1. Introduction

Non-oxide ceramics are promising materials for a range of high temperature applications. Selected current and future applications are listed in Table I. In all such applications, the ceramics are exposed to high temperature gases. Therefore it is critical to understand the response of these materials to their environment.

The variables to be considered here include both the type of ceramic and the environment to which it is exposed. Non-oxide ceramics include borides, nitrides, and carbides. Most high temperature corrosion environments contain oxygen and hence the emphasis of this chapter will be on oxidation processes of the type (Shaw et al., 1987; Jacobson and Opila, 1999):

\[
\begin{align*}
2 \text{MB} + \frac{5}{2} \text{O}_2(g) &= 2 \text{MO} + \text{B}_2\text{O}_3 \\
\text{MC} + \text{O}_2 &= \text{MO} + \text{CO}(g) \\
\text{MN} + \frac{1}{2} \text{O}_2 &= \text{MO} + \text{N}_2(g)
\end{align*}
\]

However other corrosive environments, such as water vapor and molten salts, will also be discussed.

In oxidation studies, most of the oxidation properties are determined by the oxide film. Therefore ceramic materials discussed in this chapter will be classified according to the particular oxide film. In general, long-term applications require a protective oxide with a low oxygen diffusivity. The leading protective oxides for ceramics are silica (SiO\textsubscript{2}) and alumina (Al\textsubscript{2}O\textsubscript{3}). Other protective oxides, such as boria (B\textsubscript{2}O\textsubscript{3}), may be acceptable for short-term applications. Figure 1 illustrates the range of oxidation rates for different protective oxides (Shaw et al., 1987).

Most of the emphasis here will be on the silica forming ceramics—silicon carbide (SiC) and silicon nitride (Si\textsubscript{3}N\textsubscript{4}). Due to the proposed applications of silica-forming ceramics, behavior of these materials in complex corrosive environments will also be discussed. Next alumina forming
ceramics, aluminum nitride (AlN) and aluminum carbide (Al₄C₃), will be examined. The third class of ceramics to be considered are boria forming materials, such as boron nitride (BN). Next, transition metal carbides and nitrides which form the corresponding transition metal oxides will be discussed. Ceramic materials which form mixed oxide scales will also be considered. Finally, the oxidation of ceramic matrix composites (CMCs), which involves several parallel processes, is an important current area of research and will conclude the chapter.

High temperature corrosion of metals and alloys has been more extensively studied than high temperature corrosion of ceramics. Many of the same experimental techniques and basic principles apply to both metals and ceramics. In the case of metals, our fundamental knowledge is derived from studies on high purity metals and alloys. Oxidation/corrosion studies on metals are typically done on small coupons. For ceramics, high purity chemically vapor deposited (CVD) samples have recently become available for such studies. CVD coupons are available for most ceramics, however in some cases only ceramic powders are available. Experimental approaches and analyses of data must be modified for powders. Finally, an important chemical difference between metals and ceramics is the formation of gases in the case of ceramic oxidation, as illustrated in reactions (2) and (3). The escape of gases through the oxide film is an important issue.

2. Experimental Techniques

There are two objectives in high temperature corrosion studies: understanding the interaction of the ceramic and corrodent on a fundamental level and assessing the behavior of the ceramic in an actual application environment. The former requires a clean, well-defined environment designed to vary one or two parameters methodically (e.g. temperature and/or oxygen pressure). The latter requires a complex environment which accurately models the actual application. The emphasis of this chapter will be on fundamental oxidation mechanisms. However, when data are available, more complex environments, which come closer to the actual application, will also be discussed.

Most fundamental oxidation studies involve high purity oxygen in a well-controlled laboratory furnace at constant temperature. The amount of oxide can be monitored as a function of time via several techniques. The most common are optical measurements of oxide thickness and thermogravimetric monitoring of total weight gain. A thermogravimetric apparatus is illustrated in Figure 2. It is also possible to monitor oxidation kinetics via sensitive pressure gauges to monitor oxygen gas uptake and/or product gas release. Recently Gozzi and co-workers (1997) have developed a zirconia sensor to precisely monitor oxygen uptake in refractory carbide oxidation studies.
Kinetic information obtained via the above techniques is very useful in determining the rate controlling step of the reaction process. Scale formation which is linear with time obeys the following rate law and indicates an interfacial reaction or gas phase diffusion is rate controlling (Kofstad, 1988):

\[ x = k_i t \]  

Here \( x \) is the oxide scale thickness, \( k_i \) is the linear rate constant and \( t \) is the time. Scale formation which is parabolic with time obeys the following rate law and indicates diffusion through the oxide scale is rate controlling:

\[ x^2 = k_p t \]  

Here \( k_p \) is the parabolic rate constant. Another common observation is linear weight loss, which generally indicates that gas phase diffusion of volatile products through a boundary layer is rate controlling. For a flat plate coupon, this boundary layer-controlled flux, \( J \), is given by (Geiger and Poirier, 1973):

\[ J = 0.664 (Re)^{0.5} (Sc)^{0.33} \frac{\rho v D}{L} \]  

\[ Re = \frac{\rho v L}{\eta}; \quad Sc = \frac{\eta}{\rho D} \]  

Here, the dimensionless quantities \( Re \) and \( Sc \) are the Reynolds and Schmidt numbers, respectively. In these expressions, \( \rho \) is the density of the gas atmosphere above the sample, \( \rho_v \) is the density of the volatile species at the oxide/gas interface, \( v \) is the linear gas velocity, \( \eta \) is the viscosity, \( L \) is a characteristic dimension of the coupon and \( D \) is the diffusivity of the volatile species in the boundary layer.

The above rate laws are for coupons of well-defined geometry, however they can be adapted to powders as well. Jander (Hulbert, 1969) first developed a 'shrinking sphere' model for this purpose. First assume the radius of each particle, \( R \), decreases linearly with time:

\[ \alpha = 1 - \frac{R^3}{R_0^3} = 1 - \frac{(R_0 - kt)^3}{R_0^3} \]
1 - (1 - \alpha)^{1/3} = \frac{k}{R} t \quad (9)

Here \( R_0 \) is the initial radius of the particle and \( \alpha \) is the volume fraction reacted. A plot of \( 1 -(1-\alpha)^{1/3} \) vs. \( t \) gives a linear plot if an interface reaction is rate controlling. If diffusion in the solid is rate controlling then Equation (8) becomes:

\[ \alpha = 1 - \frac{R^3}{R_0^3} = 1 - \left( \frac{R_0 - k_p \sqrt{t}}{R_0} \right)^3 \quad (10) \]

\[ 1 - (1-\alpha)^{1/3} = \frac{k_p \sqrt{t}}{R} \quad (11) \]

Now a plot of \( 1-(1-\alpha)^{1/3} \) vs. \( \sqrt{t} \) will give a linear plot if diffusion is rate controlling. In theory, boundary-layer controlled diffusion rate laws can also be extended to powders, but diffusion through a powder bed is a more complex problem.

Once an oxide scale is established on the coupon or powder, its composition and morphology may be characterized by a variety of optical, electron optical, and x-ray diffraction techniques (XRD). In many cases the products are amorphous and other techniques are used to supplement XRD, such as chemical analyses of the scale alone. Scanning electron microscopy (SEM) has proven most useful, although transmission electron microscopy (TEM) is playing an increasingly important role as well.

The preceding discussion has dealt with studies of ceramics in pure environments at constant temperatures. In actual applications, high temperature materials are rarely at constant temperatures, but rather thermally cycled. Materials which show good isothermal resistance may not show oxidation resistance under thermal cycling. Thus thermal cycling oxidation tests are important in understanding overall oxidation behavior.

A variety of complex test rigs have been developed to simulate the behavior of high temperature materials in actual process environments. Consider first combustion conditions. Figure 3 shows calculated combustion gases for burning of a typical hydrocarbon fuel. Note that oxygen is only one constituent of the hot gas stream. It consists of \( \text{CO}_2, \text{H}_2\text{O}, \text{N}_2 \) and \( \text{O}_2 \) in fuel-lean environments and \( \text{CO}, \text{H}_2, \) and \( \text{N}_2 \) in fuel-rich environments. Burner rigs are very useful in simulating turbine engine conditions (Deadmore et al., 1978). Figure 4 shows a schematic of a burner rig. It is critical to control the parameters as tightly as possible in such a rig in order to allow comparison of one burner to another. Similar large scale rigs have been developed in order to simulate the environments in coal slagging (Ferber and Tennery, 1983) and various chemical process environments.
3. Oxidation of Silica Formers

To begin this section on the oxidation of silica formers, SiC and Si₃N₄, some of the properties of silica are briefly reviewed. Next, the oxidation of silicon is considered. While silicon is not a ceramic, its oxidation behavior nevertheless forms the basis for discussion of other silica formers. The section then proceeds from the oxidation mechanisms and rates obtained under the most ideal, clean conditions in pure oxygen to those conditions closer to actual application environments. These include effects of small amounts of impurities, sintering aids, effects of other oxidants, cyclic oxidation, effects of oxidation on strength, volatility effects, and hot corrosion.

3.1. Properties of Silica

Silica scales can exist in either an amorphous or a crystalline form. Both are composed of SiO₄ tetrahedra. The crystalline phase generally observed is cristobalite. Tridymite is sometimes observed when impurities are present, however, this phase will not be discussed further. Cristobalite can be envisioned as a regular lattice of six member rings composed of alternating silicon and oxygen atoms as shown in Figure 5a. The amorphous structure is an open and disordered structure in which seven or eight member rings can also be found as shown schematically in Figure 5b.

Under many conditions, transport of oxidant through the silica scale limits the growth rate of silica and the corresponding recession of the matrix material. Understanding the transport of oxidant or gaseous products through silica is therefore crucial to predicting the oxidation rates of silica forming materials. Oxygen transport in silica has been reviewed by Lamkin et al. (1992) in detail and is summarized briefly here. Two transport mechanisms can occur. First, molecular oxygen can move through the interstitial rings in the silica structure by a permeation mechanism. Second, charged oxygen species can move from oxygen lattice site to lattice site. The second process is expected to have a higher enthalpy since bond breaking is required.

While transport through silica is slower than any other oxide under ideal conditions as shown in Figure 1, and thus limits the growth rate of silica scales, the transport properties of this oxide are very susceptible to the environment and under any real application conditions these slow rates of silica growth are seldom observed.

3.2. Basic Oxidation Mechanisms of Silica Formers
3.2.1. Oxidation of Silicon

The oxidation of silicon has been studied extensively due to the need for formation of dielectric silica layers on silicon by thermal oxidation for semiconductor devices. Deal and Grove (1965) have written a classic paper on this subject in which gas phase diffusion rates, linear reaction rates of oxygen with silicon, as well as oxygen transport rates through the growing silica scale are all considered. It was found that silicon oxidation can be described by a linear parabolic rate law:

\[ x_o^2 + Ax_o = B(t + \tau) \]  

where \( x_o \) is the oxide thickness, \( t \) is the time, and \( \tau \) is a time correction for any pre-existing oxide. \( B \) is the parabolic rate constant, often referred to as \( k_p \) in the oxidation literature. \( B/A \) is the linear rate constant referred to as \( k_l \) in the oxidation literature. At short times Equation 12 reduces to \( x_o = (t + \tau)B/A \) which is the simple case for reaction controlled oxidation. At long times Equation 12 reduces to \( x_o^2 = Bt \), which is the simple case for solid state diffusion controlled oxidation.

In addition Deal and Grove (1965) show that oxygen transport through the silica scale occurs by molecular permeation of \( O_2 \) since \( B \) is directly proportional to the oxygen partial pressure. Also, the parabolic rates measured by Deal and Grove (1965) were in good agreement with rates measured by Norton for permeation of molecular oxygen through silica.

3.2.2. Oxidation of Pure SiC in Dry Oxygen

Since SiC oxidation also results in the formation of a silica scale, many analogies can be made with the oxidation of silicon. First, the oxidation of SiC can be modeled using the Deal and Grove treatment of linear parabolic kinetics. Ramberg et al. (1996) have studied the oxidation of pure SiC in the temperature range of 800-1100°C and shown that the linear reaction rate limits the oxidation rate in this temperature range. At temperatures of 1200°C and higher, the oxidation kinetics are generally considered purely parabolic (Costello and Tressler, 1986; Zheng et al., 1990a; Ogbuji and Opila, 1995). As first noted by Motzfeldt (1964), the parabolic rates of SiC should be 1.5 to 2 times slower than Si due to the additional oxygen consumed in the formation of gaseous products by the reactions:

\[ \text{SiC} + \frac{3}{2} O_2(g) = \text{SiO}_2 + \text{CO}(g) \]  
\[ \text{SiC} + 2 \text{O}_2(g) = \text{SiO}_2 + \text{CO}_2(g) \]
This has been shown to be true experimentally (Motzfeldt, 1964; Costello and Tressler, 1986; Ogbuji and Opila, 1995). Parabolic oxidation rates for pure materials, Si and SiC, as well as Si$_3$N$_4$ discussed below, are compared in Figure 6 on an Arrhenius plot.

Oxidation of silicon is typically studied to 1200°C, whereas SiC is useful to much higher temperatures. Thus, additional high temperature phenomena for SiC oxidation must be discussed. First, Zheng et al. (1990a,b) have shown by oxygen isotope tracer studies that significant exchange with the oxygen lattice occurs at temperatures of 1300°C and higher. In addition it was found that the dependence of the oxidation rate of SiC on oxygen partial pressure decreases as temperature is increased (Zheng et al. 1990b). Both these results suggest that at higher temperatures, increased transport of oxygen through the silica scale by motion of charged oxygen species on the oxygen network sites occurs. The overall oxidation rate, however, is still dominated by molecular permeation through interstitial sites up to temperatures of 1500°C (Ogbuji and Opila, 1995).

An additional complication occurring at higher temperatures is the crystallization of the silica scale. Thermally grown silica crystallites are spherulitic (Ogbuji, 1981) as shown in Figure 7. For very pure SiC, crystallization occurs at temperatures of 1300°C and higher (Ogbuji and Opila, 1995). Several studies report decreases in weight change with time and attribute this to crystallization of the silica scale (Narushima et al., 1989; Costello and Tressler, 1981; Fergus and Worrell, 1990). Recent work (Ogbuji, 1997) shows that oxygen transport through cristobalite is slower than through amorphous silica by about a factor of 30 at 1300°C. Under clean oxidation conditions, however, a continuous film of amorphous silica is found under the cristobalite. In addition, a film of amorphous silica is found between the cristobalite spherulites (Ogbuji, 1981). Thus, the extent of crystallization is not complete enough to result in a slowing in the observed oxidation rate of SiC at least for times less than one hundred hours (Ogbuji and Opila, 1995).

An upturn in the oxidation rate and activation energy of SiC at temperatures of about 1350 to 1400°C has been observed in several studies (Costello and Tressler, 1981, 1986; Fergus and Worrell, 1990; Fillipuzzi et al., 1992; Frisch et al., 1988; Ramberg, 1997; Singhal, 1976a; Zheng et al., 1990a) This has been variously attributed to crystallization of the scale (Costello and Tressler, 1981; Singhal, 1976a; Frisch et al., 1988) or a change in oxygen transport mechanism (Costello and Tressler, 1986; Fergus and Worrell, 1990; Zheng et al., 1990a; Fillipuzzi et al., 1992). However, it has also been shown that this increase in rate and activation energy is not observed for very pure materials in very clean conditions (Ogbuji and Opila, 1995, Ramberg, 1997) even when crystallization occurs. It has been suggested that this upturn in rates and energies may be
attributed to impurities either in the oxidation environment (Ogbuji and Opila, 1995) or in the SiC (Ramberg, 1997).

Oxidation of SiC at temperatures greater than 1500°C has been studied by Narushima et al. (1989) and Schiroky et al. (1986). Narushima et al. (1989) showed that the oxidation kinetics were parabolic at temperatures as high as 1675°C. The oxide scales formed were initially amorphous and after four hours primarily crystalline. After the scale crystallized the oxidation rates were about an order of magnitude slower. Schiroky et al. (1986) found that at temperatures of 1700°C and higher the oxide scale, being near its melting point, became amorphous. Under these conditions bubbles formed in the silica scale due to the presence at the interface of the gaseous products CO and or CO₂ at pressures greater than one atmosphere.

3.2.3. Oxidation of Pure Si₃N₄ in Dry Oxygen

The oxidation of silicon nitride typically occurs with parabolic kinetics (Choi et al. (1989), Du et al., 1989a; Schiroky et al. (1986); Ogbuji and Opila, 1995), is dependent on oxygen partial pressure, and is independent of nitrogen partial pressure (Du et al., 1989a). These three observations show that the oxidation kinetics are controlled by oxygen transport across a growing scale. However, it is also well known that the oxidation of pure Si₃N₄ in clean conditions at temperatures less than 1500°C are much lower than the oxidation rates of Si and SiC (Choi et al., 1989; Du et al., 1989a; Ogbuji and Opila, 1995) as shown in Figure 6. This has been attributed to the formation of a silicon oxynitride layer which forms below the silica scale during oxidation of Si₃N₄. The composition of the oxynitride layer has been described as both a discrete phase, Si₂ON₂, (Du et al., 1989a) and as a solid solution of SiO₂ and Si₃N₄ of variable composition, SiOₙNₙ (Ogbuji, 1995). This leads to the formulation of the oxidation reactions in two possible ways. For the discrete oxynitride phase the oxidation proceeds by a two-step reaction:

\begin{align}
\text{Si}_3\text{N}_4 + 3/4\text{O}_2(g) &= 3/2 \text{Si}_2\text{ON}_2 + 1/2 \text{N}_2(g) \\
\text{Si}_2\text{ON}_2 + 3/2 \text{O}_2(g) &= 2\text{SiO}_2 + \text{N}_2(g)
\end{align}

(15) (16)

While for the solution oxynitride, the oxidation reaction has been formulated by Sheldon as:

\[ \text{Si}_3\text{N}_4(1-x)\text{O}_6x = 3 8x \text{O}_2(g) = \text{Si}_3\text{N}_4(1-x-8x)\text{O}6x(8x) + 2 8x \text{N}_2(g) \]

(17)

The evidence that supports a discrete oxynitride phase is the discontinuity in the etching rate of the oxide scale formed on silicon nitride (Du et al., 1989a) as well as the contrast observed in TEM and SEM images of scale cross-sections (Ogbuji and Smialek, 1991). An SEM image of the
oxynitride/oxide formed on Si$_3$N$_4$ is shown in Figure 8. However, for a diffusion controlled oxidation process, it is assumed that equilibrium is achieved at the oxide/matrix interface. This assumption results in impossibly high nitrogen pressures at the Si$_2$O$_x$/N$_2$ interface (Du et al., 1989b). Other work by Ogbuji and Bryan (1995) shows an amorphous oxynitride layer with a continuous gradation in oxygen and nitrogen composition from SiO$_2$ to Si$_3$N$_4$. The accompanying model (Ogbuji, 1995) proposes that oxidation proceeds by a substitution of oxygen for nitrogen across the graded SiO$_x$N$_y$ network rather than at an interface. This is illustrated in Figure 9. The contrast observed in images of scale cross-sections is attributed to the interface between the amorphous subscale and the crystalline silica top layer.

The high temperature oxidation of pure Si$_3$N$_4$ has been studied by Hirai et al. (1980) as well as Narushima et al. (1993a). As for SiC, as the temperature is increased toward the melting point of silica, the scales formed on Si$_3$N$_4$ become more amorphous and the oxidation kinetics change from parabolic to linear.

3.3. Oxidation of SiC and Si$_3$N$_4$ in the Presence of Low Level Impurities

The oxidation rates of SiC and especially Si$_3$N$_4$ in clean environments are quite low. However, for any real application, impurities in the oxidizing environment are likely to be present. The effect of impurities on the oxidation rate of these materials can be quite significant. This issue has been explored experimentally in a number of ways. First, one of the most common contaminants, sodium, was ion implanted in both SiC (Zheng et al., 1992a) and Si$_3$N$_4$ (Zheng et al., 1992b) at levels between 50-1000 ppm. Upon oxidation, the Na was incorporated in the oxide scale. The oxidation kinetics remained parabolic for both implanted materials. The oxidation rates were enhanced by a factor of two in the case of SiC, but by an order of magnitude for the case of Si$_3$N$_4$. Nevertheless, the oxidation rate of Si$_3$N$_4$ remained lower than that of SiC. The activation energy for oxidation remained the same for SiC, but decreased for Si$_3$N$_4$. The enhanced oxidation in the presence of Na is attributed to the accelerated transport of oxygen through the silica scale which is due to incorporation of Na in the silica network and the formation of non-bridging oxygens. The incorporation of Na into the silica network is shown schematically in Figure 5c. In the case of Si$_3$N$_4$ the more dramatic enhancement in oxidation rates is attributed to the modification of the silicon oxynitride structure by sodium. While this phase is still present, in this case, its protective properties diminish substantially with incorporation of sodium.

Enhanced parabolic oxidation rates of SiC and Si$_3$N$_4$ due to contamination from furnace tubes used in oxidation experiments have also been observed (Choi et al., 1989; Opila, 1995; Ogbuji and Opila, 1995; Fox, 1998). This enhancement is also attributed to Na impurities from impurities in alumina tubes (Opila, 1995). In one case (Fox, 1998), the oxidation rates of SiC and Si$_3$N$_4$,
both very pure CVD materials, were found to follow parabolic kinetics, but at essentially the same rate. It was proposed that impurities from the alumina furnace tubes modified the structure of the silicon oxynitride so much that it no longer offered any additional barrier to oxygen transport than the silica scale. Backhaus-Ricault and Gogotsi (1995) have also shown that silicon oxynitride is absent in hot isostatically pressed additive-free silicon nitride. This result may also be due to minor amounts of impurities resulting from the HIP process.

Finally, moderate amounts of alkali impurities have been intentionally introduced into silica scales growing on both SiC and Si₃N₄ by vapor techniques (Pareek and Shores, 1991; McNallan et al., 1994; Sun et al., 1994). Alkali halides or alkali salts are vaporized in one portion of a furnace and transported with a carrier gas to the test specimen in another zone of the furnace. Depending on the activity of the alkali vapor species in the test, the oxide scale composition varied from 0.4 (Pareek and Shores, 1991) to 30 (Sun et al., 1994) mole percent alkaline oxide. Oxidation kinetics for the lowest levels of alkali impurity in the scale were parabolic, but elevated over rates observed for pure silica formers (Pareek and Shores, 1991). As the alkali content increased, oxidation rates increased and changed from parabolic to linear kinetics (Pareek and Shores, 1991; McNallan et al., 1994; Sun et al., 1994). The linear oxidation kinetics were ascribed to gas phase mass transfer controlled kinetics (McNallan et al., 1994) or interfacial corrosion reaction limited kinetics (Sun et al., 1994; Pareek and Shores, 1991).

In combustion environments, the deposition of alkali impurities in combination with impurities in common fuels, called hot corrosion, results in dramatically increased corrosion rates. This more severe corrosion process is discussed in Section 3.9.

3.4. Oxidation of SiC and Si₃N₄ Containing Sintering Aids

Sintering aids are often added to SiC and Si₃N₄ to promote densification of these materials. Sintering aids have several noteworthy effects on the oxidation of silica forming materials including increased oxidation rates, change in rate controlling mechanism, and alteration of the oxide scale structure.

First, for SiC materials, typical sintering aids are B, C, and/or Al. B and C are typically added to the SiC in the form of B₄C (Browning et al., 1987). For commercially available SiC containing B₄C at levels of about 0.5 wt% levels (Browning et al., 1987), parabolic oxidation rates are the same as those observed for CVD SiC (comparing results from Costello and Tressler, 1986 to Ogbuji and Opila, 1995). It has been shown that the boron segregates into the oxide leaving a depleted region in the SiC near the interface (Costello et al., 1981). Boron in turn is depleted near the oxide surface due to volatilization losses. While boron at this concentration has no apparent
affect on the oxidation rate, bubbles are formed in the oxide scale at temperatures between 1200 and 1300°C as partial pressures of volatile boron species approach one atm. (Fergus and Worrell, 1990).

Aluminum additions to SiC are made in the form of alumina additions typically in the amount of 4 wt%. This additive has been studied by Hinze et al. (1975a), Singhal and Lange (1975), Singhal (1976a), and Costello and Tressler (1981). It has been found that the parabolic oxidation rates increase with alumina content (Singhal and Lange, 1975) as shown in Figure 10. The alumina additives tend to enhance crystallization of the scale (Singhal, 1976a; Costello and Tressler, 1981). Glassy areas of the oxide scale are found to be enriched in Al as well as other impurities. Also, for high aluminum contents, mullite can be found in the oxide scale. Apparent activation energies for oxidation are higher than for pure SiC. This has been variously attributed to a mechanism controlled by outdiffusion of the oxidation product CO (Singhal, 1976a), crystallization of the silica scale (Costello and Tressler, 1981), or to incorporation of aluminum in the glass-forming network (Hinze et al., 1975a). Work by Costello and Tressler (1981) indicates that the oxidation rate is controlled by oxygen transport through the Al2O3-altered scales.

Densification of silicon nitride by conventional hot pressing also requires the use of sintering aids, typically MgO and Y2O3. As for SiC, the oxidation rates of these materials increase as the silica scale becomes less pure. Singhal (1976b) showed that silicon nitride, 1 wt% MgO oxidizes parabolically but that oxidation rates are higher than for pure Si3N4. The oxide was composed of cristobalite with segregated magnesium silicate (MgSiO3) grains on the surface. Similar yttrium silicate phases are shown in Figure 11. It was also demonstrated that oxidation rates were independent of both oxygen and nitrogen partial pressure. It was suggested that out-diffusion of the Mg cation to the silica surface controls the rate. In a series of elegant oxidation/reoxidation experiments, Cubiciotti and Lau (1978, 1979) demonstrated that diffusion of the sintering aid cation, Mg or Y, to the surface controlled the oxidation rate of these materials. The weight change kinetics for an oxidation/reoxidation experiment are shown in Figure 12 for Si3N4 containing MgO. After oxidizing and removing the scale, reoxidation follows the original kinetic curve showing that the oxidation rate is not dependent on silica thickness. Clarke (1983) has modeled this process assuming gradients in chemical potential of these cations are maintained by continuous formation of SiO2 and the consumption of cations in the scale to form silicates at the oxide surface.

It should be noted that if too much sintering aids are added to these materials, cation-rich silicates form in the grain boundaries during oxidation. The volume expansion that accompanies this phase formation results in destruction of the material (Patel and Thompson, 1988). Finally, for large amounts of phases added to SiC or Si3N4 the resulting material may be considered a composite
material. Oxidation of these materials may result in multiphase oxide scales. This type of behavior is considered in more detail in Section 7 below.

3.5. Oxidation of Silica Formers in Other Oxidants

3.5.1. Oxidation of Si, SiC, and Si₃N₄ in Water Vapor

Return now to the behavior of very pure materials and consider the effect of changes in the gaseous environment from dry oxygen to those environments more commonly found for high temperature applications of silica forming ceramics. Many proposed high temperature applications for SiC and Si₃N₄ are in combustion environments. Water vapor is always present in combustion environments since it is a product of the burning of hydrocarbon fuels. Water vapor is typically found at a level of about 10% of the total gas pressure, independent of fuel-to-air ratio, as shown in Figure 3 (Jacobson 1993). The effects of water vapor on oxidation of SiC and Si₃N₄ are therefore of significant interest and can be classified in three areas. First, water vapor increases the transport of impurity species to the silica scale. As already demonstrated, these impurities increase the oxidation rate of silica formers. Second, water vapor, as the primary oxidant species, intrinsically increases the oxidation rate of silica formers. Finally, silica scales are volatile in water vapor due to the formation of gaseous silicon hydroxide species. Under these conditions SiC and Si₃N₄ oxidize to form silica scales which simultaneously volatilize. The net result is paralinear kinetics and enhanced rates of matrix recession. Each of these effects will now be considered in more detail.

3.5.1.1. Water Vapor-Impurity Effects

Water vapor enhances the transport of impurities because of the stability of gaseous hydroxides. Many impurities found in real application environments, such as Na and K, readily form hydroxides at partial pressures greater than other alkali-containing vapor species. This is documented in several of the studies already referenced in Section 3.3 on low level impurities above. For example, Pareek and Shores (1991) have calculated that KOH formation from K₂CO₃ in water vapor increases the activity of potassium-containing species in the vapor phase as shown in Figure 13. The resulting increased activity of potassium in the silica scale increases the oxidation rate of SiC. Opila (1994) also noted that sodium impurities from furnace tubes contaminated SiO₂ scales on SiC at an increased level in the presence of water vapor presumably to due sodium hydroxide formation. Higher oxidation rates were observed for samples oxidized in 99.8% pure alumina tubes than in high-purity fused quartz tubes. Thus, oxidation rates of silica formers are enhanced in water vapor due to the formation of impurity-carrying hydroxide species which contaminate the silica scale.
3.5.1.2. Intrinsic Effects of Water Vapor on Oxidation

Water vapor also intrinsically increases the oxidation rate of silica formers compared to results obtained in dry oxygen. To begin this discussion, consider again the results of Deal and Grove (1965) for silicon. As in oxygen containing environments, linear-parabolic oxidation kinetics are again observed for silicon in water vapor. Deal and Grove (1965) show that the parabolic rate constant, \( B \) is proportional to the product of the effective diffusivity of the oxidant in the silica scale, \( D_{\text{eff}} \), and the solubility of the oxidant in the silica scale, \( C^* \), or:

\[
B = 2D_{\text{eff}}C^*/N
\]  

where \( N \) is the number of oxidant molecules incorporated in a unit volume of oxide. Based on work of Norton (1961) and Moulson and Roberts (1961) for the diffusivities and solubilities of oxygen and water in silica, as well as the parabolic oxidation rates of silicon measured by Deal and Grove (1965), the following summary can be made. The effective diffusivity of water in silica is about an order of magnitude less than oxygen as shown in Figure 14. The solubility of water in silica, however, is two to three orders of magnitude higher than oxygen. Thus, the parabolic rate constant for silicon in water vapor is one to two orders of magnitude higher than in oxygen as can be seen in Figure 15. Since the solubilities of water and oxygen in silica are relatively temperature independent, (Deal and Grove, 1965; Moulson and Roberts, 1961) the temperature dependence of the parabolic rate constant in each environment reflects the enthalpy of diffusion for that oxidant species in silica. Simply put, the enhanced oxidation rates for silica formers in water vapor are due to the high solubility of water in silica.

As already discussed, it is generally concluded that the oxidation of silicon in dry oxygen is controlled by the transport of molecular oxygen through the silica scale. Likewise in water vapor, Deal and Grove (1965) concluded that transport of molecular water in silica is the rate controlling mechanism for parabolic oxidation of silicon. Molecular transport limited parabolic oxidation for silicon is demonstrated in Figure 16. The plot of parabolic rate constant versus oxidant partial pressure for both oxygen and water vapor results in a line of slope one. In separate experiments, Deal and Grove (1965) also conclude that the oxidation of silicon by water vapor is independent of the carrier gas, based on results with oxygen and argon carrier gases, demonstrating that water vapor is the primary oxidant in water vapor/oxygen mixtures.

Again, comparisons between the oxidation of SiC and Si\(_3\)N\(_4\) in water vapor and the oxidation of Si can be made since they are all silica formers. The oxidation of SiC in water vapor has been studied extensively, however, in many of these studies, it is difficult separate effects from
impurities and those from water vapor. Several studies have been conducted for pure materials in clean water vapor-containing environments (Fung and Kopanski, 1984; Lu et al., 1984; Suzuki et al., 1982; Tressler et al., 1985). The results of each of these studies shows that SiC oxidation is accelerated in water vapor compared to dry oxygen. However, Opila (in press) has examined the oxidation kinetics of CVD SiC over a range of water vapor partial pressures, mainly in an oxygen carrier gas, but also in an argon carrier gas. These results show a number of interesting trends. First, at 1200°C bubbles are found in the amorphous scales for water vapor contents of 25% H₂O/O₂ and higher. The density of bubbles increases with water vapor partial pressure. An example of bubbles formed in a silica scale from oxidation of SiC in 90% H₂O/10% O₂ at 1200°C for 64h is shown in Figure 17. These bubbles are not observed at 1100°C nor are they observed for crystalline scales formed at 1400°C. Again, these bubbles are an effect of gaseous oxidation products formed by the reaction:

\[ \text{SiC} + 3 \text{H}_2\text{O(g)} = \text{SiO}_2 + \text{CO(g)} + 3 \text{H}_2(\text{g}) \]  

(19)

Also, Figure 18 (Opila, in press) shows that the variation of parabolic rate constant with water vapor partial pressure yields a power law exponent less than one. This indicates that some ionic oxidant species contributes to the oxidation process in addition to molecular water. Finally, crystallization of the silica scale is shown to be only weakly dependent on water vapor and more strongly dependent on an impurity nucleation mechanism.

Similarly, the oxidation of Si₃N₄ has primarily been studied for materials containing sintering aids. The only study where pure Si₃N₄ has been studied in a clean water vapor-containing environment is that by Choi et al. (1989). In this study, the oxidation rate is found to depend on carrier gas and in addition, a complicated dependence on water vapor partial pressure is observed. Oxidation rates in high water vapor partial pressures are lower than rates observed in lower water vapor partial pressures. This result may be due to silica volatility which is discussed in detail below. In addition, it is not yet established whether the silicon oxynitride layer is protective in water vapor environments.

To summarize, both SiC and Si₃N₄ oxidize more rapidly in water vapor very likely due to the increased solubility of water in the silica scale. The apparent dependence of the parabolic rate constant on water vapor partial pressure is not the same as for silicon oxidation at lower temperatures, indicating more complex oxidation mechanisms are operative. This is an area that requires additional work before an in depth understanding of the oxidation mechanisms in water vapor are understood.
3.5.1.3. Silica Scale Volatility in Water Vapor

Another important effect of water vapor on oxidation of silica-forming materials is its ability to readily react with the silica scale to form volatile silicon hydroxide species. There is a substantial amount of work in the 1950’s and 60’s that indicated silica was volatile in water vapor containing environments although the identity of the volatile species was not known (Glemser and Wendlandt, 1963). Silica volatility was also measured more recently by Cheng and Cutler, (1979). However, it is not until recently the following formation reaction has been identified in a transpiration study by Hashimoto (1992) and confirmed by Opila et al. (1997) using mass spectrometry:

$$\text{SiO}_2 + 2\text{H}_2\text{O}(g) = \text{Si(OH)}_4(g)$$  (20)

The effect of silica volatility on the oxidation of silica-forming materials has also gone largely unnoticed until recently. Silicon oxidation literature contains only a few comments that there is a slight deficiency of silica formed during thermal oxidation of silicon in water vapor (Deal, 1963; Ligenza, 1962). This can be attributed to the interest in the short-term oxidation behavior of silicon.

The importance of this effect has only recently been documented for silicon carbide (Opila and Hann, 1997; Robinson et al., 1999; and Opila et al., 1999). SiC exposed to water vapor-containing environments will oxidize to form a silica scale which simultaneously volatilizes to form Si(OH)$_4$ by the above reaction. The resulting kinetics are paralinear in nature. The change in oxide thickness for a paralinear process is shown in Figure 19a and is described by the following equation:

$$\frac{dx}{dt} = \frac{k_P}{2x} - k_I$$  (21)

Similar expressions are developed for surface recession and weight change (Opila and Hann, 1997; Tedmon, 1966). Surface recession kinetics are shown in Figure 19a while paralinear weight change is shown in Figure 19b. At long times a steady state process is achieved in which silica is consumed at the same rate it forms resulting in a linear recession rate of SiC. Paralinear oxidation kinetics of Si$_3$N$_4$ have also recently been documented by Fox et al. (1997b). A comparison of kinetics in 50% H$_2$O/O$_2$ for SiO$_2$, CVD SiC, and CVD Si$_3$N$_4$ is shown in Figure 20. Silica shows a simple linear weight loss. SiC shows a larger weight gain before the weight loss than Si$_3$N$_4$ since the ratio of $k_P/k_I$ is larger for SiC.
The volatility rate of silica, $k_v$, is shown by Opila and Hann (1997) in furnace studies, and Robinson and Smialek (1999) in burner rig studies, to be controlled by transport of the volatile silicon hydroxide species through a gaseous laminar boundary layer. Equation 6 gives the following expression for the volatility rate, $k_v$:

$$k_v = 0.664 \frac{Re^{1/2} Sc^{1/3} D_{PV}}{L}$$

where, in this case, $D$ is the interdiffusion coefficient of Si(OH)$_4$ in the boundary layer, $\rho_v$ is the density of the volatile species Si(OH)$_4$ in the boundary layer, and $L$ is the characteristic length of the test coupon. Silica volatility is dependent on pressure, temperature, and gas velocity. Rewriting Equation 22 in terms of these parameters results in:

$$k_v \propto \frac{V^{1/2}}{P_{total}^{3/2}} P_{Si(OH)_4}$$

The temperature dependence of the volatility rate is included in the $P_{Si(OH)_4}$ term and arises from the enthalpy for reaction of Equation 20. Also, it can be seen from Equation 20 that the partial pressure of Si(OH)$_4$ increases with the square of the water vapor partial pressure. This pressure dependence has been demonstrated by Hashimoto (1992) for silica. Inserting this relationship in Equation 23, while remembering that in combustion environments the water vapor partial pressure scales with the total pressure, results in:

$$k_v \propto V^{1/2} P_{total}^{3/2}$$

The pressure dependence, $P_{total}^{3/2}$, has been demonstrated by Robinson and Smialek (1999) for SiC recession as shown in Figures 21 and 22. In addition, Figure 22 shows the temperature and velocity dependence for silica volatility in combustion conditions. Many proposed applications of SiC and Si$_3$N$_4$ involve high pressure and high velocity combustion environments, thus silica volatility and the corresponding substrate recession become critical life limiting processes especially for thin components.

3.5.2. Oxidation of SiC in CO$_2$

Another common component of gaseous environments in which silica forming materials can be used is carbon dioxide. Again this gaseous species is a product of combustion of hydrocarbon fuels. The oxidation reaction for SiC in CO$_2$ can be given by the following:
SiC + 3CO₂(g) = SiO₂ + 4CO(g)  \hspace{1cm} (25)

Oxidation of SiC in this environment is much less studied than in oxygen and water vapor. Nevertheless, Antill and Warburton, (1970), Fitzer and Ebi, (1974) and Opila and Nguyen, (1998) all find that oxidation rates of SiC in pure CO₂ are much lower than in oxygen. This is shown graphically in Figure 23. While the rate limiting mechanism is unclear, it can be concluded that passive oxidation of SiC in CO₂ is negligible compared to the rates observed in oxygen or water vapor.

3.6. Cyclic Oxidation of Silica-Formers

Cyclic oxidation tests simulate the temperature cycles of actual engine mission cycles for applications such as aircraft engines. These tests subject the samples to cooling stresses which form due to thermal expansion mismatches between the substrate and the growing oxide scale. Thermal expansion curves for SiC, Si₃N₄, amorphous SiO₂, and cristobalite are shown in Figure 24. Assuming the scale is stress-free at temperature, then during cooling, compressive stresses would develop in the scale for amorphous silica, while tensile stresses would develop in the scale for cristobalite. Since most real applications of silica-formers result in at least partially crystalline scales, it is expected that the silica scale will be in tension on cooling. Cracks in the oxide may form. On reheating, the stress is relieved, the cracks heal at temperature resulting in protective behavior of the oxide. In contrast, for amorphous silica on silica-formers and typical superalloys, the scale buckles and spalls on cooling due to compressive stress formation. Thus, one expects for silica-forming ceramics, in contrast to superalloys, little difference between isothermal and cyclic oxidation rates. This has in fact been shown to be the case by Andrews and Riley, (1990), Maeda et al., (1989), Opila et al., (1993), and Park et al. (1998). Figure 25 shows typical cyclic oxidation behavior for a number of silica-forming ceramics exposed to 5h cycles at 1300°C for a total of 1000h. In general the silica growth shows fairly protective behavior. In addition, little difference in flexural strength between isothermally and cyclically exposed samples is observed (Andrews and Riley, 1990; Maeda et al., 1989).

3.7. Effect of Oxidation on Fracture Strength of Silica Formers

Oxidation of SiC and Si₃N₄ affects the fracture strength of these materials. After short oxidation times, increased fracture strengths are reported for SiC (Easier et al., 1981; Becher, 1983) and Si₃N₄ (Jakus et al., 1984). The initial formation of amorphous silica rounds existing surface flaws. At longer oxidation times, surface pits, bubble formation at the oxide/matrix interface, attack of secondary phases and grain boundary attack may all lead to flaws larger than those originally
present. Strength reductions are then observed (Easler et al., 1981; Becher, 1983, Jakus et al., 1984). The formation of these oxidation flaws is dependent on the sample purity and the oxidation temperature. Thus pure materials oxidized at lower temperatures exhibit less oxidation-induced strength degradation.

Tressler (1990) has also summarized the oxidation effects on crack growth and creep in SiC and Si₃N₄. Crack growth is generally controlled by grain boundary phases. Materials containing glassy grain boundary phases are more susceptible to crack growth than single phase materials. When oxidation increases the glassy grain boundary phase, the threshold stress intensity above which crack growth occurs decreases. Likewise, if oxidation results in a more creep resistant grain boundary phase due to crystallization or depletion of additive cations, the threshold stress intensity increases.

3.8. Volatility

3.8.1. Direct Vaporization

SiC and Si₃N₄ ceramics form vapor species via a variety of mechanisms. The simplest mechanism is direct vaporization of SiC or Si₃N₄. There are numerous vaporization studies of these compounds (Hincke and Brantley, 1930; Drowart et al., 1958; Zmbov et al., 1971). SiC vaporizes to Si(g), Si₂C(g), SiC₂(g), and Si₂(g) in vacuum; Si₃N₄ decomposes to Si(l) and N₂(g) in vacuum. However, in most applications there is some oxygen present in the atmosphere and the interaction of this oxygen with SiC and Si₃N₄ is important in determining the volatile species.

Volatility diagrams are the best way to represent the vapor species over a condensed phase (Wagner, 1958; Kohl et al., 1977; Heuer and Lou, 1990). Figure 26(a) is a simple volatility diagram for the vapor species over Si in varying oxygen potentials. Note the condensed phases in Figure 26(a). The vapor species on the left hand side are in equilibrium with Si(l) and the vapor species on the right hand side are in equilibrium with SiO₂(s). SiO(g) is the primary vapor species above Si(l) or SiO₂(s) for P(O₂) from about 10⁻²⁴ to 10⁻⁴ bar. Below 10⁻²⁴ bar Si(g) becomes dominant; above 10⁻⁴ bar SiO₂(g) becomes dominant.

Heuer and Lou (1990) also develop volatility diagrams for SiC and Si₃N₄, shown in Figures 26(b) and 26(c), respectively. The basic difference between these and Figure 26(a) is that the activity of Si is now less than unity. The Si₃N₄ diagram also includes a region of stability for SiN₂O₂. For clarity, these diagrams only show the dominant vapor species.
3.8.2. Oxide/Substrate Reactions

As the volatility diagrams indicate, in most cases a stable SiO$_2$ film is present on a SiC or Si$_3$N$_4$ substrate in all but the lowest oxygen potentials. At high temperatures, the silicon dioxide scale reacts with the SiC or Si$_3$N$_4$ substrate, generating volatile species as:

$$\text{SiC(s)} + 2 \text{SiO}_2(\text{s}) = 3 \text{SiO(g)} + \text{CO(g)} \quad (26)$$

$$\text{Si}_2\text{N}_2\text{O(s)} + \text{SiO}_2(\text{s}) = 3 \text{SiO(g)} + \text{N}_2(\text{g}) \quad (27)$$

Note that the reaction for Si$_3$N$_4$ has been written with the Si$_2$N$_2$O, although this intermediate layer may be variable in composition.

The formation of vapor species at the oxide/ceramic interface can be displayed with a predominance diagram and overlaid equilibrium vapor species. These are illustrated in Figures 27(a) and 27(b) for the Si-C-O and Si-N-O systems, respectively. At temperatures above 1800 K, substantial pressures can be generated at the SiC/SiO$_2$ interface (Jacobson et al., 1992). This is shown in a plot of P(Total) vs. temperature in Figure 28. Note that the pressure is quite significant for carbon-saturated SiC/SiO$_2$ and can be reduced with silicon-saturated SiC/SiO$_2$.

Figure 29 is a photograph of a coupon of SiC oxidized at 2073 K for 1 hr. Note the extensive bubbling. Some of this is due to the rapid oxidation rates and formation of CO(g) bubbles as an oxidation product. However some of the bubbling is very likely due to the SiC/SiO$_2$ interactions described above.

3.8.3. Active Oxidation

The most common route for silica-forming ceramics to vaporize is active oxidation. Silica is a unique protective oxide in that it has a stable gaseous suboxide. Its stability is comparable to that of SiO$_2$(g) (Chase et al., 1986):

$$\frac{1}{2} \text{Si} + \text{O}_2(\text{g}) = \frac{1}{2} \text{SiO}_2(\text{g}) \quad \Delta G^\circ_{1500 \text{K}} = -322 \text{ kJ/mol} \quad (28)$$

$$\text{Si} + \frac{1}{2} \text{O}_2(\text{g}) = \text{SiO(g)} \quad \Delta G^\circ_{1500 \text{K}} = -227 \text{ kJ/mol} \quad (29)$$

Consider a bare silicon surface. At low partial pressures of oxygen, there is sufficient oxygen to form SiO(g), but insufficient oxygen to form a stable SiO$_2$(s) film. This is termed 'active oxidation'
and can lead to substantial material consumption. At higher partial pressures of oxygen, a stable SiO$_2$(s) film forms. This situation is termed 'passive oxidation'.

The theoretical and experimental aspects of active oxidation have been nicely summarized in a recent overview (Narushima, et. al, 1997). There are several approaches to calculating the oxygen partial pressure at which the active-to-passive transition occurs. Before discussing this issue it is important to recognize that the active-to-passive transition and the passive-to-active transition are different. The former is attained by gradually increasing oxygen pressure over a bare Si, SiC, or Si$_3$N$_4$ to the point where SiO(g) is no longer dominant and a stable SiO$_2$ film forms; the latter is attained by gradually decreasing oxidant pressure over a SiO$_2$-protected surface to the point where SiO(g) formation dominates. These transitions occur at partial pressures of oxygen that differ by several orders of magnitude.

Wagner (1958) was the first to approach these transitions for pure silicon. Wagner's (1958) approach is based on the condition for Si/SiO$_2$ equilibria. This concept has been extended to SiC and Si$_3$N$_4$ (Singhal, 1976c). In the active region, a flux of oxygen atoms strikes the SiC surface and reacts to form SiO(g) and CO(g).

\[
\text{SiC} + O_2(g) = \text{SiO}(g) + \text{CO}(g)
\]  

(30)

As pressure is increased, the resultant flux of SiO(g) and CO(g) increase. Eventually, there is sufficient SiO(g) and CO(g) so that SiC and SiO$_2$ can co-exist in equilibrium as shown in the equations below. The exact condition for SiC and SiO$_2$ equilibrium has been an issue of some controversy. There are three possibilities:

\[
\begin{align*}
\text{SiC}(s) + 2 \text{SiO}_2(s) & = 3 \text{SiO}(g) + \text{CO}(g) \\
\text{SiC}(s) + \text{SiO}_2(s) & = 2 \text{SiO}(g) + \text{C}(s) \\
2 \text{SiC}(s) + \text{SiO}_2(s) & = 3 \text{Si}(l,s) + 2 \text{CO}(g)
\end{align*}
\]

(31)  
(32)  
(33)

In order to determine when the necessary P(SiO) and/or P(CO) is attained, the flux of O$_2$(g) striking the surface must be examined. The relevant diagram is Figure 30(a). The flux of oxygen is written in terms of oxygen atom flux for clarity in mass conservation.

\[
J_{o, \text{incident}}^{\text{incident}} = 2D_{O_2} \frac{P_{O_2}}{\delta_{O_2}RT}
\]

(34)

The resultant fluxes of SiO(g) and CO(g) emerging from the surface are given by:

\[
J_{SiO}^{\text{reaction}} = D_{SiO} \frac{P_{SiO}}{\delta_{SiO}RT}
\]

(35)
\[
J_{\text{co}}^{\text{reaction}} = D_{\text{CO}} \frac{p_{\text{CO}}^*}{\delta_{\text{CO}} RT}
\]

(36)

Here \( J \) is the flux of the subscripted species, \( D \) is the diffusion coefficient of the subscripted species, \( p \) is the pressure of the subscripted species, \( \delta \) is the thickness of the boundary layer, \( R \) is the gas constant and \( T \) is the absolute temperature. In the case of \( p^* \), this is the pressure calculated from the above equilibria. Thus the transition pressure is:

\[
p_{\text{O}}^{\text{transition}} = \left( \frac{D_{\text{SiO}}}{D_{\text{O}}^2} \right) \left( \frac{\delta_{\text{SiO}}}{\delta_{\text{SiO}}} \right) p_{\text{SiO}}^* = \left( \frac{D_{\text{SiO}}}{D_{\text{O}}^2} \right)^{1/2} p_{\text{SiO}}^*
\]

(37)

\[
p_{\text{CO}}^{\text{transition}} = \left( \frac{D_{\text{CO}}}{D_{\text{O}}^2} \right) \left( \frac{\delta_{\text{CO}}}{\delta_{\text{CO}}} \right) p_{\text{CO}}^* = \left( \frac{D_{\text{CO}}}{D_{\text{O}}^2} \right)^{1/2} p_{\text{CO}}^*
\]

(38)

The relationship between diffusivities and boundary layer thickness allows the above simplification. It is generally agreed that the second equation more closely describes the active-to-passive transition for SiC.

Next consider Wagner's (1958) calculation of the oxygen partial pressure at the passive-to-active transition. The passive-to-active transition begins with a stable SiO\(_2\) film. As oxygen potential is lowered, this SiO\(_2\) film decomposes and eventually active oxidation begins. Wagner (1958) derives the passive-to-active transition from the decomposition of SiO\(_2\):

\[
\text{SiO}_2(s) = \text{SiO}(g) + \frac{1}{2} \text{O}_2(g)
\]

(39)

When the partial pressure of oxygen is low enough, SiO\(_2\) decomposes and is no longer stable. This marks the beginning of active oxidation. For SiC, the active-to-passive and passive-to-active transition differ by about four orders of magnitude.

More recently, a purely thermodynamic approach has been developed, based on the volatility diagrams illustrated in Figures 26(a-c). Heuer and Lou (1990) modify the original volatility diagram of Wagner (1958) in this approach. The modifications include the representation of only the dominant volatile species and inclusion of an equimolar \( p(\text{O}_2) = \frac{1}{2} p(\text{SiO}) \) line. A lower and upper limit for the active-to-passive transition can then be calculated. (Heuer and Lou, 1990 and Schneider, 1995). The lower limit is from a modification of Turkdogan's theory of enhanced vaporization of metals in reactive gas streams (Turkdogan, 1963) and is illustrated in Figure 30(b). The Turkdogan theory is based on a metal vaporizing into a gas stream, reacting to form a metal oxide 'fog' in the gas, and condensing. In the case of Si, SiC, or Si\(_3\)N\(_4\) the vapor pressure of Si(g), Si-C gas species or Si-N gas species are much too low to have much effect; however the SiO(g) can form appreciable vapor pressures at low oxygen potentials. If a bare SiC surface is actively
oxidizing, SiO(g) diffuses outward from the sample. At the point when \( P(O_2) = \frac{1}{2} P(SiO) \), then the following reaction in the gas phase occurs:

\[
SiO(g) + \frac{1}{2} O_2(g) = SiO_2(\text{smoke}) \quad (40)
\]

The point at which \( P(O_2) = P(SiO) \) can be determined from the intersection of the equimolar line with the \( P(SiO) \) line in Figures 26(b-c). Some of the smoke will condense on the SiC surface and this is the onset of passive oxidation. This lower-limit transition is very similar to the passive-to-active transition derived by Wagner (1958), but does not contain the small corrections due to mass transport.

An upper limit for the active-to-passive transition can also be derived from a volatility diagram (Schneider, 1995). Going back to Wagner's criteria that enough SiO(g) must form to satisfy the equilibria:

\[
Si + SiO_2 = 2SiO(g) \quad (41)
\]

The underline indicates the activity of silicon is unity (pure silicon) or less (SiC or Si3N4). This equilibrium pressure of SiO is given by point T in the volatility diagrams of Figures 26(a-c). The amount of \( O_2(g) \) necessary to produce SiO(g) for the above equilibrium is determined from:

\[
Si + \frac{1}{2} O_2 = SiO(g) \quad (42)
\]

This upper limit for active oxidation is thus found at the intersection of the equimolar line (due to the stoichiometry of the above equation) and a horizontal line drawn from the SiO(g) pressure at point T. Note that this particular point has no physical meaning. It is simply a convenient way of reading the proper \( P(O_2) \) from the graph. Again, this is similar to the Wagner active-to-passive transition without the small correction for mass transport effects.

Figure 31 summarizes these approaches and some experimental data (Narushima et al., 1997). The experimental data exhibit considerable spread, likely due to morphology and mass-transport effects (Hinze and Graham, 1976) which are not fully incorporated into the current models.

The above discussion has focused on oxygen as the primary oxidant in determining the active-to-passive and passive-to-active transition. There are studies in the measurement of the active-to-passive transition in \( H_2O/H_2 \) (Kim and Readey, 1987) and \( CO_2/CO \) (Narushima et al., 1993b). It has been shown that these are equivalent to the oxygen cases with \( H_2O \) and \( CO_2 \) acting as the primary oxidant, respectively (Opila and Jacobson, 1995).
3.9. Molten Salt Corrosion

In combustion environments sodium impurities are very common. These can lead to highly corrosive salt deposits. An example is ingested NaCl in combustion air, which reacts with sulfur impurities in the fuel to form highly stable Na$_2$SO$_4$ (Jacobson, 1989):

$$2 \text{NaCl} (v) + \text{SO}_3(g) + \text{H}_2\text{O}(g) = \text{Na}_2\text{SO}_4(l) + 2 \text{HCl}(g) \quad (43)$$

The sodium may come from a marine environment in a gas turbine, process chemicals in an industrial furnace, or even fuel impurities. Under the appropriate conditions, Na$_2$SO$_4$ forms as a condensed phase, depositing on parts. Figure 32 shows the calculated dew points for Na$_2$SO$_4$ deposition. The rates of deposition are also critical and have been treated in detail (Stearns et al., 1983).

Deposition of corrosive salts is a continuous process. It is difficult to simulate this in a laboratory furnace. For this reason many salt corrosion studies are done by seeding a flame in a burner rig. Figure 33 is a comparison of SiC treated in a burner without salt and with a flame seeded with NaCl. Note the extensive corrosion in the latter case.

These processes are best interpreted with the acid/base theory of oxides (Rapp, 1986; Jacobson, 1993). Na$_2$SO$_4$ decomposes to Na$_2$O, which is the key reactant:

$$\text{Na}_2\text{SO}_4(l) = \text{Na}_2\text{O}(s) + \text{SO}_3(g) \quad (44)$$

The partial pressure of SO$_3(g)$ [or SO$_2(g) + \text{O}_2(g)$] effectively sets the activity of Na$_2$O(s). A high activity of Na$_2$O is a basic molten salt; a low activity of Na$_2$O is an acidic molten salt. Since SiO$_2$ is an acidic oxide, it is readily attacked by a basic oxide:

$$2 \text{SiO}_2(s) + \text{Na}_2\text{O} = \text{Na}_2\text{O}_2\text{SiO}_2(l) \quad (45)$$

This reaction describes the process which leads to the thick, glassy product in Figure 33. The solid protective SiO$_2$ layer has been replaced by a liquid, non-protective sodium silicate layer. The liquid layer permits rapid diffusion of oxygen inward and carbon monoxide outward, resulting in accelerated oxidation. This provides more silica for the above reaction and leads to rapid consumption of SiC.
The necessary conditions for the silica reaction with Na$_2$O follow from basic chemical thermodynamics (Jacobson, 1989). Assuming unit activities for SiO$_2$ and Na$_2$O$_2$(SiO$_2$), the threshold Na$_2$O activity [a(Na$_2$O)] for SiO$_2$ dissolution can be calculated:

$$a(\text{Na}_2\text{O}) = \exp(-\Delta G^\circ/RT) \quad (46)$$

Here $\Delta G^\circ$ is the free energy for reaction (45) above. At 900°C, the minimum $a(\text{Na}_2\text{O})$ is $10^{-11}$, meaning any activity greater than this will lead to silicate formation. Since the $a(\text{Na}_2\text{O})$ is set by the partial pressure of SO$_3$(g), this means that a P(SO$_3$) of $5 \times 10^{-5}$ bar or less will lead to silicate formation.

Some of these predictions have been verified in a burner using different sulfur levels in the fuel to set a partial pressure of SO$_3$(g). These results are shown in Figure 34 for quartz coupons heated in a low sulfur fuel (0.05%) and high sulfur fuel (0.5%). The first set of coupons was exposed to a low partial pressure of SO$_3$ and hence a more basic salt. Figure 34(a) shows there is clear evidence of more corrosion and chemical analysis revealed sodium silicate. The second set of coupons was exposed to a higher partial pressure of SO$_3$ and hence a more acidic salt. Figure 34(b) indicates less product formation and chemical analysis indicated no sodium silicate formation.

The preceding example illustrates the chemistry for Na$_2$O reactions with pure silica. Actual applications are more complex since additional elements may alter the basicity of the oxide. In the case of silica-forming ceramics, carbon is especially important. Carbon is a constituent of SiC, excess carbon may be added as a sintering aid, and carbon is often used as a fiber coating in fiber reinforced composites. In order to understand the effect of carbon, the following electrochemical cell was used to measure $a(\text{Na}_2\text{O})$ (Watt et al., 1981 and Jacobson, 1989):

$$\text{Pt}, \text{O}_2(r)|\text{ZrO}_2|\text{Na}_2\text{SO}_4, \text{SO}_3|\text{mullite}|10 \text{m/o Ag}_2\text{SO}_4, \text{Na}_2\text{SO}_3, \text{Ag} \quad (47)$$

Zirconia acts as a membrane for oxide anions and mullite acts as a membrane for sodium cations. The $\text{O}_2(r)$ is the oxygen reference and the $\text{Ag}_2\text{SO}_4/\text{Na}_2\text{SO}_3$ is the Na$_2$O reference. The net reaction is thus:

$$\text{Na}_2\text{O} + \text{Ag}_2\text{SO}_4 = 0.5 \text{O}_2(\text{g}) + \text{Na}_2\text{SO}_4 + 2 \text{Ag} \quad (48)$$

Using the Nernst Equation, $a(\text{Na}_2\text{O})$ can be determined from the cell voltage (E) and the free energy change for the above reaction. At 900°C, for example, $a(\text{Na}_2\text{O})$ is given by:
\[ E = 1.498 + 0.116 \log a(\text{Na}_2\text{O}) \] (49)

Figure 35 shows \( a(\text{Na}_2\text{O}) \) measurements taken for \( \text{Na}_2\text{SO}_4/(0.01 \ \text{SO}_2/\text{O}_2) \) and \( \text{Na}_2\text{SO}_4/(0.01 \ \text{SO}_2/\text{O}_2) + \text{carbon} \). Note, addition of carbon makes the melt swing strongly basic. This very likely occurs due to the formation of sodium sulfide and the subsequent formation of excess \( \text{Na}_2\text{O} \):

\[ \text{Na}_2\text{SO}_4(\text{l}) + \text{C} = \text{Na}_2\text{S} + 2 \text{CO}_2 \] (50)

\[ \text{Na}_2\text{S} + 3 \text{Na}_2\text{SO} = 4 \text{Na}_2\text{O} + 4 \text{SO}_2 \] (51)

The fact that carbon drives \( \text{Na}_2\text{SO}_4 \) more basic explains why \( \text{SiC} \) corrodes more readily in molten salts than \( \text{Si}_3\text{N}_4 \). It also explains why excess carbon in a ceramic promotes corrosion.

Once basic \( \text{Na}_2\text{O} \) begins to react with \( \text{SiO}_2 \), this process continues until all of the available \( \text{Na}_2\text{O} \) is consumed. In a combustion situation where sodium is ingested from the air, there is an essentially infinite supply of sodium and it is possible for a large amount of the ceramic to be consumed. In a situation where the supply of sodium is limited, the reaction will stop at some composition given by the liquidus \( \text{Na}_2\text{O}-\text{SiO}_2/\text{SiO}_2 \) line on the \( \text{Na}_2\text{O}-\text{SiO}_2 \) phase diagram. A layered \( \text{Na}_2\text{O}-\text{SiO}_2(\text{l})-\text{SiO}_2(\text{s}) \) structure will result (Mayer and Riley, 1978 and Jacobson, 1986). This is shown in Figure 36.

The reactive species, \( \text{Na}_2\text{O} \), can form from other salts as well. \( \text{Na}_2\text{CO}_3 \) readily decomposes:

\[ \text{Na}_2\text{CO}_3 = \text{Na}_2\text{O} + \text{CO}_2 \] (52)

In fact, \( \text{Na}_2\text{CO}_3 \) decomposes more readily than \( \text{Na}_2\text{SO}_4 \) and creates a more basic molten salt. McNallan et al. (1990) have shown that \( \text{Na}_2\text{O} \) can also form from the reaction of \( \text{NaCl} \) and water vapor:

\[ 2 \text{NaCl} + \text{H}_2\text{O} = \text{Na}_2\text{O} + 2 \text{HCl} \] (53)

In addition to \( \text{Na}_2\text{O} \), other oxides can potentially attack the protective \( \text{SiO}_2 \) film. Sea salt leads to \( \text{MgSO}_4 \) and \( \text{CaSO}_4 \) deposits as well as \( \text{Na}_2\text{SO}_4 \) deposits. This mixed sulfate can be described by molten salt solution models (Pelton, 1983; Bornstein, 1996) and leads to magnesium, calcium, and sodium silicates (Fox, 1998). Many fuels contain vanadium impurities. \( \text{V}_2\text{O}_3 \) is an acidic oxide and thus does not react directly with \( \text{SiO}_2 \). However, it accelerates \( \text{Si}_3\text{N}_4 \) oxidation, possibly due to the solubility of \( \text{SiO}_2 \) in \( \text{V}_2\text{O}_3 \) (Lamkin, et. al, 1990).
Similar principles apply to molten slags, which are complex mixtures of oxides. Here it is more difficult to define basicity, however a useful index has been the weight percent ratio of basic to acidic oxides. These are the basic and acidic oxides commonly encountered (Ferber and Tennery, 1983, 1984):

Basic Oxides: $\text{Na}_2\text{O}, \text{K}_2\text{O}, \text{MgO}, \text{CaO}, \text{Fe}_2\text{O}_3$

Acidic Oxides: $\text{SiO}_2, \text{Al}_2\text{O}_3, \text{TiO}_2$

The basic slags clearly degrade silica-forming materials more readily than the acidic slags. In a locally low partial pressure of oxygen, metal silicides may form (Ferber and Tennery, 1984).

It has been shown that molten salt-induced corrosion can lead to substantial attack of SiC and Si$_3$N$_4$. It is important to discuss the influence of this corrosion on microstructure and mechanical properties. Figure 37 illustrates the surface of as-received SiC (with B and C additives), the glassy coating after corrosion in Na$_2$SO$_4$, and then after cleanly removing the glassy coating with hydrofluoric acid. The last step reveals a highly pitted structure as opposed to the dense starting material. The reasons for this non-uniform attack are very likely due to structural discontinuities (e.g. grain boundaries, dislocations, etc.) in the SiC and bubble formation in the corrosion reactions (Jacobson and Smialek, 1986). Microstructural examination clearly indicates enhanced attack near grain boundaries, but also a correlation between the larger pits and bubbles in the silicate. A proposed mechanism for this is illustrated in Figure 38.

The strengths of monolithic ceramics are highly dependent on surface finish. Thus the pitting observed in Figure 37(c) is expected to lead to a strength decrease. Depending on the salt chemistry, which in turn determines the degree of pitting, room temperature strength reductions of up to 50% have been observed (Smialek and Jacobson, 1986; Jacobson et al., 1986; Fox and Smialek, 1990). In nearly all cases the fracture origins were the corrosion pits.

3.10. Refractory Oxide Coatings on Silica-Forming Ceramics

One possible solution to the corrosion problems of silica-forming ceramics is the application of a refractory oxide coating. Refractory oxides are generally more chemically inert than SiC and Si$_3$N$_4$ and this approach should combine the best properties of both materials. There are a number of important considerations in the application of a coating. The first is a good match between coefficients of thermal expansion (CTEs) and the SiC or Si$_3$N$_4$ substrate. Table II lists CTEs for SiC, Si$_3$N$_4$, and several refractory oxides. Note the remarkably close match of SiC and mullite.
Mullite coatings on SiC were first developed by researchers at Solar Turbines (Price, 1992) and further refined at the NASA Lewis Research Center (Lee et al., 1988, 1989). Critical issues are substrate surface roughening for adherence and application of fully crystalline mullite. Figure 39 shows a mullite coating on a SiC substrate.

Mullite-coated SiC exhibits substantially improved behavior in Na₂SO₄-induced corrosion as compared to uncoated SiC (Jacobson et al., 1996). This is due to the formation of higher melting sodium-alumino-silicates, as opposed to the low melting sodium silicates discussed in the section 3.9. above. However, in time, sodium penetrates the mullite coating and forms sodium silicate at the mullite/SiC interface. In the case of water vapor containing environments, mullite coatings offer limited benefits as the silica activity is only 0.4. Hence silica is readily volatilized from mullite by water vapor (Lee et al., 1989).

Other refractory oxide coatings are also candidates, however the problem of CTE mismatch must be considered (Table II). Graded coatings may be a possibility. Stability of the coating/substrate interface is also a critical issue.

4. Oxidation of Alumina-Forming Non-oxides

From Figure 1 it can be seen that after silica forming materials, those materials with the next most protective oxide scale would be alumina formers. Since alumina scales are less sensitive to impurities, AIN and Al₄C₃ appear, at first glance, to be ideal high temperature oxidation resistant materials. However, the recession rates in this figure are based on those observed for alumina forming metal alloys. Oxidation of both of these ceramic materials results in the formation of gaseous products which can alter the protective qualities of the alumina scale.

4.1. Oxidation of AIN

As previously mentioned the oxidation of AIN is complex since nitrogen gas is generated, as shown in the following reaction:

\[ 2\text{AIN} + \frac{3}{2} \text{O}_2(g) = \text{Al}_2\text{O}_3 + \text{N}_2(g) \]  \hspace{1cm} (54)

The literature contains many discrepant reports on both the rate controlling oxidation mechanism as well as the oxidation rate. Including only the results for dense bulk AIN, the oxidation behavior of AIN can be summarized as follows. In general, the oxidation kinetics of AIN in dry air (Sato et al., 1987) or dry oxygen (Jones and Scott, 1989; Kuromitsu et al., 1992; Opila et al., 1998) were parabolic. The oxidation rate was found to be similar to that for alumina formation on NiAl (Opila
et al., 1998) as shown in Figure 40. Thus the rate controlling mechanism was concluded to be the same, i.e., oxidation rates are limited by transport of aluminum and/or oxygen through the \( \alpha \)-alumina scale by short circuit paths such as grain boundaries. In laboratory air (Bellosi et al., 1993; Robinson and Dieckmann, 1994), wet air (Sato et al., 1987; Kim and Moorhead, 1994) or wet oxygen (Kuromitsu et al., 1992; Opila et al., 1998), however, the oxidation kinetics were much faster and generally linear in nature at temperatures below about 1200 to 1300°C. Typical weight change kinetics for AlN in 10% \( \text{H}_2\text{O} \)/90% \( \text{O}_2 \) are shown in Figure 41. At higher temperatures in wet air, kinetics were again parabolic (Sato et al., 1987; Kim and Moorhead, 1994). Oxidation rates were found to increase with water vapor content (Kuromitsu et al., 1992; Sato et al., 1987; Kim and Moorhead, 1994). It was concluded that oxidation in even small amounts of water vapor by the reaction

\[
2\text{AlN} + 3\text{H}_2\text{O(g)} = \text{Al}_2\text{O}_3 + \text{N}_2(g) + 3\text{H}_2(g)
\]  

results in the formation of micropores (Sato et al., 1987; Opila et al., 1998; Azema et al., 1991) 20-100 nm in diameter, at temperatures below 1200-1300°C. A typical fracture section of the alumina scale from AlN exposed in 10% \( \text{H}_2\text{O} \)/90% \( \text{O}_2 \) at 1200°C for 48h is shown in Figure 42. For these high porosities, transport of the oxidant through the micropores is rapid and oxidation is limited by the surface reaction between AlN and the oxidant, hence the observed linear kinetics. At higher temperatures the scale becomes denser due to sintering of the micropores and the observed oxidation kinetics become parabolic once again. Thus, the oxidation rate of AlN under very dry conditions is parabolic and quite low, but at water vapor contents typical of normal humidity levels in air, rapid linear oxidation occurs.

4.2. Oxidation of Al\(_4\)C\(_3\)

While the oxidation of Al\(_4\)C\(_3\) has not been studied nearly as much as AlN there are indications its oxidation resistance is similar. LeFort and Marty (1989) report that a nonprotective \( \alpha \)-\( \text{Al}_2\text{O}_3 \) scale is formed on Al\(_4\)C\(_3\) powder. The oxide scale was found to be severely cracked. Gas generation during oxidation may again explain the nonprotective nature of this oxide.

5. Boria forming Ceramics

Most borides actually form a mixed boria-metal oxide scale and will be discussed in Section 7.1.1. However, boron nitride (BN) and boron carbide (B\(_4\)C) form a pure boria scale. BN has been used for many years as a crucible material in vacuum and is of recent interest as a fiber coating for continuous fiber reinforced ceramic matrix composites. The most common form of BN is a
hexagonal, graphite-like structure. This discussion will be confined to oxidation of hexagonal BN.

In pure oxygen BN oxidizes as:

\[ 2 \text{BN}(s) + \frac{3}{2} \text{O}_2(g) = \text{B}_2\text{O}_3(l) + \text{N}_2(g) \]  

(56)

Some NO\(_2\)(g) has been detected (Pechentkoyskaya and Nazarchuk, 1981, Jacobson et al. 1999), but N\(_2\) is likely the primary gaseous product. In pure oxygen, oxidation of BN becomes significant above about 800°C (Podobeda et al., 1976).

Boria is not a desirable protective oxide. It has a low melting point (410°C) and tends to be very reactive with water vapor, forming a number of stable H\(_x\)B\(_y\)O\(_z\)(g) compounds (Chase et al. 1986). There are a number of oxidation studies of BN in the literature and the results indicate complex behavior (Lavrenko and Alexeev, 1986; Oda and Yoshio, 1993; Jacobson et al., 1999a). A recent study of BN indicates that oxidation behavior is quite dependent on structure of the starting BN material and water vapor content in the gas stream (Jacobson et al., 1999a). The first point is illustrated in Figure 43. Three types of BN—a CVD material deposited at a low temperature (Figure 43a), a CVD material deposited at a high temperature (Figure 43a), and a hot pressed material (Figure 43b)—are shown. Note the very different oxidation behavior. This is attributed to the differences in porosity and oxygen content of the various types of BN. These factors may provide rapid access for oxygen into the internal surfaces of the BN material, leading to more extensive oxidation. It also important to note that oxidation kinetics of BN are sensitive to crystallographic orientation, as observed by several investigators (Oda and Yoshio, 1993; Jacobson et al., 1999a).

Figure 44 shows the vapor species above B\(_2\)O\(_3\) in an oxygen/10% H\(_2\)O environment. Note that the hydroxide species exhibit substantial vapor pressures. Thus in water vapor containing environments, BN is expected to volatilize via B\(_2\)O\(_3\) and H\(_x\)B\(_y\)O\(_z\)(g) species.

Boron carbide (B\(_4\)C) is an interesting ceramic due to its high hardness and strength. It also forms a B\(_2\)O\(_3\) scale on oxidation and thus behaves similarly to BN (Gogotsi and Lavarenko, 1992). As with BN, oxidation of B\(_4\)C becomes significant above about 600°C and rapid oxidation occurs above 1200°C (Lavrenko et al., 1976). During oxidation, carbon is depleted in the B\(_4\)C region near the oxide/ceramic interface. This contributes to an oxidation-induced decrease in hardness after the B\(_2\)O\(_3\) scale is removed (Lavrenko et al., 1976).

6. Transition Metal Carbides and Nitrides
The transition metal carbides and nitrides have a number of important properties including hardness, wear resistance, and high temperature strength. However, these carbides and nitrides do not form oxides with low oxygen diffusivities. There is also generally a large thermal expansion mismatch between metal carbides and their corresponding oxides, leading to scale cracking and the formation of easily accessible paths for oxygen inward. In addition, porous oxide scales may form with rapid gas diffusion paths inward. These factors all lead to rapid oxide growth rates. The major features in the oxidation of these materials will be discussed.

Often there is a stable oxycarbide or oxynitride formed between the oxide layer and the corresponding carbide or nitride. In the case of the carbides, the carbon may be liberated in a number of different forms. In some cases, free carbon is observed in the oxide scale. In other cases, all the carbon is oxidized to CO or CO₂. In the case of the nitrides, the liberated nitrogen escapes as NOₓ or N₂.

6.1. Transition Metal Carbides

Now consider some specific carbides. Table III lists the stable transition metal carbides (Storms, 1967). The carbides to the right of these tend to form carbon/metal triple bonds, whereas the carbides to the left of these are of very low stability. Currently there is the most interest in ZrC, HfC and TaC as high temperature materials, due to their extremely high melting points. These materials are under consideration for rocket thrusters (Patterson et al., 1996) and ZrC is important in nuclear applications (Rama Rao and Venugopal, 1994).

Shimada and co-workers (1990, 1995, 1997) have extensively studied the oxidation of ZrC powders and single crystals at temperatures from 380-600°C at reduced oxygen pressures. Consider powder oxidation first. In the early stages of oxidation, oxycarbide forms. ZrC can have up to 60% of its interstitial positions occupied by oxygen. This is followed by nucleation of ZrO₂ on the oxycarbide. At higher temperatures, ZrO₂ grows rapidly and the associated volume expansion leads to cracking. Single crystals show a similar mechanism, shown schematically in Figure 45. At short times, an amorphous layer of ZrO₂₋ₓ grows parabolically (diffusion control). This is oxygen-deficient zirconia and the low oxygen potential allows free carbon to be deposited in the scale. Then the outer layers are oxidized to cubic zirconia and zone (2) starts to form linearly with time (interface control). This zone contains less free carbon and numerous cracks.

There are two important points from the above observations—the formation of free carbon in the oxide scale and the formation of cubic zirconia. In order for free carbon to precipitate, the oxygen pressure must be very low, as determined from the following reaction:
ZrC(s) + O₂(g) = ZrO₂(s) + C \tag{57}

Thus at 600°C, for the above reaction to be satisfied, P(O₂) is 4 \times 10^{-44} \text{ bar}. If oxygen-deficient ZrO₂ₓ is formed, these extremely low oxygen potentials are possible. The formation of cubic zirconia at these low temperatures is surprising since it is not expected to be stable until 1590°C in a pure Zr-O system (Levin and McMurdie, 1975). However it has been shown that the cubic phase is stabilized in oxygen-deficient zirconia with carbon (Wang and Liang, 1998).

At higher temperatures, ZrC oxidizes at much more rapid rates. Kuriakose and Margrave (1964) report that the type of ZrC they examined is destructively oxidized above 652°C. However, Voitovich and Pugach (1973) observe little weight change until about 700°C. Clearly the form of the starting ZrC material must play a role in its oxidation behavior. At 800°C Voitovich and Pugach (1973) observe a two phase scale with a distribution of carbon in ZrO₂. At 1000°C the scale becomes more porous, at 1100°C a thin layer of carbon is observed at the ZrC/ZrO₂ interface. They also report the presence of cubic zirconia, which apparently is stabilized by carbon. At 1200°C, sintering of the scale is observed, which is surprising at these temperatures. At this high temperature, no free carbon is observed in the scale.

Currently, there is a good deal of interest in HfC for rocket thrusters (Patterson, 1996). Shimada et al. (1997) have examined the oxidation of HfC from 600 to 900°C. There are many similarities to ZrC oxidation. A two zone structure is formed where the inner layer forms with diffusion-controlled kinetics and the outer layer forms with interface-controlled kinetics. Unlike ZrC oxidation, only the monoclinic form of the product oxide (HfO₂) is observed. Carbon is observed in both layers (Shimada, 1999). At temperatures of 1400-2100°C Bargeron et al. (1993) also observed a multi-layer product structure. In this case the outer layer contained no carbon.

Courtright et al. (1991) have observed the formation of porous oxide scales on HfC and PrC at high temperatures. At temperatures of 1200 to 1530°C, Holcomb and St. Pierre (1993) have applied a counter-current diffusion model for carbon oxidation through this porous oxide. This is based on the two-step oxidation of carbon in the HfC, as shown in Figure 46 and the following reactions:

\begin{align*}
\text{HfC(s)} + 3 \text{ CO}_2(\text{g}) & = \text{HfO}_2(\text{s}) + 4 \text{ CO(}\text{g}) \tag{58} \\
4 \text{ CO(}g) + 2 \text{ O}_2(\text{g}) & = 4 \text{ CO}_2(\text{g}) \tag{59}
\end{align*}

This two step oxidation is necessary due to the thermodynamic incompatibility of CO and O₂, which requires CO₂ to be the oxidant. They set up the necessary flux equations for each region and derive concentration profiles and parabolic rate constants. As expected for gas-phase
diffusion control, the temperature dependence is weak. Agreement between the model and measured rate constants is good.

The oxidation of other carbides, such as TiC, NbC, and TaC have also been studied. A two-layer oxide scale has also been observed on TiC (Voitovich, 1997; Shimada, 1996; and Shimada, 1999). Shimada (1996) has examined the oxidation of TiC powders as a function of temperature. Oxidation begins at about 300°C with the formation of TiC,O,1-x. At higher temperatures anatase is formed and CO₂ is evolved. As temperatures approach 900°C, oxidation of the powder is nearly complete and rutile is observed. More recently Shimada (1998) has studied oxidation of single crystal TiC isothermally and observed a two layer product scale with carbon in the oxide film. Voitovich (1997) reports several reaction stages with initial parabolic kinetics followed by linear kinetics. He reports an inner layer composed of TiO, Ti₂O₃, and TiO₂ and an outer layer of TiO₂.

Recently Desmaison-Brut and co-workers (1997) have examined the oxidation of TaC and Ta₂C between 750 and 850°C. The TaC oxidizes directly to Ta₂O₅, whereas the Ta₂C forms Ta₂O₅ and an intermediate layer of tantalum oxycarbide. They attribute the different reaction rates and activation energies observed for Ta₂C as opposed to TaC to the formation of oxycarbide on Ta₂C.

As discussed in the Section 2, Gozzi and co-workers (1997) have reported on a novel method of following oxygen uptake in a refractory carbide oxidation reaction with a yttria-stabilized zirconia sensor. They use low oxygen pressures (~10⁻⁵ bar) and their approach is extremely sensitive to small changes in oxygen. They are able to follow oxygen incorporation into a carbide lattice at low temperatures. Interestingly at temperatures of about 400°C, the refractory carbides (Table II) behave very similarly, indicating behavior is controlled by the carbon sublattice.

6.2. Transition Metal Nitrides

Formation of oxynitride layers between transition metal nitrides and their corresponding oxides is common (Douglass et al., 1996). Voitovich and Pugach (1975) attribute the better oxidation resistance of HfN as compared to ZrN to differences in the structure of the oxynitride layer. The oxynitride formed on ZrN contains more vacancies than that formed on HfN. They also observe cubic ZrO₂ at low temperatures (500°C)—apparently stabilized by the oxygen-deficient zirconia and nitrogen in the system. Desmaison et al. (1976) have also observed a mixture of monoclinic zirconia and a mixture of cubic or tetragonal zirconia. Their kinetics and morphology observations indicate the oxidation of ZrN is controlled by an interfacial chemical reaction.

There is interest in TiN both as a wear-resistant material and as a diffusion barrier in electronics.
TiN exhibits a range of non-stoichiometry, with different compositions showing different oxidation mechanisms (Desmaison et al., 1979). However, it is reported to have better oxidation resistance than ZrN or HfN (Voitovich and Pugach, 1975). The TiO₂ scale forms in layers with pores and tends to exfoliate (Desmaison et al., 1979), leading to complex oxidation kinetics. Tampieri and co-workers (1991) report varying oxidation kinetics, depending on the temperature range. At temperatures of 700-800°C parabolic kinetics suggest diffusion control and at higher temperatures diffusion through micro-cracks in the scale becomes important.

There is interest in TaN as a thin film resistor in the electronics industry. A range of oxidation kinetics and morphologies have been reported (Ibidunni, 1993; Bouzouita et al., 1983, 1987, 1988). In one case tantalum oxynitride is reported, in some cases porous scales are reported. These differences are partially due to oxidizing gases (O₂, CO₂/CO, or air) and may also be due to differences in the starting materials.

7. Oxidation of Ceramics Forming Multi-Component Scales

Thus far, the oxidation/corrosion behavior of materials which are nominally single phase, and form scales which are nominally single phase oxides has been discussed. In practice, composite materials are often desired to enhance mechanical properties or to improve other properties of the materials such as thermal conductivity, electrical conductivity, as well as oxidation resistance. In these cases, multi-phase materials are developed to enhance the desired properties. When these materials are subjected to high temperature oxidizing environments, the resulting oxidation products are also more complex. In general, oxidation results in three classes of scales: mixed oxides, compound oxides, or solution oxides. Examples of each will be discussed below. In addition, fiber reinforced composites are a special class of composite materials in which an interphase material is required between the fiber and the matrix to provide the required fracture toughness. The unique issues with these materials will also be discussed.

7.1. Mixed Oxide Scales

7.1.1. Transition Metal Borides

This section reviews the oxidation behavior of HfB₂, ZrB₂, and TiB₂. HfB₂ and ZrB₂ as well as their oxides, HfO₂ and ZrO₂, are ceramics with capability for short-time use at ultra-high temperatures for applications such as rocket engine components and re-entry shields. TiB₂ is of interest for its high hardness, high melting point, and especially its good electrical and thermal conductivity. Applications of TiB₂ include coatings resistant to liquid metals and wear-resistant parts. Unlike transition metal carbides and nitrides, the oxidation of transition metal borides can result in a two
phase scale consisting of boria in addition to the transition metal oxide. In general, at low
temperatures and short times the borides have lower parabolic oxidation rates because transport
of oxidant now occurs through a liquid boria phase rather than through the gas phase. The
oxidation resistance of several transition metal diborides have been ranked by Kaufmann et al.
(1967) in order of decreasing oxidation resistance as follows: HfB₂, ZrB₂, TiB₂, TaB₂, NbB₂.

The oxidation of borides is complex for a number of reasons. First, consider the formation of
liquid boria. It is well known that boria is very volatile, forming both B₂O₃(g) (Soulen et al., 1955)
and volatile hydroxide species (Margrave, 1956; Meschi et al., 1960). At temperatures where the
volatility of boria is readily observed, oxidation kinetics will therefore follow a paralinear rate law in
which oxidation weight gains occur simultaneously with boria volatilization. At temperatures below
1000°C, both Irving and Worsley (1968) and Courtois et al. (1993) find protective oxidation
behavior of TiB₂ attributed to a liquid boria film. A layered oxide is observed at these temperatures
and is shown schematically in Figure 47. At temperatures higher than 1000°C, Courtois et al.
(1993) observe paralinear behavior due to boria formation and volatilization. At these
temperatures, Irving and Worsley (1968) observe that boria is volatilized as soon at is formed and
the titania layer controls the oxidation kinetics. Tripp and Graham (1971) developed an
experimental TGA technique for oxidation of ZrB₂ which measures each component of the
paralinear kinetics: consumption of oxygen, weight change due to boria volatility, as well as total
sample weight change. Typical results are shown in Figure 48. At higher temperatures, the boria
will volatilize as soon as it is formed leaving only a porous transition metal oxide scale, just as in
the case of carbides and nitrides. Parabolic kinetics may still be observed, but the rate limiting
step is gas phase diffusion in the pores rather than diffusion through a condensed phase. In other
cases, at high temperatures and low pressures, boiling boria may rupture the transition metal
oxide scales resulting in linear oxidation rates controlled by the chemical reaction at the
scale/matrix interface.

Other complications for oxidation of transition metal borides at high temperatures, where boria no
longer offers protection, result from the properties of transition metal oxides. The monoclinic to
tetragonal phase transition for zirconia and hafnia is accompanied by a volume changes which
results in less protective scales (Berkowitz-Mattuck, 1966; Kaufmann et al., 1967). In addition, the
titania grown on TiB₂ is usually highly textured and cracked resulting in linear kinetics at longer
times (Tampieri and Bellosi, 1993).

7.1.2. Oxidation of Complex Systems Resulting in Mixed Oxide Scales
While any number of complex mixtures of non-oxide ceramics can be formulated and studied with regard to their oxidation resistance, one hopes to obtain a material whose oxidation properties are dominated by the most protective oxide in the system. However, for this to occur, a continuous layer of this oxide must be formed in the scale, otherwise the oxidation rates are controlled by less protective oxides, gas phase diffusion, or the linear reaction rate of oxidant with the matrix. Only several systems of particular note will be discussed here. In the first of these systems, Si-containing material, e.g. SiC, additions are made to less oxidation-resistant materials in hopes of improving the oxidation behavior by forming silica layers. In other systems, additions are made to SiC or Si$_3$N$_4$ to improve properties such as wear resistance. The oxidation behavior is then tested to ensure that these additions do not degrade the matrix oxidation resistance to unacceptable levels.

7.1.2.1. HfB$_2$/SiC

The first system of interest is the HfB$_2$-SiC system. While HfB$_2$ is the most oxidation resistant transition metal boride, oxidation rates are still quite high, especially at high temperatures where boria volatilizes as soon as it is formed. Claugherty et al. (1968) explored a variety of ways for improving the oxidation resistance of HfB$_2$. It was found that SiC additions increased the oxidation resistance of HfB$_2$ with best results obtained for 20 vol % SiC. The resulting scale was more adherent than that formed on pure HfB$_2$. The increased oxidation resistance of borides with SiC additions was explored further for ZrB$_2$ by Tripp et al. (1973) and for HfB$_2$ by Hinze et al. (1975b). At short times a transient rapid oxidation rate occurred as short circuit diffusion through ZrO$_2$ or HfO$_2$ controlled the oxidation rate. After this transient, parabolic kinetics were observed in both cases with the reaction rate controlled by diffusion of oxidant through a protective silica layer in the scale. Microscopy revealed a layer of silica on top of porous HfO$_2$. At temperatures below 1300°C, boria is found in the scale in the form of a borosilicate glass. Boria volatility is suppressed by the silica. At higher temperatures, selective vaporization of the boria occurs from the borosilicate glass.

7.1.2.2. Ti$_3$SiC$_2$

Another material of more recent interest is Ti$_3$SiC$_2$. This material has unique properties such as easy machinability, high thermal shock resistance, good ductility and high yield points at high temperatures. It has only recently been formulated as a single phase compound. The oxidation resistance has been studied by Barsoum et al. (1997). Parabolic oxidation is observed. Again a layered structure is observed with an inner layer of silica and titania and an outer layer of oriented rutile. The oxidation rates of this material are quite high suggesting diffusion control by the titania
Barsoum et al. (1997) suggest that it may be possible to improve the oxidation resistance of Ti$_3$SiC$_2$ by adding TiSi$_2$ to form a continuous silica layer upon oxidation.

### 7.1.2.3. TiC or TiN additions to Si$_3$N$_4$

Composite materials based on Si$_3$N$_4$ with additions of TiC or TiN have been developed to increase wear resistance. The effect of the TiC (Opsommer et al., 1998) and TiN (Bellosi et al., 1990) additions on the oxidation behavior of these Si$_3$N$_4$ composites has been studied. In both cases, these additions have adverse effects on the oxidation resistance of Si$_3$N$_4$ especially at low temperatures. At low temperatures Bellosi et al. (1990) have shown that the Si$_3$N$_4$-TiN oxidation rate is dominated by the rate of reaction of the TiN phase with oxygen. At intermediate temperatures oxidation is rapid and parabolic, dominated by diffusion in the titania. At temperatures of 1200°C and higher, rates are parabolic and transport in the silicate phase in the scale limits the oxidation rates. In all three regimes, oxidation rates increase with the amount of TiC or TiN additive. At high temperatures, incorporation of titanium into the silicate scale as well as formation of TiO$_2$ may contribute to these enhanced oxidation rates (Opsommer et al., 1998). Therefore, additions of TiC or TiN to Si$_3$N$_4$ must be optimized so that wear resistance can be increased enough while not degrading the oxidation resistance of the Si$_3$N$_4$ matrix to unacceptable levels.

### 7.2. Composites which form Compound Oxides

In the preceding discussion on oxidation of composites, each matrix phase reacted with oxygen to form its corresponding oxide. No interactions between matrix or oxide phases were reported. In this section, a few examples are discussed in which the composite scale contains oxide compounds formed from more than one component of the matrix phases, e.g. mullite, zircon, or aluminum borates. This is not intended to be a complete review, but an illustration of the kinds of effects compound formation can have on oxidation behavior. For example, formation of the compound oxide scales may result in higher oxidation rates if the compound oxide is more permeable to oxygen, or likewise, lower oxidation rates if the compound oxide is less oxygen permeable. The volatility of scale components can also be reduced by the formation of compound oxides as will be discussed below for aluminum borates.

#### 7.2.1. SiC-Reinforced Al$_2$O$_3$

SiC reinforcement of alumina results in significant improvements in mechanical properties such as toughness, creep strength, wear, and thermal fatigue compared to monolithic alumina, however,
these composites suffer from degradation due to thermal oxidation. Luthra and Park (1990) and Deqing and Lopez (1994) have studied the oxidation behavior of these composites. Parabolic oxidation is observed but at rates three orders of magnitude higher than for pure SiC indicating the reaction product is not pure silica. Oxidation rates increase with amount of SiC reinforcement. The observed reaction products are not the predicted equilibrium phases. For low SiC contents where alumina and mullite are predicted, alumina, mullite, and a glassy aluminosilicate are present in the oxidized zone. For higher SiC contents where silica and mullite are predicted, primarily mullite and glassy aluminosilicate phases are present after oxidation. Luthra and Park (1990) propose that oxygen diffusion through alumina controls the oxidation rate for low SiC concentrations. At higher concentrations the oxidation rate is controlled by oxygen diffusion in and/or diffusion of CO out through mullite or the glassy aluminosilicate phase. The composites form a layered structure in which a thin layer containing partially-reacted SiC particles is found beneath a completely reacted zone. A schematic of this structure can be found in Figure 49. TEM work on a similar composite (also containing zirconia in the alumina matrix) shows amorphous carbon and graphite (Backhaus-Ricault, 1991) are also present in the inner partially reacted SiC oxidation zone. The outer fully reacted zone contains porosity which is attributed to bubbles from CO evolution (Luthra and Park, 1990; Deqing and Lopez, 1994; Backhaus-Ricault, 1991). Thus, the formation of mullite and aluminosilicate glassy phases result in enhanced rates of transport of oxygen to the SiC reinforcement phase and corresponding rapid oxidation rates.

7.2.2. \( \text{Si}_2\text{N}_2\text{O-ZrO}_2 \)

The porous \( \text{Si}_2\text{N}_2\text{O-ZrO}_2 \) composite material is another interesting example of the effects of compound formation on oxidation behavior. O'Meara et al. (1995) and Heim et al. (1995,1997a,b) have documented the case in which \( \text{Si}_2\text{N}_2\text{O} \) in this composite oxidizes to form \( \text{SiO}_2 \) at increasing rates up to temperatures around 1300°C. At temperatures higher than this, it is proposed that zircon, \( \text{ZrSiO}_4 \), or a mixture of zircon and an amorphous zircon silicate forms which effectively acts as a barrier to further oxidation. Schematics of the oxidation structure below and above the zircon formation temperature are shown in Figure 50. Courtright et al. (1991) have summarized the oxygen permeability of a number of oxides. In this summary the oxygen permeability of zircon is reported to be higher than that of silica. Therefore, the results of Heim et al. (1995, 1997a,b) are inconsistent with the permeability of zircon reported by Courtright (1991). Nevertheless, the concept of increasing the oxidation resistance of a material by choosing a composite which forms a less oxygen permeable compound oxide is viable.

7.2.3. \( \text{TiB}_2-\text{Al}_2\text{O}_3 \)
From the discussion above on TiB₂ oxidation, it is known that the oxidation resistance of TiB₂ relies on the formation of a protective liquid boria scale. At temperatures high enough to volatilize the boria directly, linear oxidation rates are sometimes observed due to the nonprotective nature of the titania scale. The addition of Al₂O₃ or AlIN to the TiB₂ results in better oxidation resistance (Tampieri and Bellosi, 1993; Schneider et al., 1996). At intermediate temperatures, 700-900°C, Al₄B₂O₉ is formed which slows the oxidation rate (Tampieri and Bellosi, 1993). In addition, boria volatility is reduced since it is entrapped in the aluminum borate phase (Schneider et al., 1996). At higher temperatures a less boria-rich phase is formed, Al₁₈B₄O₃₃, and higher oxidation rates are observed in the case of the AlN composite (Schneider et al., 1996). This system demonstrates the possibility to suppress volatility by compound formation during oxidation.

7.3. Composites which form Single Phase Oxide Solid Solutions

While few examples of this phenomenon exist and little work has been done in this area, it is possible to affect the oxidation behavior of materials which form a single phase oxide solid solution. For example, upon oxidation, a HfC/ZrC composite would form a solution of HfO₂/ZrO₂. Use of this type of composite, and choosing the HfC to ZrC ratio might be useful to control the monoclinic to tetragonal transformation temperature of the resulting scale or other properties of the matrix or scale.

8. Ceramic Matrix Composites

Ceramic matrix composites (CMCs) are perhaps the most promising ceramics today for high temperature structural applications. These CMCs include a wide range of multiphase ceramics with the additional phase added to impart improved fracture toughness. Examples of this include whiskers, particulates, and continuous fibers in a ceramic matrix (Evans, 1990). In many cases, the additional phase and/or coating around that phase has substantially different oxidation behavior than the matrix. This leads to complex oxidation and corrosion behavior for the entire composite.

Most of the current research on CMCs is focused on SiC fibers in matrices of oxides, Si₃N₄, or SiC. The fibers impart fracture toughness by deflecting the matrix cracks. To properly deflect cracks, a fiber must not bond with the matrix. For this reason, inert fiber coatings are needed. To date, only two adequate coatings have been found—graphitic carbon and boron nitride. Thus the entire composite has a slowly oxidizing phase (fibers and matrix) and an easily oxidizable phase (fiber coating). Generally at high temperatures, SiO₂ rapidly forms on the outside of the matrix and attack of the fiber coating is minimized. Even a matrix crack may quickly seal. However at
intermediate (500-1000°C) temperatures, silica forms very slowly and oxidative attack of the fiber coating becomes a critical issue. These constituents are shown schematically in Figure 51.

Consider first carbon fiber coatings. The attack of the coating at low temperatures is evident in the TGA data shown in Figure 52 (Bhatt, 1989). Note that the lower temperatures show an initial weight loss due to the oxidation of the carbon coatings of the fibers. At higher temperatures the matrix rapidly forms a film of SiO₂ and this oxidation is not evident. However in applications, components must operate at intermediate temperatures and hence it is important to understand this problem.

The most thorough treatment of this problem is from Filipuzzi and co-workers (1994) for carbon-coated SiC fibers in an SiC matrix. Their initial treatment consists of a matrix with uni-directional fibers with one end exposed. The process consists of a two-step oxidation of carbon and concurrent oxidation of the fiber and surrounding matrix. As carbon oxidizes, an annular region around the fiber is left. A schematic of this region is shown in Figure 53. The pressure dependence of SiC oxidation indicates that most of the oxidation would occur at the mouth of pore, as shown in Figure 53. Filipuzzi and co-workers first measured oxidation properties of each of the constituent parts of the composite as well as the composite as a whole. Standard TGA methods as well as resistivity measurements give an indication of the amount of carbon consumed. These data are used in an overall description of the composite oxidation.

The constituent data for the composite is used in the following analytical model. This is based on differential equations, which describe:

1. Diffusion of oxygen into the pore.
2. Two step oxidation of carbon to CO and CO₂.
3. Diffusion of CO out of the pore.
4. Consumption of oxygen by oxidation of the matrix.
5. Consumption of oxygen by oxidation of the fiber.

Appropriate boundary conditions are introduced and the equations are solved to generate kinetic curves such as those in Figure 52. Excellent agreement between their model and experiments are obtained. Their analysis indicates that thinner carbon coatings give better performance.

This analysis has been extended to other situations. Eckel and co-workers (1995) have examined the lower temperature regime, where SiC oxidation can be ignored. At temperatures near about 500°C reaction control becomes important in the oxidation of carbon. Windisch and co-workers
(1997) have seen evidence of reaction control at higher temperatures (1073-1373 K) and reduced oxygen pressures.

Recently Halbig et al. (1998) have examined the oxidation of SiC matrices with carbon fibers. Oxidation tests indicate low temperature oxidation is kinetically controlled by reaction of oxygen with the fibers and high temperature oxidation is diffusion controlled by diffusion of oxygen to the fibers. They have developed finite difference models of this process.

The other commonly used fiber coating is hexagonal BN. There are several important phenomena in the oxidation of a silicon-based composite with BN fiber coatings. These are illustrated in Figure 54(a) and (b). As the silicon-based material and the BN oxidize, both $B_2O_3$ and the $SiO_2$ form. These react to form low melting borosilicate melts, as shown in Figure 54a. Related to this, it is well-known that boron dramatically enhances the oxidation of silica forming materials (Schlichting, 1984). As discussed in Section 5 on monolithic BN, water vapor readily reacts with the $B_2O_3$ oxidation product and forms volatile $H_xB_yO_z(g)$ species. This can lead to effective volatilization of BN, much as carbon volatilizes, illustrated in Figure 54(b). However, unlike carbon, distinct interfaces are not formed with the BN. Rather there is a 'plug' of borosilicate formed and a gradual transition from oxide to nitride (Jacobson et al., 1999b). This is illustrated in Figure 55.

A description of this has been developed based on gas phase diffusion of the $H_xB_yO_z(g)$ species outward and boron-enhanced oxidation of the SiC to form a borosilicate plug (Jacobson et al., 1999b). This is shown schematically in Figure 56. Exposures of the SiC fiber with BN fiber coatings in a SiC matrix were conducted at 700 and 800°C in 1 and 10% $H_2O/O_2$ atmospheres with one end of the test sample ground off to expose the fiber ends and coatings. Then recession distances were measured. As illustrated in Figure 55, the exact distance was difficult to determine as interfaces were not always sharp. In order to apply the model to the experiments an adjustable parameter of boron-influenced SiC oxidation rates had to be introduced. Using substantially enhanced oxidation rates, reasonable agreement between the model and experimentally measured recession distances were obtained.

9. Summary

As structural ceramics find more applications in high temperature systems, oxidation and corrosion at high temperatures becomes an important field of study. In this chapter, the critical issues in this field have been surveyed. Ceramics have been classified according to the type of protective oxide they form. These include silica formers, alumina formers, boria formers, and transition metal oxide formers. Most of the literature covers silica formers since there are a
number of near-term applications for these materials. Basic oxidation mechanisms, water vapor interactions, volatilization routes, and salt-induced corrosion were discussed for these materials. Less information is available on alumina-forming ceramics. However the rapid oxidation rate in water vapor appears to be a major problem. Boria formers show rapid oxidation rates due to the formation of a liquid oxide film and are volatile in the presence of water vapor due to highly stable H-B-O(g) species formation. Transition metal carbides and nitrides also show rapid oxidation rates due to rapid transport in the oxide scale and cracking of that scale.

Multi-component ceramics allow the optimization of various physical properties. These include ceramics which form multi-component oxides as well as fiber-reinforced ceramic matrix composites. However, the oxidation behavior of these materials is complex compared to the pure materials. The leading fiber-reinforced composites are silicon-based and contain continuous SiC fibers with coatings of graphitic carbon or hexagonal boron nitride. The oxidation of the fiber coating at intermediate temperatures is a major issue and models of this process are discussed for both carbon and boron nitride coatings.
References:


Hincke, W. B., L. R. Brantley (1930), JACS 52, 48-52.


General References:


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<th>Ceramics</th>
<th>Approximate Use Temperatures</th>
<th>Environment</th>
<th>References</th>
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<td>SiC, Si$_3$N$_4$, Composites</td>
<td>900-1400</td>
<td>Combustion gases deposits: Na, Mg, Ca sulfates, Sodium vanadates</td>
<td>Probst (1986)</td>
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<td>Combustor Liners</td>
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<td>Blades and Vanes</td>
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<td>Re-entry shields</td>
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<td>1000-1800</td>
<td>Reduced pressure N$_2$, O$_2$, CO$_2$, N, O</td>
<td>Lacombe and Bonnet (1990) Jacobson and Rapp (1995)</td>
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<td>High temperature semiconductors</td>
<td>SiC</td>
<td>Use: 600 Processing: 1200</td>
<td>Air</td>
<td>Chelnokov and Syrkin (1997), Powell et al. (1992)</td>
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Table II. Coefficients of thermal expansion for silica-forming ceramics and several refractory oxides (Lee et al., 1994)

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<tr>
<td>Si$_3$N$_4$</td>
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<tr>
<td>Mullite (3Al$_2$O$_3$2SiO$_2$)</td>
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<td>Alumina (Al$_2$O$_3$)</td>
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<td>Partially stabilized zirconia (0.09Y$_2$O$_3$ZrO$_2$)</td>
<td>10</td>
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</tbody>
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Table III. Transition Metal Carbides (Storms, 1967). The most stable carbides are shown in the box with strong metal-carbon and metal-metal bonds.

<table>
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<th>Mn</th>
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<td>Nb</td>
<td>Mo</td>
<td>Tc</td>
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<td>La</td>
<td>Hf</td>
<td>Ta</td>
<td>W</td>
<td>Re</td>
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</tbody>
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Figure 1
Figure 2
Figure 3
Figure 4
Figure 6

- Si, 119 kJ/mol, extrapolated from Deal & Grove (1965)
- CVD SiC, 118 kJ/mol Ogbuji & Grove (1995)
- CVD Si$_3$N$_4$, 363 kJ/mol Ogbuji & Grove (1995)
Figure 9
Figure 10
Figure 12
Figure 13
Figure 14
Figure 15

Wet O\textsubscript{2} data corrected to 1 bar H\textsubscript{2}O

\[ E_\text{a} = 68 \text{ kJ/mole} \]

Dry O\textsubscript{2} (1 bar)

\[ E_\text{a} = 119 \text{ kJ/mole} \]

Parabolic rate constant, \( \mu \text{m}^2/\text{h} \)

Temperature, °C

1200 \hspace{1cm} 1000 \hspace{1cm} 900 \hspace{1cm} 800 \hspace{1cm} 700 \hspace{1cm} 600 \hspace{1cm} 550

1000/T, K

0.7 \hspace{1cm} 0.8 \hspace{1cm} 0.9 \hspace{1cm} 1.0 \hspace{1cm} 1.1 \hspace{1cm} 1.2

From E-7625 fig 2
Figure 17
Figure 18
Figure 19
Figure 20

\[ \text{SiC: } k_p = 2.7 \times 10^{-3} \text{ mg}^2/\text{cm}^4 \text{ h} \]
\[ k_q = 5.3 \times 10^{-3} \text{ mg/cm}^2 \text{ h} \]

\[ \text{Si}_3\text{N}_4: \]
\[ k_p = 9.5 \times 10^{-4} \text{ mg}^2/\text{cm}^4 \text{ h} \]
\[ k_q = 5.2 \times 10^{-3} \text{ mg/cm}^2 \text{ h} \]

\[ \text{SiO}_2: \]
\[ k_p = 2.2 \times 10^{-3} \text{ mg/cm}^2 \text{ h} \]
Figure 21

SIC weight loss, mg/cm²

Exposure time, hr

6 atm, 1335°C
5 atm, 1315°C
10 atm, 1332°C
15 atm, 1329°C
Temperature, °C

$k_i (mg/cm^2-hr), T(K), P(atm), v(m/s) R(kJ/mol-K)$

\[ k_i = 2.06 \exp\left(\frac{-108}{RT}\right) p^{1.5} v^{0.5} \]

Figure 22
Figure 23
Figure 24
Figure 25

- CVD SiC
- CVD Si$_3$N$_4$
- NT 154 Si$_3$N$_4$
Figure 26(a)
Figure 26 (b)
Figure 26 (c)
Figure 27
Figure 30
Figure 31
Figure 32
Figure 35
Figure 37
(1) Initial Oxidation and Dissolution

\[ \text{SiC} + \frac{3}{2} \text{O}_2 = \text{SiO}_2 + \text{CO} \]

\[ x\text{SiO}_2 + \text{Na}_2\text{CO}_3 = \text{Na}_2\text{O} \cdot x(\text{SiO}_2) + \text{CO}_2 \]

(2) Gas Bubbles and Pit Formation

(3) Melt Saturation and Healing Layer

Figure 38
Figure 40

Temperature, °C

Log $k_0$, μm²/s

AIN
$322 \pm 113$ kJ/mol

Ni-Al$_{0.25}$
$293$ kJ/mol

$10 \ 000/\Theta, K$
Figure 41
Figure 43

(a) Specific weight change, mg/cm²

(b) Specific weight change, mg/cm²

Time, hr
Figure 44
Figure 45
Figure 46
Figure 47
Figure 48

Specific weight change, mg/cm²

Oxygen consumption
Sample weight
Vaporization

Time, hr

E-11463 Opila 8 pt/100% scan arg
Figure 49

- **SiC particles**
- **Oxidation product**
- **Unreacted matrix oxide**

Unreacted composite → Semi-acted SiC particles → Reaction product
Figure 51
Figure 52
Figure 53
Fig 55
Fiber

Flowing $O_2/H_2O$

Matrix

Growing $SiO_2$ film--closes pore

Borosilicate

Intact boron nitride

Boundary layer

$h_xb_yO_z(s)$

SiC Fiber

Matrix

(From E-10734-M fig 11) Figure 56