A MODEL FOR AXIAL MAGNETIC BEARINGS INCLUDING EDDY CURRENTS

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ABSTRACT

This paper presents an analytical method of modelling eddy currents inside axial bearings. The problem is solved by dividing an axial bearing into elementary geometric forms, solving the Maxwell equations for these simplified geometries, defining boundary conditions and combining the geometries. The final result is an analytical solution for the flux, from which the impedance and the force of an axial bearing can be derived. Several impedance measurements have shown that the analytical solution can fit the measured data with a precision of approximately 5%.

INTRODUCTION

Modelling magnetic bearings is necessary to achieve reasonable results for magnetic force and bearing losses in order to design bearings and amplifiers. Furthermore, controller design is based on a model of the plant including magnetic bearings. Magnetic bearings are usually modelled as equivalent electrical circuits. With simple models neglecting eddy currents, hysteresis, saturation and material nonlinearities, it is possible to achieve approximations for the magnetic force which are sufficient for most applications. Nevertheless, better modelling can improve bearing design and system behaviour. For some applications such as self sensing bearings a more precise model is necessary.

The absolute permeability $\mu$ describes the relationship between the magnetic flux density $B$ and the magnetic field strength $H$ (see equation (5)). The relative permeability $\mu_r$ is material dependent, the permeability of vacuum $\mu_0$ is constant. For simple models $\mu_r$ can be assumed to be constant, but for real materials $\mu_r$ depends on the magnetic field strength (i.e. $\mu_r = \mu_r(H)$). With increasing magnetic field strength the gradient $\mu$ decreases (to a limit value $\mu_0$). Therefore, the magnetic flux density is nearly constant for high field strengths (saturation). While increasing and decreasing $H$, the flux density $B$ has different values for forward or backward loops due to hysteresis of the material. There are complex mathematical models such as the Preisach model [May91] to describe hysteresis, but these do not have an analytical solution.

Changing magnetic fields inside conducting materials cause currents to flow. These eddy currents have a large influence on the behaviour of magnetic bearings. The main goal of this paper
is to achieve an analytical description of a bearing model including eddy currents. Nonlinear
material behaviour is considered in numerical calculations. Hysteresis is neglected in this paper.
For elementary geometric forms, analytical solutions for eddy currents can be found. These
solutions are sufficient to describe the behaviour of axial bearings. The elementary geometric
forms are a semi-infinite plate, a rotationally symmetric plate and a semi-infinite cylinder.

**ELEMENTARY GEOMETRIC FORMS**

Electromagnetic fields and, therefore, eddy currents also can be described using Maxwell’s
equations ([Küp90], [Jac83]). In order to achieve analytical solutions it is necessary to divide
an axial bearing into elementary geometric forms. Contrary to numerical results, analytical
solutions can be more useful for bearing and controller design. Finite Element analysis can
solve two or three-dimensional problems and more complex geometric forms. These numerical
calculations are therefore used to verify analytical results and simplifications.

![Cross section of an axial bearing](image)

Figure 1: Cross section of an axial bearing. The bearing is divided into elementary forms.

**Maxwell’s Equations**

Considering that the frequencies of the currents and fields are sufficiently low so that we can
neglect the displacement current $\partial D/\partial t$, Maxwell’s equations can be written as

$$\text{curl} E = \frac{\partial B}{\partial t} = -\frac{\partial B}{\partial t} \frac{\partial H}{\partial t}$$  (1)

$$\text{curl} H = J$$  (2)

$$\text{div} B = 0$$  (3)

Here $E$ is the electric field strength and $J$ the current density. The generalized Ohm’s law,
with the electric conductivity $\sigma$ is given by:

$$J = \sigma E$$  (4)

The material equation is:

$$B = \mu_0 \mu_r H = \mu H$$  (5)
Combining equations (1), (2), (4) and (5) we obtain:

$$\text{curl(curl } H) = -\mu \sigma \frac{\partial H}{\partial t}$$

(6)

A sinusoidal variation of the fields enables the transformation of equation (6) into the Fourier space. $\partial / \partial t$ can be replaced by $j\omega$ giving:

$$\text{curl(curl } H) = -j\omega \mu \sigma H = -\alpha^2 H$$

(7)

where:

$$\alpha = \sqrt{j\omega \mu \sigma} = (1 + j) \sqrt{\frac{\omega \mu \sigma}{2}}$$

(8)

Maxwell’s equations (2) and (3) can be integrated using the theorems of Stokes and Gauss ($\Theta$: magnetomotive force, $\Phi$: magnetic flux) to give:

$$\Theta = \int_{\partial A} H d\mathbf{s} = \sum I + \int_{A} J d\mathbf{A}$$

(9)

$$\Phi = \int_{\partial V} B d\mathbf{A}$$

(10)

In the following sections, solutions for the elementary geometric forms semi-infinite plate, rotational symmetric plate and semi-infinite cylinder are given.

Figure 2: Cross section of
(a) a semi-infinite plate, (b) a rotational plate, (c) a semi-infinite cylinder

Semi-infinite Plate

A very long plate with thickness $d$ and a magnetic field only in the $z$-direction has an eddy current flow in the $x$-direction. Equation (7) leads to the following differential equation:

$$\text{curl } \begin{pmatrix} 0 \\ 0 \\ H_x \end{pmatrix} = \begin{pmatrix} \frac{\partial H_x}{\partial y} \\ 0 \\ 0 \end{pmatrix} = \begin{pmatrix} J_x \\ 0 \\ 0 \end{pmatrix}$$

(11)
with the general solution:

\[
H_z(y) = c_1 e^{\alpha y} + c_2 e^{-\alpha y}
\]  

(14)

\[
J_x(y) = \frac{\partial H_z}{\partial y} = c_1 \alpha e^{\alpha y} - c_2 \alpha e^{-\alpha y}
\]  

(15)

Rotational Symmetric Plate

\[
\text{curl} \left( \begin{array}{c} H_r \\ 0 \\ 0 \end{array} \right) = \left( \begin{array}{c} \frac{\partial H_r}{\partial r} \\ 0 \\ 0 \end{array} \right) = \left( \begin{array}{c} 0 \\ J_\varphi \\ 0 \end{array} \right)
\]

(16)

\[
\text{curl} \left( \frac{\partial H_r}{\partial z} \right) = \left( \begin{array}{c} -\frac{\partial^2 H_r}{\partial z^2} \\ 0 \\ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial H_z}{\partial z} \right) \end{array} \right) = \left( \begin{array}{c} -\alpha^2 H_r \\ 0 \\ 0 \end{array} \right)
\]

(17)

Equation (17) represents the following two differential equations:

\[
\frac{\partial^2 H_r}{\partial z^2} = \alpha^2 H_r
\]

(18)

\[
\frac{1}{r} \frac{\partial H_r}{\partial z} + \frac{\partial^2 H_r}{\partial z \partial r} = 0
\]

(19)

The solution for the differential equations (18) and (19) can be found by separation of the variables:

\[
H_r(r, z) = R(r) Z(z)
\]

(20)

Equation (18) can be transformed to:

\[
\frac{\partial^2 Z}{\partial z^2} = \alpha^2 Z
\]

(21)

with the general solution:

\[
Z(z) = z_1 e^{\alpha z} + z_2 e^{-\alpha z}
\]

(22)

Equation (19) can be transformed to:

\[
\frac{1}{r} R + \frac{\partial R}{\partial r} = 0
\]

(23)

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with the general solution:

$$\mathcal{R}(r) = r_1 \frac{1}{r}$$  \hspace{1cm} (24)

Equation (22) and (24) lead to:

$$H_r(r, z) = \frac{1}{r} (c_1 e^{\alpha z} + c_2 e^{-\alpha z})$$  \hspace{1cm} (25)

Without regard for $r$, equation (25) is similar to equation (14) (the case of the semi-infinite plate).

**Semi-infinite Cylinder**

$$\text{curl} \begin{pmatrix} 0 \\ 0 \\ H_z \end{pmatrix} = \begin{pmatrix} 0 \\ -\frac{\partial H_z}{\partial r} \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ J_\varphi \\ 0 \end{pmatrix}$$  \hspace{1cm} (26)

$$\text{curl} \begin{pmatrix} 0 \\ -\frac{\partial H_z}{\partial r} \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ -\frac{1}{r} \frac{\partial}{\partial r} (r \frac{\partial H_z}{\partial r}) \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ -\alpha^2 H_z \end{pmatrix}$$  \hspace{1cm} (27)

$$\frac{\partial^2 H_z}{\partial r^2} + \frac{1}{r} \frac{\partial H_z}{\partial r} - \alpha^2 H_z = 0$$  \hspace{1cm} (28)

The substitution $\bar{r} = \alpha r$ transforms the differential equation (28) into Bessel's equation:

$$\bar{r}^2 \frac{\partial^2 H_z}{\partial \bar{r}^2} + \bar{r} \frac{\partial H_z}{\partial \bar{r}} - \bar{r}^2 H_z = 0$$  \hspace{1cm} (29)

with the general solution:

$$H_z(\bar{r}) = c_1 I_0(\bar{r}) + c_2 K_0(\bar{r})$$  \hspace{1cm} (30)

or written without substitution:

$$H_z(r) = c_1 I_0(\alpha r) + c_2 K_0(\alpha r)$$  \hspace{1cm} (31)

Here, $I_0$ is the modified Bessel function of first type and zeroth order and $K_0$ is the modified Bessel function of second type and zeroth order. More detailed information can be found in [AS65]. For calculations with Bessel functions, especially with complex arguments, a power series expansion is necessary. A solution without Bessel functions is therefore desirable.

It can be clearly seen that equation (31) leads to equation (14) when $R_i \gg d$ ($R_i$: inner radius of the cylinder). A numerical comparison of the results for equation (31) and equation (14) shows that the difference between these two equations is only significant for unrealistically small values of $R_i$ ($R_i < d$).

An analysis of the cases of the semi-infinite cylinder and the rotational symmetric plate shows that these cases can be described with the equations of the semi-infinite plate. The problem of eddy currents in a semi-infinite plate is well known in the literature ([Sto74], [Küp90]), because the choice of lamination thickness is important for high power transformers and electrical machines.
Solutions of differential equations require definitions of boundary conditions. Current field lines have to be closed. This is valid for eddy currents too. It is therefore assumed that eddy currents in a semi-infinite plate turn back at infinity. Eddy currents flow in one direction at one edge and in the other direction at the second edge (see figure 3 (a)). This solution is given in the literature ([Sto74], [Küp90]).

For a semi-infinite cylinder a second solution is possible. Contrary to a plate, a cylinder is geometrically closed and therefore eddy currents can flow in one direction (see figure 3 (b)). Calculations with the Finite Element program FEMAG ([Ins94]) produce this solution by default.

![Figure 3: The two different boundary conditions. Cylinders with or without intersection.](image)

**Cylinder with intersection**

The constants \( c_1 \) and \( c_2 \) of equations (14) and (15) can be found by defining two boundary conditions. The first boundary condition is given by

\[
H_z\left(-\frac{d}{2}\right) = H_0
\]  

(32)

The second boundary condition can be derived using Kirchhoff's law. The sum of all the currents has to be 0.

\[
\int_{-\frac{d}{2}}^{\frac{d}{2}} J_x(y)\,dy = 0
\]  

(33)

Equation (33) can be written as:

\[
H_z\left(\frac{d}{2}\right) = H_z\left(-\frac{d}{2}\right)
\]  

(34)
The solutions for $H_z$, $J_x$ and the flux $\Phi$ with the boundary condition of a cylinder with intersection are as follows ($A$ is the area of the cross section of a cylinder):

$$H_z(y) = H_0 \frac{\cosh(\alpha y)}{\cosh(\alpha \frac{d}{2})}$$

(35)

$$J_x(y) = \alpha H_0 \frac{\sinh(\alpha y)}{\cosh(\alpha \frac{d}{2})}$$

(36)

$$\Phi = \int_{-\frac{d}{2}}^{\frac{d}{2}} B dA = \mu \frac{A}{d} \int_{-\frac{d}{2}}^{\frac{d}{2}} H_z(y) dy = A \mu H_0 \frac{\tanh(\alpha \frac{d}{2})}{\alpha \frac{d}{2}}$$

(37)

Figure 4 shows the decrease of the $H$-field in the middle of the core and the increase of eddy currents near the surface of the core with increasing frequency. The following model parameters have been used:

- $d = 6 \cdot 10^{-3}$ m
- $A = 1.39 \cdot 10^{-3}$ m$^2$
- $\mu_0 = 4\pi \cdot 10^{-7}$ H/m
- $\mu_r = 5000$
- $H_0 = 79.6$ A/m
- $B_0 = \mu_0 \mu_r H_0 = 0.5$ T
- $\sigma = 2 \cdot 10^7$ (\Omega m)$^{-1}$

**Cylinder without intersection**

In the case of a cylinder without intersection the start and end of the plate are fixed together and therefore eddy currents can have closed field lines and flow in one direction (the opposite
direction of the coil current). The eddy current density at the outer surface of the core is zero.

c1 and c2 in equations (14) and (15) can be found by defining two boundary conditions. The first boundary condition is given by:

\[ H_z(-\frac{d}{2}) = H_0 \] (38)

The second boundary condition is given by:

\[ J_x\left(\frac{d}{2}\right) = 0 \] (39)

The solutions for \( H_z, J_x \) and \( \Phi \) with the boundary condition of a cylinder without intersection are:

\[ H_z(y) = H_0 \frac{\cosh(\alpha_1 \frac{d}{2} - \alpha y)}{\cosh(\alpha d)} \] (40)

\[ J_x(y) = \alpha H_0 \frac{\sinh(\alpha_1 \frac{d}{2} - \alpha y)}{\cosh(\alpha d)} \] (41)

\[ \Phi = A\mu H_0 \frac{\tanh(\alpha d)}{\alpha d} \] (42)

Figure 5: \(|H_z(y)| \) and \(|J_x(y)| \) for a cylinder without intersection

Figure 5 shows the decrease of the \( H \)-field near to the outer surface of the core and the increase of eddy currents near to the inner surface of the core with increasing frequency. The model parameters of the preceding section have been used.

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1The unidirectional eddy current flow reminds us of a transformer with a shorted secondary winding. With the transformer equation, a model for eddy currents would be very simple. A transformer however has two windings and one core. Currents flow inside the windings and the core is only used as a magnetic material. An axial bearing has one winding and the core has two functions (for conducting and as a magnetic material). The transformer equation therefore would lead to an unacceptable error.
Comparison of the two boundary conditions

The generalized form of the flux equation can be written as:

$$\Phi = \Phi_0 \frac{\tanh(\gamma)}{\gamma}$$  \hspace{1cm} (43)

with $\Phi_0 = A\mu H_0$ and with $\gamma = \alpha \frac{d}{2} \ldots \alpha d$.

A comparison of the flux equations (36) and (42) shows that $\gamma$ is two times larger for unidirectional than for bidirectional eddy currents. Figure 6 shows the frequency dependent flux for both cases with the same $\Phi_0$. The flux bandwidth of bidirectional eddy currents is 4 times higher (see figure 6). This does not correspond to measurements which only show a factor of 2. This difference is caused by the frequency dependence of $\Phi_0$.

In the case of bidirectional eddy currents, $\Phi_0$ is constant for all frequencies. The $H$-field only changes inside the core and the boundary conditions at the surface of the core remain the same. This leads to frequency independance of the global $H$-field.

For unidirectional eddy currents the boundary conditions change at the surface of the core. This leads to a global change of the $H$-field and, therefore, to a change of $\Phi_0$. In order to achieve analytical solutions it is necessary to expand the problem of modelling eddy currents with the problem of modelling inductances. Calculation of inductances can be done for simple geometric setups, but for complex geometries such as axial bearings analytical models are not available.

Several measurements have shown that the flux curves are similar for both eddy current cases. Therefore, the equations for bidirectional eddy currents can be used for unidirectional eddy currents too. It is clear that axial bearings with screw-holes or other gaps have uni- and bidirectional eddy currents and, therefore, the flux curve is between the two extreme eddy current cases.
ELECTRICAL MODEL AND FORCES

The last boundary value \( H_0 \) can be found by closing the magnetic circuit. In order to build the complete magnetic circuit different materials have to be combined. Solutions for the flux are independent from the depth of the core (i.e. from the z-direction). Different materials lead to different solutions. It is clear that the flux changes in a transition region from one solution to the other. Finite Element analysis is used to calculate the flux distribution in a transition region between different materials. Figure 7 shows that the flux in the air gap is homogeneous even for high frequencies and that the transition region in the iron is small. Combining several materials is possible because the flux density distribution inside one material has little influence on the flux density distribution inside the next material.

\[ \Phi = \int B_m \, dA_m = A_m \mu_0 \mu_r H_m \frac{\tanh(\gamma_m)}{\gamma_m} \]  

For calculations with eddy currents, it is useful to define the cut-off frequency \( f_g \) of the flux decrease inside the core. In the literature \( f_g \) is called the cut-off frequency of the lamination (see also [Küp90]).

\[ f_g = \frac{1}{\pi \sigma \mu d^2} \]  

Axial bearing materials normally have either a very low \( f_g \) (unlaminated iron) or a very high \( f_g \) (laminated iron, special magnetic core materials such as Corovac\(^2\)) and air \((f_{g_{air}} \rightarrow \infty))\).

\[ \oint H \, dl = \sum_{m=1} H_m x_m = NI \]  

\(^2\)Corovac is a high frequency core material with \( \sigma \approx 5 \cdot 10^5 \text{S/m} \) and \( \mu_r \approx 130 \). The \( \mu_r \) of Corovac is low compared to ferrite cores which are used for high frequency transformers. The machinability of Corovac is however much better.
Solving the system of equations (45) and (46) leads to the flux equation (for \( n \) different materials):

\[
\Phi = \frac{\mu_0 N I}{\sum_{m=1}^{n} \frac{x_m}{\mu_{r,m} A_m \tanh(\gamma_m)}}
\] (47)

The modelling of axial bearings cannot be improved using more than two terms of the sum. The first term represents all materials with low \( f_g \) and the second term includes all materials with a behaviour as air (\( f_g \) is very high and \( \tanh(\gamma_m)/\gamma_m = 1 \)). This simplification is based upon the fact that a connection of two materials with similar \( f_g \) leads to:

\[
a_1 \frac{\tanh(\gamma_1)}{\gamma_1} + a_2 \frac{\tanh(\gamma_2)}{\gamma_2} \approx a_3 \frac{\tanh(\gamma_3)}{\gamma_3}
\] (48)

The reduced flux equation is:

\[
\Phi = \frac{\mu_0 N A I}{2x + \frac{l_{fe}}{\mu_r \tanh(\gamma)}} = \frac{\mu_0 N A I}{2x + \frac{1}{1 + \frac{1}{a_{fe} \tanh(\gamma)}}}
\] (49)

with:

\[
a_{fe} = \frac{2x}{l_{fe}}
\] (50)

All areas are normalized to the area of the air gap \( A \). \( x \) is the length of the air gap (including lengths of materials with very high \( f_g \)) and \( l_{fe} \) is the length of the un laminated core.

**Saturation**

In the previous calculations \( \mu_r \) has been assumed to be constant. Saturation of the magnetic material leads to a limit value for the flux. The solution for the flux therefore gives unrealistically high values at the surface. Even with simple models for saturation (e.g. Fröhlich model [Sto74]), it is not possible to derive an analytical solution for the differential equations. Numerical calculations have shown that the behaviour of the H-field with saturation is very similar to the behaviour without saturation. It is therefore not necessary to include saturation into a model for eddy currents.

**Equivalent Electrical Circuit**

When \( x \) is assumed to be constant, the impedance of an eddy current affected coil can be written as:

\[
Z = \frac{U}{I} = \frac{N d\Phi}{dt} \frac{l}{I} = j \omega L_0 \frac{1}{1 + \frac{1}{a_{fe} \tanh(\gamma)}}
\] (51)

with:

\[
L_0 = \mu_0 N^2 A \frac{1}{2x}
\] (52)
Forces

The magnetic force is proportional to the square of the flux. Therefore a reduction of the flux due to eddy currents will also reduce the force to:

\[
F = \frac{\Phi^2}{\mu_0 A} = \frac{K I^2}{4} \frac{1}{x^2 \left(1 + \frac{1}{\alpha_{fe} \tanh(\gamma)}\right)^2}
\]

(53)

with:

\[
K = \frac{\mu_0 N^2 A}{A}
\]

(54)

Figure 8: Normalized force dynamics |F/F_{DC}| of an axial bearing (f_o \approx 0.02Hz, \alpha_{fe} \approx 200)

NONLINEAR PARAMETER IDENTIFICATION

Parameter Identification

It is clear that the impedance of a coil not only depends on the inductance affected by the eddy currents. An impedance model must also include the coil resistance R_{cu}, the coil inductance L_{cu} and the coil capacitance C_{cu} (see figure 9). The last two parameters are necessary to model the frequency range above 10kHz. Thus, the total impedance is:

Figure 9: Complete equivalent electrical circuit
\[ Z_{\text{tot}} = \frac{1}{j\omega C_{\text{cu}}} \| (Z + R_{\text{cu}} + j\omega L_{\text{cu}}) \]  

(55)

Robustness of the identification can be improved by defining the constraints. The authors used the constraint algorithm from the MATLAB Optimization Toolbox. The parameter range described below is sufficient to approximate all measured axial bearings. Less stringent constraints are possible, but computational time increases.

All parameter ranges excluding \( R_{\text{cu}} \) are written with exponential notation. The fitting algorithm therefore fits the exponent of the parameters. This method avoids the value 0 and guarantees convergence.

The ranges of the fitted parameters are given below.

\[
\begin{align*}
R_{\text{cu}} &= 0 \ldots 2 \ \Omega \\
C_{\text{cu}} &= 1 \cdot 10^{-12} \ldots 1 \cdot 10^{-9} \ F \\
L_{\text{cu}} &= 1 \cdot 10^{-6} \ldots 1 \cdot 10^{-3} \ H \\
L_0 &= 1 \cdot 10^{-3} \ldots 1 \cdot 10^9 \ H \\
\alpha_{fe} &= 1 \cdot 10^0 \ldots 1 \cdot 10^3 \\
f_g &= 1 \cdot 10^{-3} \ldots 1 \cdot 10^1 \ Hz
\end{align*}
\]

Figure 10: Example 1: Axial bearing with un laminated core and target.
– - : calculated data, - - : measured data
Identified parameters: \( R_{\text{cu}} \approx 1.5\Omega, C_{\text{cu}} \approx 800\mu F, L_{\text{cu}} \approx 120\mu H, L_0 \approx 12m H, \alpha_{fe} \approx 3, f_g \approx 7Hz \)
The fitting error $e$ is approximately 5% with the usual air gap and increases up to 10% when the air gap is small and the flux is saturated and is given by:

$$
e = \sqrt{\frac{1}{n} \sum_{i=1}^{n} \left(1 - \frac{Z_{fit}}{Z_{meas}}\right)^2} \quad (56)$$

where $n$ is the number of measured values, $Z_{fit}$ the fitted impedance and $Z_{meas}$ the measured impedance.

**Losses**

The spectral losses of an axial bearing are given by:

$$P = \text{real}(U\text{conj}(I)) = |U|^2 \text{real(} \text{conj} \left( \frac{1}{Z_{tot}} \right)) \quad (57)$$

Figure 12 shows normalized loss based on impedance measurements of an axial bearing with a Corovac core. At low frequencies losses are very high due to $R_{cu}$. With increasing frequency...
losses decrease. It is therefore useful to choose the switching frequency of an amplifier as high as possible. As mentioned in section *Comparison of the two boundary conditions*, a radial intersection of an axial bearing can increase the bandwidth of the flux and leads to lower losses. It is only sufficient to intersect unlaminated parts of the core with low $f_g$, i.e. in most cases the axial bearing target. An axial bearing without intersection suffers more than 40% higher loss than one with intersection at frequencies typically used for switching amplifiers (20 - 200 kHz).

![Graph](image)

Figure 12: upper plot: normalized losses $P_n = P/|u|^2$ of an axial bearing with Corovac core and steel target (--- $P_1$: target intersected, - - - $P_2$: target not intersected)
lower plot: relative difference between the two curves of the upper plot $P_r = (P_2 - P_1)/P_1$

**Measurement**

Measurements of axial bearing impedances are made with a LCR-meter (HP 4284A) including a power current source (HP 42841A). The LCR-meter has the possibility to measure with logarithmically swept frequencies from 20Hz up to 1MHz. A host computer controls the measurement via GPIB-interface.
The power current source allows measurements of very high inductances and can superpose the measurement signal with a DC-bias current of up to 20A.

**CONCLUSIONS**

A model for eddy currents in axial bearings has been derived which can describe the behaviour for a wide frequency range. From the analytical solution for the flux, the impedance, the magnetic force and the losses can be calculated. Measurements have shown that the model is very precise.
REFERENCES


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