OPTIMAL TRAJECTORIES
FOR THE AEROASSISTED FLIGHT EXPERIMENT,
PART 2, EQUATIONS OF MOTION IN AN INERTIAL SYSTEM

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\textsuperscript{1}This work was supported by NASA-Marshall Space Flight Center, Grant No. NAG-8-755, by Jet Propulsion Laboratory, Contract No. 956415, and by Boeing Military Airplane Company.

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Abstract. This report is the second of a series dealing with the determination of optimal trajectories for the aeroassisted flight experiment (AFE). The AFE refers to the study of the free flight of an autonomous spacecraft, shuttle-launched and shuttle-recovered. Its purpose is to gather atmospheric entry environmental data for use in designing aeroassisted orbital transfer vehicles (AOTV).

It is assumed that: the spacecraft is a particle of constant mass; the Earth is rotating with constant angular velocity; the Earth is an oblate planet, and the gravitational potential depends on both the radial distance and the latitude; however, harmonics of order higher than four are ignored; the atmosphere is at rest with respect to the Earth.

Under the above assumptions, the equations of motion for hypervelocity atmospheric flight (which can be used not only for AFE problems, but also for AOT problems and space shuttle problems) are derived in an inertial system. Transformation relations are supplied which allow one to pass from quantities computed in an inertial system to quantities computed in an Earth-fixed system and viceversa.

Key Words. Flight mechanics, hypervelocity flight, atmospheric flight, coordinate systems, equations of motion, transformation techniques, optimal trajectories, aeroassisted flight experiment, aeroassisted orbital transfer, space shuttle reentry.
I. Introduction

This report is the second of a series dealing with the determination of optimal trajectories for the aeroassisted flight experiment (AFE). The AFE refers to the study of the free flight of an autonomous spacecraft, shuttle-launched and shuttle-recovered. Its purpose is to gather atmospheric entry environmental data for use in designing aeroassisted orbital transfer vehicles (AOTV).

It is assumed that: (a) the spacecraft is a particle of constant mass; (b) the Earth is rotating with constant angular velocity; (c) the atmosphere is at rest with respect to the Earth; (d) the Earth is an oblate planet, and the gravitational potential depends on both the radial distance and the latitude; however, harmonics of order higher than four are ignored.

Under the above assumptions, the equations of motion for hypervelocity atmospheric flight (which can be used not only for AFE problems, but also for AOT problems and space shuttle problems) are derived in an inertial system. Transformation relations are supplied which allow one to pass from quantities computed in an inertial system to quantities computed in an Earth-fixed system (Ref. 1), and viceversa.

Previous Research. Previous research on the topics covered here can be found in Refs. 2-12. For the general theory of flight paths and coordinate systems, see Refs. 2-3; for the equations of flight over a spherical Earth, see Refs. 2-4; for the perturbed motion about an oblate Earth, see Ref. 5; for AFE problems, see Ref. 6; for reentry problems, see Ref. 7; for methods of orbit determination, see Refs. 8-9; for the values of the astrophysical quantities, see Ref. 10; for the values of the characteristic constants of the oblate Earth, see Refs. 11-12.
Outline. Section 2 contains the notations, and Section 3 defines the basic coordinate systems. The relations between coordinate systems are discussed in Section 4, and the angular velocity (or evolutory velocity) is introduced in Section 5. The kinematical equations for an inertial system are derived in Section 6, and the dynamical equations are obtained in Section 7. Section 8 summarizes the results, and Section 9 presents the transformation relations which allow one to pass from quantities computed in an inertial system to quantities computed in an Earth-fixed system, and viceversa.
2. **Notations**

Throughout the paper, the following notations are employed:

- \( \mathbf{a} = \text{acceleration, } \text{m/sec}^2 \);
- \( \mathbf{A} = \text{aerodynamic force, } \text{N} \);
- \( \mathbf{C}_D = \text{drag coefficient} \);
- \( \mathbf{C}_L = \text{lift coefficient} \);
- \( \mathbf{C}_Q = \text{side force coefficient} \);
- \( \mathbf{D} = \text{drag force, } \text{N} \);
- \( \mathbf{f} = \text{latitudinal component of the gravitational acceleration, } \text{m/sec}^2 \);
- \( \mathbf{g} = \text{radial component of the gravitational acceleration, } \text{m/sec}^2 \);
- \( \mathbf{J}_2 = \text{characteristic constant of the Earth's gravitational field} \);
- \( \mathbf{J}_3 = \text{characteristic constant of the Earth's gravitational field} \);
- \( \mathbf{J}_4 = \text{characteristic constant of the Earth's gravitational field} \);
- \( \mathbf{L} = \text{lift force, } \text{N} \);
- \( \mathbf{m} = \text{mass, } \text{kg} \);
- \( \mathbf{M} = \text{Mach number} \);
- \( \mathbf{Q} = \text{side force, } \text{N} \);
- \( \mathbf{r} = \text{radial distance, } \text{m} \);
- \( \mathbf{r}_e = \text{equatorial radius, } \text{m} \);
- \( \mathbf{r}_p = \text{polar radius, } \text{m} \);
- \( \mathbf{R}_e = \text{Reynolds number} \);
- \( \mathbf{S} = \text{reference surface area, } \text{m}^2 \);
- \( \mathbf{T} = \text{thrust force, } \text{N} \);
- \( \mathbf{U} = \text{Earth's gravitational potential, } \text{m}^2/\text{sec}^2 \);
\[ V = \text{velocity, m/sec}; \]
\[ W = \text{gravitational force, N}; \]
\[ x = \text{Cartesian coordinate, m}; \]
\[ y = \text{Cartesian coordinate, m}; \]
\[ z = \text{Cartesian coordinate, m}; \]
\[ \alpha = \text{angle of attack, rad}; \]
\[ \gamma = \text{path inclination, rad}; \]
\[ \theta = \text{longitude, rad}; \]
\[ \mu = \text{bank angle, rad}; \]
\[ \mu_e = \text{Earth's gravitational constant, m}^3/\text{sec}^2; \]
\[ \rho = \text{air density, kg/m}^3; \]
\[ \sigma = \text{sideslip angle, rad}; \]
\[ \phi = \text{latitude, rad}; \]
\[ \chi = \text{heading angle, rad}; \]
\[ \omega = \text{angular velocity of the Earth with respect to an inertial system, rad/sec}; \]
\[ \omega_{hi} = \text{angular velocity of the local horizon system with respect to the inertial system, rad/sec}. \]

**Subscripts**
- \( b \) = body axes system;
- \( e \) = Earth axes system;
- \( h \) = local horizon system;
- \( i \) = inertial axes system;
- \( w \) = wind axes system.

**Superscripts**
- \( \cdot \) = derivative with respect to time;
- \( \rightarrow \) = vector quantity.
3. **Basic Coordinate Systems**

The basic coordinate systems for flight over a spherical Earth are the inertial axes system $Ox_i^iY_i^iZ_i^i$, the Earth axes system $Ox_e^eY_e^eZ_e^e$, the local horizon system $Px_h^hY_h^hZ_h^h$, the wind axes system $Px_w^wY_w^wZ_w^w$, and the body axes system $Px_b^bY_b^bZ_b^b$.

3.1. **Inertial Axes System.** The inertial axes system $Ox_i^iY_i^iZ_i^i$ is a Cartesian reference frame defined as follows. Its origin $O$ is the center of the Earth; the $z_i$-axis is aligned with the axis of rotation of the Earth and is positive northward; the axes $x_i, y_i$ are orthogonal to the $z_i$-axis and are directed radially; the trihedral $Ox_i^iY_i^iZ_i^i$ is right-handed. In particular, the plane $x_i, y_i$ contains the fundamental parallel (the Equator); and the plane $x_i, z_i$ contains the fundamental meridian (the Greenwich meridian) at a particular time instant, the time instant $t = 0$. The symbols $\hat{i}_i, \hat{j}_i, \hat{k}_i$ denote the unit vectors of the inertial axes system; these unit vectors are time invariant by definition.

3.2. **Earth Axes System.** The Earth axes system $Ox_e^eY_e^eZ_e^e$ is a Cartesian reference frame which is rigidly attached to the Earth. Its origin $O$ is the center of the Earth; the $z_e$-axis is aligned with the axis of rotation of the Earth and is positive northward; the axes $x_e, y_e$ are orthogonal to the $z_e$-axis and are directed radially; the trihedral $Ox_e^eY_e^eZ_e^e$ is right-handed. In particular, the plane $x_e, y_e$ contains the fundamental parallel (the Equator); and the plane $x_e, z_e$ contains the fundamental meridian (the Greenwich meridian) at all time instants. The symbols $\hat{i}_e, \hat{j}_e, \hat{k}_e$ denote the unit vectors of the Earth axes system; these unit vectors are time dependent.

3.3. **Local Horizon System.** The local horizon system $Px_h^hY_h^hZ_h^h$ is a Cartesian reference frame defined as follows. Its origin $P$ is identical with
the instantaneous position of the spacecraft; the \(z_h\)-axis is directed radially (that is, vertical) and is positive downward; the axes \(x_h, y_h\) are orthogonal to the \(z_h\)-axis (therefore, they are tangent to the spherical surface through \(P\); they form the so-called local horizon plane); the trihedral \(P_{x_h} y_h z_h\) is right-handed. In particular, the \(x_h\)-axis is tangent to the local parallel through \(P\) and is positive eastward; the \(y_h\)-axis is tangent to the local meridian through \(P\) and is positive southward. The symbols \(\hat{i}_h, \hat{j}_h, \hat{k}_h\) denote the unit vectors of the local horizon system.

3.4. **Wind Axes System.** The wind axes system \(P_{x_w} y_w z_w\) is a Cartesian reference frame defined as follows. Its origin \(P\) is identical with the instantaneous position of the spacecraft; the \(x_w\)-axis is tangent to the flight path (relative velocity) and is positive forward; the axes \(y_w, z_w\) are orthogonal to the \(x_w\)-axis and are such that the trihedral \(P_{x_w} y_w z_w\) is right-handed. In particular, the \(z_w\)-axis is contained in the plane of symmetry of the spacecraft and is positive downward for the normal flight attitude of the spacecraft; the \(y_w\)-axis is positive rightward for the normal flight attitude of the spacecraft. The symbols \(\hat{i}_w, \hat{j}_w, \hat{k}_w\) denote the unit vectors of the wind axes system.

3.5. **Body Axes System.** The body axes system \(P_{x_b} y_b z_b\) is a Cartesian reference frame defined as follows. Its origin \(P\) is identical with the instantaneous position of the spacecraft; the \(y_b\)-axis is orthogonal to the plane of symmetry of the spacecraft and is positive rightward; the axes \(x_b, z_b\) are orthogonal to the \(y_b\)-axis, are contained in the plane of symmetry, and are such that the trihedral \(P_{x_b} y_b z_b\) is right-handed. In particular, the \(x_b\)-axis is positive forward, the \(y_b\)-axis is positive
rightward, and the $z_b$-axis is positive downward for the normal flight attitude of the spacecraft. The symbols $\mathbf{i}_b, \mathbf{j}_b, \mathbf{k}_b$ denote the unit vectors of the body axes system.
4. **Relations between Coordinate Systems**

In this section, the relationships between the different coordinate systems are derived; more specifically, attention is focused on the following system pairs: inertial axes-Earth axes; Earth axes-local horizon; inertial axes-local horizon; local horizon-wind axes; and wind axes-body axes.

4.1. **Basic Relations.** In the inertial axes system, a point P can be described via its Cartesian coordinates \(x_i, y_i, z_i\) or via its spherical coordinates \(r_i, \theta_i, \phi_i\). Here, \(r_i\) is the radial distance from the center of the Earth; \(\theta_i\) is the longitude, positive eastward; and \(\phi_i\) is the latitude, positive northward. The following relations hold between Cartesian coordinates and spherical coordinates:

\[
\begin{align*}
    x_i &= r_i \cos \theta_i \cos \phi_i, \\
    y_i &= r_i \sin \theta_i \cos \phi_i, \\
    z_i &= r_i \sin \phi_i.
\end{align*}
\]

In the Earth axes system, a point P can be described via its Cartesian coordinates \(x_e, y_e, z_e\) or via its spherical coordinates \(r, \theta, \phi\). Here, \(r\) is the radial distance from the center of the Earth; \(\theta\) is the longitude, positive eastward; and \(\phi\) is the latitude, positive northward. The following relations hold between Cartesian coordinates and spherical coordinates:

\[
\begin{align*}
    x_e &= r \cos \theta \cos \phi, \\
    y_e &= r \sin \theta \cos \phi, \\
    z_e &= r \sin \phi.
\end{align*}
\]
Note that the inertial axes system and the Earth axes system coincide at time
instant \( t = 0 \). Also note that the Earth axes system rotates with constant
angular velocity \( \omega \) with respect to the inertial axes system. Therefore,
in spherical coordinates, the following transformation relations hold:

\[
\begin{align*}
    r_i &= r, \quad (3a) \\
    \theta_i &= \theta + \omega t, \quad (3b) \\
    \phi_i &= \phi, \quad (3c)
\end{align*}
\]

implying the following inverse relations:

\[
\begin{align*}
    r &= r_i, \quad (4a) \\
    \theta &= \theta_i - \omega t, \quad (4b) \\
    \phi &= \phi_i. \quad (4c)
\end{align*}
\]

4.2. Transformation from Inertial Axes to Earth Axes. The Earth axes
system \( O_{e}x_{e}y_{e}z_{e} \) is obtained from the inertial axes system \( O_{i}x_{i}y_{i}z_{i} \) by means
of the counterclockwise rotation \( \omega t \) around the \( z_{i} \)-axis. Note that the
\( z_{e} \)-axis is the same as the \( z_{i} \)-axis.

In vector-matrix notation, the transformation leading from the inertial
axes to the Earth axes can be expressed as follows:

\[
\begin{bmatrix}
    \mathbf{i}_e \\
    \mathbf{j}_e \\
    \mathbf{k}_e
\end{bmatrix} =
\begin{bmatrix}
    \cos(\omega t) & \sin(\omega t) & 0 \\
    -\sin(\omega t) & \cos(\omega t) & 0 \\
    0 & 0 & 1
\end{bmatrix}
\begin{bmatrix}
    \mathbf{i}_i \\
    \mathbf{j}_i \\
    \mathbf{k}_i
\end{bmatrix}, \quad (5a)
\]
with the implication that

\[
\begin{bmatrix}
\hat{i}_i \\
\hat{j}_i \\
\hat{k}_i
\end{bmatrix} =
\begin{bmatrix}
\cos(\omega t) & -\sin(\omega t) & 0 \\
\sin(\omega t) & \cos(\omega t) & 0 \\
0 & 0 & 1
\end{bmatrix}
\begin{bmatrix}
\hat{i}_e \\
\hat{j}_e \\
\hat{k}_e
\end{bmatrix}.
\]

(5b)

4.3. Transformation from Earth Axes to Local Horizon. The local horizon system \( P_{hxhyhz} \) can be obtained from the Earth axes system \( O_{xeyleze} \) by means of the combination of four rotations and one translation. This requires the definition of four intermediate coordinate systems: the system \( O_{x1y1z1} \); the system \( O_{x2y2z2} \); the system \( P_{x3y3z3} \); and the system \( P_{x4y4z4} \).

The system \( O_{x1y1z1} \) is obtained from the Earth axes system \( O_{xeyleze} \) by means of the counterclockwise rotation \( \theta \) around the \( ze \)-axis. Note that the \( z1 \)-axis is the same as the \( ze \)-axis, that the axes \( x1,y1 \) are contained in the equatorial plane, and that the axes \( x1,z1 \) are contained in a meridian plane. The symbols \( \hat{i}_1,\hat{j}_1,\hat{k}_1 \) denote the unit vectors of the system \( O_{x1y1z1} \).

The system \( O_{x2y2z2} \) is obtained from the system \( O_{x1y1z1} \) by means of the clockwise rotation \( \phi \) around the \( y1 \)-axis. Note that the \( y2 \)-axis is the same as the \( y1 \)-axis, that the axes \( x2,z2 \) are contained in a meridian plane, and that the axes \( y2,z2 \) are contained in a plane parallel to the local horizon plane. The symbols \( \hat{i}_2,\hat{j}_2,\hat{k}_2 \) denote the unit vectors of the system \( O_{x2y2z2} \).

The system \( P_{x3y3z3} \) is obtained from the system \( O_{x2y2z2} \) by means of the radial translation \( r \), leading from point \( O \) to point \( P \). Since there is no rotation, the axes \( x3,y3,z3 \) are parallel to the axes \( x2,y2,z2 \); in particular, the axes \( x3,z3 \) are contained in a meridian plane, while
the axes $y_3,z_3$ are contained in the local horizon plane. The symbols $\hat{i}_3,\hat{j}_3,\hat{k}_3$ denote the unit vectors of the system $P_3Y_3Z_3$.

The system $P_4Y_4Z_4$ is obtained from the system $P_3Y_3Z_3$ by means of the counterclockwise rotation $\pi/2$ around the $z_3$-axis. Note that the $z_4$-axis is the same as the $z_3$-axis, that the axes $y_4,z_4$ are contained in a meridian plane, while the axes $x_4,z_4$ are contained in the local horizon plane. The symbols $\hat{i}_4,\hat{j}_4,\hat{k}_4$ denote the unit vectors of the system $P_4Y_4Z_4$.

The local horizon system $P_hY_hZ_h$ is obtained from the system $P_4Y_4Z_4$ by means of the clockwise rotation $\pi/2$ around the $x_4$-axis. Note that the $x_h$-axis is the same as the $x_4$-axis, that the axes $y_h,Z_h$ are contained in a meridian plane, and the axes $x_h,Y_h$ are contained in the local horizon plane.

In vector-matrix notation, the successive transformations leading from one coordinate system to another can be expressed as follows:

\[
\begin{bmatrix}
\hat{i}_1 \\
\hat{j}_1 \\
\hat{k}_1
\end{bmatrix} =
\begin{bmatrix}
\cos\theta & \sin\theta & 0 \\
-sin\theta & \cos\theta & 0 \\
0 & 0 & 1
\end{bmatrix}
\begin{bmatrix}
\hat{i}_e \\
\hat{j}_e \\
\hat{k}_e
\end{bmatrix},
\]

(6a)

\[
\begin{bmatrix}
\hat{i}_2 \\
\hat{j}_2 \\
\hat{k}_2
\end{bmatrix} =
\begin{bmatrix}
\cos\phi & 0 & \sin\phi \\
0 & 1 & 0 \\
-sin\phi & 0 & \cos\phi
\end{bmatrix}
\begin{bmatrix}
\hat{i}_1 \\
\hat{j}_1 \\
\hat{k}_1
\end{bmatrix},
\]

(6b)

and
Equations (6) imply that

\[
\begin{bmatrix}
\hat{i}_2 \\
\hat{j}_2 \\
\hat{k}_2
\end{bmatrix} =
\begin{bmatrix}
\cos\theta \cos\phi & \sin\theta \cos\phi & \sin\phi \\
-sin\theta & \cos\theta & 0 \\
-cos\theta \sin\phi & -sin\theta \sin\phi & \cos\phi
\end{bmatrix}
\begin{bmatrix}
\hat{i}_e \\
\hat{j}_e \\
\hat{k}_e
\end{bmatrix},
\] (8)

while Eqs. (7) imply that

\[
\begin{bmatrix}
\hat{i}_h \\
\hat{j}_h \\
\hat{k}_h
\end{bmatrix} =
\begin{bmatrix}
0 & 1 & 0 \\
0 & 0 & -1 \\
-1 & 0 & 0
\end{bmatrix}
\begin{bmatrix}
\hat{i}_2 \\
\hat{j}_2 \\
\hat{k}_2
\end{bmatrix}.
\] (9)
Therefore, upon combining Eqs. (8)-(9), we see that

\[
\begin{bmatrix}
\mathbf{i}_h \\
\mathbf{j}_h \\
\mathbf{k}_h
\end{bmatrix} =
\begin{bmatrix}
\sin\theta & \cos\theta & 0 \\
\cos\theta \sin\phi & \sin\theta \sin\phi & -\cos\phi \\
-\cos\theta \cos\phi & -\sin\theta \cos\phi & -\sin\phi
\end{bmatrix}
\begin{bmatrix}
\mathbf{i}_e \\
\mathbf{j}_e \\
\mathbf{k}_e
\end{bmatrix},
\]  
(10a)

with the implication that

\[
\begin{bmatrix}
\mathbf{i}_e \\
\mathbf{j}_e \\
\mathbf{k}_e
\end{bmatrix} =
\begin{bmatrix}
\sin\theta & \cos\theta \sin\phi & -\cos\theta \cos\phi \\
\cos\theta & \sin\theta \sin\phi & -\sin\theta \cos\phi \\
0 & -\cos\phi & -\sin\phi
\end{bmatrix}
\begin{bmatrix}
\mathbf{i}_h \\
\mathbf{j}_h \\
\mathbf{k}_h
\end{bmatrix}.
\]  
(10b)

4.4. **Transformation from Inertial Axes to Local Horizon.** The transformation leading from the inertial axes system to the local horizon system requires the combination of five rotations and one translation, as one surmises from Sections 4.2 and 4.3. If one combines Eqs. (5a) and (10a) and accounts for Eqs. (3), the following result is obtained:

\[
\begin{bmatrix}
\mathbf{i}_h \\
\mathbf{j}_h \\
\mathbf{k}_h
\end{bmatrix} =
\begin{bmatrix}
\sin\theta_i & \cos\theta_i & 0 \\
\cos\theta_i \sin\phi_i & \sin\theta_i \sin\phi_i & -\cos\phi_i \\
-\cos\theta_i \cos\phi_i & -\sin\theta_i \cos\phi_i & -\sin\phi_i
\end{bmatrix}
\begin{bmatrix}
\mathbf{i}_i \\
\mathbf{j}_i \\
\mathbf{k}_i
\end{bmatrix},
\]  
(11a)

with the implication that
Transformation from Local Horizon to Wind Axes. The wind axes system $\mathbf{P}_w \mathbf{Y}_w \mathbf{Z}_w$ can be obtained from the local horizon system $\mathbf{P}_h \mathbf{Y}_h \mathbf{Z}_h$ by means of the combination of three rotations. This requires the definition of two intermediate coordinate systems: the system $\mathbf{P}_x_5 \mathbf{Y}_5 \mathbf{Z}_5$ and the system $\mathbf{P}_x_6 \mathbf{Y}_6 \mathbf{Z}_6$.

The system $\mathbf{P}_x_5 \mathbf{Y}_5 \mathbf{Z}_5$ is obtained from the local horizon system $\mathbf{P}_x_5 \mathbf{Y}_5 \mathbf{Z}_5$ by means of the counterclockwise rotation $\chi$ around the $\mathbf{z}_h$-axis. Note that the $\mathbf{z}_5$-axis is the same as the $\mathbf{z}_h$-axis, that the axes $\mathbf{x}_5, \mathbf{y}_5$ are contained in the local horizon plane, and that the axes $\mathbf{x}_5, \mathbf{z}_5$ are contained in the plane $(\mathbf{OP}, \mathbf{V})$, where $\mathbf{OP}$ is the radius vector connecting the points $\mathbf{O}$ and $\mathbf{P}$ and $\mathbf{V}$ is the relative velocity vector, namely, the velocity of the spacecraft with respect to the Earth axes system. Also note that the axis $\mathbf{x}_5$ has the direction of the projected relative velocity vector $\mathbf{V}_p$; this is the projection of $\mathbf{V}$ on the local horizon. The angle $\chi$ is called the heading angle and is positive if the projected relative velocity vector $\mathbf{V}_p$ is directed outward with respect to the local parallel. The symbols $i_5, j_5, k_5$ denote the unit vectors of the system $\mathbf{P}_x_5 \mathbf{Y}_5 \mathbf{Z}_5$.

The system $\mathbf{P}_x_6 \mathbf{Y}_6 \mathbf{Z}_6$ is obtained from the system $\mathbf{P}_x_5 \mathbf{Y}_5 \mathbf{Z}_5$ by means of the counterclockwise rotation $\gamma$ around the $\mathbf{y}_5$-axis. Note that the $\mathbf{y}_6$-axis is the same as the $\mathbf{y}_5$-axis and that the axes $\mathbf{x}_6, \mathbf{z}_6$ are contained in the plane $(\mathbf{OP}, \mathbf{V})$. Also note that the $\mathbf{x}_6$-axis is positive forward and that the
The $z_6$-axis is positive downward. The angle $\gamma$ is called the path inclination and is positive if the relative velocity vector $\vec{V}$ is inclined upward with respect to the local horizon. The symbols $\hat{i}_6, \hat{j}_6, \hat{k}_6$ denote the unit vectors of the system $P_{x_6}y_6z_6$.

The wind axes system $P_{x_w}y_wz_w$ is obtained from the system $P_{x_6}y_6z_6$ by means of the counterclockwise rotation $\mu$ around the $x_6$-axis. Note that the $x_w$-axis is the same as the $x_6$-axis. Also, note that the $x_w$-axis is positive forward, the $y_w$-axis is positive rightward, and the $z_w$-axis is positive downward and is contained in the plane of symmetry of the spacecraft. The angle $\mu$ is called the angle of bank and is positive if the spacecraft is banked to the right.

In vector-matrix notation, the successive transformations leading from one coordinate system to another can be expressed as follows:

\[
\begin{bmatrix}
\hat{i}_5 \\
\hat{j}_5 \\
\hat{k}_5
\end{bmatrix} = \begin{bmatrix}
\cos \chi & \sin \chi & 0 \\
-sin \chi & \cos \chi & 0 \\
0 & 0 & 1
\end{bmatrix} \begin{bmatrix}
\hat{i}_h \\
\hat{j}_h \\
\hat{k}_h
\end{bmatrix},
\]

\[
\begin{bmatrix}
\hat{i}_6 \\
\hat{j}_6 \\
\hat{k}_6
\end{bmatrix} = \begin{bmatrix}
\cos \gamma & 0 & \sin \gamma \\
0 & 1 & 0 \\
\sin \gamma & 0 & \cos \gamma
\end{bmatrix} \begin{bmatrix}
\hat{i}_5 \\
\hat{j}_5 \\
\hat{k}_5
\end{bmatrix},
\]

\[
\begin{bmatrix}
\hat{i}_w \\
\hat{j}_w \\
\hat{k}_w
\end{bmatrix} = \begin{bmatrix}
1 & 0 & 0 \\
0 & \cos \mu & \sin \mu \\
0 & -\sin \mu & \cos \mu
\end{bmatrix} \begin{bmatrix}
\hat{i}_6 \\
\hat{j}_6 \\
\hat{k}_6
\end{bmatrix}.
\]
Equations (12) lead to
\[
\begin{bmatrix}
\hat{i}_w \\
\hat{j}_w \\
\hat{k}_w
\end{bmatrix} =
\begin{bmatrix}
\cos \gamma \cos \chi & \cos \gamma \sin \chi & -\sin \gamma \\
\sin \mu \sin \gamma \cos \chi & \sin \mu \sin \gamma \sin \chi & \sin \mu \cos \gamma \\
-\cos \mu \sin \chi & +\cos \mu \cos \chi & \sin \mu \cos \gamma
\end{bmatrix}
\begin{bmatrix}
\hat{i}_h \\
\hat{j}_h \\
\hat{k}_h
\end{bmatrix},
\tag{13a}
\]
with the implication that
\[
\begin{bmatrix}
\hat{i}_h \\
\hat{j}_h \\
\hat{k}_h
\end{bmatrix} =
\begin{bmatrix}
\cos \gamma \cos \chi & \sin \mu \sin \gamma \cos \chi & \cos \mu \sin \gamma \cos \chi \\
-\cos \mu \sin \chi & +\sin \mu \sin \chi & \cos \mu \sin \gamma \cos \chi \\
-\sin \gamma & \sin \mu \cos \gamma & \cos \mu \cos \gamma
\end{bmatrix}
\begin{bmatrix}
\hat{i}_w \\
\hat{j}_w \\
\hat{k}_w
\end{bmatrix},
\tag{13b}
\]

4.6. **Transformation from Wind Axes to Body Axes.** The body system $P_{x_b}Y_bZ_b$ can be obtained from the wind axes system $P_{x_w}Y_wZ_w$ by means of the combination of two rotations. This requires the definition of one intermediate coordinate system, the system $P_{x_7}Y_7Z_7$.

The system $P_{x_7}Y_7Z_7$ is obtained from the wind axes system $P_{x_w}Y_wZ_w$ by means of the counterclockwise rotation $\sigma$ around the $z_w$-axis. Note that the $z_7$-axis is the same as the $z_w$-axis and that the axes $x_7,z_7$ are contained in the plane of symmetry of the spacecraft. Also note that the axis $x_7$ is positive forward, the axis $y_7$ is positive rightward, and the axis $z_7$ is positive downward. The angle $\sigma$ is called the sideslip angle and is positive if the relative velocity vector $\vec{V}$ is directed leftward with respect to the plane of symmetry of the spacecraft.
The body axes system $P_x b_y b_z b$ is obtained from the system $P_x y_7 z_7$ by means of the counterclockwise rotation $\alpha$ around the $y_7$-axis. Note that the $y_b$-axis is the same as the $y_7$-axis and that the axes $x_b, z_b$ are contained in the plane of symmetry of the spacecraft. Also note that the axis $x_b$ is positive forward, the axis $y_b$ is positive rightward, and the axis $z_b$ is positive downward. The angle $\alpha$ is called the angle of attack and is positive if the relative velocity vector $\vec{V}$ is directed downward with respect to the $x_b$-axis of the spacecraft.

In vector-matrix notation, the successive transformations leading from one coordinate system to another can be expressed as follows:

\[
\begin{bmatrix}
i_7 \\
j_7 \\
k_7
\end{bmatrix} =
\begin{bmatrix}
\cos \alpha & \sin \alpha & 0 \\
-sin \alpha & \cos \alpha & 0 \\
0 & 0 & 1
\end{bmatrix}
\begin{bmatrix}
i_w \\
j_w \\
k_w
\end{bmatrix}, \quad (14a)
\]

\[
\begin{bmatrix}
i_b \\
j_b \\
k_b
\end{bmatrix} =
\begin{bmatrix}
\cos \alpha & 0 & -\sin \alpha \\
0 & 1 & 0 \\
\sin \alpha & 0 & \cos \alpha
\end{bmatrix}
\begin{bmatrix}
i_7 \\
j_7 \\
k_7
\end{bmatrix}, \quad (14b)
\]

with the implication that

\[
\begin{bmatrix}
i_b \\
j_b \\
k_b
\end{bmatrix} =
\begin{bmatrix}
\cos \alpha \cos \alpha & \cos \alpha \sin \alpha & -\sin \alpha \\
-sin \alpha & \cos \alpha & 0 \\
\sin \alpha \cos \alpha & \sin \alpha \sin \alpha & \cos \alpha
\end{bmatrix}
\begin{bmatrix}
i_w \\
j_w \\
k_w
\end{bmatrix}, \quad (15a)
\]
and that

\[
\begin{bmatrix}
    \hat{i}_w \\
    \hat{j}_w \\
    \hat{k}_w
\end{bmatrix} =
\begin{bmatrix}
    \cos \alpha \cos \sigma & -\sin \alpha & \sin \alpha \cos \sigma \\
    \cos \alpha \sin \sigma & \cos \sigma & \sin \alpha \sin \sigma \\
    -\sin \alpha & 0 & \cos \alpha
\end{bmatrix}
\begin{bmatrix}
    \hat{i}_b \\
    \hat{j}_b \\
    \hat{k}_b
\end{bmatrix},
\]

(15b)
5. **Angular Velocity**

In this section, we compute the angular velocity (or evolutory velocity) of the local horizon system with respect to the inertial axes system. To do so, consider the behavior of the spacecraft between the time instants $t$ and $t + dt$, and denote by $d\Omega_{hi}$ the infinitesimal vectorial rotation of the local horizon system with respect to the inertial axes system. This infinitesimal vectorial rotation can be decomposed into partial rotations as follows:

$$d\Omega_{hi} = d\Omega_{h4} + d\Omega_{43} + d\Omega_{32} + d\Omega_{21} + d\Omega_{1e} + d\Omega_{ei}, \quad (16)$$

with

$$d\Omega_{h4} = 0, \quad (17a)$$
$$d\Omega_{43} = 0, \quad (17b)$$
$$d\Omega_{32} = 0, \quad (17c)$$
$$d\Omega_{21} = -d\phi \hat{J}_1, \quad (17d)$$
$$d\Omega_{1e} = d\theta \hat{k}_e, \quad (17e)$$
$$d\Omega_{ei} = \omega dt \hat{k}_i. \quad (17f)$$

Here, $d\theta$ denotes the infinitesimal change of the longitude, $d\phi$ denotes the infinitesimal change of the latitude, and $\omega dt$ denotes the infinitesimal rotation of the Earth axes system with respect to the inertial axes system. The rotation $\omega dt$ occurs around the $z_i$-axis and is positive counterclockwise;
the rotation $d\theta$ occurs around the $z_e$-axis and is positive counterclockwise; and the rotation $d\phi$ occurs around the $y_i$-axis and is positive clockwise. This explains the difference in the signs appearing on the right-hand sides of Eqs. (17d), (17e), (17f).

Upon combining Eqs. (16)-(17), we see that the infinitesimal vectorial rotation of the local horizon system with respect to the inertial axes system can be written as

$$d\Omega_{hi} = w dt k_i + d\theta k_e - d\phi j_1.$$  \hspace{1cm} (18)

As a consequence, the angular velocity of the local horizon system with respect to the inertial axes system is given by

$$\omega_{hi} = d\Omega_{hi}/dt = \omega k_i + \theta k_e - \phi j_1.$$  \hspace{1cm} (19)

In the light of Eqs. (5)-(11), the unit vectors $k_i$, $k_e$, $j_1$ can be expressed in terms of the unit vectors of the local horizon system as follows:

$$k_i = -\cos\phi \ j_h - \sin\phi \ k_h,$$  \hspace{1cm} (20a)

$$k_e = -\cos\phi \ j_h - \sin\phi \ k_h,$$  \hspace{1cm} (20b)

$$j_1 = \ i_h,$$  \hspace{1cm} (20c)

so that

$$\omega_{hi} = -\phi \ i_h - (\theta \cos\phi + \omega \cos\phi) \ j_h - (\theta \sin\phi + \omega \sin\phi) \ k_h.$$  \hspace{1cm} (21)

Next, we invoke Eqs. (3),
\[
\begin{align*}
  r_i &= r, \\
  \theta_i &= \theta + \omega t, \\
  \phi_i &= \phi,
\end{align*}
\]

and observe that
\[
\begin{align*}
  \dot{r}_i &= \dot{r}, \\
  \dot{\theta}_i &= \dot{\theta} + \dot{\omega}, \\
  \dot{\phi}_i &= \dot{\phi}.
\end{align*}
\]

As a consequence, the angular velocity (21) can be rewritten as
\[
\omega_{hi} = -\phi_i \hat{i}_h - \dot{\theta}_i \cos \phi_i \hat{j}_h - \dot{\theta}_i \sin \phi_i \hat{k}_h.
\]

Next, Poisson's formulas are employed to compute the derivatives of the unit vectors of the local horizon system with respect to time:
\[
\begin{align*}
  \frac{d\hat{i}_h}{dt} &= \omega_{hi} \times \hat{i}_h, \\
  \frac{d\hat{j}_h}{dt} &= \omega_{hi} \times \hat{j}_h, \\
  \frac{d\hat{k}_h}{dt} &= \omega_{hi} \times \hat{k}_h.
\end{align*}
\]

Upon combining Eqs. (24)-(25), we obtain the relations
\[
\frac{d\hat{i}_h}{dt} = - (\dot{\theta}_i \sin \phi_i) \hat{j}_h + (\dot{\theta}_i \cos \phi_i) \hat{k}_h.
\]
\[
\begin{align*}
\frac{d\dot{i}_h}{dt} &= (\dot{\theta}_i \sin \phi_i) \dot{i}_h - \phi_i \dot{k}_h, \quad (26b) \\
\frac{d\dot{k}_h}{dt} &= -(\dot{\theta}_i \cos \phi_i) \dot{i}_h + \phi_i \dot{j}_h, \quad (26c)
\end{align*}
\]
whose vector-matrix form is the following:
\[
\begin{bmatrix}
\frac{d\dot{i}_h}{dt} \\
\frac{d\dot{j}_h}{dt} \\
\frac{d\dot{k}_h}{dt}
\end{bmatrix}
= \begin{bmatrix}
0 & -\dot{\theta}_i \sin \phi_i & \dot{\theta}_i \cos \phi_i \\
\dot{\theta}_i \sin \phi_i & 0 & -\dot{\phi}_i \\
-\dot{\theta}_i \cos \phi_i & \dot{\phi}_i & 0
\end{bmatrix}
\begin{bmatrix}
\dot{i}_h \\
\dot{j}_h \\
\dot{k}_h
\end{bmatrix}.
\quad (27)
\]
It is interesting to note that Eq. (27) can also be obtained by taking the time derivative of Eq. (11a) and using Eq. (11b).
6. **Kinematical Equations**

In this section, we derive the scalar relationships corresponding to the vectorial equation

\[
d\mathbf{OP}/dt = \mathbf{V}_i.
\]  

(28)

Here, \(\mathbf{V}_i\) denotes the velocity of the spacecraft with respect to the inertial axes system and \(\mathbf{OP}\) denotes the position vector joining the center of the Earth \(O\) with the spacecraft position \(P\).

First, we observe that the position vector \(\mathbf{OP}\) is given by

\[
\mathbf{OP} = -r_i \mathbf{k}_h,
\]

(29)

where \(r_i\) is the radial distance from the center of the Earth and \(\mathbf{k}_h\) is the third unit vector of the local horizon system. As a consequence, the time derivative of \(\mathbf{OP}\) can be written as

\[
d\mathbf{OP}/dt = -r_i \dot{\mathbf{k}}_h - r_i (d\mathbf{k}_h/dt),
\]

(30)

where \(d\mathbf{k}_h/dt\) is given by Eq. (26c). Therefore, upon combining Eqs. (26c) and (30), we obtain the relation

\[
d\mathbf{OP}/dt = \mathbf{\theta}_i r_i \cos \phi_i \mathbf{i}_h - \mathbf{\phi}_i r_i \mathbf{j}_h - r_i \mathbf{k}_h.
\]

(31)

We recall that the relative velocity vector \(\mathbf{V}\) (velocity of the spacecraft with respect to the Earth axes system) can be identified via three elements: the relative velocity modulus \(V\), the relative path inclination \(\gamma\), and the relative heading angle \(\chi\) (see Section 4.5). Analogously, the inertial velocity vector \(\mathbf{V}_i\) (velocity of the spacecraft with respect to the inertial axes system) can be identified via three elements: the inertial velocity
modulus $V_i$, the inertial path inclination $\gamma_i$, and the inertial heading angle $\chi_i$.

Let $\hat{V}_p$ and $\hat{V}_{ip}$ denote the projections of the vectors $\hat{V}$ and $\hat{V}_i$ on the local horizon. The relative path inclination $\gamma$ is the angle between the vectors $\hat{V}$ and $\hat{V}_p$; analogously, the inertial path inclination $\gamma_i$ is the angle between the vectors $\hat{V}_i$ and $\hat{V}_{ip}$. The relative heading angle $\chi$ is the angle between the vectors $\hat{V}_p$ and $\hat{i}_h$; analogously, the inertial heading angle $\chi_i$ is the angle between the vectors $\hat{V}_{ip}$ and $\hat{i}_h$. The conventions used for the signs of $\gamma$, $\chi$.

With the above understanding, just as the relative velocity vector $\hat{V}$ can be written in terms of its components on the local horizon system,

$$\hat{V} = V\cos\gamma \cos\chi \hat{i}_h + V\cos\gamma \sin\chi \hat{j}_h - V\sin\gamma \hat{k}_h,$$

the inertial velocity vector $\hat{V}_i$ can be written in terms of its components on the local horizon system,

$$\hat{V}_i = V_i\cos\gamma_i \cos\chi_i \hat{i}_h + V_i\cos\gamma_i \sin\chi_i \hat{j}_h - V_i\sin\gamma_i \hat{k}_h.$$

Finally, upon combining Eqs. (28), (31), (33), and upon projecting the resulting vectorial equation on the axes of the local horizon system, we obtain the following scalar form of the kinematical equations:

$$\dot{\theta}_i = V_i\cos\gamma_i \cos\chi_i / r_i \cos\phi_i,$$

$$\dot{r}_i = V_i\sin\gamma_i,$$

$$\dot{\chi}_i = -V_i\cos\gamma_i \sin\chi_i / r_i,$$
6.1. Relations between Inertial Velocity and Relative Velocity. We employ the theorem of composition of velocities, which states that the inertial velocity $\vec{V}_i$ is the sum of the relative velocity $\vec{V}$ and the transport velocity $\vec{V}_t$,

$$\vec{V}_i = \vec{V} + \vec{V}_t,$$

with

$$\vec{V}_t = \omega \times \vec{OP}.$$  \hspace{1cm} (36)

Owing to the fact that

$$\vec{\omega} = -\omega \cos \phi \hat{j}_h - \omega \sin \phi \hat{k}_h = -\omega \cos \phi_i \hat{i}_h - \omega \sin \phi_i \hat{k}_h,$$

$$\vec{OP} = -r_\kappa \hat{k}_h = -r_i \hat{k}_h,$$

the transport velocity (36) can be rewritten as

$$\vec{V}_t = \omega r \cos \phi \hat{i}_h = \omega r_i \cos \phi_i \hat{i}_h.$$ \hspace{1cm} (38)

Therefore, upon combining Eqs. (32), (33), (35), (38), and upon projecting the resulting vectorial equation on the axes of the local horizon system, the following scalar equations are obtained:

$$V_i \cos \gamma_i \cos \chi_i - \omega r_i \cos \phi_i = V \cos \gamma \cos \chi,$$
\hspace{1cm} (39a)

$$V_i \cos \gamma_i \sin \chi_i = V \cos \gamma \sin \chi,$$ \hspace{1cm} (39b)

$$V_i \sin \gamma_i = V \sin \gamma.$$ \hspace{1cm} (39c)
Laborious manipulations, omitted for the sake of brevity, lead to the following transformation relations:

\[ V_i = \sqrt{V^2 + 2w_V \cos \gamma \cos \phi + (\omega \cos \phi)^2}, \] (40a)

\[ \tan \gamma_i = \frac{V \sin \gamma}{\sqrt{(V \cos \gamma)^2 + 2w_V \cos \gamma \cos \phi + (\omega \cos \phi)^2}}, \] (40b)

\[ \tan \chi_i = \frac{V \cos \gamma \sin \chi_i}{(V \cos \gamma \cos \chi_i + \omega \cos \phi)}. \] (40c)

Equations (40) imply the following inverse relations:

\[ V = \sqrt{V_i^2 - 2w_i V_i \cos \gamma_i \cos \chi_i \cos \phi_i + (\omega_i \cos \phi_i)^2}, \] (41a)

\[ \tan \gamma = \frac{V_i \sin \gamma_i}{\sqrt{(V_i \cos \gamma_i)^2 - 2w_i V_i \cos \gamma_i \cos \chi_i \cos \phi_i + (\omega_i \cos \phi_i)^2}}, \] (41b)

\[ \tan \chi = \frac{V_i \cos \gamma_i \sin \chi_i}{(V_i \cos \gamma_i \cos \chi_i - \omega_i \cos \phi_i)}. \] (41c)
7. **Dynamical Equations**

In this section, we derive the scalar relationships corresponding to the vectorial equation

\[ \dot{T} + \dot{A} + \dot{W} = ma_i = m(dV_i/dt), \]  

(42)

where \( \dot{T} \) is the thrust, \( \dot{A} \) is the aerodynamic force, \( \dot{W} \) is the gravitational force, \( m \) is the mass of the spacecraft, and \( \dot{a}_i \) is the inertial acceleration. We consider the case where the engine is shut-off, so that

\[ \dot{T} = 0, \]  

(43)

and the mass of the spacecraft is constant. Hence, Eq. (42) is written as

\[ \dot{A} + \dot{W} = ma_i = m(dV_i/dt). \]  

(44)

We now compute the components of the vectors appearing in Eq. (44) on the axes of the local horizon system.

7.1. **Aerodynamic Force.** The components of the aerodynamic force on the wind axes are the drag \( D \), the side force \( Q \), and the lift \( L \). Therefore, the aerodynamic force can be written as

\[ \dot{A} = -D\dot{t}_w - Q\dot{\ell}_w - Lk_w. \]  

(45)

No special significance is implied in the signs appearing on the right-hand side of Eq. (45). These signs merely reflect the conventions adopted in this report with regard to the positive values for the drag, the side force, and the lift.
In vector-matrix form, Eq. (45) can be rewritten as follows:

\[
\mathbf{\dot{A}} = \begin{bmatrix}
D & 0 & 0 & i_w \\
0 & Q & 0 & j_w \\
0 & 0 & L & k_w
\end{bmatrix}
\]

(46)

Therefore, upon combining Eqs. (13a) and (46), we obtain the relation

\[
\mathbf{A} = \begin{bmatrix}
-D \cos \gamma \cos \chi & -D \cos \gamma \sin \chi & D \sin \gamma & i_h \\
-Q \sin \mu \sin \gamma \cos \chi & -Q \sin \mu \sin \gamma \sin \chi & -Q \sin \mu \cos \gamma & j_h \\
+Q \cos \mu \sin \chi & -Q \cos \mu \cos \chi & -L \cos \mu \sin \chi & k_h
\end{bmatrix}
\]

(47)

The next step consists of rewriting the elements of the matrix in Eq. (47) in terms of inertial velocity elements \( V_i, \gamma_i, \chi_i \), instead of relative velocity elements \( V, \gamma, \chi \). For this purpose, let the following functions of the inertial velocity elements be defined:

\[
M_1 = M_1(V_i, \gamma_i, \chi_i) = V_i \cos \gamma_i \cos \chi_i - \omega_r \cos \phi, \quad (48a)
\]

\[
M_2 = M_2(V_i, \gamma_i, \chi_i) = V_i \cos \gamma_i \sin \chi_i, \quad (48b)
\]

\[
M_3 = M_3(V_i, \gamma_i, \chi_i) = V_i \sin \gamma_i, \quad (48c)
\]

and

\[
N_1 = N_1(V_i, \gamma_i, \chi_i) = \sqrt{(M_1^2 + M_2^2 + M_3^2)}, \quad (49a)
\]

\[
N_2 = N_2(V_i, \gamma_i, \chi_i) = \sqrt{M_1^2 + M_2^2}/\sqrt{M_1^2 + M_2^2 + M_3^2}, \quad (49b)
\]
\[ N_3 = N_3(V_1, \gamma_1, x_1) = M_3/\sqrt{M_1^2 + M_2^2 + M_3^2}, \quad (49c) \]
\[ N_4 = N_4(V_1, \gamma_1, x_1) = M_1/\sqrt{M_1^2 + M_2^2}, \quad (49d) \]
\[ N_5 = N_5(V_1, \gamma_1, x_1) = M_2/\sqrt{M_1^2 + M_2^2}. \quad (49e) \]

In the light of Eqs. (39), the following relations can be established:

\[ M_1 = V \cos \gamma \cos x, \quad (50a) \]
\[ M_2 = V \cos \gamma \sin x, \quad (50b) \]
\[ M_3 = V \sin \gamma, \quad (50c) \]

and

\[ N_1 = V, \quad (51a) \]
\[ N_2 = \cos \gamma, \quad (51b) \]
\[ N_3 = \sin \gamma, \quad (51c) \]
\[ N_4 = \cos x, \quad (51d) \]
\[ N_5 = \sin x. \quad (51e) \]

Therefore, upon combining Eqs. (47) and (50)-(51), we obtain the relation

\[ A = \begin{bmatrix}
-DN_2 N_4 & -DN_2 N_5 & DN_3 & i_h \\
-Qsin \mu N_3 N_4 & -Qsin \mu N_3 N_5 & -Qsin \mu N_2 & j_h \\
+Qcos \mu N_5 & -Qcos \mu N_4 & & k_h \\
-Lcos \mu N_3 N_4 & -Lcos \mu N_3 N_5 & -Lcos \mu N_2 & \\
-Lsin \mu N_5 & +Lsin \mu N_4 & & \\
\end{bmatrix}. \quad (52) \]
7.2. **Gravitational Force.** Here, we assume that the Earth is an oblate planet and that its mass has radial symmetry with respect to the axis of rotation. Because the equatorial radius \( r_e \) is larger than the polar radius \( r_p \), the gravity force \( \vec{W} \) has two components: the radial component \( m g \), directed toward the center of the Earth, and the latitudinal component \( m f \), tangent to the local meridian and directed toward the Equator. Therefore, the gravity force can be written as

\[
\vec{W} = m f \hat{j}_h + m g \hat{k}_h. \tag{53}
\]

The radial component \( g \) and the latitudinal component \( f \) of the acceleration of gravity are related to the Earth's gravitational potential \( U \) by the expressions

\[
g = \partial U / \partial r, \quad f = (1/r) \partial U / \partial \phi, \tag{54}
\]

where

\[
U = -\left(\mu_e / r_i^2\right) \left[ 1 + J_2 \left( r_e / r_i \right)^2 H_2 + J_3 \left( r_e / r_i \right)^3 H_3 + J_4 \left( r_e / r_i \right)^4 H_4 \right], \tag{55a}
\]

\[
H_2 = 1/2 - (3/2) \sin^2 \phi, \tag{55b}
\]

\[
H_3 = (3/2) \sin \phi - (5/2) \sin^3 \phi, \tag{55c}
\]

\[
H_4 = -(3/8) + (30/8) \sin^2 \phi - (35/8) \sin^4 \phi. \tag{55d}
\]

Here, \( \mu_e \) is the Earth's gravitational constant, \( r_e \) is the equatorial radius, and \( J_2, J_3, J_4 \) denote the characteristic constants of the Earth's gravitational field. Note that the expression for \( U \) is approximate, since harmonics of order higher than four are ignored.
Upon combining Eqs. (54)-(55), we see that the components of the acceleration of gravity can be written as

\[ g = \left( \frac{\mu_e}{r_i^2} \right) \left[ 1 + J_2 \left( \frac{r_e}{r_i} \right)^2 G_2 + J_3 \left( \frac{r_e}{r_i} \right)^3 G_3 + J_4 \left( \frac{r_e}{r_i} \right)^4 G_4 \right]. \]  
\[ (56a) \]

\[ G_2 = \frac{3}{2} - \left( \frac{9}{2} \right) \sin^2 \phi_i, \]
\[ (56b) \]

\[ G_3 = 6 \sin \phi_i - 10 \sin^3 \phi_i, \]
\[ (56c) \]

\[ G_4 = -\frac{15}{8} + \left( \frac{150}{8} \right) \sin^2 \phi_i - \left( \frac{173}{8} \right) \sin^4 \phi_i, \]
\[ (56d) \]

and

\[ f = \left( \frac{\mu_e}{r_i^2} \right) \left[ J_2 \left( \frac{r_e}{r_i} \right)^2 F_2 + J_3 \left( \frac{r_e}{r_i} \right)^3 F_3 + J_4 \left( \frac{r_e}{r_i} \right)^4 F_4 \right], \]
\[ (57a) \]

\[ F_2 = 3 \sin \phi_i \cos \phi_i, \]
\[ (57b) \]

\[ F_3 = -\left( \frac{3}{2} \right) \cos \phi_i + \left( \frac{15}{2} \right) \sin^2 \phi_i \cos \phi_i, \]
\[ (57c) \]

\[ F_4 = -\left( \frac{15}{2} \right) \sin \phi_i \cos \phi_i + \left( \frac{35}{2} \right) \sin^3 \phi_i \cos \phi_i. \]
\[ (57d) \]

7.3. **Inertial Acceleration.** Let \( V_{ixh}, V_{iyh}, V_{izh} \) denote the components of the inertial velocity on the local horizon system,

\[ V_{ixh} = V_i \cos \gamma_i \cos \chi_i, \]
\[ (58a) \]

\[ V_{iyh} = V_i \cos \gamma_i \sin \chi_i, \]
\[ (58b) \]

\[ V_{izh} = -V_i \sin \gamma_i. \]
\[ (58c) \]

With this understanding, the inertial velocity (33) can be rewritten as

\[ \dot{V}_i = V_{ixh} \dot{h} + V_{iyh} \dot{h} + V_{izh} \dot{h}. \]  
\[ (59) \]
Therefore, the inertial acceleration is given by

\[
\frac{dV_i}{dt} = \frac{dV_{ix_h}}{dt} + \frac{dV_{iy_h}}{dt} + \frac{dV_{iz_h}}{dt}
\]

\[
+ V_{ix_h}(\frac{di_h}{dt}) + V_{iy_h}(\frac{dj_h}{dt}) + V_{iz_h}(\frac{dk_h}{dt}).
\]  

(60)

If we combine Eqs. (26) and (34), the time derivatives of the unit vectors of the local horizon system can be written as

\[
\frac{di_h}{dt} = -(V_{ic} \cos \gamma_i \cos \chi_i \tan \delta_i / r_i)k_h + (V_{ic} \cos \gamma_i \cos \chi_i / r_i)k_h,
\]  

(61a)

\[
\frac{dj_h}{dt} = (V_{ic} \cos \gamma_i \cos \chi_i \tan \delta_i / r_i)i_h + (V_{sc} \cos \gamma_i \sin \chi_i / r_i)k_h,
\]  

(61b)

\[
\frac{dk_h}{dt} = -(V_{ic} \cos \gamma_i \cos \chi_i / r_i)i_h - (V_{ic} \cos \gamma_i \sin \chi_i / r_i)j_h.
\]  

(61c)

Upon combining Eqs. (58), (60), (61), the inertial acceleration becomes

\[
\frac{dV_i}{dt} = (V_{ix_h} + V_{ic} \cos \gamma_i \cos \chi_i \sin \chi_i \tan \delta_i / r_i + V_{ic} \cos \gamma_i \sin \gamma_i \cos \chi_i / r_i)i_h
\]

\[
+ (V_{iy_h} - V_{ic} \cos \gamma_i \cos \gamma_i \cos \gamma_i \sin \gamma_i \cos \gamma_i / r_i)j_h
\]

\[
+ (V_{iz_h} + V_{ic} \cos \gamma_i \cos \chi_i / r_i)k_h.
\]  

(62)

7.4. Scalar Equations. Next, we combine Eqs. (44), (52), (53), (62).

Upon projecting the resulting vectorial equation on the axes of the local horizon system, we obtain the following scalar equations:

\[
V_{ixh} = -DN_2N_4/m + (QN_5/m)\cos \mu - (QN_3N_4/m)\sin \mu
\]

\[
- (LN_5/m)\sin \mu - (LN_3N_4/m)\cos \mu
\]

\[
- (V_i^2 / r_i)(\cos \gamma_i \cos \gamma_i \sin \chi_i \tan \delta_i + \cos \gamma_i \sin \gamma_i \cos \chi_i),
\]  

(63a)
\[ V_{1y h} = -DN_2 N_5 / m - (QN_4 / m) \cos \mu - (QN_3 N_5 / m) \sin \mu \]
\[ + (LN_4 / m) \sin \mu - (LN_3 N_5 / m) \cos \mu \]
\[ + f + (V_i^2 / r_i) (\cos^2 \gamma_i \cos^2 \chi_i \tan \phi_i - \cos \gamma_i \sin \gamma_i \sin \chi_i), \] (63b)

\[ V_{iz h} = DN_3 / m - (QN_2 / m) \sin \mu - (LN_2 / m) \cos \mu \]
\[ + g - (V_i^2 / r_i) \cos^2 \gamma_i. \] (63c)

We recall that the components of the inertial velocity on the axes of the local horizon system are given by [see Eqs. (58)]

\[ V_{ix h} = V_i \cos \gamma_i \cos \chi_i, \] (64a)

\[ V_{iy h} = V_i \cos \gamma_i \sin \chi_i, \] (64b)

\[ V_{iz h} = -V_i \sin \gamma_i, \] (64c)

with the implication that

\[
\begin{bmatrix}
\dot{V}_{ix h} \\
\dot{V}_{iy h} \\
\dot{V}_{iz h}
\end{bmatrix}
= \begin{bmatrix}
\cos \gamma_i \cos \chi_i & -\sin \gamma_i \cos \chi_i & -\sin \gamma_i \\
\cos \gamma_i \sin \chi_i & -\sin \gamma_i \sin \chi_i & \cos \chi_i \\
-\sin \gamma_i & -\cos \gamma_i & 0
\end{bmatrix}
\begin{bmatrix}
\dot{V}_i \\
\dot{V}_i \gamma_i \\
\dot{V}_i \cos \gamma_i \chi_i
\end{bmatrix},
\] (65)

and that

\[
\begin{bmatrix}
\dot{V}_i \\
\dot{V}_i \gamma_i \\
\dot{V}_i \cos \gamma_i \chi_i
\end{bmatrix}
= \begin{bmatrix}
\cos \gamma_i \cos \chi_i & \cos \gamma_i \sin \chi_i & -\sin \gamma_i \\
-\sin \gamma_i \cos \chi_i & -\sin \gamma_i \sin \chi_i & -\cos \gamma_i \\
-\sin \chi_i & \cos \chi_i & 0
\end{bmatrix}
\begin{bmatrix}
\dot{V}_{ix h} \\
\dot{V}_{iy h} \\
\dot{V}_{iz h}
\end{bmatrix}.
\] (66)
The final step consists of combining Eqs. (63) and (66). This leads to the following scalar form of the dynamical equations:

\[
V_i = -DN_7/m + (DN_6/m)\cos\gamma_i \cos x_i \\
\quad - (QN_4N_6/m)\sin\mu \sin\gamma_i + (QN_6/mN_2)\cos\mu \cos\gamma_i \sin x_i \\
\quad - (LN_4N_6/m)\cos\mu \sin\gamma_i - (LN_6/mN_2)\sin\mu \cos\gamma_i \sin x_i \\
+ f\cos\gamma_i \sin x_i - g\sin\gamma_i, 
\]  

(67a)

\[
V_i \dot{\gamma}_i = - (DN_6/m)\sin\gamma_i \cos x_i \\
\quad + (Q/mN_2)\sin\mu \cos\gamma_i - (QN_6N_7/mN_2)\sin\mu \sin^2\gamma_i \cos x_i - (QN_6/mN_2)\cos\mu \sin\gamma_i \sin x_i \\
\quad + (L/mN_2)\cos\mu \cos\gamma_i - (LN_6N_7/mN_2)\cos\mu \sin^2\gamma_i \cos x_i + (LN_6/mN_2)\sin\mu \sin\gamma_i \sin x_i \\
\quad - f\sin\gamma_i \sin x_i + (v_i^2/r_i - g)\cos\gamma_i, 
\]  

(67b)

\[
V_i \cos\gamma_i \dot{x}_i = -(DN_6/m)\sin x_i \\
\quad - (QN_6N_7/mN_2)\sin\mu \sin\gamma_i \sin x_i - (QN_7/mN_2)\cos\mu \cos\gamma_i + (QN_6/mN_2)\cos\mu \cos x_i \\
\quad - (LN_6N_7/mN_2)\cos\mu \sin\gamma_i \sin x_i + (LN_7/mN_2)\sin\mu \cos\gamma_i - (LN_6/mN_2)\sin\mu \cos x_i \\
\quad + f\cos x_i + (v_i^2/r_i)\cos^2\gamma_i \cos x_i \tan\phi_i, 
\]  

(67c)

where

\[
N_6 = \omega_i \cos\phi_i / N_1, 
\]  

(68a)

\[
N_7 = V_i / N_1. 
\]  

(68b)

Equations (67) can be rewritten as
\[ V_i = -\frac{D \gamma_7}{m} + \frac{(D \gamma_6/m)}{m} \cos \gamma_i \cos x_i \]
- \( (Q \gamma_4 N_6/m) \sin \mu \sin \gamma_i + (Q \gamma_6/m N_2) \cos \mu \cos \gamma_i \sin x_i \)
- \( (L \gamma_4 N_6/m) \cos \mu \sin \gamma_i - (L \gamma_6/m N_2) \sin \mu \cos \gamma_i \sin x_i \)
+ \( f \cos \gamma_i \sin x_i - g \sin \gamma_i \), \hspace{1cm} (69a)

\[ \dot{\gamma}_i = -\frac{D \gamma_6}{m V_i} \sin \gamma_i \cos x_i \]
+ \( (Q/m N_2 V_i) \sin \mu \cos \gamma_i - (Q \gamma_6 N_7/m N_2 V_i) \sin \mu \sin^2 \gamma_i \cos x_i - (Q \gamma_6/m N_2 V_i) \cos \mu \sin \gamma_i \sin x_i \)
+ \( (L/m N_2 V_i) \cos \mu \cos \gamma_i - (L \gamma_6 N_7/m N_2 V_i) \cos \mu \sin^2 \gamma_i \cos x_i + (L \gamma_6/m N_2 V_i) \sin \mu \sin \gamma_i \sin x_i \)
- \( (f/V_i) \sin \gamma_i \sin x_i + (V_i/r_i - g/V_i) \cos \gamma_i \), \hspace{1cm} (69b)

\[ \dot{x}_i = -\frac{D \gamma_6}{m V_i} \sin x_i / \cos \gamma_i \]
- \( (Q \gamma_6 N_7/m N_2 V_i) \sin \mu \tan \gamma_i \sin x_i - (Q \gamma_7/m N_2 V_i) \cos \mu + (Q \gamma_6/m N_2 V_i) \cos \mu \cos x_i / \cos \gamma_i \)
- \( (L \gamma_6 N_7/m N_2 V_i) \cos \mu \tan \gamma_i \sin x_i + (L \gamma_7/m N_2 V_i) \sin \mu - (L \gamma_6/m N_2 V_i) \sin \mu \cos x_i / \cos \gamma_i \)
+ \( (f/V_i) \cos x_i / \cos \gamma_i + (V_i/r_i) \cos \gamma_i \cos x_i \tan \phi_i \). \hspace{1cm} (69c)
8. **Summary of Results**

In this report, we have derived the equations of motion of a spacecraft under the following assumptions: (a) the spacecraft is a particle of constant mass; (b) the Earth is rotating with constant angular velocity; (c) the atmosphere is at rest with respect to the Earth; (d) the Earth is an oblate planet, and the gravitational potential depends on both the radial distance and the latitude; however, harmonics of order higher than four are ignored.

An inertial axes system has been used, and the following kinematical and dynamical equations have been obtained:

\[ \dot{\theta}_i = V_i \cos \gamma_i \cos x_i / r_i \cos \phi_i, \quad (70a) \]
\[ \dot{\phi}_i = -V_i \cos \phi_i \sin x_i / r_i, \quad (70b) \]
\[ \dot{r}_i = V_i \sin \gamma_i, \quad (70c) \]

and

\[ \dot{V}_i = -D_{N_7}/m + (D_{N_6}/m) \cos \gamma_i \cos x_i \]
\[ - (Q_{N_6}/m) \sin \mu \sin \gamma_i + (Q_{N_6}/m) \cos \mu \cos \gamma_i \sin x_i \]
\[ - (L_{N_6}/m) \cos \mu \sin \gamma_i - (L_{N_6}/m) \sin \mu \cos \gamma_i \sin x_i \]
\[ + f \cos \gamma_i \sin x_i - g \sin \gamma_i, \quad (71a) \]
\[ \dot{\gamma}_i = -(D_{N_6}/mV_i) \sin \gamma_i \cos x_i \]
\[ + (Q/mN_2V_i) \sin \mu \cos \gamma_i - (Q_{N_7}/mN_2V_i) \sin \mu \sin^2 \gamma_i \cos x_i - (Q_{N_6}/mN_2V_i) \cos \mu \sin \gamma_i \sin x_i \]
\[ + (L/mN_2V_i) \cos \mu \cos \gamma_i - (L_{N_7}/mN_2V_i) \cos \mu \sin^2 \gamma_i \cos x_i + (L_{N_6}/mN_2V_i) \sin \mu \sin \gamma_i \sin x_i \]
\[ - (f/V_i) \sin \gamma_i \sin x_i + (V_i/r_i - g/V_i) \cos \gamma_i, \quad (71b) \]
\[ x_i = -\frac{(DN_6/mV_i)\sin\chi_i}{\cos\gamma_i} \]

\[ - \frac{(QN_7/mN_2V_i)\sin\mu \tan\gamma_i}{\sin x_i} - \frac{(QN_6/mN_2V_i)\cos\mu}{\cos x_i} \]

\[ - \frac{\cos\mu \tan \gamma_i}{\sin x_i} + \frac{(LN_7/mN_2V_i)\sin\mu}{\cos x_i} \]

\[ + \frac{f/V_i}{\cos x_i} \tan \phi_i + \frac{(V_i/r_1)}{\cos \gamma_i} \cos x_i \tan \phi_i. \] 

(71c)

In the dynamical equations (71), the quantities \( N_i \) depend on the inertial velocity elements \( V_i, \gamma_i, x_i \) and are given by

\[ N_1 = \sqrt{M_1^2 + M_2^2 + M_3^2}, \] 

(72a)

\[ N_2 = \sqrt{M_1^2 + M_2^2} / \sqrt{M_1^2 + M_2^2 + M_3^2}, \] 

(72b)

\[ N_3 = M_3 / \sqrt{M_1^2 + M_2^2 + M_3^2}, \] 

(72c)

\[ N_4 = M_1 / \sqrt{M_1^2 + M_2^2}, \] 

(72d)

\[ N_5 = M_2 / \sqrt{M_1^2 + M_2^2}, \] 

(72e)

\[ N_6 = \omega r_1 \cos \phi_i / N_1, \] 

(72f)

\[ N_7 = V_i / N_1, \] 

(72g)

with

\[ M_1 = V_i \cos \gamma_i \cos x_i - \omega r_1 \cos \phi_i, \] 

(73a)

\[ M_2 = V_i \cos \gamma_i \sin x_i, \] 

(73b)

\[ M_3 = V_i \sin \gamma_i. \] 

(73c)
8.1. **Aerodynamic Force.** In Eqs. (71), the drag, the side force, and the lift are given by

\[
D = \frac{1}{2} C_D \rho S (V_i^2 - 2 \omega r_i V_i \cos \gamma_i \cos \phi_i + \omega^2 r_i^2 \cos^2 \phi_i),
\]

\[
Q = \frac{1}{2} C_Q \rho S (V_i^2 - 2 \omega r_i V_i \cos \gamma_i \cos \phi_i + \omega^2 r_i^2 \cos^2 \phi_i),
\]

\[
L = \frac{1}{2} C_L \rho S (V_i^2 - 2 \omega r_i V_i \cos \gamma_i \cos \phi_i + \omega^2 r_i^2 \cos^2 \phi_i),
\]

where \( C_D \) is the drag coefficient, \( C_Q \) is the side force coefficient, \( C_L \) is the lift coefficient, \( \rho \) is the air density, and \( S \) is a reference surface area. In turn, the aerodynamic coefficients are functions of the form

\[
C_D = C_D(\alpha, \sigma, M, R_e),
\]

\[
C_Q = C_Q(\alpha, \sigma, M, R_e),
\]

\[
C_L = C_L(\alpha, \sigma, M, R_e),
\]

where \( \alpha \) is the angle of attack, \( \sigma \) is the sideslip angle, \( M \) is the Mach number, and \( R_e \) is the Reynolds number.

8.2. **Gravitational Force.** In Eqs. (71), the radial component and the latitudinal component of the acceleration of gravity are given by

\[
g = \left( \mu_e / r_i^2 \right) \left[ 1 + J_2 (r_e / r_i)^2 G_2 + J_3 (r_e / r_i)^3 G_3 + J_4 (r_e / r_i)^4 G_4 \right],
\]

\[
G_2 = 3/2 - (9/2) \sin^2 \phi_i,
\]

\[
G_3 = 6 \sin \phi_i - 10 \sin^3 \phi_i,
\]

\[
G_4 = -15/8 + (150/8) \sin^2 \phi_i - (175/8) \sin^4 \phi_i,
\]
and

\[ f = \left( \frac{u_e}{r_i^2} \right) \left\{ J_2 \left( \frac{r_e}{r_i} \right)^2 F_2 + J_3 \left( \frac{r_e}{r_i} \right)^3 F_3 + J_4 \left( \frac{r_e}{r_i} \right)^4 F_4 \right\}, \quad (77a) \]

\[ F_2 = 3 \sin \phi_i \cos \phi_i, \quad (77b) \]

\[ F_3 = -(3/2) \cos \phi_i + (15/2) \sin^2 \phi_i \cos \phi_i, \quad (77c) \]

\[ F_4 = -(15/2) \sin \phi_i \cos \phi_i + (35/2) \sin^3 \phi_i \cos \phi_i. \quad (77d) \]

8.3. Physical Constants. The major physical constants appearing in the system (70)-(77) have the following values:

\[ \omega = 0.729211595 \text{ E-04 rad/sec}, \quad (78a) \]

\[ u_e = 0.39860064 \text{ E+15 m}^3/\text{sec}^2, \quad (78b) \]

\[ J_2 = 0.10826271 \text{ E-02}, \quad (78c) \]

\[ J_3 = -0.25358868 \text{ E-05}, \quad (78d) \]

\[ J_4 = -0.1624618 \text{ E-05}, \quad (78e) \]

\[ r_e = 0.6378164 \text{ E+07 m}, \quad (78f) \]

\[ r_p = 0.6356755 \text{ E+07 m}. \quad (78g) \]

Here, \( \omega \) is the Earth's angular velocity; \( u_e \) is the Earth's gravitational constant; \( J_2, J_3, J_4 \) are the characteristic constants of the Earth's gravitational field; \( r_e \) is the Earth's equatorial radius; and \( r_p \) is the Earth's polar radius. Note that the Earth's sea-level radius \( r_{se} \) varies with the latitude \( \phi_i \) according to the relation
\[ r_{s\omega} = \frac{1}{2}(r_e + r_p) + \frac{1}{2}(r_e - r_p)\cos(2\phi_i). \] (79)

8.4. **Spacecraft Data.** For the AFE vehicle, it is assumed that

\[ m = 0.16782918 \times 10^4 \text{ kg}, \] (80a)
\[ S = 0.14314 \times 10^2 \text{ m}^2, \] (80b)
\[ \alpha = 0.17000 \times 10^2 \text{ deg}, \] (80c)
\[ C_L = -0.370696 \times 10^0, \] (80d)
\[ C_D = 0.131452 \times 10^1. \] (80e)

Here, \( m \) is the spacecraft mass at atmospheric entry; \( S \) is the reference surface area; \( \alpha \) is the angle of attack; \( C_L \) is the lift coefficient; and \( C_D \) is the drag coefficient. Note that, for the aeroassisted flight experiment, the angle of attack is kept constant; the aerodynamic coefficients are assumed to be independent of the Mach number and the Reynolds number; and the spacecraft is controlled via the angle of bank.
9. **Transformation Relations**

In this section, we summarize the transformation relations which allow one to pass from (i) quantities computed in an Earth-fixed system to (ii) quantities computed in an inertial system, and vice versa.

9.1. Spacecraft Position. Let \( r, \theta, \phi \) denote the spherical coordinates of the spacecraft \( P \) in the Earth-fixed system \( O_{x_e y_e z_e} \). Let \( r_i, \theta_i, \phi_i \) denote the spherical coordinates of the same spacecraft in the inertial system \( O_{x_1 y_1 z_1} \). Assume that the axes of the Earth-fixed system coincide with the axes of the inertial system at time instant \( t = 0 \). Then, the following transformation relations hold:

\[
\begin{align*}
    r_i &= r, \\
    \theta_i &= \theta + \omega t, \\
    \phi_i &= \phi.
\end{align*}
\]

Equations (81) imply the following inverse relations:

\[
\begin{align*}
    r &= r_i, \\
    \theta &= \theta_i - \omega t, \\
    \phi &= \phi_i.
\end{align*}
\]

9.2. Spacecraft Velocity. Let \( V, \gamma, \chi \) denote the velocity modulus, the path inclination, and the heading angle in the Earth-fixed system \( O_{x_e y_e z_e} \). Let \( V_i, \gamma_i, \chi_i \) denote the velocity modulus, the path inclination, and the heading angle in the inertial system \( O_{x_1 y_1 z_1} \). The following transformation relations hold:
\[ V_i = \sqrt{V^2 + 2 \omega_r V \cos \gamma \cos \phi + (\omega_r \cos \phi)^2}, \quad (83a) \]
\[ \tan \gamma_i = \frac{V \sin \gamma}{\sqrt{(V \cos \gamma)^2 + 2 \omega_r V \cos \gamma \cos \phi + (\omega_r \cos \phi)^2}}, \quad (83b) \]
\[ \tan \chi_i = \frac{V \cos \gamma \sin \chi}{(V \cos \gamma \cos \chi + \omega_r \cos \phi)}. \quad (83c) \]

Equations (83) imply the following inverse relations:
\[ V = \sqrt{V_i^2 - 2 \omega_r V_i \cos \gamma_i \cos \phi_i + (\omega_r \cos \phi_i)^2}, \quad (84a) \]
\[ \tan \gamma = \frac{V_i \sin \gamma_i}{\sqrt{(V_i \cos \gamma_i)^2 - 2 \omega_r V_i \cos \gamma_i \cos \phi_i + (\omega_r \cos \phi_i)^2}}, \quad (84b) \]
\[ \tan \chi = \frac{V_i \cos \gamma_i \sin \chi_i}{(V_i \cos \gamma_i \cos \chi_i - \omega_r \cos \phi_i)}. \quad (84c) \]

9.3. Cartesian Coordinates. After the spacecraft position is known in spherical coordinates, the corresponding Cartesian coordinates can be computed. The following transformation relations hold:
\[ x_e = r \cos \theta \cos \phi, \quad (85a) \]
\[ y_e = r \sin \theta \cos \phi, \quad (85b) \]
\[ z_e = r \sin \phi, \quad (85c) \]

and
\[ x_i = r_i \cos \theta_i \cos \phi_i, \quad (86a) \]
\[ y_i = r_i \sin \theta_i \cos \phi_i, \quad (86b) \]
\[ z_i = r_i \sin \phi_i. \quad (86c) \]
9.4. Cartesian Velocity Components. After the spacecraft velocity elements \( V, \gamma, \chi \) or \( V_i, \gamma_i, \chi_i \) are known, the Cartesian velocity components \( x_e, y_e, z_e \) or \( x_i, y_i, z_i \) can be computed. The following transformation relations hold:

\[
\begin{align*}
\dot{x}_e &= -V \sin \theta \cos \gamma \cos \chi + V \cos \theta \sin \phi \cos \gamma \sin \chi + V \cos \theta \cos \phi \sin \gamma, \\
\dot{y}_e &= V \cos \theta \cos \gamma \cos \chi + V \sin \theta \sin \phi \cos \gamma \sin \chi + V \sin \theta \cos \phi \sin \gamma, \\
\dot{z}_e &= -V \cos \phi \cos \gamma \sin \chi + V \sin \phi \sin \gamma,
\end{align*}
\]  

(87a) (87b) (87c)

and

\[
\begin{align*}
\dot{x}_i &= -V_i \sin \theta_i \cos \gamma_i \cos \chi_i + V_i \cos \theta_i \sin \phi_i \cos \gamma_i \sin \chi_i \\
&\quad + V_i \cos \theta_i \cos \phi_i \sin \gamma_i, \\
\dot{y}_i &= V_i \cos \theta_i \cos \gamma_i \cos \chi_i + V_i \sin \theta_i \sin \phi_i \cos \gamma_i \sin \chi_i \\
&\quad + V_i \sin \theta_i \cos \phi_i \sin \gamma_i, \\
\dot{z}_i &= -V_i \cos \phi_i \cos \gamma_i \sin \chi_i + V_i \sin \phi_i \sin \gamma_i.
\end{align*}
\]  

(88a) (88b) (88c)
10. Conclusions

This report is the second of a series dealing with the determination of optimal trajectories for the aeroassisted flight experiment (AFE). The AFE refers to the study of the free flight of an autonomous spacecraft, shuttle-launched and shuttle-recovered. Its purpose is to gather atmospheric entry environmental data for use in designing aeroassisted orbital transfer vehicles (AOTV).

It is assumed that: the spacecraft is a particle of constant mass; the Earth is rotating with constant angular velocity; the Earth is an oblate planet, and the gravitational potential depends on both the radial distance and the latitude; however, harmonics of order higher than four are ignored; the atmosphere is at rest with respect to the Earth.

Under the above assumptions, the equations of motion for hypervelocity atmospheric flight (which can be used not only for AFE problems, but also for AOT problems and space shuttle problems) are derived in an inertial system. Transformation relations are supplied which allow one to pass from quantities computed in an inertial system to quantities computed in an Earth-fixed system, and vice versa.
References


