Ionisation of Atomic Hydrogen by Positron Impact

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1 Experimental Setup

With our crossed beam apparatus [1] we measure the relative impact-ionisation cross section of atomic hydrogen by positron impact. A layout of the scattering region is given in fig. 1.

2 Data Taking

Because of the H₂-molecules in the atomic beam and the residual gas both H⁺ and H²⁺ ions are produced (see table 1). Positrons ionize atoms or molecules through two different processes: (a) impact ionisation: leaving an ion, a free electron and the projectile, or (b) positronium formation: leaving an ion and a positronium ‘atom’. Impact ionisation leads to time-correlated signals on both detectors. The positron and the ion signals are processed by a time-to-pulse-height converter. If the ion follows the projectile in less than 4 µs this event is stored in a multichannel analyzer, H⁺ and H²⁺ ions can be distinguished by their different flight times. Because of a relative high background on the ion-detector (∼ 10 s⁻¹) produced by La-photons we analyze time-correlated signals only. In our set-up the detection of atomic ions produced via dissociative ionisation plays a minor role (less than 1% of the total ionisation signals). As long as the detection probabilities for the correlated positron-ion pairs and the overlap of projectile and target beam are unknown we can only determine relative ion-formation probabilities. By switching the polarity of the optical elements for the primary beam transport we can also measure the respective values for electron impact. To obtain σ⁻¹⁺(H₁) and σ⁻¹⁺(H₂) the ion-formation probabilities are normalized at 100 eV to the data of Shah et al. [2] and Rapp, Englander-Golden [3], respectively. The same normalization factors are also used for the normalization of the positron impact ionisation data on H₁ and H₂, respectively. Fortunately we can check this procedure by comparing our e⁺-H₂ results with those obtained in a different apparatus [4]. The energy of the projectiles can be varied between 10 eV and 600 eV; the intensities are in the order of 3000 s⁻¹. The observed ion-formation probabilities are rather low (∼< 5 × 10⁻⁶ ions/projectile), so automated around-the-clock measurements were performed for more than 100 days to obtain the presented data.

3 Results

In figure 2 the first measurements on the ionisation of atomic hydrogen by positron impact are shown.
Below 400 eV the e⁺-H₁ impact ionisation cross section (fig. 2) is significantly higher than the respective e⁻ cross section, at 50 eV about a factor of two. Qualitatively the cross section shows the shape predicted by the theoretical estimates, but at maximum all calculated values are too low. As mentioned above we measure the positron impact ionisation cross section for molecular hydrogen simultaneously (fig. 3). Our values agree excellently with those from Fromme et al. [4]. In order to check the performance of the apparatus we measure the number of time-correlated H₂⁺ and H₂⁺-ions for electron impact ionisation. The comparisons (fig. 4, 5) show a good agreement.

Note: We detect only those projectiles that are scattered into a angular sphere of ± 30°. This may cause two errors in the detection of the correlated ion-projectile-pairs, depending on (a) the energy or, (b) the charge of the scattering projectile. To correct for the effect of error (a) we will form the ratio of the ion-formation probabilities produced by e⁺ and e⁻ and multiplied them with the e⁻ cross sections from literature. So far there is only incomplete information on the effect of error (b).

### Table 1: The most important processes for the positron impact ionisation of atomic and molecular hydrogen

<table>
<thead>
<tr>
<th>Process</th>
<th>Cross Section</th>
<th>Threshold</th>
</tr>
</thead>
<tbody>
<tr>
<td>e⁺ + H₁ → Ps + H₁⁺</td>
<td>σₚₛ(H₁)</td>
<td>6.8 eV</td>
</tr>
<tr>
<td>e⁺ + H₁ → e⁺ + e⁻ + H⁺</td>
<td>σₐₗₒₙ(H₁)</td>
<td>13.6 eV</td>
</tr>
<tr>
<td>e⁺ + H₂ → Ps + H₂⁺</td>
<td>σₚₛ(H₂)</td>
<td>8.6 eV</td>
</tr>
<tr>
<td>e⁺ + H₂ → e⁺ + e⁻ + H₃⁺</td>
<td>σₐₗₒₙ(H₂)</td>
<td>15.4 eV</td>
</tr>
<tr>
<td>e⁺ + H₂ → Ps + H₁ + H⁺⁺</td>
<td>σₚₛ,Diss(H₂)</td>
<td>11.1 eV</td>
</tr>
<tr>
<td>e⁺ + H₂ → e⁺ + e⁻ + H₁ + H⁺</td>
<td>σₐₗₒₙ,Diss(H₂)</td>
<td>17.9 eV</td>
</tr>
</tbody>
</table>

Figure 2: Positron impact ionisation cross sections: I. σ⁺ₐₗₒₙ(H₁).

Figure 3: Positron impact ionisation cross sections: II. σ⁺ₐₗₒₙ(H₂).

4 Future / Acknowledgements

In a collaboration of members of the Brookhaven National Laboratory, the City College of CUNY and the University of Bielefeld this experiment will be continued at BNL. The apparatus will be modified and it will be possible to achieve more precise data with higher positron intensities, especially at ener-
gies near the threshold. The work at Bielefeld was supported by the Deutsche Forschungsgemeinschaft (DFG). Our Brookhaven project is being supported by the Bundesministerium für Forschung und Technologie (BMFT).

References

    Atomic Physics with Positrons


    J. Chem. Phys. 43 (1965), 1464 - 1479


[8] Wetmore, A.E., Olson, R.E.
    Phys. Rev. 34 (1986), 2822 - 2829