Infrared Spectroscopy of Matrix-Isolated Polycyclic Aromatic Hydrocarbon Cations. 2. The Members of the Thermodynamically Most Favorable Series through Coronene

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Gaseous, ionized polycyclic aromatic hydrocarbons (PAHs) are thought to be responsible for a very common family of interstellar infrared emission bands. Here the near- and mid-infrared spectra of the cations of the five most thermodynamically favored PAHs up to coronene: phenanthrene, pyrene, benzo[ghi]pyrene, benzo- [ghi]perylene, and coronene, are presented to test this hypothesis. For those molecules that have been studied previously (pyrene, pyrene-d10, and coronene), band positions and relative intensities are in agreement. In all of these cases we report additional features. Absolute integrated absorbance values are given for the phenanthrene, perdeuteriophenanthrene, pyrene, benzo[ghi]perylene, and coronene cations. With the exception of coronene, the cation bands corresponding to the CC modes are typically 2–5 times more intense than those of the CH out-of-plane bending vibrations. For the cations, the CC stretching and CH in-plane bending modes give rise to bands that are an order of magnitude stronger than those of the neutral species, and the CH out-of-plane bends produce bands that are 5–20 times weaker than those of the neutral species. This behavior is similar to that found in most other PAH cations studied to date. The astronomical implications of these PAH cation spectra are also discussed.

I. Introduction

Polycyclic aromatic hydrocarbons (PAHs) form an important class of organic molecules, having long been known to play prominent roles as environmental toxins, and in soot formation. During the past decade, evidence has been mounting that they are important components of the interstellar medium as well. Under the harsh conditions of many interstellar environments, PAHs are expected to be predominantly ionized. For this reason we have undertaken a study of the infrared spectroscopic properties of matrix-isolated PAH cations. While the primary goal of this research is to investigate the role of PAHs in the interstellar medium, the results are of general interest. PAH cations are considered to be important intermediates in combustion, and PAH intermediates and primary reaction products are directly related to the carcinogenicity of some members of this hydrocarbon family.

Ionized and neutral PAH molecules have been proposed as the carriers of a very common family of interstellar infrared fluorescence bands at 3050, 1615, 1310, 1150, and 885 cm⁻¹ (3.3, 6.2, 7.7, 8.7, and 11.3 μm). The ubiquity of this spectrum suggests that free molecular PAHs are common throughout interstellar space, and the intensity of the bands indicates that they are as abundant as the most abundant polyatomic interstellar molecules known. As much as 30–40% of the radiant energy from some of these sources is emitted in this family of bands. Reviews of this field can be found in ref 3.

While the interstellar PAH hypothesis rests on the general resemblance of the interstellar emission spectra to the absorption spectra of aromatic hydrocarbons, the available data are mainly for neutral PAHs suspended in salt pellets or organic solvents. Under such conditions, interactions with the surrounding medium and with other PAH molecules (clusters) strongly perturb the transitions.8 While this body of information has been invaluable in establishing the interstellar PAH hypothesis, there are important differences in details between the laboratory and interstellar spectra. Differences in band position and relative band intensities, as well as the fact that most interstellar PAHs are expected to be isolated and ionized in the regions where they emit, have hindered the verification of the hypothesis and the exploitation of PAHs as probes of the interstellar medium. Thus, the infrared spectra of neutral, isolated PAHs and of ionized, isolated PAHs are of fundamental importance to further progress in this field.

There has been substantial progress recently in experimentally measuring the emission spectra of vibrationally excited, gas phase aromatic molecules. Elegant UV-pumped IR emission experiments on gas phase neutral PAHs have been carried out.9–13 The results of these experiments, in conjunction with absorption studies of PAHs in the gaseous state at various temperatures,14–18 show that (i) the emission peak frequency of the aromatic CH stretch in gas phase PAHs shifts as a function of internal energy content and matches that of the interstellar band at 3050 cm⁻¹ at reasonable levels of excitation; (ii) the bandwidth increases with energy and matches that of the interstellar band at similarly appropriate internal energies; and (iii) the plateau that underlies the interstellar feature, and extends out to about 2740 cm⁻¹, is reproduced as well. Schlemmer et al.,12 in reporting the gas phase emission from naphthalene and pyrene, show that the relative intensities of the emission bands differ in absorption from those for the same species and from the relative intensities of the interstellar emission features. Thus, they conclude that small neutral PAHs cannot be the carriers of the interstellar bands. This is, of course, consistent with PAHs being ionized in the emission zones. These emission experiments on neutral PAHs are inherently very difficult, and it will likely be some time before this type of measurement will be carried out on ionized PAHs. These results raise important, fundamental questions regarding the molecular physics of the infrared emission process from vibrationally excited molecules.
There is also a need for reliable, quantitative infrared absorption data on neutral and ionized PAHs at the low temperatures characteristic of the interstellar medium. We have initiated a systematic program to obtain this information. In the first paper in this series, we reported our experimental approach to the problem and presented results from 4000 to 200 cm\(^{-1}\) on the smallest PAH, naphthalene \((C_{10}H_{8})\), and its fully deuterated analog, \((C_{10}D_{8})\).\(^{19}\) We now have the infrared spectra for a number of neutral PAHs and their cations. Here the spectroscopic properties of the members of the thermodynamically most favored PAH species containing between 6 and 24 carbon atoms are presented. Stein and co-workers\(^{20}\) have extensively investigated the thermodynamic properties of aromatic species, showing that, of the hundreds of PAH isomers possible in this size range, those in Figure 1 are thermodynamically most favored under high-temperature PAH growth conditions. Thus, these are likely to be important in the circumstellar environments in which PAHs are produced and able to survive the harsh conditions of many of the interstellar emission regions. In subsequent papers in this series, different aspects of PAH cation structure, including tetracyclic PAHs, polyacenes, and PAHs incorporating 5-membered rings, will be explored. The neutral spectra of these PAHs will be presented elsewhere.\(^{21,22}\) Our motivation is to understand how structure and size influence the infrared spectra of small PAHs in their neutral and cationic forms.

Efforts to provide this type of information are also underway at several other institutions. Defrees and Miller carried out theoretical calculations on the expected infrared spectral properties of the naphthalene and anthracene cations and predicted surprising relative band strength differences between the neutral and ionized forms of these PAHs.\(^{23-25}\) Subsequently, Defrees et al.\(^{25}\) and Pauzat et al.\(^{24}\) have greatly expanded this work. Professor Vala and co-workers at the University of Florida have published the infrared spectroscopic properties for cations of the PAHs naphthalene, anthracene, pyrene, perylene, and coronene,\(^{25-28}\) and d'Hendecourt et al. have reported the spectrum of the coronene cation.\(^{29}\)

This paper is laid out as follows. The experimental techniques are summarized in section II. In section III are presented the near- and mid-infrared spectra of the phenanthrene, phenanthrene-\(d_{10}\), pyrene, pyrene-\(d_{10}\), benzo[e]pyrene, benzo[ghi]perylene, and coronene cations. These results are discussed in section IV, where trends and implications regarding the interstellar question are pointed out.

II. Experimental Section

The matrix isolation technique is employed to isolate individual polycyclic aromatic hydrocarbon (PAH) molecules in an argon matrix, where they are photoionized and probed spectroscopically. Argon matrices are known to be suitable for vibrational studies, typically causing small shifts in the 0–15 cm\(^{-1}\) range. For electronic transitions, a neon matrix is required to avoid serious perturbations. With the exception of the sample deposition procedure, our experimental procedure has been described in detail previously.\(^{19}\) A schematic cross section of the sample chamber is shown in Figure 2. Briefly, an infrared-transmitting substrate (CsI) is suspended inside a high-vacuum chamber and cooled by a closed-cycle helium refrigerator. The ultra-high vacuum chamber \((p \approx 10^{-8} \text{ mTorr})\) is equipped with multiple inlet ports, and the cooler is mounted in such a way that the infrared window can be rotated to face any of these ports without breaking the vacuum. Typically, the CsI window was cooled to 10 K and positioned to face the sample deposition inlets. Samples were prepared by co-deposition of a gaseous PAH with a large overabundance of argon to a thickness appropriate for generation of the cation and measurement of its infrared spectrum. After sample deposition was complete, the cold head was rotated to face the beam of an infrared spectrometer so that a prephotolysis spectrum could be recorded. The sample was then rotated to face a vacuum ultraviolet lamp for photolysis. The sample was finally returned to the scanning position where the postphotolysis spectrum was recorded. Comparison of this spectrum with the neutral spectrum permitted identification of the ion features that appeared upon photolysis.

A. Sample Preparation. PAHs containing three or more rings do not have sufficient vapor pressure at room temperature to permit their preparation as a gaseous mixture with argon, as
was the case in our previous study involving naphthalene. This necessitated the vaporization of the PAH of interest and subsequent co-deposition of the gaseous PAH molecules with excess argon. To this end, PAH samples were placed in resistively heated Pyrex tubes (1/2 in. diameter) and mounted on the sample chamber through a stainless steel Cajon Ultratorr fitting. The temperature of the PAH was monitored using a chromel/alumel thermocouple mounted on the exterior of the tube with aluminum tape. Such an arrangement should potentially be useful at temperatures as high as 316 °C (the failure temperature of the Viton O-ring) and with samples as small as a few milligrams. Argon was admitted through a second port (Figure 2). The argon deposition line was liquid nitrogen trapped to minimize contamination.

Sample quality was found to be optimal for PAH vapor pressures in the range 10−30 mTorr. Higher vapor pressures required higher argon deposition rates that exceeded the thermal conductivity of the CsI window, warming the matrix. The annealing that resulted produced a matrix that was highly scattering at short wavelengths, crippling the Lyman-α photoionization efficiency and the signal-to-noise ratio of the near-IR spectrum. Conversely, lower vapor pressures required longer deposition times, which necessarily increased the contaminants in the matrix and, in turn, reduced the ionization efficiency. Thermochemical data were used to establish the approximate temperature necessary for each PAH investigated. Optimum tube temperatures were as follows: phenanthrene, 10 °C; pyrene, 60 °C; benzo[e]pyrene, 97 °C; benzo[ghi]perylene, 137 °C; and coronene, 158 °C. The optimal argon flow rate was estimated to be between 0.5 and 1.0 mmol/h.

As was the case in our earlier study of the naphthalene cation, photoionization of PAH/Ar matrices was accomplished using the Lyman-α emission from a microwave-powered discharge lamp optimized by the use of a 10% H2 in He gas mixture at ~75 mTorr. A MgF2 vacuum window allowed transmission of this radiation to the sample. For those experiments incorporating the electron acceptor CCl4 in the matrix, lower energy photolysis was necessary. Lyman-α photons (10.35 eV) are sufficiently energetic to ionize matrix-isolated CCl4 (IEccl4 = 11.28 eV) molecules, negating their usefulness as electron acceptors. For those experiments, UV radiation was generated by a discharge in pure H2 gas at ~150 mTorr. This optimized the broad molecular hydrogen emission band centered around 160 nm (7.77 eV) with respect to the Lyman-α line. A CaF2 vacuum window (cutoff λ ≈ 150 nm) served as a filter to exclude the residual Lyman-α radiation.

This technique results in ion yields ranging from 5 to 10%. Specifically, the yields for the spectra presented here were as follows: phenanthrene, 9 ± 4%; phenanthrene-d10, 7 ± 4%; pyrene, 10 ± 5%; benzo[e]pyrene, 6 ± 3%; benzo[ghi]perylene, 10 ± 5%; and coronene, 11 ± 6%. These yields were determined as previously described.19

Pure argon samples having pressures in the range 75−225 Torr were prepared as described previously.19

The PAHs used in this investigation include phenanthrene (Aldrich Chemical Co., 98.4% purity), phenanthrene-d10 (Aldrich, 97% purity), pyrene (Aldrich, 99% purity), pyrene-d10 (Cambridge Isotope Laboratories, 98% purity), benzo[e]pyrene (Aldrich, 99% purity), benzo[ghi]perylene (Aldrich, 98% purity), and coronene (Pfaltz and Bauer, Inc., 97% purity). All samples were used without further purification. Matheson prepurified argon (99.998% minimum) was used in these studies.

B. Spectroscopy. The sample chamber, now equipped with CsI vacuum windows, is suspended in the sample compartment of an FTIR spectrometer ( Nicolet Analytical Instruments, Model 740). All spectra reported here were measured at 0.5 cm−1 resolution. This resolution is critical in detecting ion bands that fall near the position of a neutral band. Spectra typically were generated through the coaddition of five blocks of 200 scans, a number that optimized both the signal-to-noise ratio and the time requirements of each experiment.

Mid-infrared spectra (7000−500 cm−1) were collected using an MCT-B detector/KBr beam splitter combination. Near-infrared spectra (15000−8800 cm−1) were collected using a silicon detector/quartz beam splitter combination. Conversion between the two spectral regions allowed the measurement of both electronic and vibrational band intensities from a single matrix.

C. Cation Band Identification. Three criteria must be satisfied before a photoproduct band is attributed to a particular PAH cation. First, the bands must appear only when the associated neutral PAH is present in the matrix. An extensive number of control experiments have been carried out to ensure that this is the case. Second, the bands attributed to the cation must be markedly enhanced when CCl4, an electron acceptor, is present in the matrix at a concentration of 1 part in 200. This behavior establishes that the bands arise from a positive ion. This effect is demonstrated in Figure 3, which shows a portion of the spectrum of the benzo[ghi]perylene cation generated both with and without CCl4 in the matrix. Third, all the bands attributed to the cation must evolve in a similar fashion. Spectra collected as a function of photolysis time show that all of the bands assigned to the respective cation grow and decline together in a manner distinct from other features produced by photolysis. These band correlation plots are presented on a molecule by molecule basis along with the appropriate spectroscopic data. In general, the intensities of the bands corresponding to the PAH cation (plotted showing data points) peak after 4−8 min of photolysis and then remain essentially constant or fall off slightly upon further photolysis. This behavior is consistent with that observed previously for the naphthalene cation and discussed in ref 19. Conversely, the bands of other photoproducts (plotted without data points shown) tend to increase steadily with photolysis time. Rogue photoproduct bands were not common, typically numbering 2−4 in each case. Most common were...
the thermodynamically most important vibrations of the following reasons. First, as little is known about the infrared properties of PAH cations, the smallest members of the PAH family were chosen to minimize the number of fundamental vibrations and spectral complexity. Second, as this work is motivated by the astrophysical question, the PAHs most likely to be important in this context were given first preference. Thus, the thermodynamically most favorable members of the high-temperature PAH formation route containing 6–24 carbon atoms were studied (Figure 1). These spectra are presented here in the second paper of the series. Third, as the effects of structure and size on the infrared spectra of simple PAH cations are not known, related PAHs with limited, well-defined structural differences have been studied together. Finally, as size increases, PAH samples tend to become increasingly refractory and increasingly difficult and expensive to obtain. Thus, in the interest of experimental practicality, smaller PAHs were the logical starting point. Small PAHs are also relevant from the astrophysical point of view. Indeed, a QRRK analysis of the interstellar spectra indicates that molecules containing 20–30 carbon atoms dominate the fluorescence at the shortest wavelengths.\textsuperscript{2a,3a,4}

The positions and relative intensities of the mid-infrared cation bands for the thermodynamically most favorable series of PAHs up to coronene are shown in Table 1. In addition, the spectra of the perdeuterated phenanthrene (C\textsubscript{14}D\textsubscript{10}\textsuperscript{+}) and pyrene (C\textsubscript{16}D\textsubscript{16}\textsuperscript{+}) cations have also been measured. Band positions and relative intensities are compared with those of the fully hydrogenated cations and with those of naphthalene (C\textsubscript{10}H\textsubscript{8}\textsuperscript{+}) and phenanthrene (C\textsubscript{14}H\textsubscript{10}\textsuperscript{+}) in Table 2. The infrared spectra of the cations of the thermodynamically most favorable series of PAHs are presented in the order of increasing molecular size. Only those regions of the spectrum where cation bands appear are shown.

III. Results

There are many PAHs that one can study. Those studied at the NASA Ames Research Center have been selected for the following reasons. First, as little is known about the infrared properties of PAH cations, the smallest members of the PAH family were chosen to minimize the number of fundamental vibrations and spectral complexity. Second, as this work is motivated by the astrophysical question, the PAHs most likely to be important in this context were given first preference. Thus, the thermodynamically most favorable members of the high-temperature PAH formation route containing 6–24 carbon atoms were studied (Figure 1).\textsuperscript{19} These spectra are presented here in the second paper of the series. Third, as the effects of structure and size on the infrared spectra of simple PAH cations are not known, related PAHs with limited, well-defined structural

### Table 1: Infrared Frequencies and Relative Intensities for the Cations of the Most Thermodynamically Favorable PAHs

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<th>( I_\text{nl} )</th>
<th>( I_\text{nu} )</th>
<th>( I_\text{nl} )</th>
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### Table 2: Comparison of the Infrared Frequencies and Relative Intensities for Isotopomeric Forms of the Cations of Naphthalene, Phenanthrene, and Pyrene

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Infrared Spectroscopy of PAH Cations

The Phenanthrene Cation, \( C_{14}H_{10}^+ \). The mid-infrared spectrum of the phenanthrene cation is shown in Figure 4. The band frequencies and relative intensities are listed in Table 1. The evolution of the bands assigned to the cation by Lyman-\( \alpha \) photolysis is plotted in Figure 5. This band correlation, one of the criteria of ion assignment, is looser for the weaker bands whose signal-to-noise ratios lie in the 2-3 range. Nonetheless, their overall behavior of a sharp rise and leveling off warrants their assignment to the cation.

Only weak new bands appear in the normally very strong CH out-of-plane bending region between 900 and 600 cm\(^{-1}\). Most of the detectable new bands fall between 1570 and 1250 cm\(^{-1}\) and, thus, are assigned principally to the CC stretching and CH in-plane bending modes of the phenanthrene cation. There undoubtedly are other cation bands that are discernible in the normal CH out-of-plane bending region between 900 and 600 cm\(^{-1}\) from the CH stretching and CH in-plane bending modes of the phenanthrene cation. Their assignment to the analogous bands of the neutral counterparts.23a,b,24 The spectrum of the naphthalene cation, the first multiring aromatic of the series, has been presented elsewhere.19,25 Additional discussions of the pyrene and coronene cations studied to date are listed in Table 1. The spectrum of neutral phenanthrene is richest between 1600 and 580 cm\(^{-1}\), the cation band frequencies lie between 1570 and 582 cm\(^{-1}\). The mid-infrared spectrum of the naphthalene cation is shown in Figure 4. In Figure 5, the cation frequency spectrum, with bands between 1570 and 582 cm\(^{-1}\) and (\( a_{135} \)) = (0.09)(0.2493) = 0.229 cm\(^{-1}\). If the absorption strengths of the CH out-of-plane bends in the ion were the same as those of the neutral molecule, we would expect to see the growth of ion bands with integrated intensities \( N_b \text{ cm}^{-1} \) and \( N_b \text{ cm}^{-1} \) and \( N_b \text{ cm}^{-1} \) and \( N_b \text{ cm}^{-1} \) and \( N_b \text{ cm}^{-1} \).

The Perdeuterated Phenanthrene Cation, \( C_{14}D_{10}^+ \). The mid-infrared spectrum of the fully deuterated phenanthrene cation intensity pattern is just the opposite, with the bands in the CC stretching and CH in-plane bending region being stronger than those in the CH out-of-plane bending region. Interestingly, the strongest cation bands fall at 1565 and 1277 cm\(^{-1}\), frequencies within the envelopes of the strongest interstellar emission features.

Two cation bands are discernible in the normal CH out-of-plane bending region between 900 and 680 cm\(^{-1}\), namely, those at 836.0 and 756.2 cm\(^{-1}\). They fall close to the regions expected for two and four adjacent hydrogen atoms per aromatic ring, respectively, the types of peripheral H atoms on phenanthrene. Note that both bands (836.0 and 756.2 cm\(^{-1}\)) are shifted to higher frequencies by 15-20 cm\(^{-1}\) from the very intense bands corresponding to the CH out-of-plane bending vibrations in the neutral molecule (813 and 735 cm\(^{-1}\)). Furthermore, one can estimate the intensity of these features relative to those of the neutral molecule. Let \( n\text{ cm}^{-1} \) be the number of neutral molecules in the (prephotolysis) sample. \( n\text{ cm}^{-1} \) is the number of ions produced by photolysis. The ionization fraction, \( I(=n^+/n)\), has been measured directly as 0.09 ± 0.04 (9 ± 4%) by the fractional decrease in the neutral band areas with photolysis for this experiment. The integrated intensities (\( f d\nu \equiv a \)) of the bands from the CH out-of-plane bending modes of the neutral phenanthrene molecules are \((a_{135}) = 0.165 \text{ cm}^{-1}\) and \((a_{735}) = 0.249 \text{ cm}^{-1}\). If the absorption strengths of the CH out-of-plane bends in the ion were the same as those of the neutral molecule, we would expect to see the growth of ion bands with integrated intensities 9 ± 4% those of the prephotolysis neutral bands or \((a_{135}) = (0.09)(0.165) = 1.5 \times 10^{-2} \text{ cm}^{-1}\) and \((a_{735}) = (0.09)(0.2493) = 2.2 \times 10^{-2} \text{ cm}^{-1}\).

In reality, the intensities of the 836 and 756 cm\(^{-1}\) ion bands are measured to be \((a_{135}) = 1.123 \times 10^{-3} \text{ cm}^{-1}\) and \((a_{735}) = 1.406 \times 10^{-3} \text{ cm}^{-1}\), respectively, or 13\( x \) and 16\( x \) weaker, on average, than those in the neutral molecule. This is in general agreement with our previous results for the naphthalene molecule, where the CH out-of-plane mode was found to be suppressed by a factor of 18 in the cation.

The two bands at 694.5 and 582 cm\(^{-1}\) probably arise from CCC plane bending motions.
deuterated PAHs have the is shown in Figure 6. The evolution of the positions. Due to its low intensity and separation from other features, the cation band at 1013.4 cm⁻¹ is not shown.

Figure 6. Mid-infrared bands of the perdeuterated phenanthrene cation isolated in an argon matrix at 10 K: (a) before photolysis; (b) after 8 min of in situ photolysis. The cation bands are labeled with their positions. Due to its low intensity and separation from other features, the cation band at 1013.4 cm⁻¹ is not shown.

is shown in Figure 6. The evolution of the bands assigned to the cation with Lyman-α photolysis is plotted in Figure 7. These bands, together with the corresponding hydrogenated ion bands, are listed in Table 2.

Comparisons between the spectra of the hydrogenated and deuterated PAHs have been used to shed light on the nature of a particular band. For phenanthrene, however, without the benefit of calculations that take molecular symmetry into account, this technique alone is of limited use. The 1565, 1558.2, 1551, and 1513 cm⁻¹ bands all fall in the CC stretching region, with a $\nu_{CC}^\text{H}/\nu_{CC}^\text{D}$ ratio of 1.01–1.03 to their presumed deuterated counterparts at 1547.4, 1528/1525, 1507, and 1459.5 cm⁻¹. This assignment to CC stretching modes does not require deuteration.

While the positions of the bands in the 1300–1260 cm⁻¹ region of perhydrophenanthrene are suggestive of CC stretching vibrations, it is conceivable that they might have some CH in-plane bend character as well. Simply taking the bands between 1300 and 1250 cm⁻¹ and comparing them to the frequencies of the perdeuterated molecule between 1296.5 and 1201 cm⁻¹, one obtains $\nu_{CH}^{\text{H}}/\nu_{CH}^{\text{D}}$ ratios in the 1–1.05 range, again consistent with modes that are primarily CC stretching in nature. The difficulty here is that bands also appear upon deuteration between 1194.1 and 856.2 cm⁻¹, suggesting that there are active modes in the hydrogenated species corresponding to CH in-plane bends. By assuming a $\nu_{CH}^{\text{H}}/\nu_{CH}^{\text{D}}$ ratio of 1.15, this would predict CH in-plane bands between about 1370 and 980 cm⁻¹, encompassing the bands between 1300 and 1250 cm⁻¹. As the spectrum of neutral phenanthrene has several absorptions in this region, it is plausible that the corresponding cation bands are screened. It has also been demonstrated that surprisingly large changes in intensity in a given vibrational mode can result from deuteration. Thus, it may also be that the features observed in this region for the perdeuterated cation are too weak to be visible in the fully hydrogenated cation. It is clear that deuteration alone is not sufficient to make band assignments in the larger PAH molecules.

The Pyrene Cation, $C_{16}H_{10}^+$. The mid-infrared absorption bands of the pyrene cation are shown in Figure 8. The cation band frequencies and relative intensities are listed in Table 1. These compare favorably with those reported by Vala et al. for this species. We have also observed seven additional bands attributable to the pyrene cation. These bands lie at frequencies of 1556.0, 1356.1, 1255.7, 1253.7, 1188.7, 953.8, and 867.0 cm⁻¹. The bands at 1255.7/1253.7, 1188.7, 953.8, and 867.0 cm⁻¹ are weaker than those previously reported, while the bands at 1556.0, and 1356.1 cm⁻¹ overlap with previously reported bands and may arise through site effects.

The mid-infrared spectrum is characterized by moderate to strong bands that are rather evenly distributed between 1600 and 600 cm⁻¹. Of the twelve bands between 1600 and 1200 cm⁻¹, the CC stretching region, three are very strong. We have detected two bands in the CH in-plane bend region, at 1102 and 976 cm⁻¹.

The 861 cm⁻¹ frequency of the detected out-of-plane CH bend is interesting. Pyrene has two rings with three adjacent H atoms and two rings with two adjacent H atoms. The 861 cm⁻¹ band falls at the high limit of the normal range expected for two adjacent H atoms. As the 690 cm⁻¹ band falls nearly 60 cm⁻¹ below the range for neutral PAHs with three adjacent H atoms on a ring, it is attributed to a CCC bend. Apparently, due to either screening or the inherent weakness of the feature, the absorption corresponding to the triply adjacent CH out-of-plane bend remains undetected. Using an analysis similar to that presented for the phenanthrene cation, we find that the integrated
Figure 8. Mid-infrared spectra of the pyrene cation isolated in an argon matrix at 10 K: (a) before photolysis; (b) after 8 min of in situ photolysis. The cation bands are labeled with their positions. Due to its low intensity and separation from other features, the cation band at 1102.0 cm⁻¹ is not shown.

The Perdeuterated Pyrene Cation, C₁₆D₁₀⁺. The mid-infrared spectrum of the fully deuterated pyrene cation is shown in Figure 9. These bands, together with the corresponding hydrogenated ion bands, are also listed in Table 2. The frequencies and assignments are discussed in depth by Vala et al.²⁷ We have found five additional bands attributable to the perdeuterated pyrene cation. These lie at frequencies of 1540.1, 1521.1, 1366.4, 1288.0, and 952.5 cm⁻¹.

The Benzo[e]pyrene Cation, C₂₀H₁₂⁺. The mid-infrared absorption spectrum of the benzo[e]pyrene cation is shown in Figure 10. The cation band frequencies and relative intensities are listed in Table 1. The evolution of the bands assigned to C₂₀H₁₂⁺ with Lyman-α photolysis is plotted in Figure 11. As with the phenanthrene cation, although this band correlation is looser for the weaker bands, the overall behavior of a sharp rise and leveling off warrants their assignment to the cation. The lines plotted without data points present the behavior of two photoproduct bands not associated with the benzo[e]pyrene cation.

There are two cation bands in the CH out-of-plane bending region. Three are expected since there are three different classes of chemical groups. The intensity of the doubly adjacent CH out-of-plane bending mode in the pyrene cation is reduced by a factor of 5 relative to the neutral molecule.

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region screened by the stronger neutral absorptions, these results indicate that the force field for the cation is sufficiently different that these out-of-plane motions not only are reduced dramatically in intensity but that the bond strengths and interactions are altered as well. If the 873.3 cm$^{-1}$ band does indeed arise from the doubly adjacent CH out-of-plane bending vibration, its intensity is suppressed by a factor of 9 from that of the neutral molecule. Furthermore, if the 848 cm$^{-1}$ band arises from the triply adjacent mode, it is remarkable both for its large blue shift from the neutral position ($\nu = 772$ cm$^{-1}$, $\Delta \nu = 76$ cm$^{-1}$) and for the fact that it is only suppressed by a factor of 2 relative to the neutral, much less than the 10–20× reduction typically observed in other PAHs.

The band at 685.7 cm$^{-1}$ is assigned to a CCC out-of-plane bend since, in contrast to the other ion CH out-of-plane bending bands, it falls 62 cm$^{-1}$ lower in frequency than the lowest frequency, strong neutral band and about 40 cm$^{-1}$ below the region normally associated with four adjacent H atoms, in the range normally associated with five adjacent H atoms.

The Benzo[ghi]perylene Cation, C$_{22}$H$_{12}$. The mid-infrared absorption bands of the benzo[ghi]perylene cation are shown in Figure 12. The cation band frequencies and relative intensities are listed in Table 1. The evolution of the bands assigned to C$_{22}$H$_{12}^+$ with Lyman-α photolysis is plotted in Figure 13. The two lines plotted without data points present the behavior of two photoproduct bands not associated with the benzo[ghi]perylene cation. Although the benzo[ghi]perylene cation spectrum is similar to that of the benzo[e]pyrene cation in that it is rich in the 1600–1200 cm$^{-1}$ range, it has a very different intensity pattern. This spectrum is dominated by four strong bands in a grouping, which is similar to that found for the interstellar emission spectrum, rather than by one strong band. The strongest band falls at 1578 cm$^{-1}$, followed by a gap of nearly 200 cm$^{-1}$, at which point three bands become important, namely, those at 1401.3, 1369, and 1324.4 cm$^{-1}$. As with phenanthrene and benzo[e]pyrene,
The band at 863.6 cm\(^{-1}\) has been detected in the CH out-of-plane bending region. This band does not fall in the CH in-plane bend region. This is puzzling as the force field for the cation is quite different from that of the neutral molecule. As a doubly adjacent CH out-of-plane bend, the intensity of this band is suppressed by a factor of 4 from that of the 846 cm\(^{-1}\) band of the neutral molecule.

The band at 639.6 cm\(^{-1}\) is assigned to a CCC out-of-plane bend producing a band slightly stronger than that arising from the CH bend of the cation. The ionization percentage calculated for this experiment is 11%, comparable to that obtained for the other PAHs in the series. Furthermore, the intensity of the detected CH out-of-plane bending feature is also comparable to that measured for the other cations. Thus, we conclude that the CC stretching modes in the coronene cation are anomalously weak, giving rise to the atypical intensity pattern.

Two bands fall in the CC stretching range (1579 and 1378.6 cm\(^{-1}\)), none are in the CH in-plane bend region, and one falls in the CH out-of-plane bend region (874.5 cm\(^{-1}\)). These three bands are labeled with their positions. The coronene cation spectrum is remarkable for its simplicity and unusual intensity pattern. Only three cation bands have been reported, and the intensity pattern is opposite those of all other PAH cations studied to date, with the CH out-of-plane bend producing a band slightly stronger than that arising from the CC stretching modes. It should be pointed out that the ionization percentage calculated for this experiment is 11%, comparable to that obtained for the other PAHs in the series. Furthermore, the intensity of the detected CH out-of-plane bending feature is also comparable to that measured for the other cations. Thus, we conclude that the CC stretching modes in the coronene cation are anomalously weak, giving rise to the atypical intensity pattern.

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TABLE 3: Near-Infrared Frequencies and Relative Intensities of Several Cations from the Thermodynamically Most Favored PAHs

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<th>pyrene</th>
<th>benzo[e]pyrene</th>
<th>benzo[ghi]perylene</th>
<th>coronene</th>
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<td>ν (cm⁻¹)</td>
<td>l_int</td>
<td>ν (cm⁻¹)</td>
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<td>10687</td>
<td>0.255</td>
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</table>

Figure 15. Near-infrared spectrum of the naphthalene cation isolated in an argon matrix at 10 K. This is the 0–0 band of the D₂ – D₀ transition.

Figure 16. Near-infrared spectrum of the phenanthrene cation isolated in an argon matrix at 10 K. This is the 0–0 transition. The mid-infrared spectrum of the same sample is shown in Figure 4.

Figure 17. Near-infrared spectrum of the pyrene cation isolated in an argon matrix at 10 K. This is the 0–0 band of the D₂ – D₀ transition. The mid-infrared spectrum of the same sample is shown in Figure 8 (argon matrix, 10 K).

Figure 18. Near-infrared spectrum of the benzo[e]pyrene cation isolated in an argon matrix at 10 K. This is the 0–0 band of the D₂ – D₀ transition. The mid-infrared spectrum of the same sample is shown in Figure 10.

were measured on the same matrices for which the mid-IR bands listed in Table 1 were measured.

The Naphthalene Cation, C₁₀H₈⁺. The near-infrared spectrum of the naphthalene cation is shown in Figure 15. The electronic spectrum of this cation has been described in detail previously.²⁵,²⁴

The Phenanthrene Cation, C₁₄H₁₀⁺. The near-infrared spectrum of the phenanthrene cation is shown in Figure 16, with frequencies and relative intensities listed in Table 3. The near-infrared spectrum of the phenanthrene cation isolated in an argon matrix has been discussed by Andrews et al.³⁵ and that in a neon matrix by Salama et al.³⁶ The near-infrared spectrum of the phenanthrene-d_{10} cation is not shown since it resembles that of perhydrophenanthrene. The frequencies of phenanthrene-d_{10} are listed in Table 2.

The Pyrene Cation, C₁₆H₁₀⁺. The near-infrared spectrum of the pyrene cation is shown in Figure 17, with frequencies and relative intensities listed in Table 3. The near-infrared spectra of the pyrene cation isolated in argon and neon matrices have been described previously.²⁷,³⁷

The Benzo[e]pyrene Cation, C₂₀H₁₇⁺. The near-infrared spectrum of the benzo[e]pyrene cation is shown in Figure 18, with frequencies and relative intensities listed in Table 3.

The Benzo[ghi]perylene Cation, C₂₂H₁₂⁺. The near-infrared spectrum of the benzo[ghi]perylene cation is shown in Figure 19, with frequencies and relative intensities listed in Table 3.
Infrared Spectroscopy of PAH Cations

Figure 19. Near-infrared spectrum of the benzo[ghi]perylene cation isolated in an argon matrix at 10 K. This is the $D_0 \rightarrow D_0$ transition. The strongest band corresponds to the 0–0 transition. The mid-infrared spectrum of the same sample is shown in Figure 12.

The Coronene Cation, $C_{24}H_{12}^+$. The near-infrared spectrum of the coronene cation is shown in Figure 20, with frequencies and relative intensities listed in Table 3. The near-infrared spectrum of the coronene cation isolated in an argon matrix has been reported by d’Hendecourt et al. 29

C. Integrated Absorbance Values. Integrated absorbance values [A (kilometers/mole)] for the mid-infrared cation bands of the PAHs phenanthrene, pyrene, benzo[ghi]perylene, and coronene are listed in Table 4. These A-values were determined using the near-infrared spectra presented earlier in the following way.

The integrated absorbance value, $A$, is given by the relation

$$ A = \int \tau \, dv/N^+ $$

where $N^+$ is the number density of absorbers, in this case PAH cations, and $\tau \, dv$ is the integrated area of the strongest near-infrared band (in centimeters$^{-1}$). $N^+$ (in moles centimeters$^{-2}$) is related to the area of this band by the equation

$$ N^+ = \tau \Delta \nu / \epsilon_\text{max} \Delta \nu $$

where $\epsilon_\text{max} \Delta \nu$ is the integrated molar absorption coefficient in liters mole$^{-1}$. Combination of eqs 1 and 2 gives

$$ A_{\text{NR}} = \epsilon_\text{max} \Delta \nu $$

where $\Delta \nu$ is the measured full width at half-height (FWHH) of the near-infrared band, and the value of the molar absorption coefficient $\epsilon_\text{max}$ was determined from the oscillator strength ($f$) using

$$ n_{\text{A}} f = 4.32 \times 10^{-9} \epsilon_\text{max} \Delta \nu $$

A straightforward unitary conversion gives the integrated absorbance value in units of kilometers/mole. The ratio of the integrated area of the mid-infrared bands to that of the near-infrared band in the same matrix could then be used to scale $A_{\text{NR}}$, giving the integrated absorbance values presented in Table 4. All integrated band areas were calculated in base 10 with software provided by Nicolet. The specific oscillator strengths, molar absorption coefficients, and near-IR/mid-IR band ratios used for our calculations are given in the following.

It should also be emphasized that our experiments indicate that the presence of CCl$_4$ significantly enhances the near-infrared band intensities with respect to the mid-infrared band intensities. For example, without CCl$_4$ in the matrix, the ratio of the integrated intensity of the 9005 Å phenanthrene electronic band to the 1565 cm$^{-1}$ vibrational band is 66. With CCl$_4$, the ratio is 139. The presence of CCl$_4$ enhances the ratio by about 2 times. For benzo[ghi]perylene, without CCl$_4$ in the matrix, the ratio of the integrated intensity of the 9181 Å electronic band to the 1349.2 cm$^{-1}$ vibrational band is 43. With CCl$_4$, the ratio is 59, an enhancement of about 40%. For the 7663 Å and 1401.3 cm$^{-1}$ bands of benzo[ghi]perylene, the enhancement is about 6 times ($\tau \, dv_{\text{NR}} / \tau \, dv_{\text{MR}} = 9$ without CCl$_4$ and 55 with CCl$_4$).

In ref 19, we showed that CCl$_4$ did not influence the absolute strength of naphthalene’s mid-infrared bands, but did not investigate its influence on any electronic transitions.

In view of the near- to mid-infrared intensity ratio alterations induced by CCl$_4$, A-values were determined for matrices without CCl$_4$ present. Ultimately, the accuracy of these values rests on the molar absorption coefficient used to determine the number of neutral PAH molecules lost upon photolysis, the technique generally used to determine the molar absorption coefficients and oscillator strengths of the near-infrared cation bands. 36

As this information is not available for all of the PAHs studied here, we have been unable to determine A-values for all the cations reported. The integrated A-values for the other PAH cation bands will be evaluated as this information becomes available.

This is an extremely important issue as there is wide variance in the literature on this subject. As with the naphthalene cation, and as discussed at length in ref 19, while the relative intensities of the cation bands reported here agree well in those cases where independently measured data are available (pyrene and coronene), there are discrepancies between the reported absolute intensities. However, while discrepancies still exist between our results and those from Professor Vala’s group at the University of Florida, they are much more modest than was the case for the naphthalene cation (2–4× as compared to 50×). Such variances may well be within the experimental uncertainties associated with the measurements, conservatively estimated to be less than a factor of 2. By far the largest source of error in these measurements stems from the lack of accurate molar absorbances for the near-IR bands. In order to reduce the relative uncertainty in these values, and to eliminate the possibility of matrix effects influencing the electronic transitions differently from the vibrational transitions, we are developing an alternative method to determine A-values. In view of the fundamental importance of these values, we present the details of our A-value calculations.

Naphthalene Integrated A-Values. The analysis described earlier was applied to the naphthalene cation in order to check
the technique. We had previously derived the naphthalene cation $A$-values using an entirely independent analysis.\textsuperscript{19} Use of the amended molar absorption coefficient for the 6754 cm$^{-1}$ band of the naphthalene cation,\textsuperscript{19} $\epsilon^*_{\text{max}} = 565$ L mol$^{-1}$ cm$^{-1}$, and FWHH = 45 cm$^{-1}$ in eq 3 gives $A_{\text{NR}} = 254$ km/mol. The ratio of the 1218 cm$^{-1}$ band of the naphthalene cation, the strongest mid-infrared feature, to the 6754 cm$^{-1}$ band of the naphthalene cation, is 1/66, giving an $A$-value for this band of $A_{1218} = 3.9$ km/mol. This is in excellent agreement with the value of 4.1 km/mol presented in ref 19, supporting the validity of this technique.

**Phenanthrene Integrated $A$-Values.** The oscillator strength for the 9005 Å band of phenanthrene is $6 \times 10^{-5}$.\textsuperscript{36} Use of eq 4 and a 53 cm$^{-1}$ FWHH gives $\epsilon^*_{\text{max}} = 338$ L mol$^{-1}$ cm$^{-1}$. Equation 3 gives $A_{\text{NR}} = 1.79 \times 10^2$ km/mol. Finally, the ratio of the strongest mid-infrared band, 1565 cm$^{-1}$, to the 9005 Å near-infrared band is 1/66, giving $A_{1565} = 2.7$ km/mol. The other mid-infrared $A$-values were scaled using the relative intensities shown in Table 1.

**Phenanthrene-d$_{10}$ Integrated $A$-Values.** The oscillator strength for the 8973 Å band of phenanthrene-d$_{10}$ is assumed equal to that for phenanthrene, i.e., $6 \times 10^{-5}$.\textsuperscript{36} This gives $\epsilon^*_{\text{max}} = 188$ L mol$^{-1}$ cm$^{-1}$ (FWHH = 95 cm$^{-1}$) and $A_{\text{NR}} = 179$ km/mol. The ratio of the 1528 cm$^{-1}$ mid-infrared band to the 8973 Å band in phenanthrene-d$_{10}$ is 1/42, giving $A_{1528} = 4.2$ km/mol. The other mid-infrared $A$-values were scaled using the relative intensities shown in Table 2.

**Pyrene Integrated $A$-Values.** The oscillator strength for the 7837 Å band of pyrene is assumed to be about $1 \times 10^{-3}$.\textsuperscript{39} This gives $\epsilon^*_{\text{max}} = 1420$ L mol$^{-1}$ cm$^{-1}$ (FWHH = 210 cm$^{-1}$) and $A_{\text{NR}} = 2.99 \times 10^2$ km/mol. The ratio of the 1356.1 cm$^{-1}$ mid-infrared band to the 7837 Å band of pyrene is 1/11, giving $A_{1356} = 282$ km/mol. The other mid-infrared $A$-values were scaled using the relative intensities shown in Table 1.

**Benzo[ghi]perylene Integrated $A$-Values.** The oscillator strength for the 7663 Å band of benzo[ghi]perylene is also taken as $1 \times 10^{-3}$.\textsuperscript{39} This corresponds to $\epsilon^*_{\text{max}} = 3830$ L mol$^{-1}$ cm$^{-1}$, based on a 78 cm$^{-1}$ FWHH. The integrated absorbance value is $A_{\text{NR}} = 2.99 \times 10^2$ km/mol. The ratio of the 1578.2 cm$^{-1}$ mid-infrared band to the 7663 Å band of benzo[ghi]-perylene is 1/8, giving $A_{1578} = 374$ km/mol. The other mid-infrared $A$-values were scaled using the relative intensities shown in Table 1.

**Coronene Integrated $A$-Values.** The oscillator strength for the 9607 Å band of coronene is $1.8 \times 10^{-3}$.\textsuperscript{40} This corresponds to $\epsilon^*_{\text{max}} = 5.84 \times 10^2$ L mol$^{-1}$ cm$^{-1}$, based on a 92 cm$^{-1}$ FWHH. This, in turn, gives $A_{\text{NR}} = 5.38 \times 10^3$ km/mol. The ratio of the 874.5 cm$^{-1}$ mid-infrared band to the 9607 Å band of coronene is 1/30, giving $A_{874} = 179$ km/mol. The other mid-infrared $A$-values were scaled using the relative intensities shown in Table 1.

**IV. Astrophysical Implications**

Comparisons between the interstellar emission spectra and the absorption spectra of PAHs in KBr pellets have shown important differences, which are used by some to draw conclusions about the nature of interstellar PAHs and by others to raise doubts as to the validity of the PAH hypothesis. One important difference concerns relative band intensities. As shown in Figure 21, which presents the IR emission spectrum from the Great Nebula in Orion,\textsuperscript{41} the interstellar 1610 cm$^{-1}$ (6.2 μm) feature and broad emission envelope centered near 1300 cm$^{-1}$ (7.7 μm), which are assigned to PAH CC stretching modes, are much more intense than the interstellar feature near 885 cm$^{-1}$ (11.3 μm), which is attributed to the CH out-of-plane bending mode. This behavior is just the opposite of that of PAHs in KBr pellets and in solvents where the strongest absorptions by far are those due to the out-of-plane CH bending vibrations (often 5–8 times greater than the average band strength in the CC stretch and CH in-plane bending region between about 1650 and 1100 cm$^{-1}$ (6 and 9 μm)). The relative intensities of the interstellar features have been taken by some as evidence that interstellar PAHs are severely dehydrogenated, in spite of the fact that atomic hydrogen is on the order of 10 000 times more abundant in most of the emission zones than the PAHs.\textsuperscript{42} To invoke such a high degree of dehydrogenation under these conditions seems unreasonable.

This difficulty is removed if the PAHs in space are largely ionized. With the exception of coronene, we have shown the following: First, the intensities of the bands in the 1650–1100 cm$^{-1}$ range, the CC stretching and CH in-plane bending regionüb
of PAH cations, are typically 2–5 times greater than those of the CH out-of-plane bending modes, which fall between 900 and 550 cm\(^{-1}\). Second, the CH out-of-plane bending modes of the cations are typically 5–20 times weaker than the out-of-plane bending modes of the neutral molecules, and third, the bands in the CH stretching and CH in-plane bending regions are generally an order of magnitude stronger than those for the corresponding transitions in neutral PAHs. The coronene exception is notable. The 874.5 cm\(^{-1}\) CH out-of-plane bend of the coronene cation is slightly stronger than the strongest CC stretch at 1579 cm\(^{-1}\).

This work and work presented by the University of Florida group on other PAH cations\(^{26-28}\) confirm the behavior predicted theoretically by Defrees et al.\(^{23}\) and Pauzat et al.\(^{24}\) and resolve one of the more troubling aspects of the interstellar PAH model. Other impediments to the full exploitation of the interstellar PAH hypothesis raised by comparing KBr pellet spectra to the theoretically generated by \textit{in situ} photoysis.

As illustrated by Figure 21, these are consistently the broadest and most intense members of the interstellar emission band family. The observation that the most intense interstellar emission features fall in the PAH CC stretching and CH in-plane bending regions is completely consistent with the composite emission spectrum expected from a mixture dominated by free, ionized, small PAHs.

V. Conclusions

The near- and mid-infrared spectra of the phenanthrene, pyrene, benzo[\(e\)]pyrene, benzo[\(ghij\)]perylenes, and coronene cations, isolated in argon matrices, are reported. Ions were generated by \textit{in situ} photoysis.

With the exception of coronene, the strongest mid-infrared absorption bands of these PAH cations fall between 1600 and 1200 cm\(^{-1}\), the CC stretching and CH in-plane bending region. These tend to be 2–5 times more intense than the bands due to the CH out-of-plane bending modes. The strongest bands tend to fall in groupings between 1600 and 1550 cm\(^{-1}\) and between 1400 and 1200 cm\(^{-1}\).

On the other hand, integrated absorbance values for the CH modes in the cation are strongly depressed with respect to the neutral molecules. Only weak new bands grow in the normally very strong CH out-of-plane bending region about 900 and 600 cm\(^{-1}\). The extent to which these out-of-plane motions are reduced in intensity and the bond strengths and interactions are altered is far more than would be expected from first principles by the removal of one electron from the \(\pi\) system. Similarly, as has been the case with previous PAH cations studied, no new features were found in the CH stretching region between 3200 and 2900 cm\(^{-1}\), implying that the CH stretching modes of the cation are comparable to or weaker than those for the neutral species. Theory predicts them to be weaker.\(^{24}\)

Most of the detectable new bands fall between 1570 and 1250 cm\(^{-1}\) and, thus, are assigned principally to the CC stretching modes. Those at the lower frequencies in this range may have some in-plane CH bending character as well.

The observation that the strongest cation bands coincide with CC stretching modes, behavior just the opposite of that of the neutral species, has important ramifications on the interpretation of the observed astronomical emission spectrum, which is dominated by very broad features in this region at 1310 and 1620 cm\(^{-1}\). Since this seems to be a general characteristic of PAH cations, it resolves one of the greatest discrepancies with the interstellar PAH hypothesis.

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**Note Added in Proof:** Recent unpublished calculations on PAH cation integrated absorbance values (\(A\)-values) by Langhoff show considerable PAH to PAH variation when compared to the values in Table 4. The experimental values tend to be no more than a factor of 2 greater than the theoretical values for pyrene, but they are 2–3 orders of magnitude smaller for phenanthrene. This points out the need for a better understanding of \(A\)-values in PAH cations.

**References and Notes**


(21) Hudgins, D. M.; Sandford, S. A. manuscript in preparation.
(22) Hudgins, D. M., Sandford, S. A., manuscript in preparation.
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