Reduced Dimer Production in Solar-Simulator-Pumped Continuous Wave Iodine Lasers Based on Model Simulations and Scaling and Pumping Studies

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Abstract

A numerical rate equation model for a continuous wave iodine laser with longitudinally flowing gaseous lasant is validated by approximating two experiments that compare the perfluoroalkyl iodine lasants n-C₃F₇I and t-C₄F₉I. The salient feature of the simulations is that the production rate of the dimer (C₄F₉)₂ is reduced by one order of magnitude relative to the dimer (C₃F₇)₂. The model is then used to investigate the kinetic effects of this reduced dimer production—especially how it improves output power. Related parametric and scaling studies are also presented. When dimer production is reduced, more monomer radicals (t-C₄F₉) are available to combine with iodine ions, thus enhancing depletion of the laser lower level and reducing buildup of the principal quencher, molecular iodine. Fewer iodine molecules result in fewer downward transitions from quenching and more transitions from stimulated emission of lasing photons. Enhanced depletion of the lower level reduces the absorption of lasing photons. The combined result is more lasing photons and proportionally increased output power.
\( Q \)  
lasant heating from pump, W \( \cdot \) m\(^{-3}\)

\( q_i \)  
quenching rate coefficient, cm\(^3\) \( \cdot \) sec\(^{-1}\)

\( R \)  
gas constant, \( 8.314 \) J \( \cdot \) mol\(^{-1}\) \( \cdot \) K\(^{-1}\)

\( R^* \)  
\[ R^* = \frac{R}{M \times 10^{-3}} \] J \( \cdot \) kg\(^{-1}\) \( \cdot \) K\(^{-1}\)

\( R \)  
radical n-C\(_3\)F\(_7\), t-C\(_4\)F\(_9\), or f-C\(_3\)F\(_7\)

\( R_a \)  
reflectivity of highly reflecting mirror, 1.0

\( R_p \)  
reflectivity of output mirror, 0.75

\( r \)  
radial coordinate in cylindrical frame, m

\( r_N \)  
normalized radial coordinate, \( r / r_i \)

\( r_i \)  
inner radius of laser tube, m

\( S \)  
distance along inner perimeter of laser tube, m

\( s \)  
arc length, m

\( SC \)  
solar constant, \( 1.35 \) kW \( \cdot \) m\(^{-2}\)

\( T \)  
lasant temperature, K

\( w \)  
lasant flow speed

\( Z \)  
avxial distance downstream of pump entrance

\( Z_{G+} \)  
avxial distance on which inversion density \([I^*] - [I] / 2 > 0\), positive gain length, cm

\( Z_L \)  
active length of laser tube, 33 cm

\( Z_p \)  
length of pump, 15 cm

\( z \)  
normalized axial distance downstream of pump entrance, \( Z / Z_L \)

\( z_{G+} \)  
normalized positive gain length, \( Z_{G+} / Z_L \)

\( z_p \)  
normalized pump length, \( Z_p / Z_L = 0.45 \)

\( \alpha \)  
constant for \( C_v \) (eq. (26)), J \( \cdot \) mol\(^{-1}\) \( \cdot \) K\(^{-1}\)

\( \beta \)  
constant for \( C_v \) (eq. (26)), K\(^{-1}\)

\( \gamma \)  
distance along light ray, m

\( \delta \)  
absorption length for pumping photons, m

\( \eta \)  
density of lasant gas, kg \( \cdot \) m\(^{-3}\)

\( \theta \)  
angle, rad

\( \nu_L \)  
frequency of laser beam, Hz

\( \nu_p \)  
pump frequency, Hz

\( \xi_1 \)  
photodissociation rate of [RI], sec\(^{-1}\)

\( \xi_2 \)  
photodissociation rate of [I\(_2\)], sec\(^{-1}\)

\( \rho \)  
number density of lasing photons, \( \rho_+ + \rho_-, \) photons \( \cdot \) cm\(^{-3}\)

\( \rho_\pm \)  
number density of lasing photons moving along Z-axis, photons \( \cdot \) cm\(^{-3}\)

\( \rho_p \)  
number density of pumping photons, photons \( \cdot \) cm\(^{-3}\)

\( \rho_{pN} \)  
normalized pumping photon density

\( \sigma \)  
cross section for laser-beam absorption and stimulated emission, cm\(^2\)

\( \psi \)  
angle, rad

\([ \cdot ] \)  
number density, particles \( \cdot \) cm\(^{-3}\)

\( \langle \cdot \rangle \)  
average over cross section of laser tube

Subscript:

\( 0 \)  
value at \( Z = 0 \)

**Model Description**

**Laser Geometry**

The CW laser used in references 14 and 15 is shown schematically in figure 1. The Z-axis of the 1-D mathematical model is parallel to the optical axis of the laser cavity and points in the direction of the lasant flow. The origin is located where the lasant enters the elliptical pump chamber. Upstream of this point, the lasant is undissociated and does not interact with the laser beam, provided the lasant is free of absorbing or scattering impurities. The pumping region spans the distance \( 0 < Z < Z_p \), where \( Z_p = 15 \) cm. We assume that the incident pumping radiation at the laser tube is axisymmetric. Measurements of the actual incident pumping radiation are given in appendix D. Downstream of the pump, the lasant has a nonzero inversion density and continues to interact with the laser beam until the lasant is withdrawn at the end of the tube \( Z_L \), where \( Z = Z_L = 33 \) cm. In the computation, the active length of the laser tube \((0 < Z \leq Z_L)\) is normalized to unity \((0 < z \leq z_{G+})\), and the pump spans the normalized distance \( 0 < z < z_p \), where \( z_p = 0.45 \).

**Photochemical Reactions**

The kinetic reactions included in the laser model are as follows.
Introduction

A solar power station advantageously placed in space could beam power to other spacecraft and to planetary surfaces, including the surface of the Earth, as discussed in references 1-3. Experiments related to this concept include tests of solar-simulator-pumped iodine lasers, as discussed in references 4-7. These experiments are supported by modeling efforts reported in references 8-13. An important finding is that the gaseous perfluoroalkyl iodine lasants n-C₃F₇I and t-C₄F₉I have markedly different production rates for the dimers (C₃F₇)₂ and (C₄F₉)₂. Our modeling effort is devoted principally to understanding the effects of this differing dimer production on laser performance.

Lee (ref. 14) and Lee et al. (ref. 15) present two experimental comparisons of the lasants n-C₃F₇I and t-C₄F₉I when flowed longitudinally in a continuous wave (CW) laser. Laser output power is measured in the first comparison as the lasant flow speed is varied and in the second comparison as the intensity of the solar-simulator pump is varied. In both comparisons, the output power \( P_{out} \) for t-C₄F₉I is found to be about three times greater than that for n-C₃F₇I. This increase in \( P_{out} \) is not explained by an increased utilization of the pump spectrum. For the solar simulator used, the pump spectrum utilization is only 20 percent greater for t-C₄F₉I than for n-C₃F₇I.

The dimer density is represented by \([R₂]\), where \( R \) represents either of the perfluoroalkyl radicals n-C₃F₇ or t-C₄F₉ and brackets denote the number density. Ershov, Zaleskij, and Sokolov (ref. 16) have shown experimentally that the rate of \([R₂]\) production is much less for t-C₄F₉I than for n-C₃F₇I, although a numerical value for the ratio of these production rates is not given. Lee (ref. 14) notes that the reduced \([R₂]\) production increases t-C₄F₉I recyclability. He also speculates that the reduced \([R₂]\) production for t-C₄F₉I would make more monomer radicals \([R]\) available to combine with iodine atoms \([I]\). Consequently, more of the laser lower level \([I]\) would be depleted and less \([I₂]\), the principal quencher of excited iodine \([I^*]\), would be formed. This reduction of \([I₂]\) would also reduce the lasant flow speed.

Our purpose is to examine these speculations and especially to determine the kinetic effects of reduced \([R₂]\) production on \( P_{out} \). Our approach is to use a one-dimensional (1-D) numerical rate equation model. This model is described in the next section. The model is tuned by using the data given in references 14 and 15; a discussion of that process follows the model description. Thereafter, general properties of the tuned model and its solutions are given. The solutions indicate that the experimental power curves obtained in references 14 and 15 could be improved by optimizing the lasant flow speed \( w \) so that molecular iodine does not build up within the pump region of the laser. This optimization is based on parametric and scaling studies that are presented in appendixes A-D. The flow speed optimization is modeled in the section “Diagnostic Plots,” and the n-C₃F₇I and t-C₄F₉I lasants are again compared. A fictitious lasant (identical to n-C₃F₇I except for a reduced \([R₂]\) production rate) is modeled in the section, “Optimization of Laser Performance.” The purpose is to isolate the kinetic effects of reduced \([R₂]\) production. Concluding remarks are given next and the appendixes follow.

Symbols

\( B \)  
Brightness of pump lamp image, \( W \cdot m^{-2} \cdot rad^{-1} \)

\( C \)  
Carbon atom

\( C_1, C_2 \)  
Constants of integration

\( C_p \)  
Specific heat at constant pressure, \( J \cdot mol^{-1} \cdot K^{-1} \)

\( C_p^* \)  
\( \frac{C_p}{M \times 10^{-3}} \cdot J \cdot kg^{-1} \cdot K^{-1} \)

\( C_v \)  
Specific heat at constant volume, \( J \cdot mol^{-1} \cdot K^{-1} \)

\( c \)  
Speed of light in vacuum

\( c_i \)  
Three-body, reaction-rate coefficients, cm\(^3\) sec\(^{-1}\)

\( e \)  
Energy density of pump radiation, \( J \cdot m^{-3} \)

\( F \)  
Incident flux of pump radiation, W/m\(^2\)

\( F \)  
Fluorine atom

\( f-C₃F₇I \)  
Fictitious lasant identical to n-C₃F₇I, except that \( k_3 \) and \( k_4 \) are reduced by factor of 0.1

\( h \)  
Planck’s constant, \( 6.626 \times 10^{-34} \cdot J \cdot sec \)

\( I \)  
Iodine atom

\( I_p \)  
Intensity of pump radiation, SC

\( I^* \)  
Iodine atom in excited state

\( I₂ \)  
Molecular iodine

\( k_i \)  
Two-body reaction rate coefficients, cm\(^{-3}\) sec\(^{-1}\)

\( M \)  
Molecular weight, g mo1\(^{-1}\)

\( P_{out} \)  
Laser output power, W

\( p \)  
Lasant pressure, Pa

\( p_{out} \)  
Laser output power density, W cm\(^{-2}\)
\[
d\frac{d(w[I^*])}{dZ} = \xi_1[R1] + 0.51\xi_2[I_2] - k_1[R1][I^*] \\
- q_2[I_2][I^*] - c\sigma_p\left([I^*] - \frac{1}{2}[I]\right) \\
- q_3[R][I^*] - q_4[R_1][I^*] \\
- q_5[I^*][I] - k_2[R1][I^*] \\
- c_1[R1][I^*][I] - q_1[R1][I^*] \tag{6e}
\]

where \(c\) is the speed of light in vacuum and \(\rho\) is the number density of lasing photons.

The latter density is given by

\[
\rho = \rho_+ + \rho_-
\tag{7}
\]

where the symbols + and − indicate the direction of photon motion along the Z-axis. For steady-state CW operation, these photon densities satisfy the equations

\[
d\frac{d\rho_+}{dZ} = \rho_+\sigma \left([I^*] - \frac{1}{2}[I]\right) \tag{8a}
\]

\[
d\frac{d\rho_-}{dZ} = -\rho_-\sigma \left([I^*] - \frac{1}{2}[I]\right) \tag{8b}
\]

and the boundary condition at \(Z = 0\)

\[
\rho_+(0) = R_a\rho_-(0) \tag{9a}
\]

At \(Z = Z_L\)

\[
\rho_-(Z_L) = R_b\rho_+(Z_L) \tag{9b}
\]

where \(R_a\) and \(R_b\) are the mirror reflectivities at \(Z = 0\) and \(Z = Z_L\), respectively. The quantity \(\rho_-(Z)\) may be eliminated from the formulation because equations (8a) and (8b) are satisfied by

\[
\rho_+(Z)\rho_-(Z) = \text{Constant} \tag{10}
\]

The boundary conditions (eqs. (9a) and (9b)) then give

\[
\rho_-(Z) = \frac{\rho_+(0)^2}{R_a\rho_+(Z)} \tag{11a}
\]

and from equation (7)

\[
\rho_+(Z) = \rho_+(0) + \frac{\rho_+(0)^2}{R_a\rho_+(Z)} \tag{11b}
\]

The boundary conditions (eqs. (9a) and (9b)) also reduce to

\[
\rho_+(Z_L) = \rho_+(0)(R_aR_b)^{-1/2} \tag{12}
\]

The formulation is now complete in terms of \(\rho_+(Z)\), and \(\rho_-(Z)\) is determined from equation (11a). The output power density \(P_{\text{out}}\) in W cm\(^{-2}\) is given by

\[
P_{\text{out}} = \rho_+(Z_L)(1-R_b) \tag{13}
\]

Compressible Fluid Dynamics

Part of the incident pump power is dissipated by heat. The resulting dependence on \(Z\) (m) of temperature \(T\) (K), pressure \(p\) (Pa), and speed \(w\) (m sec\(^{-1}\)) is determined approximately from the 1-D, steady-state fluid dynamic equations for an inviscid, nonconducting gas (ref. 22) as follows.

**Continuity equation.** Under all these conditions the continuity equation becomes

\[
d\frac{d\rho}{dZ} = 0 \tag{14}
\]

where \(\eta\) is the density (kg m\(^{-3}\)). Hence,

\[
\eta = \frac{C_1}{w} \tag{15}
\]

where the constant \(C_1\) is given by

\[
C_1 = \eta_0w_0 \tag{16}
\]

where \(\eta_0\) is the density and \(w_0\) is the speed at \(Z = 0\).

**Momentum equation.** The corresponding momentum equation is

\[
\frac{dw}{dZ} = \frac{1}{\eta}\frac{dp}{dZ} \tag{17a}
\]
**Photodissociation reactions.** The parent molecule RI is irradiated in the pump region, and the following reaction occurs:

$$RI + h\nu_p \rightarrow R + I^*$$

(1a)

where \(h\) is Planck's constant, \(h\nu_p\) is the energy of pumping photons, \(\xi_1\) is the photodissociation rate of RI, and \(I^* = 2P_{1/2}\), which is the laser upper level. A sequence of reactions produces \(I_2\), which can also be photodissociated in the pump according to the reactions

$$I_2 + h\nu_p \rightarrow \begin{cases} I + I^* & (51\text{-percent probability}) \\ 2I & (49\text{-percent probability}) \end{cases}$$

(1b)

where \(\xi_2\) is the photodissociation rate of \(I_2\), and \(I = 2P_{3/2}\), which is the laser lower level. Photodissociation rates for \(\xi_1\) and \(\xi_2\) are listed in table I with other laser rate coefficients. Data in this table were obtained from references 17–21.

**Absorption and stimulated emission reactions.** Photons in the laser beam can be absorbed by I atoms or undergo stimulated emission by excited \(I^*\) atoms according to the reactions

$$I^* + h\nu_L \leftrightarrow I + 2h\nu_L$$

(1c)

where \(h\nu_L\) is the energy of the lasing photons and \(\sigma\) is the cross section for absorption and stimulated emission.

**Two-body reactions.**

$$\begin{align*}
R + I^* & \rightarrow RI \\
R + I & \rightarrow R I \\
R + R & \rightarrow R_2 \\
R + RI & \rightarrow R_2 + I \\
R + I_2 & \rightarrow RI + I \\
RI + I^* & \rightarrow I_2 + R \\
I + RI & \rightarrow I_2 + R \\
\end{align*}$$

(2)

where \(k_i\) are two-body reaction-rate coefficients.

**Pyrolysis.** The lasant RI and dimer \(R_2\) can also undergo thermal dissociation, especially at high temperatures. Thus,

$$\begin{align*}
RI & \rightarrow R + I \\
R_2 & \rightarrow R + R \\
\end{align*}$$

(3)

**Three-body reactions.**

$$\begin{align*}
I^* + I + RI & \rightarrow I_2 + RI \\
I + I + RI & \rightarrow I_2 + RI \\
I + I_2 & \rightarrow I_2 + I_2 \\
I + I + R_2 & \rightarrow I_2 + R_2 \\
\end{align*}$$

(4)

where \(c_i\) are the three-body reaction-rate coefficients.

**Quenching reactions.** The following reactions quench the excited state of iodine atoms:

$$\begin{align*}
1^* + RI & \rightarrow I + RI \\
I^* + I_2 & \rightarrow I + I_2 \\
I^* + R & \rightarrow I + R \\
I^* + R_2 & \rightarrow I + R_2 \\
I^* + I & \rightarrow I + I \\
\end{align*}$$

(5)

where \(q_i\) are the quenching reaction-rate coefficients.

**Rate Equations**

For steady-state CW operation, the 1-D rate equations are purely functions of \(Z\) as shown:

$$\frac{d}{dZ}(w[RI]) = k_1[R][I^*] + k_2[R][I] + k_3[R][I_2] - k_4[R][I] - k_5[R][RI]$$

(6a)

$$\frac{d}{dZ}(w[R]) = \xi_1[RI] - k_6[R][I] - k_7[R][I_2]$$

(6b)

$$\frac{d}{dZ}(w[R_2]) = k_8[R][I] + k_9[R][RI] + k_{10}[R_2]$$

(6c)

$$\frac{d}{dZ}(w[I_2]) = c_1[RI][I^*][I] + c_2[RI][I]^2 + c_3[I_2][I]^2 - \xi_2[I_2]$$

(6d)
experimental conditions with these most favorable rate coefficients). Therefore, some adjustment of the rate coefficients was necessary for the theoretical study. Most rate coefficients are not measured directly but are inferred semiempirically through a combination of measurements and models. Thus, uncertainties in published values of rate coefficients depend not only on measurement uncertainties but also on model uncertainties. The latter uncertainties often are not evaluated.

Wishing to depart from the published values as little as possible, we performed a sensitivity study with the model. This study showed that under the given experimental conditions, the model was most sensitive to $k_2$, $k_3$, $k_5$, and $q_2$ and the output power was increased by increasing $k_2$ and $k_5$ and by decreasing $k_3$ and $q_2$. Accordingly, we multiplied $k_2$ and $k_5$ and also the pump photodissociation rates $\xi_1$ and $\xi_2$ by a factor of 2.55 and divided $k_3$ and $q_2$ by the same factor to achieve lasing and obtain the n1-model curve of figure 2(a). These adjusted values for n-C$_3$F$_7$I are listed in table I, where temperature dependence is included where known.

The r1-model curve of figure 2(a) was obtained next by using the t-C$_4$F$_9$I data given in table II(a) and the corresponding rate coefficients given in table I. As shown in table I, the principal difference assumed for t-C$_4$F$_9$I is that the production rate of the dimer $[R_2]$ is reduced by a factor of 0.1 (reactions $k_3$ and $k_4$). The other rate coefficients and parameters are taken to be the same as those for n-C$_3$F$_7$I, except for $M$, $\alpha$, $\beta$, and $\xi_1$. Various factors for reduced $[R_2]$ production were also tried, but the factor 0.1 gave the best match, subject to the values assumed for the other rate coefficients of t-C$_4$F$_9$I.

We had originally found that the n2- and r2-model curves almost overlapped. Recalling that chemical analysis of the n-lasant had revealed a 20-percent dimer impurity, as shown in appendix A, we incorporated this feature into the model. This change succeeded in increasing somewhat the vertical separation of the theoretical curves in figure 2(b) without much affecting those in figure 2(a). Chemical analysis of the t-C$_4$F$_9$I lasant (appendix A) indicated an even higher amount of an unknown impurity; however, a 20-percent dimer contaminant for both n-C$_3$F$_7$I and t-C$_4$F$_9$I gave the best match in figure 2. This adjustment completed our tuning of the model.

Diagnostic Plots

Besides $P_{\text{out}}$, the laser model computes and plots the following densities as functions of $z$: the parent molecule $[R]$, the monomer $[R]$, the dimer $[R_2]$, molecular iodine $[I_2]$, excited iodine $[I^*]$, ground-state iodine $[I]$, the inversion density $[I^*] - [I]/2$, and the lasing photon densities $\rho_+$ and $\rho_-$. Also computed and plotted are $T$, $\eta$, $p$, and $w$. The plots of these computed quantities (diagnostic plots) are useful in the analysis of laser performance. Figures 3-6 present diagnostic plots for the four model curves of figure 2.

Diagnostic plots for both end points of the n1-model curve of figure 2(a) are shown in figure 3. These diagnostic plots show that the laser is not operating optimally. In figure 3(a), $[I_2]$ rapidly builds up at the midpoint of the pump region (positive normalized gain length $z_{G+}$) and remains high downstream of this point until, at the exit of the pump region ($z_{p}$), it reaches an even higher plateau. The effects of this $[I_2]$ buildup on $[I^*]$, $[I]$, and the inversion density $[I^*] - [I]/2$ are shown in figures 3(b)-3(d). In particular, the inversion density becomes negative midway through the pump region and remains so downstream of this point, although it becomes somewhat less negative at the exit of the pump region.

These modeling results are consistent with the well-known observations (ref. 19) that $[I_2]$ is a strong quencher of $[I^*]$ and can make the inversion density become negative. This negative inversion density results in a net absorption of lasing photons on the downstream half of the pump section, as shown by the curves for the lasing photon densities $\rho_+$ and $\rho_-$ in figure 3(e). In particular, $\rho_+$ reaches its peak value in the middle of the pump section and thereafter decreases all the way to the output mirror at $z = 1$. Because $P_{\text{out}}$ is proportional to $\rho_+$ at the output mirror, the model laser is not operating optimally. For optimal operation, $\rho_+$ would continue to increase until the flow exits the pump at $z = 0.45$.

Other features of this model run warrant explanation as well. Because $[R]$ reacts much more readily with $[I]$ than with $[I^*]$ to form $[RI]$ (i.e., $k_2 >> k_1$), $[R]$ decreases rapidly as $[I]$ increases at the midpoint of the pump section, as shown in figure 3(f). The distribution of $[R_2]$ is nonzero at $z = 0$ because of the 20-percent initial $[R_2]$ contaminant, as shown in figure 3(g).

Figures 3(h)-3(k) show plots of $T$, $w$, $\eta$, and $p$ (almost constant). These four plots are representative of all the model runs and are subsequently omitted. The moderate increases in $T$ and $w$ from solar-simulator-pump heating generally agree with the experiments in which a water-cooled quartz laser tube was used. (See fig. 1.)

The r1-model curves of figure 4 and one n2-model curve ($I_p = 1100$ SC) of figure 5 also show similar non-optimal laser performance. However, in the other n2-model curve ($I_p = 450$ SC) of figure 5 and in the r2-model curves of figure 6, the buildup of $[I_2]$ occurs only downstream of the pump, which indicates optimal operation.
From equation (16), this equation may be written as

\[
\frac{d}{dZ}(p + C_1 w) = 0 \quad (17b)
\]

Upon integration, we obtain

\[
p = C_2 - C_1 w \quad (18)
\]

where

\[
C_2 = p_0 + C_1 w_0 \quad (19)
\]

**Equation of state.** If the lasant is idealized as a perfect gas, the equation of state is given by

\[
p = \eta R^* T \quad (20)
\]

Here, \( R^* \) (J \cdot kg\(^{-1}\) \cdot K\(^{-1}\)) is given by

\[
R^* = \frac{R}{M \times 10^{-3}} \quad (21)
\]

where the gas constant \( R = 8.314 \) J \cdot mol\(^{-1}\) \cdot K\(^{-1}\), and \( M(\) g \cdot mol\(^{-1}\) ) is the molecular weight of the lasant. Solving equation (20) for \( T \), we obtain

\[
T = \frac{w}{R^*} \left( \frac{C_2}{C_1} - w \right) \quad (22)
\]

From equation (20), \( C_1 \) may also be written

\[
C_1 = \frac{p_0 w_0}{R^* T_0} \quad (23)
\]

where \( p_0, w_0, \) and \( T_0 \) are all measured quantities.

**Energy equation.** The 1-D, steady-state energy equation for an inviscid, nonconducting perfect gas is given by

\[
C_p^* \eta w \frac{dT}{dZ} - w \frac{dp}{dZ} = Q \quad (24)
\]

where \( C_p^* \) (J \cdot kg\(^{-1}\) \cdot K\(^{-1}\)) is the specific heat at constant pressure divided by \( M \times 10^{-3} \), and \( Q \) (W \cdot m\(^{-3}\)) is the heat from the incident pump radiation. From equations (15), (18), and (22), this equation may be rewritten as

\[
\frac{d w}{dZ} = \frac{R^* Q}{R^* C_1 w + C_p^* (C_2 - 2C_1 w)} \quad (25)
\]

**Specific heats.** In equation (25), \( C_p^* \) may be approximated from data in reference 22 as follows. The specific heat at constant volume \( C_v \) (J \cdot mol\(^{-1}\) \cdot K\(^{-1}\)) is closely approximated by

\[
C_v = \alpha \exp \left[ \beta (T - 300) \right] \quad (26)
\]

where \( \alpha \) and \( \beta \) are constant for each lasant as given in table I. The specific heat at constant pressure \( C_p \) in J \cdot mol\(^{-1}\) \cdot K\(^{-1}\) is

\[
C_p = C_v + R \quad (27)
\]

Thus, we have (in J \cdot kg\(^{-1}\) \cdot K\(^{-1}\))

\[
C_p^* = \frac{\alpha}{M \times 10^{-3}} \exp \left\{ \frac{\beta}{R^*} \left[ \left( \frac{C_2}{C_1} - w \right) - 300 \right] \right\} + R^* \quad (28)
\]

and from equation (22)

\[
C_p^* = \frac{\alpha}{M \times 10^{-3}} \exp \left\{ \frac{\beta}{R^*} \left[ \left( \frac{C_2}{C_1} - w \right) - 300 \right] \right\} + R^* \quad (29)
\]

Equation (29), which for a given lasant gas gives \( C_p^* \) as a function of \( w \), is appropriate for substituting back into equation (25).

**Numerical integration procedure.** For a given heating rate \( Q(Z) \) and values for the constants \( C_1 \) and \( C_2 \) (obtained by measuring \( p_0, T_0, \) and \( w_0 \) at \( Z = 0 \)), equation (25) can be numerically integrated to give \( w(Z) \). The other fluid dynamic fields, \( \eta(Z), p(Z), \) and \( T(Z) \) follow from equations (15), (18), and (22), respectively. The flow speed \( w(Z) \) appears explicitly in the rate equations (6a)-(6f). The density \( \eta(Z) \) determines the parent molecule number density \( [R]\), and \( T(Z) \) enters the temperature-dependent rate coefficients.

**Model Tuning**

The rate equation model was tuned by matching as closely as possible the experimental curves for the lasants n-C\(_3\)F\(_7\)I and t-C\(_4\)F\(_9\)I, as given in references 14 and 15. The laboratory data for these curves are given in tables II(a) and II(b), where the original lasant flow units in standard cubic centimeters per second have been converted to meters per second. The wide range of pressure and pump \( I_p \) intensity values makes these data sets appropriate for tuning the model. The best match achieved with the model is shown in figure 2. This match, although imperfect, was obtained with difficulty as described in the discussion that follows.

The data in table II(a) for the n1 experiment curve of figure 2(a) gave a theoretical \( P_{out} = 0 \) W when used with the published rate coefficients for n-C\(_3\)F\(_7\)I that are the most favorable for lasing (i.e., the model did not even reach the threshold for lasing under the given
reduced [R₂] production are highly beneficial to efficient operation of a CW iodine laser.

Concluding Remarks

Two experiments which compare the performance of the lasants n-C₃F₇I and t-C₄F₉I can be approximated by a one-dimensional numerical rate equation model. In this model, the principal difference assumed for the lasants is that the dimer production rate for t-C₄F₉I is one order of magnitude less than for n-C₃F₇I. The model results indicated that laser output power could be increased by optimizing the lasant flow speeds so that the principal quencher, molecular iodine, does not build up within the pump region. This optimization method was based on parametric and scaling studies that are also presented. Such optimized model runs showed that t-C₄F₉I had a larger output power than n-C₃F₇I, although the power increment depended on the operating conditions.

Optimization of the flow speeds improved the basis for comparing n-C₃F₇I and t-C₄F₉I; however, the kinetic effects of reduced dimer production were still not clear because the two lasants had different molecular weights, pump spectrum utilizations, densities, and flow speed. To clarify these kinetic effects and determine how reduced dimer production results in greater output power, we also modeled a fictitious lasant t-C₃F₇I, which differed from n-C₃F₇I only by a reduced dimer production rate. The results of this theoretical comparison confirmed earlier speculations; that is, simply by reducing dimer production, we produce more monomer radicals (C₃F₇). These radicals then combine with iodine atoms to enhance depletion of the laser lower level and reduce the growth of the principal quencher, molecular iodine. Both of these effects tend to increase the lasing photon density and, hence, the output power.

This theoretical study also found that an order-of-magnitude decrease in the dimer production rate halves the lasant flow speed required to prevent the buildup of molecular iodine in the pump region. Also, more lasant molecules are recovered after they pass through the laser tube and less iodine molecules are produced; hence, reduced dimer production improves lasant recyclability. If other properties are equal, reduced dimer production is clearly a desirable feature for the flowing lasant in a continuous wave iodine laser.

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Optimization of Laser Performance

Although we cannot prevent the buildup of \([I_2]\) downstream of the pump, we can prevent it inside the pump region, where it does the most harm, by increasing the lasant flow speed \(w\). This effect of \(w\)—together with other factors that could be used to control the buildup of \([I_2]\)—is discussed in appendix B. In this context, we take the optimal \(w\) to be the minimum flow speed for which the buildup of \([I_2]\) occurs only downstream of the pump.

We found that optimizing the theoretical power curves in figure 2 provides a better way to compare the lasants \(n\)-C\(_3\)F\(_7\)I and \(t\)-C\(_4\)F\(_9\)I. Optimizing the \(n\)-model and \(t\)-1-model curves of figure 2(a) by increasing \(w\) and by eliminating the initial dimer contaminant (i.e., setting \([R_2(0)] = 0\)) gives the increased \(P_{\text{out}}\) shown in figure 7. These curves show that under optimal conditions, \(t\)-C\(_4\)F\(_9\)I gives about twice as much output power as \(n\)-C\(_3\)F\(_7\)I at about half the flow speed. Figures 8 and 9 present diagnostic plots for both end points of each optimized curve shown in figure 7. These plots confirm that the buildup of \([I_2]\) occurs downstream of the pump and that the inversion density \([I^+]-[I]/2\) is positive throughout the pump region.

Similarly, optimization of the \(n\)-2- and \(t\)-2-model curves of figure 2(b) gives the increased \(P_{\text{out}}\) shown in figure 10. Again, \(t\)-C\(_4\)F\(_9\)I has higher output power than \(n\)-C\(_3\)F\(_7\)I, although the results are less dramatic. We conclude, for optimized flows, that \(t\)-C\(_4\)F\(_9\)I is superior to \(n\)-C\(_3\)F\(_7\)I; however, the difference in output power depends on the operating conditions.

Kinetic Analysis of Reduced Dimer Production

Although optimization of the power curves, as shown in figures 7 and 10, facilitates comparison of the lasants \(n\)-C\(_3\)F\(_7\)I and \(t\)-C\(_4\)F\(_9\)I, it does not clarify the kinetic effects of reduced dimer production because the optimized \(t\)-model curves—besides having reduced dimer production relative to the optimized \(n\)-model curves—also have different pump-spectrum utilizations, molecular weights, specific heats, densities, and flow speeds.

To analyze unambiguously the effects of reduced \([R_2]\) production, we considered a fictitious lasant \(f\)-C\(_3\)F\(_7\)I in which the only difference from \(n\)-C\(_3\)F\(_7\)I is reduced values for \(k_3\) and \(k_4\). We start with the \(n\)-1-model optimal curve of figure 7. Because \([I_2]\) builds up only downstream of the pump, the same will be true for the \(f\)-1-model (obtained from the \(n\)-1-model optimal curve by reducing \(k_3\) and \(k_4\) by a factor of 0.1 without changing \(w\) or anything else).

Reduced dimer production (without other changes) gives the increased \(P_{\text{out}}\) shown by the \(f\)-1-model curve relative to the \(n\)-1-model optimal curve in figure 11. Diagnostic plots for the \(f\)-1-model are presented in figure 12 for comparison with the diagnostic plots for the \(n\)-1-model optimum, as presented in figure 8. These diagnostic plots are supplemented by detailed plots of \([I_2]\) and \([RI]\) for the two models in figure 13.

These diagnostic and detailed plots show clearly that a reduced \([R_2]\) production rate decreases \([R_2]\), \([I_2]\), \([I^+]\), and \([I^*-I]/2\) in the pump region and increases \([R]\), \([RI]\), \(p_+\), and \(p_-\). The results confirm the speculations made in reference 1; that is, the decrease in \([R_2]\) and \([I_2]\) and the increase in \([RI]\) improve lasant recyclability. The reduced \([R_2]\) production results in a greater density of monomer radicals \([R]\) in the pump region that can combine with iodine atoms \([I]\) and thus enhance depletion of the laser lower level and reduce the buildup of the principal quencher \([I_2]\). The lower value of \([I_2]\) increases \(p_+\) and \(p_-\) because fewer downward transitions occur by quenching and more by the stimulated emission of lasing photons. The lower value of \([I]\) reduces the absorption of \(p_+\) and \(p_-\). Hence, the reductions in \([I_2]\) and \([I]\) both help increase the output power \(P_{\text{out}}\), which is proportional to \(p_+\) at the output mirror.

After repeating this reduced dimer analysis with the \(n\)-2-model optimal curve of figure 10 and the fictitious lasant \(f\)-C\(_3\)F\(_7\)I, we obtain the increased \(P_{\text{out}}\) shown by the fictitious \(f\)-2-model curve in figure 14. Figures 11 and 14 clearly show that simply by reducing the dimer production rate, we were able to increase output power, although the amount of increase depends on the operating conditions.

Although not shown, the diagnostic plots for the \(n\)-2-model optimum and the \(f\)-2-model are similar to those given in figures 8 and 12 for the \(n\)-1-model optimum and the \(f\)-1-model, respectively. The corresponding detailed plots are given in figure 15, which again shows that the reduction of \([R_2]\) production increases \([RI]\) and decreases \([I_2]\). This second comparison reinforces the previous conclusions about the benefits of reduced \([R_2]\) production.

Finally, we investigate the speculation that reduced \([R_2]\) production allows the lasant flow speed \(w\) to be reduced. In figure 11, the values of \(w_0\) for the fictitious \(f\)-1-model are the same as for the \(n\)-1-model optimum. We now optimize the \(f\)-1-model by reducing \(w_0\) to the minimum values that still prevent the buildup of \([I_2]\) inside the pump. We find that \(w_0\) can be halved and that \(P_{\text{out}}\) is not significantly affected, as shown in figure 16. Diagnostic plots for the \(f\)-1-model optimum (halved \(w_0\)) are shown in figure 17. These plots confirm that \([I_2]\) builds up only downstream of the pump. Although not shown, similar results are obtained when the fictitious \(f\)-2-model of figure 14 is optimized by halving \(w_0\). These effects of
Appendix A

Gas Chromatographic-Mass Spectrometric Analysis of Lasants n-C₃F₇I and t-C₄F₉I

A gas chromatographic-mass spectrometric (GCMS) device was used to examine samples of the two lasants. The findings are summarized below.

Analysis of n-C₃F₇I Sample

No significant difference was noted between the pre- and postlased n-C₃F₇I material. The analysis revealed the following compounds:

- \((C₃F₇)_2\) = 20 percent
- C₂F₃I: Trace amounts
- n-C₃F₇I: > 80 percent
- Unknown perfluorocarbon: Trace amounts
- C₂F₄ClI: Trace amounts

The presence of I₂ was not detected in the GCMS analysis, even though I₂ was visually detectable in the postlased sample. Several explanations can be given for the absence of I₂ in the postlased sample: the I₂ peak was masked by either the \((C₃F₇)_2\) dimer peak or the n-C₃F₇I peak, the I₂ was trapped on the gas chromatograph column, or the I₂ concentration was so small that it was not detected on the gas chromatograph.

Analysis of t-C₄F₉I Sample

More than 20 gas chromatographic peaks were detected. Only the ethyl alcohol and t-C₄F₉I peaks could be identified. The remaining peaks contained perfluoroalkyl iodine compounds and perfluoroalkanes. The major peaks are listed below and are approximate.

<table>
<thead>
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<th>Peak</th>
<th>Identification</th>
<th>Sample, percent</th>
</tr>
</thead>
<tbody>
<tr>
<td>15</td>
<td>Ethyl alcohol</td>
<td>3</td>
</tr>
<tr>
<td>62-83</td>
<td>t-C₄F₉I</td>
<td>35</td>
</tr>
<tr>
<td>444-480</td>
<td>Unknown</td>
<td>60</td>
</tr>
<tr>
<td>1154</td>
<td>Unknown</td>
<td>2</td>
</tr>
</tbody>
</table>
Appendix B

Parametric and Scaling Studies

As shown in figure 3, the inversion density \([I^*] - [I]/2 > 0\) for only a fraction of the pump length \(Z_p\). We termed this fractional distance the positive gain length \(Z_{G+}\). We performed a parametric study to determine how \(Z_{G+}\) varies with flow speed \(w\), number density \([RI]\), and pump intensity \(I_p\). The result is expressed by the proportionality

\[
Z_{G+} \propto \frac{w}{[RI] \sqrt[2]{I_p}} \quad (B1)
\]

An increase in \(w\) retards the buildup of \([I_2]\), whereas an increase in \([RI]\) and \(I_p\) increases the level of photoionization and accelerates the buildup of \([I_2]\). This proportionality was used to optimize the laser model (to achieve positive gain throughout the pump region) by increasing \(w\) until \(Z_{G+} = Z_p\).

A scaling study was also performed to determine how output power density \(P_{out} (W \cdot cm^{-2})\) scales for similar lasers—lasers that have the same ratio of positive gain length to pump length \(Z_{G+}/Z_p\). For such similar lasers, \(P_{out}\) was found to scale according to the proportionality

\[
P_{out} \propto Z_{G+} I_p [RI] \quad (B2)
\]

The output power \(P_{out}\) in W scales with the inner radius \(r_i\) of a circular cylindrical laser tube according to the proportionality

\[
P_{out} \propto \langle \rho_{pN} \rangle r_i^2 \quad (B3)
\]

where \(\langle \rho_{pN} \rangle\) is the normalized density of pumping photons averaged over the cross section. A derivation of \(\langle \rho_{pN} \rangle\) as a function of the ratio \(r_i/\delta\), where \(\delta\) is the pumping photon absorption length, is given in appendix C.

The scaling study may be summarized as follows: Two optimal CW iodine lasers are similar if the lasant flow speeds satisfy the relation

\[
\frac{w_2}{w_1} = \left(\frac{Z_{p2}}{Z_{p1}}\right)^{1/2} \left(\frac{I_{p2}}{I_{p1}}\right)^{1/2} \left(\frac{[RI]_2}{[RI]_1}\right)^{1/2} \quad (B4)
\]

For such similar lasers, the output power \(P_{out}\) scales as

\[
\frac{P_{out2}}{P_{out1}} = \left(\frac{Z_{p2}}{Z_{p1}}\right)^{1/2} \left(\frac{I_{p2}}{I_{p1}}\right)^{1/2} \left(\frac{[RI]_2}{[RI]_1}\right)^{1/2} \left(\frac{\langle \rho_{pN2} \rangle}{\langle \rho_{pN1} \rangle}\right)^{1/2} \quad (B5)
\]

A plot of \(\langle \rho_{pN} \rangle\) as a function of \(r_i/\delta\) is given in appendix C.
Appendix C

Derivation of Average Pumping Photon Density in Absorbing Laser With Circular Cylindrical Symmetry

This appendix provides a derivation of the quantity \( \langle P_p N \rangle \) used to scale the output powers of equation (B5) in appendix B.

Two-Dimensional Relation of Incident Flux to Transmitting Surface Brightness

Assume that the laser tube and the incident pump radiation are axially symmetric and independent of \( Z \). (Appendix D gives a discussion of how this idealization compares with theoretical and actual elliptical pump chambers.) We also assume that the optical image of the pump source fills the interior of the laser tube, as shown in figure 18, where the curved wall of the circular glass tube acts as a negative lens. Furthermore, we assume that in the absence of absorption the pump intensity is uniform and isotropic on a cross section of the pump image. Then, if we take a point on the perimeter, as shown in figure 19, the image brightness \( B \) (\( \text{W.m}^{-2} \cdot \text{rad}^{-1} \), interpreted as watts per meter perimeter per meter depth per radian) can be integrated to determine the incident flux of pump radiation \( F \) (\( \text{W.m}^{-2} \)):

\[
F = \int_{-\pi/2}^{\pi/2} B \cos \theta \, d\theta = 2B \tag{C1}
\]

where the angle \( \theta \) is in radians. A comparable 3-D relation is given in reference 23 on page 23.

Pumping Photon Density Transmitted by Surface Element

We now determine the differential pumping energy density \( de \) (\( \text{J.m}^{-3} \)) at an interior point at depth \( \gamma \) due to incident radiation passing through length \( dS \) of the perimeter, as shown in figure 20. Initially, assume that the laser tube is evacuated so that absorption is not a factor. Then, the power within ray \( d\theta \) (\( \text{W.m}^{-1} \cdot \text{depth} \)) is given by \( B \cos \theta \, dS \, d\theta \). Because this power results from photons moving at the speed of light in vacuum \( c \), the pumping energy (\( \text{J.m}^{-2} \) (per meter depth per meter length)) within \( d\theta \) is given by \( (B/c) \cos \theta \, dS \, d\theta \). But \( d\theta = ds/\gamma \), hence, \( de \) (\( \text{J.m}^{-3} \)) due to the incident radiation that penetrates \( dS \) is given by

\[
de = \frac{B}{c\gamma} \cos \theta \, dS \tag{C2}
\]

and the corresponding pumping photon density \( d\rho_p \) (photons \( \cdot \text{m}^{-3} \)) is given by

\[
d\rho_p = \frac{B}{h\nu_p c\gamma} \cos \theta \, dS \tag{C3}
\]

For an absorbing laser gas, this photon density becomes

\[
d\rho_p = \frac{B}{h\nu_p c\gamma} \exp\left(-\frac{\gamma}{\delta}\right) \cos \theta \, dS \tag{C4}
\]

where \( \delta \) is the absorption length at the pump frequency.

Total Pumping Photon Density at Interior Point Due to Axisymmetric Incident Flux

Substitution of equation (C1) into (C4) gives \( de \) at point \( H \) of figure 21 when \( F \) is transmitted through \( dS \):

\[
d\rho_p = \frac{F}{2h\nu_p c\gamma} \exp\left(-\frac{\gamma}{\delta}\right) \cos \theta \, dS \tag{C5}
\]

However, by introducing the angle \( \psi \) shown in figure 21 we may write

\[
d\psi = \frac{\cos \theta}{\gamma} \, dS \tag{C6}
\]

and equation (C5) becomes

\[
d\rho_p = \frac{F}{2c\nu_p} \exp\left(-\frac{\gamma}{\delta}\right) d\psi \tag{C7}
\]

The total pumping photon density (photons \( \cdot \text{m}^{-3} \)) at point \( H \) in an absorbing gas may then be written

\[
\rho_p = \frac{F}{c\nu_p} \exp\left(-\frac{\gamma}{\delta}\right) \int_0^{\pi} d\psi \tag{C8}
\]

and by the law of cosines

\[
\gamma = -r_1 \sin \psi + \sqrt{\frac{2}{r_1^2} \sin^2 \psi} \tag{C9}
\]
Hence, the pumping photon density (photons · m\(^{-3}\)) at radius \(r\) in an absorbing laser is given by

\[
p_{p}(r, r_{t}, \delta) = \frac{F}{c h \nu_{p}} \int_{0}^{\pi} \exp \left[ \frac{1}{\delta} \left( r \cos \psi \right. \right.
\]
\[\left. \left. - \sqrt{\frac{r_{t}^{2} - r^{2} \sin^{2} \psi}{1 - r_{N}^{2} \sin^{2} \psi}} \right) \right] d\psi \tag{C10}
\]

The average pumping photon density over the cross section is given by

\[
\langle p_{p}(r, \delta) \rangle = \frac{1}{\pi r_{t}^{2}} \int_{0}^{r_{t}} 2\pi r p_{p}(r, r_{t}, \delta) \tag{C11}
\]

**Normalized Pumping Photon Densities**

In the limit as \(\delta \rightarrow \infty\), equation (C10) gives the uniform photon density for a nonabsorbing laser as

\[
p_{p}(\delta \rightarrow \infty) = \frac{\pi F}{c h \nu_{p}} \tag{C12}
\]

Using this quantity to normalize the photon density in equation (C10), we obtain

\[
p_{pN}(r_{N}, r_{t}/\delta) = \frac{1}{\pi} \int_{0}^{\pi} \exp \left[ \frac{r_{t}}{\delta} \left( r_{N} \cos \psi \right. \right.
\]
\[\left. \left. \left. - \sqrt{1 - r_{N}^{2} \sin^{2} \psi} \right] \right] d\psi \tag{C13}
\]

where \(r_{N} = r/r_{t}\) is the normalized radius. Plots of this normalized pumping photon density versus the normalized radius are given in figure 22 for various values of \(r_{t}/\delta\).

The average pumping photon density in an absorbing laser, as given by equation (C11), can also be normalized by

\[
\langle p_{pN}(r_{t}/\delta) \rangle = 2 \int_{0}^{1} dr_{N} r_{N} p_{pN}(r_{N}, r_{t}/\delta) \tag{C14}
\]

This normalized average pumping photon density is plotted versus \(r_{t}/\delta\) in figure 23. This quantity is an important factor in the scaling relation (eq. (B5) from appendix B) for calculating how output power scales with the radius of the laser tube.
Appendix D

Theoretical and Experimental Performance of Elliptical Pump Chamber

A ray tracing for the elliptical cross section of the pump chamber used in the laser experiments is shown in figure 24(a). The rays are assumed to emanate isotopically from a line source (perpendicular to the page) at the right focus. However, the rays incident on the laser tube at the left focus are found to be nonisotropic. As shown in figure 24(a), they are more concentrated on the side toward the source. A decrease in the distance between the foci in the elliptical pump chamber makes the incident radiation at the laser tube more isotropic, as shown in figure 24(b).

An experimental polar plot of the incident pumping radiation measured inside the laser tube is shown in figure 25. Differences from the theoretical plots result from the finite size (11 mm diameter) of the pump lamp, from obstruction by the lamp start-up wire, and from imperfections in the curvature and finish of the optical surfaces, including the elliptical cylinder and the flat end plates.
References


Table I. Lasant Rate Coefficients and Other Parameters

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<th>Parameter</th>
<th>Unit</th>
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<th>t-C₄F₉I</th>
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<td>Same</td>
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</tr>
<tr>
<td>[R₂(0)]</td>
<td>cm⁻³</td>
<td>0.2 [R₁(0)]</td>
<td>Same</td>
<td></td>
</tr>
</tbody>
</table>

* Cited value multiplied by 2.55.
† Cited value divided by 2.55.
‡ Semiempirical value multiplied by 2.55.
Table II. Experimental Data for Comparison of Lasants n-C₃F₇I and t-C₄F₉I

(a) Based on reference 14 and plotted in figure 2(a);
$\bar{I}_p = 1000$ SC; $T_0 = 300$ K

<table>
<thead>
<tr>
<th>$w_o$, m · sec⁻¹</th>
<th>$p_0$, torr</th>
<th>$P_{out}$, W</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Lasant n-C₃F₇I</td>
<td></td>
</tr>
<tr>
<td>5.75</td>
<td>12.0</td>
<td>4.0</td>
</tr>
<tr>
<td>6.96</td>
<td>14.0</td>
<td>6.0</td>
</tr>
<tr>
<td>7.40</td>
<td>17.0</td>
<td>7.2</td>
</tr>
<tr>
<td></td>
<td>Lasant t-C₄F₉I</td>
<td></td>
</tr>
<tr>
<td>5.07</td>
<td>14.0</td>
<td>11.1</td>
</tr>
<tr>
<td>6.45</td>
<td>17.0</td>
<td>13.8</td>
</tr>
<tr>
<td>6.72</td>
<td>14.5</td>
<td>12.5</td>
</tr>
</tbody>
</table>

(b) From reference 15 and plotted in figure 2(b)

<table>
<thead>
<tr>
<th>$I_p$, SC</th>
<th>$w_o$, m · sec⁻¹</th>
<th>$p_0$, torr</th>
<th>$P_{out}$, W</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Lasant n-C₃F₇I</td>
<td></td>
<td></td>
</tr>
<tr>
<td>450</td>
<td>6.6</td>
<td>5.6</td>
<td>1.4</td>
</tr>
<tr>
<td>600</td>
<td>6.6</td>
<td>5.6</td>
<td>1.8</td>
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<tr>
<td>740</td>
<td>6.2</td>
<td>6.0</td>
<td>2.2</td>
</tr>
<tr>
<td>925</td>
<td>6.4</td>
<td>5.8</td>
<td>2.4</td>
</tr>
<tr>
<td>1100</td>
<td>5.8</td>
<td>6.4</td>
<td>2.7</td>
</tr>
<tr>
<td></td>
<td>Lasant t-C₄F₉I</td>
<td></td>
<td></td>
</tr>
<tr>
<td>450</td>
<td>5.5</td>
<td>9.0</td>
<td>5.8</td>
</tr>
<tr>
<td>600</td>
<td>7.3</td>
<td>4.5</td>
<td>6.5</td>
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<tr>
<td>740</td>
<td>7.3</td>
<td>4.5</td>
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</tr>
<tr>
<td>925</td>
<td>7.1</td>
<td>4.2</td>
<td>8.4</td>
</tr>
<tr>
<td>1100</td>
<td>7.1</td>
<td>4.2</td>
<td>9.8</td>
</tr>
</tbody>
</table>
Figure 1. Laboratory schematic of CW iodine laser with longitudinally flowing lasant gas and continuous pumping by argon arc lamp.
Figure 2. Experimental comparison of output power from n-C$_3$F$_7$I and t-C$_4$F$_9$I with corresponding model-tuning curves.

(a) Data from reference 14.

(b) Data from reference 15.
Figure 3. Theoretical plots of diagnostic quantities versus $z$ for both end points of $n1$-model curve of figure 2(a).

(a) Molecular iodine density $[I_2]$ versus $z$.

(b) Excited iodine density $[I^+]$ versus $z$. 
(c) Ground-state iodine density [I] versus z.

(d) Inversion density [I'] - [I]/2 versus z.

Figure 3. Continued.
(e) Lasing photon densities $\rho_+$ and $\rho_-$ versus $z$.

(f) Perfluoroalkyl radical density $[R]$ versus $z$.

Figure 3. Continued.
(g) Perfluoroalkyl dimer density \([R_2]\) versus \(z\).

(h) Lasant temperature \(T\) versus \(z\).

Figure 3. Continued.
(i) Lasant flow speed $w$ versus $z$.

(j) Lasant density $\eta$ versus $z$.

Figure 3. Continued.
(k) Lasant pressure $p$ versus $z$.

Figure 3. Concluded.
Figure 4. Theoretical plots of diagnostic quantities versus $z$ for both end points of $t_1$-model curve of figure 2(a).

(a) Molecular iodine density $[I_2]$ versus $z$.

(b) Excited iodine density $[I^+]$ versus $z$.
(c) Ground-state iodine density [I] versus $z$.

(d) Inversion density $[I^*] - [I]/2$ versus $z$.

Figure 4. Continued.
(e) Lasing photon densities $p_+$ and $p_-$ versus $z$.

(f) Perfluoroalkyl radical density $[R]$ versus $z$.

Figure 4. Continued.
Figure 4. Concluded.
Figure 5. Theoretical plots of diagnostic quantities versus $z$ for both end points of $n_2$-model curve of figure 2(b).
(c) Ground-state iodine density $[I]$ versus $z$.

(d) Inversion density $[I^*] - [I]/2$ versus $z$.

Figure 5. Continued.
(e) Lasing photon densities $\rho_+$ and $\rho_-$ versus $z$.

(f) Perfluoroalkyl radical density $[R]$ versus $z$.

Figure 5. Continued.
(g) Perfluoroalkyl dimer density $[R_2]$ versus $\zeta$.

Figure 5. Concluded.
Figure 6. Theoretical plots of diagnostic quantities versus $z$ for both end points of $t_2$-model curve of figure 2(b).

(a) Molecular iodine density $[I_2]$ versus $z$.

(b) Excited iodine density $[I^+]$ versus $z$. 
(c) Ground-state iodine density $[I]$ versus $z$.

(d) Inversion density $[I^*] - [I]/2$ versus $z$.

Figure 6. Continued.
(e) Lasing photon densities $\rho_+$ and $\rho_-$ versus $z$.

(f) Perfluoroalkyl radical density $[R]$ versus $z$.

Figure 6. Continued.
Figure 6. Concluded.

(g) Perfluoroalkyl dimer density [R$_2$] versus $z$. 

Figure 6. Concluded.
Figure 7. Theoretical effects on output power of optimized lasant flow speeds in n1- and t1-models of figure 2(a).
Figure 8. Theoretical plots of diagnostic quantities versus z for both end points of n1-model optimal curve of figure 7.
(c) Ground-state iodine density $[I]$ versus $z$.

(d) Inversion density $[I^*] - [I]/2$ versus $z$.

Figure 8. Continued.
Figure 8. Continued.

(e) Lasing photon densities $\rho_+$ and $\rho_-$ versus $z$.

(f) Perfluoroalkyl radical density $[R]$ versus $z$. 

40
(g) Perfluoroalkyl dimer density $[R_2]$ versus $z$.

Figure 8. Concluded.
Figure 9. Theoretical plots of diagnostic quantities versus $z$ for both end points of $t1$-model optimal curve of figure 7.
(c) Ground-state iodine density $[I]$ versus $z$.

(d) Inversion density $[I^+] - [I]/2$ versus $z$.

Figure 9. Continued.
(e) Lasing photon densities $\rho_+$ and $\rho_-$ versus $z$.

(f) Perfluoroalkyl radical density $|R|$ versus $z$.

Figure 9. Continued.
Figure 9. Concluded.
Figure 10. Theoretical effects on output power of optimized laser flow speeds in n2- and t2-models of figure 2(b).

Figure 11. Theoretical comparison of output power from fictional f1-model and n1-model optimum of figure 7.
Figure 12. Theoretical plots of diagnostic quantities versus \( z \) for both end points of fictional model curve of figure 11.

(a) Molecular iodine density \([\text{I}_2]\) versus \( z \).

(b) Excited iodine density \([\text{I}^*]\) versus \( z \).
Figure 12. Continued.

(c) Ground-state iodine density $[I]$ versus $z$.

(d) Inversion density $[I^*] - [I]/2$ versus $z$. 
(e) Lasing photon densities $\rho_+$ and $\rho_-$ versus $z$.

(f) Perfluoroalkyl radical density $[R]$ versus $z$.

Figure 12. Continued.
Figure 12. Concluded.
Figure 13. Detailed theoretical plots of fictitious $f_1$-model and $n_1$-model optimum of figure 11 at $w_0 = 19.5 \text{ m \cdot sec}^{-1}$. 

(a) Molecular iodine density $[I_2]$ versus $z$. 

(b) Lasant number density $[R]$ versus $z$. 

Figure 14. Theoretical comparison of output power from fictitious $f_2$-model and $n_2$-model optimum of figure 10.
(a) Molecular iodine density $[I_2]$ versus $z$.

(b) Lasant number density $[RI]$ versus $z$.

Figure 15. Detailed theoretical plots for fictitious f2-model and n2-model optimum of figure 14 at $I_p = 740$ SC.
Figure 16. Theoretical comparison of output power from $n_1$-model optimum and fictitious $f_1$-model of figure 11 with $f_1$-model optimized by halving the lasant flow speeds.
Figure 17. Theoretical plots of diagnostic quantities versus \( z \) for both end points of fictitious \( f1 \)-model optimal curve of figure 16.
(c) Ground-state iodine density $[I]$ versus $z$.

(d) Inversion density $[I^*] - [I]/2$ versus $z$.

Figure 17. Continued.
(c) Lasing photon densities $\rho_+$ and $\rho_-$ versus $z$.

(f) Perfluoroalkyl radical density $[R]$ versus $z$.

Figure 17. Continued.
(g) Perfluoroalkyl dimer density $[R_2]$ versus $z$.

Figure 17. Concluded.
Figure 18. Cross section of elliptical pump chamber.

Figure 19. Cross section of laser tube with pump image brightness $B$ and incident flux $F$. 
Figure 20. Cross section of laser tube with ray geometry.

Figure 21. Cross section of laser tube with geometry for computing total pumping photon energy density at radius $r$ in absorbing lasant.
Figure 22. Normalized pumping photon density $\rho_{pN}$ versus normalized radius $r_N$ for various values of $r_\delta/\delta$. 
Figure 23. Normalized cross-sectional average pumping photon density $\langle p_{pN} \rangle$ versus $r_f/\delta$. 
Figure 24. Theoretical ray tracing for elliptical pump chamber with line source at right focus.
Figure 25. Polar plot of normalized incident pumping intensity $I_p$, measured at inner radius of laser tube for elliptical pump chamber used in laser experiments. Corresponding theoretical ray tracing is shown in figure 24(a).
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13. ABSTRACT (Maximum 200 words)  
A numerical rate equation model for a continuous wave iodine laser with longitudinally flowing gaseous lasant is validated by approximating two experiments that compare the perfluoroalkyl iodine lasants n-C₃F₇I and t-C₄F₉I. The salient feature of the simulations is that the production rate of the dimer (C₄F₉)₂ is reduced by one order of magnitude relative to the dimer (C₃F₇)₂. The model is then used to investigate the kinetic effects of this reduced dimer production—especially how it improves output power. Related parametric and scaling studies are also presented. When dimer production is reduced, more monomer radicals (t-C₄F₉) are available to combine with iodine ions, thus enhancing depletion of the laser lower level and reducing buildup of the principal quencher, molecular iodine. Fewer iodine molecules result in fewer downward transitions from quenching and more transitions from stimulated emission of lasing photons. Enhanced depletion of the lower level reduces the absorption of lasing photons. The combined result is more lasing photons and proportionally increased output power.  

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