensional case, the primary difference between the cost functions for the no-noise and noisy cases was a shift in the cost function. In that instance, the noisy case was shifted so that the minimum was at a higher cost and a more negative value of $L_p$. In the two-dimensional case, the no-noise and noisy cost functions exhibit a similar shift. For two dimensions the shift is in both the $L_p$ and $L_6$ directions. The shift is small enough that the difference between the two cost functions is not visible at the scale shown in Fig. 14 or from the perspective of Fig. 15. Figure 18 shows the isoclines of constant cost for the noisy case. The figure looks much like the isoclines for the no-noise case shown in Fig. 16. The difference between Figs. 16 and 18 is a shift in $L_p$ of about 0.1. This is the difference in the value of $L_p$ at the minimum for the no-noise and noisy cases. Heuristically, one can see that the same would be true for cases with more than two unknowns. The primary difference between the two cost functions is near the minimum.

The next logical part of the cost function to examine is near the minimum. Figure 19 shows the same view of the cost function for the noisy case as was shown in Fig. 17 for the no-noise case. The shape is roughly the same as that shown in Fig. 17, but the surface is shifted such that its minimum lies over $L_p = -0.3540$ and $L_6 = 10.24$, and is shifted upward to a cost function value of approximately 3.3.

To get a more precise idea of the cost of the noisy case near the minimum, we once again need to examine the isoclines. The isoclines (Fig. 20) in this region are much more like ellipses than they are in Figs. 16 and 18. We can follow the path of the minimization example used before by including the results from Table 4 on Fig. 20. The first iteration ($L = 1$) brought the values of $L_p$ and $L_6$ very close to the values at the minimum. The next iteration essentially selected the values at the minimum when viewed at this scale. One of the reasons the convergence is so rapid in this region is that the isoclines are nearly elliptical, demonstrating that the cost is very nearly quadratic in this region. If we had started the Gauss-Newton algorithm at a point where the isoclines are much less elliptical (as in some of the border regions in Fig. 18), the
convergence would have been much slower initially, but much the same as it entered the nearly quadratic region of the cost function.

Before concluding our examination of the two-dimensional case, we need to examine the Cramér-Rao bound. Figure 21 shows the uncertainty ellipsoid, which is based on the Cramér-Rao bounds defined in an earlier section. The relationships between the Cramér-Rao bound and the uncertainty ellipsoid are discussed in Ref. 16. The uncertainty ellipsoid almost includes the correct value of \( L_p \) and \( L_q \). The Cramér-Rao bound for \( L_p \) and \( L_q \) can be determined from the projection of the uncertainty ellipsoid onto the \( L_p \) and \( L_q \) axes, and compared with the values given earlier, which were 0.1593 and 1.116 for \( L_p \) and \( L_q \), respectively.

**ESTIMATION USING FLIGHT DATA**

In the previous several sections we examined the basic mechanics of obtaining maximum likelihood estimates from computed examples with one or two unknown parameters. Now that we have a grasp of these basics, we can explore the estimation of stability and control derivatives from actual flight data. For the computationally much more difficult situation usually encountered using actual flight data, we will obtain the maximum likelihood estimates with the Iliff-Maine code (MLE3 program) described in Ref. 17. The equations of motion that are of interest are given in the ATCRRAFT EQUATIONS OF MOTION section of this paper; the remainder of the equations are given in Ref. 17.

In general, flight data estimation is fairly complex, and codes such as the Iliff-Maine code must usually be used to assist in the analysis. However, one must still be cautious about accepting the results; that is, the estimates must fit the phenomenology, and the match between the measured and computed time histories must be acceptable. This is true in all flight regimes, but one must be particularly careful in potential problem situations such as (1) in separated flow at high Mach numbers or high angle of attack, (2) with unusual aircraft configurations such as the oblique wing (Ref. 18), or (3) with modern high-performance aircraft with high-gain feedback loops. In any of the above cases, one should be particularly careful where there are even small anomalies in the match. These anomalies may indicate ignored terms in the equations of motion, separated flow, nonlinearities, sensor problems, insufficient resolution (Ref. 1), sensor location (Ref. 1), time or phase lags (Refs. 1 and 9), or any of a long list of other problems.

The following brief examples are intended to show how the above caveats and the computed examples of previous sections can be used to assist in the analysis. In the computed example, the desirability of low-noise sensors, an adequate model, and several maneuvers at a given flight condition is shown.

**Hand Calculation Example**

Sometimes evaluation of a fairly complex flight maneuver can be augmented with a simple hand calculation. One example of this can be found for the space
shuttle. The space shuttle is a large double-delta-winged vehicle designed to
enter the atmosphere from space and land horizontally. The entry control system
consists of 12 vertical reaction-control-system (RCS) jets (six up-firing and
six down-firing), 8 horizontal RCS jets (four left-firing and four right-firing)
4 elevon surfaces, a body flap, and a split rudder surface. The locations of
these devices are shown in Fig. 22. The vertical jets and the elevons are used
for both pitch and roll control. The jets and elevons are used symmetrically
for pitch control and asymmetrically for roll control. The space shuttle con-
trol system is described briefly in Ref. 6.

The shuttle example used here is from a maneuver obtained at a Mach number
of approximately 21 and an angle of attack of approximately 40°. The controls
being used for this lateral-directional maneuver are the differential elevons
and the side-firing jets (yaw jets). The maneuver is shown in Fig. 23.
Equations (15) to (31) describe the equations of motion. A simplified approach
can be used to determine some of the derivatives by hand. The approach is one
that has been used since the beginning of dynamic analysis of flight maneuvers.
In particular, for this maneuver the slope of the rates can be used to determine
the yaw jet control derivatives. This is possible for this example, even with a
high-gain feedback system, because the yaw jets are essentially step functions,
and the slope of the rates p and r can be determined before the vehicle and the
differential elevon (aileron) responses become significant. The rolling moment
due to yaw jet ($L_{YJ}$) is particularly important for the shuttle (Ref. 6 discusses
the essential nature of flight-determined $L_{YJ}$ in the redefinition of entry
maneuvers) and is, in general, more difficult to obtain than the more dominant
yawing moment due to yaw jet. Therefore, as an illustrative example, $L_{YJ}$ is
determined by hand. Figure 24 shows yaw jet activity and smoothed roll rate
plotted at expanded scales. The equation for $L_{YJ}$ is given by

$$L_{YJ} = \frac{\dot{p}I_X}{\text{(Number of yaw jets)}}$$

(42)

$$\dot{p} = \frac{\Delta p}{\Delta t} = \frac{0.07}{57.3} \pm (0.1)$$

(43)

Therefore, given that $I_X \approx 900,000$ slug-ft$^2$, and the number of yaw jets is 4,
$L_{YJ} \approx 2750$ ft-lb.

The same maneuver was analyzed with WMLE3, and the resulting match is shown
in Fig. 25. The match is very good except for a small mismatch in p at about
6 sec. This small mismatch was studied separately with WMLE3 and found to be
caused by a nonlinearity in the aileron derivative. The value from WMLE3 for
$L_{YJ}$ is 2690 ft-lb, which for the accuracy used here is essentially the same
value as obtained by the simplified method. The aileron derivatives would be
difficult to determine as accurately as the yaw jet derivatives. Although good
estimates can seldom be obtained with the slope method discussed here, rough
estimates can usually be obtained to gain some insight into values obtained with
WMLE3 (or any other maximum likelihood program). These rough estimates can then
be used to help explain unexpected values of estimates from an estimation
program.

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Sometimes a flight example becomes too complex to allow anything other than qualitative estimates to be determined by hand. The example shown in Fig. 26 is the determination of the rudder derivative for the F-8 aircraft with the yaw augmentation system on. This example, taken from Ref. 20, includes an aileron pulse and a rudder pulse. Although an independent pilot rudder pulse is input during the maneuver, the rudder is largely responding to the lateral acceleration feedback. When the rudder is moving, several other variables are also moving, thus making it difficult to use the simplified approach just discussed. However, \( \text{C}_{\text{n}R} \) can be roughly determined when the rudder moves, approximately 1.7 sec from the start of the maneuver. Most of the slope of yaw rate is probably caused by the rudder, but a poor estimate would be obtained using the hand calculation.

**Cost Function for Full Aircraft Problem**

The analysis of a lateral-directional maneuver obtained in flight typically has from 15 to 25 unknown parameters (as shown in Eqs. (15) and (31)), in contrast to the one or two in the simple aircraft example. This makes detailed examples unwieldy and any graphic presentation of the cost function impossible. Therefore, in this section we are primarily examining the estimation procedure and the process of the minimization.

For our flight example, we have chosen a lateral-directional maneuver, with both aileron and rudder inputs, that has 17 unknown parameters. The data are from the oblique wing aircraft (Ref. 18) with the wing unskewed during the maneuver. This example was chosen because it is a typical maneuver. The time history of the data and the subsequent output of MMLE3 have been published in Ref. 21. Some results of the analysis are shown in Table 7. The match between the measured time history (solid lines) and the estimated (calculated) time history (dashed lines) is shown as a function of iteration in Fig. 27. Figures 27(a) to (e) are for iterations 0 to 4, respectively. Table 7 shows that the cost remains unchanged after four iterations. A similar result was obtained for the two-dimensional simple aircraft example in Fig. 6 and Table 4.

Of the many things the analyst must consider in obtaining estimates, the two most important ones are how good is the match and how good is the convergence. A satisfactory match and monotonous convergence are necessary, but not sufficient, conditions for a successful analysis. Figure 27(e), although not perfect, is a very good match. The convergence can best be evaluated by looking at the normalized cost in the last row of Table 7. The cost has converged rapidly and monotonically in four iterations, and it remains at the converged cost. These factors are convincing evidence that the convergence is complete. Therefore, the criteria of match and convergence are satisfied in our example. In some cases we might encounter cost that does not converge rapidly (in four to six iterations) or monotonically, or stay "exactly" at the minimum value. These situations usually indicate at least a small problem in the analysis. These problems, if found, are usually traced to an instrumentation or data acquisition problem, an inadequate mathematical model, or a maneuver that contains a marginal amount of information.
Table 7 also shows that the startup values of all the coefficients are zero for the control and bias variables. Wind tunnel estimates could have been used for starting values, but the convergence of the algorithm is not very dependent on the startup values. As part of the startup algorithm, the MMLE3 program normally holds the derivatives of the state variables constant until after the first iteration, as is evident in Table 7.

Figure 27(a) shows the match between the measured and computed data for the startup values. The match is very poor because the startup values for the control derivatives are all zero, so the only motion is in response to the initial conditions. The control derivatives and bias are determined on the first iteration, resulting in the much improved match shown in Fig. 27(b). The match after two iterations, shown in Fig. 27(c), is improved as the program further modifies the control derivatives and, for the first time, adjusts the derivatives affecting the natural frequency (CnB and CnB). By the third iteration (Fig. 27(d)), the improvement in the match is almost complete, because minor adjustments to the frequency are made and the damping derivatives are changed. Fig. 27(e) shows the match when all but the most minor derivatives have ceased to change.

Several general observations can be made based on this well behaved example. The strong or most important coefficients have essentially converged in three iterations. The same effect was seen in the simple example—that is, Lk converged faster than Lk (Table 4). Some of the less important or second-order coefficients have only converged to two places after three iterations and are still changing by one digit in the fourth place at the end of six iterations. Another observation is that for some coefficients (Ct, Cn, and Cn) even though the sign is wrong after the first iteration, the algorithm quickly selects their correct values once the important derivatives have stabilized.

In general, if the analysis of a maneuver has gone well, we do not need to spend much time inspecting a table analogous to Table 7. However, if there have been problems in convergence or in the quality of the fit, a detailed inspection of such a table may be necessary. The data may show an important coefficient going unstable at an early iteration, which could cause problems later. If the starting values are grossly in error, the algorithm is driven a long way from reasonable values and then for many reasons does not behave well. Occasionally the algorithm alternately selects from two diverse sets of values of two or more coefficients on successive iterations, behaving as if the shape of the cost function were a narrow multidimensional valley analogous to but more extreme than the two-dimensional valley shown in Figs. 18 and 20.

Cramér-Rao Bounds

The earlier sections regarding the computed example have shown that the Cramér-Rao bound is a good indicator of the accuracy of an estimated parameter. The Cramér-Rao bounds can be used in a similar, but somewhat more qualitative, fashion on flight data. The Cramér-Rao bounds that are included in MMLE3 (as well as many other maximum likelihood estimation programs) have been useful in
determining whether estimates are good or bad. The aircraft example discussed here has been reported previously (for example, in Refs. 1 and 16). However, this example of the use of the Cramér-Rao bound in the assessment of flight-derived estimates is pertinent to the thrust of this paper. Figure 28 shows estimates of $C_{np}$ as a function of angle of attack for the PA-30 twin-engine general aviation aircraft (Ref. 22) at three flap settings. There is a significant amount of scatter, which makes the reliability of the information on $C_{np}$ questionable. The data shown are the estimates from the MMLE3 program, which also provides the Cramér-Rao bounds for each estimate. Past experience (Ref. 1) has shown that if the Cramér-Rao bound is multiplied by a scale factor (the result sometimes being called the uncertainty level (Refs. 1 and 16)) and plotted as a vertical bar with the associated estimate, it helps in the interpretation of flight-determined results. Figure 29 shows the same data as Fig. 28, with the uncertainty levels now included as vertical bars. The estimates with small uncertainty levels (Cramér-Rao bounds) are the best estimates, as was discussed earlier in the section on Cramér-Rao bounds for the one-dimensional case. The fairing shown in Fig. 29 goes through the estimates with small Cramér-Rao bounds and ignores the estimates with large bounds. One can have great confidence in the fairing of the estimates, because the fairing is well defined and consistent when the Cramér-Rao bound information is included. In this particular instance, the estimates with small bounds were from maneuvers where the aileron forced the motion, and the large bounds were from maneuvers where the rudder forced the motion. Therefore, in addition to aiding in the fairing of the estimates, the Cramér-Rao bounds help show that the aileron-forced maneuvers are superior for estimating $C_{np}$ for the PA-30 aircraft.

This example illustrates that the Cramér-Rao bounds are a useful tool in assessing flight-determined estimates, just as they were found useful for the simple aircraft example with computed data.

**Atmospheric Turbulence (State Noise)**

Atmospheric turbulence (state noise) cannot always be avoided in flight; therefore, it is desirable to be able to obtain stability and control derivatives in the presence of turbulence. In addition, an estimate of the turbulence time history can be of interest, particularly in the implementation of turbulence suppression systems.

Many years ago it was demonstrated that the stability and control derivatives can be adequately determined with maximum likelihood estimation techniques for maneuvers performed in smooth air. If these techniques, which do not account for turbulence, are applied to data obtained in turbulence, not only are the resulting matches of the time histories unsatisfactory but the estimated coefficients are unacceptable (Refs. 23 to 25). The technique described in Refs. 14, 23, and 25 can account for the effect of turbulence. With this technique, maximum likelihood estimates of the stability and control derivatives as well as estimates of the turbulence time histories are obtained by minimizing the cost function given by Eq. (11). Results of the application of the technique to longitudinal maneuvers obtained in turbulence have been reported previously (Refs. 23 to 25).
The lateral-directional equations (Eqs. (15), (16), (17), (18), and (29)) can be modified in a manner similar to that used to modify the longitudinal equations in Refs. 23 to 25. The turbulence (state noise) model is the Dryden expression, which is described in Ref. 26. The Iliff-Maine code (Ref. 17) can be used to obtain the maximum likelihood estimates where state noise is present.

Thirty-eight seconds of data from the PA-30 aircraft flying in turbulence was analyzed at 50 samples/sec. The best match that could be obtained with the maximum likelihood estimation method that does not account for turbulence is shown in Fig. 30. The match is unacceptable and resulted in poor estimates of the stability and control derivatives. Figure 31 shows the match obtained with the maximum likelihood estimation technique that accounts for turbulence (Re's 14 and 17). The match is excellent and the maneuver provided acceptable estimated stability and control derivatives. It is also of interest to compare the power spectra of the estimated turbulence time histories. The power spectrum of the turbulence component affecting angle of sideslip, $\dot{\theta}_y$, is shown in Fig. 32. Figure 33 presents the power spectrum of the turbulence component affecting roll rate, $\dot{\phi}_g$. The slopes of the asymptotes shown in Figs. 32 and 33 are those defined by the Dryden expression given in Ref. 26. Good agreement is shown between the power spectra and the asymptotes for $\dot{\theta}_y$ and $\dot{\phi}_g$.

The algorithm used here is based on a linearized system described by Eqs. (5) to (7) and solved by minimizing the cost function given by Eq. (11). The system need not resemble that for the aircraft stability and control problem other than in the requirement for linearity. Therefore, many formulations for the structural problem are written in the form of Eqs. (5) to (7), and the algorithm under discussion can be directly applied with these formulations.

ESTIMATION FOR SIMPLE STRUCTURAL PROBLEM

The problem of the flexible space structure is most fully characterized as a distributed parameter system with its associated distributed system control laws. The model will vary depending upon changes in its configuration or its environment, such as solar heating. As in most cases, the preferred solution is the simplest successful approach. The lumped system approach is much simpler and computationally far more efficient than the fully distributed parameter system approach. For example, structural mode control based on current state-of-the-art approaches has proved very successful. Admittedly, the aircraft structure is heavier than most spacecraft, but many aircraft structures are highly complex, consisting of many substructures within the main structure. To the novice, many of the space structures currently being investigated appear simpler than modern, large aircraft. If the lumped parameter system approach used for the aircraft problem is found to be inadequate, it seems likely that distributed parameter estimation codes will evolve to whatever complexity is necessary to solve the flexible space structure problem.
This paper has discussed some of the experience gained from the application of aircraft stability and control analysis to flight data. The codes used for this analysis are for lumped parameter systems in the time domain. The codes have been used successfully for structural problems and are fully adaptable to the frequency domain if that is found to be preferable.

Although few results have been obtained for time-domain structural analysis at the Ames Dryden Flight Research Facility, some superficial experience in structural time-domain analysis has been obtained. The following two examples show how the techniques being used for stability and control analysis can be applied to simple structural problems. The preceding section discussed the incorporation of state noise in the model. The following examples do not include the use of state noise, but state noise, if warranted, could easily be incorporated in the types of examples to be discussed.

Estimation of Structural Characteristics

All aircraft have observable structural modes. These modes usually cause no difficulty in estimating stability and control derivatives because the structural frequencies are higher than the aerodynamic frequencies. In general, if the structural frequencies are higher than the highest aerodynamic frequency by more than a factor of 5 to 10, they can be neglected unless their amplitude is so large as to mask measurements desired for the aerodynamic analysis. However, if one or more structural modes are affecting the aerodynamic modes, as may occur in large aircraft, these structural modes must be included in the mathematical model being analyzed.

Even though no completely satisfactory practical results are available that account for structural modes and their interactions with the aerodynamics, it is interesting to assess the time-domain maximum likelihood analysis of the structural modes independent of any interaction. This can be done where a structural mode is observed and no significant coupling is apparent.

Figure 34 shows a structural mode on the lateral acceleration of an aircraft where little effect was observed for structural-aerodynamic coupling. The frequency of the mode is high enough that the mode does not interact with the aerodynamic modes. Therefore, the stability and control derivatives were obtained separately and held constant for the succeeding analysis. The analysis consisted of using the maximum likelihood estimation program MLE 3 (Ref. 17) with a sixth-order model that included the lateral-directional aerodynamic modes plus one structural mode. The dynamic pressure and the velocity were allowed to vary in the analysis. The structural mode frequency and damping were estimated as linear functions of dynamic pressure. The initial conditions were also estimated. A structural mode frequency of 7.94 Hz was chosen to start the estimation process. The comparison between the original data and the match obtained with the maximum likelihood estimation method is shown in Fig. 35. The two time histories are in good agreement at the beginning of the maneuver and at the end of the maneuver, but they are 180° out of phase at a time of approximately
0.3 sec. The match shown in Fig. 35 suggests that the maximum likelihood estimator has reached a local minimum but not the global minimum. Multiple minima are not normally a problem when obtaining the stability and control derivatives of aircraft with the maximum likelihood estimation method.

The reason for the multiple minima is demonstrated by the following simple scalar example. Let the noiseless measured response be \( z(t) = \sin (w_0 t) \) and the estimated response be \( \tilde{z}_t = \sin (\omega t) \), where \( \omega \) is the only unknown coefficient. Then, by Eq. (4), the cost function becomes

\[
J(\omega, T) = \int_0^T [\sin (w_0 t) - \sin (\omega t)]^2 \, dt
\]

\[
= t - \frac{1}{4w_0} \sin (2w_0 T) - \frac{1}{4\omega} \sin (2\omega T) - \frac{2\omega}{\omega^2 - w_0^2} \left( \frac{\omega}{\omega} \sin (\omega T) \cos (w_0 T) - \cos (\omega T) \sin (w_0 T) \right)
\]

If \( T \) is chosen to represent 10 cycles, as shown in Fig. 35, then for an \( \omega_0 \) of 1 rad/sec, \( T \) equals 20\( \pi \). In Fig. 36, the cost function \( J(\omega, 20\pi) \) is shown as a function of \( \omega \). The global minimum is at an \( \omega \) of 1 rad/sec, as it should be, but there are many local minima at increments of approximately 0.05 rad/sec. If a value of less than 0.97 or greater than 1.03 were chosen for a starting estimate of \( \omega \), the algorithm would converge to a local minimum. If a value of between 0.98 and 1.02 were chosen, it would converge to the global minimum. Therefore, for this example where 10 cycles were observed, the starting value of \( \omega \) must be less than 3 percent from the correct answer to converge to the global minimum.

Figure 37 shows a sine wave for the global minimum along with a sine wave with a frequency that varies 10 percent from the global minimum. The sine waves are in phase at the beginning and end, and 180° out of phase in the middle. These data appear similar to those shown for flight data in Fig. 35. If only one or two cycles were used for the analysis, the problem illustrated in Fig. 37 would be minimized. This is apparent in Fig. 38 where only the first cycle of Fig. 37 is shown.

If \( T \) is chosen to represent only one cycle and \( \omega_0 \) remains equal to 1 rad/sec (as in Fig. 38), then \( T \) equals 2\( \pi \). The cost function \( J(\omega, 2\pi) \) is presented as a function of \( \omega \) in Fig. 39. The global minimum is correctly at an \( \omega \) of 1 rad/sec, but now the algorithm converges to the global minimum if \( \omega \) is started within approximately 25 percent of the correct value.

Knowing the sensitivity of the algorithm when a record with many lightly damped cycles is being analyzed, the data of Fig. 34 can be reanalyzed starting closer to the observed frequency. Starting the maximum likelihood estimation method with an \( \omega \) of 3.0 results in the fit shown in Fig. 40. This is an acceptable fit of the data.
Based on the preceding results, if data are to be analyzed where many cycles of a structural mode are present, the structural mode frequency, \( \omega \), must be closely approximated before starting the estimation process.

**Structural Modes in Space**

In the process of analyzing aircraft flight data, the authors have frequently observed results that clearly exhibit unmodeled dynamics. The unmodeled dynamics could be caused by many phenomena, such as higher-order aerodynamic modes or structural modes. These modes can usually be ignored and left unmodeled because they have no effect on the results of primary interest in the analysis. If the unmodeled modes cannot be ignored, then the system equations must be revised to include the unmodeled modes.

The authors have not yet found it necessary to model structural modes for data obtained in space in the process of obtaining control derivatives for the space shuttle. However, the structural modes have been observed. Figure 41 shows the response of the space shuttle to the firing of a roll jet and a yaw jet at an altitude of 430,000 ft. The space shuttle configuration and the location of the RCS jets are shown in Fig. 22. The changes in the rigid-body rates and lateral acceleration caused by the jet firings are apparent in Fig. 41. The structural modes are also excited by the jets, as evidenced by the increased ringing in each signal at the time of the jet firings. The roll jet firing has little effect on the rigid-body response for the yaw rate and lateral acceleration; however, the yaw jet results in a rigid-body response for all the signals chosen. This maneuver was analyzed to obtain control derivatives for the rigid-body response described by Eqs. (15) to (31). The resulting match between the measured and computed response is shown in Fig. 42. The estimated control derivatives are in good agreement with those obtained from the maneuvers. The unmodeled structural dynamic modes are evident, but it is apparent that the modes will have little effect on the rigid-body control derivatives. The differences between the measured and computed rigid-body responses (the residuals) for the time close to when the jets were fired are shown in Fig. 43. The data shown here are for a sample interval of 0.006 sec. Some persistent structural ringing is shown for the two rates and the lateral acceleration. However, when a jet is fired, the increased structural response is evident. The structural coefficients can be extracted directly from the residual as they were for the example in the previous section. It appears that there may be some contamination caused by the rigid-body response at the instant the jets fire. If so, this contamination can be eliminated in one of two ways: either analyze the portion of the maneuver a tenth of a second after the jet fires, or adapt the equations of motion to include the structural dynamics in addition to the rigid dynamics. The structural dynamics depicted in Fig. 43 have not been analyzed, but the procedure is straightforward. The procedure used on this case was the same as that used on the example in the preceding section. It is apparent, however, that more than one structural mode would need to be included in the model.

All the analysis techniques discussed in this paper apply to the analysis of this space shuttle example. If state noise is included in the mathematical model, then the linear form of Eqs. (5) to (7) would be required. In general,
if the structural partial differential equation can be expressed in the linear form of Eqs. (5) to (7) (with or without state noise), the structural modes can be analyzed readily, with the NLME 3 program (Ref. 17) in the time domain. If the analyst prefers, the problem can be expressed in the linear constant coefficient form and analyzed in the frequency domain, as described in Ref. 12. The relative advantages and disadvantages of time-domain analysis as compared with frequency-domain analysis are also discussed in that reference. If the equations are nonlinear, but in the form of Eqs. (1) to (3), then maximum likelihood estimates can be obtained in the time domain.

CONCLUDING REMARKS

The computed simple aircraft example showed the basics of minimization and the general concepts of cost functions themselves. In addition, the example demonstrated the advantage of low measurement noise, multiple estimates at a given condition, and the Cramèr-Rao bounds, and the quality of the match between the measured and computed data. The flight data showed that many of these concepts still hold true even though the dimensionality of the cost function makes it impossible to plot or visualize. In addition, the techniques used for the aircraft problem were shown to be applicable to the flexible structure problem.

REFERENCES


### Table 1: Values of computed time history with no measurement noise

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### Table 2: Pertinent values as a function of iteration

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### Table 3: Values of computed time history with added measurement noise

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Table 4 Pertinent values as a function of iteration

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<th>( \Phi(L) )</th>
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Table 5 Mean and standard deviations for estimate \( \hat{L}_p \)

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Table 6 Estimate of \( L_p \) and Cramer-Rao bound as a function of the square root of noise power

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**Fig. 1** Maximum likelihood estimation concept.

**Fig. 2** Aircraft axis system.

(a) Measurement and control surface definitions.  (b) Flow angle definitions.
(c) Simplified aircraft nomenclature.

Fig. 2 Concluded.

Fig. 4 Comparison of measured and computed data for each of the first three iterations.

Fig. 3 Time history with no measurement noise.

Fig. 5 Time history with measurement noise.
**Fig. 6** Comparison of measured and computed data for each iteration.

**Fig. 7** Comparison of estimated roll rate from no-noise and noisy cases.

**Fig. 8** Cost function $J(L_p)$ as a function of $L_p$ for no-noise case.

**Fig. 9** Cost function as a function of $L_p$ for noisy case.
Fig. 10 Comparison of the cost functions for the no-noise and noisy cases.

Fig. 11 Gradient of $J(L_p)$ as a function of $L_p$ for noisy case.

Fig. 12 Comparison of Newton-Raphson and Gauss-Newton values of the second gradient for the noisy case.

Fig. 13 Comparison of first iteration step size for the Newton-Raphson and Gauss-Newton algorithms for the noisy case.
Fig. 14 Large-scale view of cost function surface.

Fig. 15 Restricted view of cost function surface.
Fig. 16 Isoclines of constant cost of $L_p$ and $L_d$ for the no-noise case.

Fig. 17 Detailed view of cost function surface for no-noise case.

Fig. 18 Isoclines of constant cost in $L_p$ and $L_d$ for the noisy case.

Fig. 19 Detailed view of cost function surface for noisy case.
Fig. 20 Isoclines of constant cost for region near minimum for noisy case.

Fig. 21 Isoclines and uncertainty ellipsoid of the cost function for the noisy case.
**Fig. 22** Space shuttle configuration.

**Fig. 23** Lateral-directional space shuttle maneuver at a Mach number of 21.

**Fig. 24** Examples of obtaining $I_{yj}$ by simple calculations for the shuttle data from Fig. 23.
Fig. 25 MNLE3 match of maneuver shown in Fig. 23.

Fig. 26 Lateral-directional maneuver from F-8 aircraft with augmentation on.
Fig. 27 Match between measured and computed time histories as a function of iteration.
(c) Two iterations.

Fig. 27 Continued.

(d) Three iterations.
Fig. 28 Variations of $C_{np}$ with angle of attack without uncertainty levels.

(e) Four iterations.

Fig. 27 Concluded.
Fig. 29 Variations of $C_{np}$ with angle of attack with uncertainty levels.

Fig. 30 Match of flight data obtained in turbulence (state noise) and computed data obtained from maximum likelihood estimator that does not account for turbulence.

Fig. 31 Match of flight data obtained in turbulence (state noise) and computed data obtained from a maximum likelihood estimator that accounts for turbulence.
Fig. 32 Power spectral density of $\hat{\phi}_g$ obtained from maneuver shown in Fig. 31.

Fig. 33 Power spectral density of $\hat{\phi}_g$ obtained from maneuver shown in Fig. 31.

Fig. 34 Structural mode oscillation observed on the lateral acceleration.

Fig. 35 Match of measured and computed lateral acceleration obtained when maximum likelihood estimator converged to local minimum.

Fig. 36 Cost functional for 10 cycles of data as function of frequency, showing close proximity of local minima to global minimum.

Fig. 37 Simple scalar example illustrating a local minimum similar to that shown for flight data in Fig. 35.
Fig. 38 Simple scalar example showing only the first cycle.

Fig. 39 Cost function for one cycle of data as function of frequency, showing wide region of convergence for global minimum.

Fig. 40 Acceptable match of measured and computed lateral acceleration.
Fig. 41 Dynamic response of space shuttle to firing of roll and yaw jets at an altitude of 430,000 ft.

Fig. 42 Maximum likelihood match of rigid-body response of the space shuttle.
Fig. 43 Differences between measured and computed rigid-body response (residual) for space shuttle. Altitude = 430,000 ft; dynamic pressure = 0.