

EAS SPECTRUM IN THE PRIMARY ENERGY REGION ABOVE 10^{15} eV BY THE AKENO AND THE YAKUTSK ARRAY DATA

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ABSTRACT

The EAS spectrum on scintillation density ρ_{600} in primary energy region $E \approx 10^{15}-10^{20}$ eV on the Yakutsk array data and recent results of the Akeno is given.

1. Introduction

At present the EAS observations at sea-level take the widest energy range of primaries. The observed EAS spectra on particle number N in a shower at $E \leq 10^{18}$ eV and on particle density ρ_{600} at a distance $R=600$ m from axis at $E_0 > 3 \cdot 10^{17}$ eV are obtained. Either for the recovery of the spectrum on E_0 or for the comparison it is reasonable to obtain these results in a form of "corrected" spectra where effects of the development fluctuations (different for N and ρ_{600}) and N and ρ_{600} measurement dispersions (different for various arrays) are taken into account. To consider the EAS spectrum on the whole it is required also to use in the analysis a common basic unit of measurement of the shower particle number (density) and a common parameter of the shower size. Yet it is reasonable and possible only on the basis of ρ_{600} : there is the experimental estimation of ρ_{600} and E_0 relationship and only in the Akeno array data there is the possibility of transition from N to ρ_{600} .

2. Results

a) Yakutsk. In the central part of the array [4] the registration of showers was triggered by a small master (SM) and on the whole array - by a big master (BM). For the analysis the shower events were selected with an axis within fixed receiving areas (different for various ranges of ρ_{600}) and for those periods of the array operation T , when $\sim 100\%$ - efficiency of registration and levels of $\delta\rho_{600}$ summary relative deviations of fluctuations of the shower development and their measurement dispersions obtained from a total measurement simulation [5] and accepted for the analysis [1,2] were provided. Each shower was individually treated as follows: 1) from approximation of measured particle densities

by $\rho(R)_i \propto R^{-n_i}$ [5] n_i and $\rho_{600,i}$ were determined; 2) $\rho_{600,i}$ was reduced to the zenith angle $\theta = 0^\circ$, atmospheric temperature 240 K and pressure 1006 mb (ρ_{600} and E_0 relationship at the atmospheric depth $X=1025 \text{ g.cm}^{-2}$ at these parameters was found) using the absorption length measured in the experiment $\lambda(\rho_{600}) = (218 \pm 15) + (172 \pm 15) \cdot \sec \theta$, g.cm^{-2} , $\theta < 60^\circ$, a barometric coefficient $\alpha_p = -0,25 \pm 0,03 \%$ per mb and temperature coefficient $\alpha_T(\rho_{600}) = 0,30 \pm 0,11 \%$ per K. For $-0,35 < \lg \rho_{600} < 0,6$ as an intermediate parameter of shower size the $\rho_{300,i}$ having the absorption length $\lambda(\rho_{300}) = 251 \pm 21 \text{ g.cm}^{-2}$, $\theta < 40^\circ$ and $\rho_{600} = (0,14 \pm 0,01) \cdot \rho_{300}^{0,89 \pm 0,02}$ were used.

Data used in spectrum construction on the whole have following common characteristics:

	$\lg[\rho_{600}, \text{m}^{-2}]$	$\delta\rho_{600}$	$ST\Omega, \text{m}^2 \cdot \text{s} \cdot \text{sr}$	Number of events
SM	$-0,35 \pm 1$	$0,40 \pm 0,17$	$(0,16 \pm 4,33) \cdot 10^{13}$	534
BM	$1 \pm 1,5$	$0,22 \pm 0,21$	$(1,88 \pm 4,40) \cdot 10^{15}$	109
BM	$> 1,5$	$0,21$	$5,69 \cdot 10^{15}$	79

Introducing into the observed intensities the corrections for the summary effect of the development fluctuations and measurement dispersions with the correction factor [2] $K = 0,98 [1 + \delta\rho_{600}^2]^{-0,5} \alpha(\alpha - 1)$ the differential $f_0(\rho_{600})$ and the integral $F_0(>\rho_{600})$ corrected EAS spectra (see Figure) were obtained.

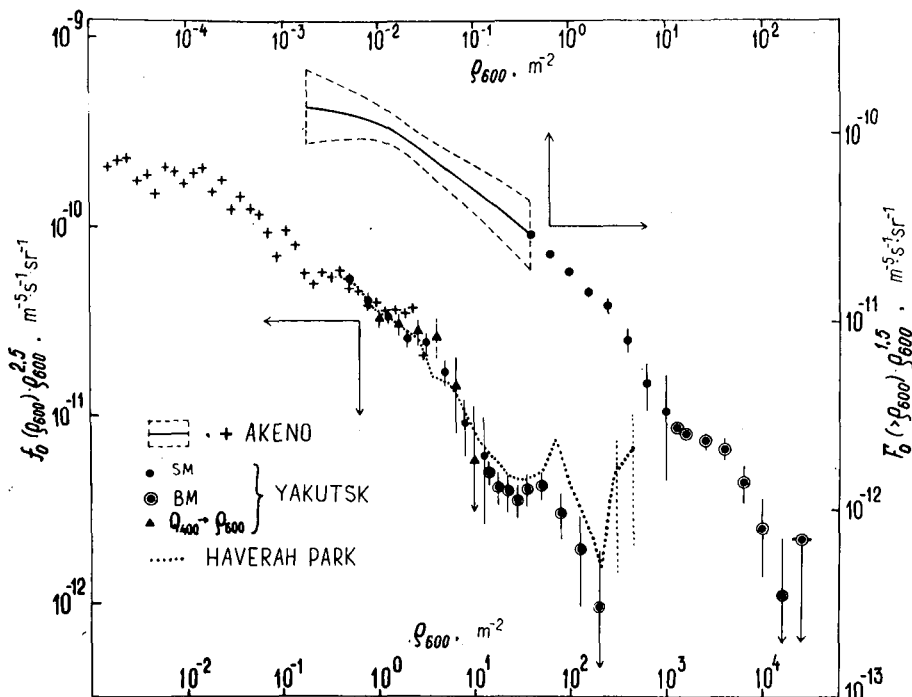
The differential spectrum for $-0,3 < \lg \rho_{600} < 1,7$ displays significant irregularities and at the description by $f_0(\rho_{600}) = A(\rho_{600}/10)^{-\alpha-1}$ has the following parameters:

$\lg \rho_{600}$	$-0,3 \pm 0,5$	$0,5 \pm 1,2$	$1,2 \pm 1,7$	$1,7 \pm 2,3$
$\lg A$	$-13,37 \pm 0,04$	$-13,63 \pm 0,02$	$-13,92 \pm 0,05$	$-13,20 \pm 0,09$
$\alpha+1$	$2,95 \pm 0,04$	$3,58 \pm 0,05$	$2,45 \pm 0,10$	$3,43 \pm 0,11$

The spectrum on ρ_{600} obtained by the relationship $\rho_{600} = (2,05 \pm 0,11) \cdot (Q_{400}/10^7)^{0,99 \pm 0,02}$ from the transformation of the density spectrum of the shower atmospheric Cerenkov light Q_{400} [1] and having the form of the spectrum of loss in atmosphere confirms the change for $-0,3 < \lg \rho_{600} < 1$.

In the Figure a dashed line corresponds to the observed spectrum on Haverah Park data [6] reduced by us to the scintillation density ρ_{600} due to [7]. In this case according

to [8] the effect of $\delta\rho_{600}$ at $\lg \rho_{600} \lesssim 1$ is small ($\lesssim 10\%$ on intensities) and at $\lg \rho_{600} \gtrsim 1$ somewhat increases. Taking into account this fact we find a satisfactory agreement of the results of both arrays. It is remarkable that the Haverah Park spectrum reveals also the steepening tendency for $1.8 \lesssim \lg \rho_{600} \lesssim 2.3$.



b) Akeno. The observed EAS spectrum at sec $\theta = 1.1$ (at the depth 1011 g.cm^{-2}) is given by $f(N_e) dN_e = A(N_e/10^6)^{-\alpha} N_e^{-1}$ with $A = (1.2 \pm 0.2) \cdot 10^{-13} \text{ m}^{-2} \text{ s}^{-1} \text{ sr}^{-1} \text{ part}^{-1}$, $\alpha_N = 1.49 \pm 0.17$ for $5 < \lg N_e < 6$ and $\alpha_N = 1.80 \pm 0.12$ for $6 < \lg N_e < 8$. Some corrections were made: the spectrum is reduced to the Yakutsk level 1025 g.cm^{-2} with absorption length $\lambda(N_e) = 235 \text{ g.cm}^{-2}$; the effect of the shower development fluctuations was taken into account on [9] with average correction factor $\bar{K}_\delta = 0.89 \cdot [1 + \delta N_e^2]^{-0.5 \alpha_N (\alpha_N - 1)} = 0.77$ where the deviations were taken according to [10] to be 0.7 for $5 < \lg N_e < 6$ and 0.44 for $6 < \lg N_e < 8$.

From [11,12] we find $\lg \rho_{600}^* = \lg [\rho_{600,e}^* + \rho_{600,\mu}^*] = 0.961 \lg N_e - 7.46$ at the depth 966 g.cm^{-2} at $T = 279 \text{ K}$.

Recounting ρ_{600} to depth 1025 g.cm^{-2} with $\lambda(\rho_{600}) = 390 \text{ g.cm}^{-2}$ at $T = 240 \text{ K}$ with $\alpha_T = 0.3 \%$ per K and to the Yakutsk basic unit of muon equivalent having the relationship $u_\mu/u_e = 1.15$ with the electron equivalent unit [3] the relationship $\lg \rho_{600} = 0.96 \cdot \lg N_e - 7.534$ is obtained.

In the Figure the differential and integral corrected EAS spectra on ρ_{600} from the Akeno data are given. For $\lg \rho_{600} < 0.12$ we obtain: $f_0(\rho_{600}) d\rho_{600} = A_0 (\rho_{600} / 10^{-1.80})^{-\alpha-1} d\rho_{600}$ with $A_0 = (1.2 \pm 0.2) \cdot 10^{-5.297} \text{ m}^{-2} \cdot \text{s}^{-1} \cdot \text{sr}^{-1} (\text{part./m}^2)^{-1}$, $\alpha = 1.55 \pm 0.18$ for $-2.76 < \lg \rho_{600} < -1.80$ and $\alpha = 1.88 \pm 0.13$ for $\lg \rho_{600} > -1.80$.

3. Conclusion

In the considered 5-decade energy range the EAS spectrum on ρ_{600} reveals significant irregularities. For $\lg \rho_{600} < 1.2$ the steepening (rather consecutive) of inclination of $f(\rho_{600})$ with increase of ρ_{600} occurs: $-\alpha-1 = -2.55 \pm 0.18$; -2.88 ± 0.12 ; -2.95 ± 0.04 and -3.58 ± 0.05 for $\Delta \lg \rho_{600} = -2.76 \pm 1.8$; -1.8 ± 0.3 ; -0.3 ± 0.5 and 0.5 ± 1.2 , respectively. At $\lg \rho_{600} > 1.2$ the irregularity is observed: $-\alpha-1 = -2.45 \pm 0.1$ at $1.2 < \lg \rho_{600} < 1.7$ and $-\alpha-1 \leq -3$ at $1.7 < \lg \rho_{600} < 2.3$. We assume that four shower events with $\lg \rho_{600} > 2.3$ from the spectrum of [6] if to eliminate the effects of methodical character could indicate the possible existence of the other irregularity in the range out of the control of the Yakutsk array.

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