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## A method for assessing the quality of model-based estimates of ground temperature and atmospheric moisture using satellite data

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**Abstract.** A method is developed for validating model-based estimates of atmospheric moisture and ground temperature using satellite data. The approach relates errors in estimates of clear-sky longwave fluxes at the top of the Earth-atmosphere system to errors in geophysical parameters. The fluxes include clear-sky outgoing longwave radiation (CLR) and radiative flux in the window region between 8 and 12  $\mu\text{m}$  (RadWn). The approach capitalizes on the availability of satellite estimates of CLR and RadWn and other auxiliary satellite data, and multiple global four-dimensional data assimilation (4-DDA) products. The basic methodology employs off-line forward radiative transfer calculations to generate synthetic clear-sky longwave fluxes from two different 4-DDA data sets. Simple linear regression is used to relate the clear-sky longwave flux discrepancies to discrepancies in ground temperature ( $\delta T_g$ ) and broad-layer integrated atmospheric precipitable water ( $\delta pw$ ). The slopes of the regression lines define sensitivity parameters which can be exploited to help interpret mismatches between satellite observations and model-based estimates of clear-sky longwave fluxes. For illustration we analyze the discrepancies in the clear-sky longwave fluxes between an early implementation of the Goddard Earth Observing System Data Assimilation System (GEOS2) and a recent operational version of the European Centre for Medium-Range Weather Forecasts data assimilation system. The analysis of the synthetic clear-sky flux data shows that simple linear regression employing  $\delta T_g$  and broad layer  $\delta pw$  provides a good approximation to the full radiative transfer calculations, typically explaining more than 90% of the 6 hourly variance in the flux differences. These simple regression relations can be inverted to "retrieve" the errors in the geophysical parameters. Uncertainties (normalized by standard deviation) in the monthly mean retrieved parameters range from 7% for  $\delta T_g$  to ~20% for the lower tropospheric moisture between 500 hPa and surface. The regression relationships developed from the synthetic flux data, together with CLR and RadWn observed with the Clouds and Earth Radiant Energy System instrument, are used to assess the quality of the GEOS2  $T_g$  and  $pw$ . Results showed that the GEOS2  $T_g$  is too cold over land, and  $pw$  in upper layers is too high over the tropical oceans and too low in the lower atmosphere.

### 1. Introduction

While much progress has been made to improve the climate characteristics of general circulation models (GCMs), the hydrological cycle stands out as a major component of the Earth-Atmosphere system which is still poorly modeled [e.g., Gates *et al.* 1999; Lau *et al.*, 1995]. Major advances in modeling the hydrological cycle are hampered by inadequate and/or incomplete measurements of such quantities as precipitation, latent heating, clouds, atmospheric and soil moisture, and ground temperature. Data obtained by methods of four-dimensional data assimilation (4-DDA) suffer from errors in these quantities, both as a result of errors in the assimilating models and a lack of observations to directly constrain the hydrological cycle.

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In this study we focus on methods to verify model-based (including 4-DDA) moisture profiles ( $q$ ) and ground temperature ( $T_g$ ) using satellite data. We focus on these two parameters because they are highly dependent on the parameterizations of sub-grid-scale processes, and they are difficult to validate because there are few reliable observations of these quantities with global coverage. The quality of the ground temperature directly reflects on the quality of the land surface formulation, while the moisture profiles are important constraints on the behavior of the boundary layer and convection schemes. The methodology described here was developed as the result of efforts to validate these quantities in a global data assimilation system, but the methodology should also prove useful for assessing systematic errors of global atmospheric models.

The basic methodology is as follows: First, sensitivity parameters are obtained by applying off-line forward radiative trans-

fer calculations to two different assimilated data products to relate differences in clear-sky longwave fluxes at the top of the Earth-atmosphere system to errors in geophysical parameters. The clear-sky longwave fluxes include clear-sky outgoing longwave radiation (CLR) and radiative flux in the window region between 8 and 12  $\mu\text{m}$  (RadWn). The assimilated data sets should provide physically reasonable and global estimates of the geophysical fields. Any discrepancies in the fields from the two different systems measure our uncertainty and provide a tool for assessing the sensitivity of the CLR to discrepancies in these fields. The sensitivity parameters together with satellite estimates of clear-sky longwave fluxes (and other auxiliary satellite data) are then used to relate the clear-sky longwave fluxes to errors in the geophysical fields. For purpose of illustration, sensitivity parameters are obtained by comparing ground temperature and atmospheric temperature and moisture profiles from an early implementation of version 2 of the Goddard Earth Observing System (GEOS2) data assimilation system [Data Assimilation Office (DAO), 1996] with those quantities from a recent operational version of the European Centre for Medium-Range Weather Forecasts (ECMWF) data assimilation system [Courtier *et al.*, 1998].

The key satellite data used for this study are the CLR and RadWn estimates from the Clouds and Earth Radiant Energy System (CERES) [Wielicki *et al.*, 1996] instrument onboard the Tropical Rainfall Measuring Mission (TRMM).

This approach to validation is viewed as a compromise between a simple comparison of model-generated and observed CLR and RadWn and full retrieval algorithms that attempt to exploit information in multiple spectral bands to obtain estimates of the geophysical quantities. The method provides insight into the nature of the errors in the model-based CLR and RadWn consistent with the broadband nature of the CLR and RadWn satellite measurements.

Section 2 describes the satellite data and radiative transfer model used for the forward calculation. The methodology is described in section 3. The sensitivity analysis, based on the GEOS2 and ECMWF assimilated data, is described in section 4. A comparison with the CERES/TRMM clear CLR and RadWn is presented in section 5. Discussion and conclusions are given in section 6.

## 2. Satellite Data and Radiative Transfer Model

The primary quantity analyzed in this study is the CLR and RadWn. The availability of high-quality satellite estimates of CLR and RadWn (see below) has made these important quantities for validating global models. Relating and understanding differences between model-computed and observed CLR and RadWn is complicated by a number of factors. Some of the differences arise from sampling differences and fundamental differences between satellite measurements and model "grid-scale" fields. In this study we focus on the problem of relating the errors (computed minus observed) in CLR and RadWn to deficiencies in the model-based estimates of the radiatively important geophysical quantities (e.g., atmospheric moisture and temperature and ground temperature). The methodology is developed with synthetic flux data to circumvent the sampling and representativeness problems. These issues are beyond the scope of this initial study, though they will need to be addressed when developing quantitative estimates of errors in the geophysical quantities based on satellite measurements. We note that in this study the model-based CLR and RadWn

are computed in an off-line mode (see below). Many institutions now routinely compute CLR on-line during the course of model integrations.

In the following we describe the CLR and RadWn data, various other auxiliary data sets used for this study, and the radiative transfer scheme used for the off-line calculation of CLR and RadWn.

### 2.1. CLR and RadWn Satellite Data

While the sensitivity parameters will be developed with synthetic flux data (see next section), our application of the methodology to validating GEOS2 (see section 5) will require satellite observations as input. For the latter we have obtained CERES/TRMM data [Wielicki *et al.*, 1996]. The quality of the CERES data is comparable to the quality of the Earth Radiation Budget Experiment (ERBE) data for instantaneous radiance, fluxes, and scene type. Generally, radiance uncertainties are at the 1% level or less. Some differences between CERES/TRMM and ERBE-ERBS are as follows: the field of view resolution, the spectral response of the instruments, and the tropical only coverage of TRMM [Wielicki *et al.*, 1996].

### 2.2. Radiative Transfer Scheme

For the sensitivity calculations we employ the radiative transfer scheme developed by Chou and Suarez [1996]. This scheme is used in the GEOS2 model. The longwave (LW) calculation has nine bands (band 10 is added to compute flux reduction due to  $\text{N}_2\text{O}$  in the 15  $\mu\text{m}$  region). The transmission and absorption of  $\text{H}_2\text{O}$ ,  $\text{H}_2\text{O}$  continuum,  $\text{CO}_2$ ,  $\text{O}_3$ ,  $\text{CF}_4$ ,  $\text{CH}_4$ , and  $\text{N}_2\text{O}$  is modeled using  $k$  distribution. In the LW, multiplication approximation for bands are assumed. The LW scheme compares well with detailed line-by-line calculations. The root mean square (rms) errors of CLR are between 1 and 3  $\text{Wm}^{-2}$  [Chou and Kouvaris, 1991; Chou *et al.*, 1995; Kratz *et al.*, 1998]. In addition, the code has been participated in the Intercomparison of Radiation Codes in Climate Models (ICRCCM) [Ellingson *et al.*, 1991].

Surface emissivity depends on surface property or vegetated surface. For a given surface type, its values are also different for different wavelengths in the LW spectrum. The global distribution of surface type and the associated emissivity that we used were taken from CERES. The Advanced Spaceborne Thermal Emission and Reflection Radiometer created an easily accessible data set on the basis of extensive measurements of the spectral reflectances of surface material in the 2–16  $\mu\text{m}$  region by Salisbury and D'Aría [1992]. Wilber *et al.* [1999] derived spectrum emissivity from the spectrum reflectance on the basis of scene type. The scene type was determined by using a 1 km map of the International Geosphere Biosphere Program scene types as supplied by the U.S. Geological Survey (USGS). A scene type of tundra was added to separate it from desert resulting in 18 surface types. We calculate the fraction of each scene type in a  $2^\circ \times 2.5^\circ$  longitude and latitude box from the  $1/6^\circ$  equal angle data. Radiative fluxes are computed in each of the  $2^\circ \times 2.5^\circ$  boxes several times on the basis of the number of scene types in a box (maximum is 18). The final mean radiative flux is an area mean of the fractional type. Surface emissivity effect on CLR is large over the desert regions, where off-line computations show it can be as high as 5–8  $\text{Wm}^{-2}$  (mostly from window bands).

### 2.3. Ozone

The analyzed ozone fields from the Goddard Earth Observing System (GEOS) ozone data assimilation system (DAS)

[Riishøjgaard *et al.*, 1999; Štajner *et al.*, 1999] are used in this study. They are obtained by assimilating total column ozone observations from the Total Ozone Mapping Spectrometer (TOMS) and ozone profile observations from the solar back-scattered ultraviolet/2 (SBUV/2) instruments into an off-line transport model driven by the GEOS-DAS assimilated winds. While the statistical analysis of ozone is performed at every model time step (15 min) using a global physical space-based analysis scheme, the analyzed ozone is written out at the same frequency and resolution as the other GEOS-DAS fields. The high quality of the ozone fields is illustrated by two examples from their validation for January 1998. The average rms difference between TOMS observations and total column ozone forecast is about 11 Dobson units or 3.7% of the average total column ozone. The comparison of analyzed ozone profiles with independent ozone observations measured by the Halogen Occultation Experiment (HALOE) onboard the NASA Upper Atmosphere Research Satellite shows that the mean profiles agree within 3.2% between the pressures of 30 and 1 hPa. In this region the HALOE measurements agree within 5% with ozonesonde, lidar, balloon, rocketsonde, and other satellites ozone measurements [Bruhl *et al.*, 1996].

While the ozone contribution to CLR is relatively small, the accuracy of its distribution is critical for the CLR computation. In particular, the high ozone concentrations in the stratosphere and its presence in the water vapor window region are important for the clear-sky flux.

### 3. Methodology

To help understand the methodology, we start with the radiative transfer equations for clear-sky conditions in the LW region of the spectrum given

$$\begin{aligned} \text{CLR}(T, q, T_g, \dots) = & \int_n^{\infty} \pi \epsilon_\nu B_\nu(T_g) \tau_\nu(u_\nu) d\nu + \left( \int_n^{\infty} \right. \\ & \left. \epsilon_\nu (1 - \epsilon_\nu) \tau_\nu(u_\nu) \int_{u_\nu}^{\infty} \pi B_\nu(T(u')) \frac{d\tau_\nu(u' - u)}{du'} du' d\nu \right. \\ & \left. + \int_n^{\infty} \int_0^{u_\nu} \pi B_\nu(T(u')) \frac{d\tau_\nu(u - u')}{du'} du' d\nu, \right) \quad (1) \end{aligned}$$

where  $\nu$  is frequency,  $B_\nu$  is the Planck parameter,  $\epsilon_\nu$  is surface emissivity,  $\tau_\nu$  is the transmission parameter at frequency  $\nu$ ,  $u$  is the path length, and  $u_\nu$  is the total path length. In (1) the first terms on the right-hand side (RHS) is the contributions from the surface. The second term is the downwelling atmospheric radiation reflected from the surface. This term tends to be small compared to the other terms. The third term is the contribution from the atmosphere.

The first step is to develop sensitivity parameters relating CLR differences to differences in the geophysical quantities. This is carried out here using data from two different four-dimensional data assimilation systems. The simulated radiances are computed using the  $T_g$ ,  $q$  and temperature ( $T$ ) profiles as input to the radiative transfer calculations using the scheme described in the previous section. The aim of this step is to try to understand which aspect of the input data is most sensitive to the CLR and how the sensitivity varies over the globe.

We consider a change in the CLR in terms of the following linearization:

$$\begin{aligned} \delta \text{CLR}(T, q, T_g) = & \frac{\partial \text{CLR}}{\partial T_g} \delta T_g + \sum_i \frac{\partial \text{CLR}}{\partial q_i} \delta q_i \\ & + \sum_i \frac{\partial \text{CLR}}{\partial T_i} \delta T_i, \quad (2) \end{aligned}$$

where  $\delta q_i$  and  $\delta T_i$  are the differences in the two moisture and temperature products in layer  $i$ . We consider the partial derivatives in (2) as sensitivity parameters relating changes in the geophysical quantities to changes in the CLR. The validity of this approximation will be examined in the following section by comparing results from the full radiative transfer calculations (referred to as the "true" values in the following discussions) with results from "best fit" linear regression equations based on (2).

### 4. Sensitivity Analysis

In this section we shall compute the sensitivity parameters defined previously using output from two different 4-DDA systems (GEOS2 and ECMWF). Note that since we are only interested in differences, there is no need in these calculations to assume one or the other 4-DDA systems is correct. In section 5 we will use the sensitivity parameters together with satellite data to validate the GEOS2 system.

GEOS2 represents a major upgrade to the baseline GEOS1 system employed in the NASA first reanalysis effort [Schubert *et al.*, 1993]. The final version of GEOS2 includes a physical-space three-dimensional variational analysis algorithm (Physical-Space Statistical Analysis System (PSAS) [Cohn *et al.*, 1998]) and numerous improvements to the general circulation model. The model improvements include a more accurate dynamics [Takacs and Suarez, 1996], a gravity wave drag scheme [Zhou *et al.*, 1996], improved diagnostic clouds, the mosaic land surface scheme of Koster and Suarez [1994], a level 2.5 moist turbulence scheme [Helfand *et al.*, 1999], and new SW and LW radiation code [Chou and Suarez, 1994]. The system also includes the capability to assimilate TOVS moisture and Special Sensing Microwave/Imager (SSM/I) total precipitable water. (Wentz, 1994)

The version of GEOS2 used here was run with 70 vertical layers extending to 0.01 hPa and with a spatial resolution of  $2^\circ \times 2.5^\circ$  latitude-longitude. The  $T_g$  output was saved every 3 hours, while the upper air  $T$  and  $q$  were saved every 6 hours. The GEOS2 data are compared with operational analyses generated by the ECMWF for January 1998. These data are available every 6 hours and have been interpolated to a  $2^\circ \times 2.5^\circ$  grid to be consistent with the GEOS2 data.

As described in the previous section,  $T_g$ ,  $T$ , and  $q$  from both analyses are input to a radiative transfer model to compute the CLR and RadWn. The primary interest for this study is in the sensitivity to  $T_g$  and  $q$ . Our initial analysis showed that  $T$  differences between GEOS2 and ECMWF are small, and the effect on the computation of CLR is around  $\pm 1.5 \text{ Wm}^{-2}$ .

In the following, CLR computed by using all GEOS2 data is referred to as CLRGEOS2, and the CLR computed using all ECMWF data is referred to as CLRECMWF. Sensitivities are estimated by recomputing the CLR after replacing selected input fields. For example, to estimate  $\partial \text{CLR} / \partial T_g$ , we compute the CLRGEOS2 as before but replacing the GEOS2  $T_g$  with

the  $T_g$  from the ECMWF data. We will refer to this as  $CLR_{GEOS2}(\delta T_g)$ , indicating the calculation is carried out with the GEOS2 ground temperature perturbed by the amount  $\delta T_g = T_g(\text{ECMWF}) - T_g(\text{GEOS2})$ . The difference  $\delta CLR(\delta T_g) = CLR_{GEOS2}(\delta T_g) - CLR_{GEOS2}$  which divided by  $\delta T_g$  is a measure of the sensitivity of CLR to ground temperature. For example, the change in CLR due to differences in  $T_g$  is then denoted by  $\delta CLR(\delta T_g)$ . An estimate of the sensitivity of CLR to  $T_g$  changes is

$$\frac{\partial CLR}{\partial T_g} \approx \frac{\delta CLR(\delta T_g)}{\delta T_g} \quad (3)$$

Since the sensitivity of CLR to moisture varies with height, sensitivity ratios are estimated separately for different layers in the atmosphere. For these calculations we used ECMWF data as the basis for determining the moisture sensitivity since we felt that the ECMWF moisture values were more reasonable especially at upper levels. In principle, it should not matter which one is used (GEOS2 or ECMWF), since we are only after the sensitivity. In fact, tests showed similar results using GEOS2, though we felt the analysis using ECMWF provided results that were somewhat easier to interpret for the moisture. The layer-integrated moisture is referred to as the precipitable water ( $\rho w$ ). For example, we compute  $\delta CLR(\delta \rho w_h) = CLR_{ECMWF}(\delta \rho w_h) - CLR_{ECMWF}$ . Here  $\delta \rho w_h$  is equal to  $\rho w$  from ECMWF minus  $\rho w$  from GEOS2, where  $\rho w_h$  is the moisture in the layer between 200 and 500 hPa. Similar calculations are done for moisture between 200 and 700 hPa ( $\rho w_{700}$ ), between 200 and 900 hPa ( $\rho w_{900}$ ), between 200 hPa and surface ( $\rho w_{sfc}$ ), between 500 and surface ( $\rho w_s$ ), and the full moisture profile (tpw). The sensitivities to  $\rho w$  are given by equations analogous to (3). For example,

$$\frac{\partial CLR}{\partial \rho w_h} \approx \frac{\delta CLR(\delta \rho w_h)}{\delta \rho w_h} \quad (4)$$

estimates the sensitivity of CLR to the  $\rho w$  between 200 and 500 hPa.

We use the period January 1998 to compute the sensitivity parameters, since we have in-house operational analyses from ECMWF for this time period. Unless noted otherwise, all calculations are based on 6 hourly data on the GEOS2  $2^\circ \times 2.5^\circ$  grid. We estimate the sensitivities locally using simple regression to relate the 6 hourly differences in CLR (or  $\delta CLR$ ) to differences ( $\delta x$ ) in the input quantities. For a single predictor the regression equation takes the form

$$\delta CLR(\delta x) = \alpha_x \delta x + \epsilon, \quad (5)$$

where  $x$  is the quantity being varied (for example,  $T_g$ ,  $\rho w_h$ , or  $\rho w_s$ ),  $\alpha_x$  is the sensitivity parameter to be estimated, and  $\epsilon$  is the component of the  $\delta CLR$  not explained by  $\delta x$ .

Before computing the sensitivity parameters, we first examine the differences in the CLR from the two assimilated data sets and look in some detail at the various band contributions to the total CLR. Figure 1 shows the mean  $\delta CLR = CLR_{GEOS2} - CLR_{ECMWF}$  for January 1998 computed from the 6 hourly data. Over land the differences can be as large as 15–20  $Wm^{-2}$  and over the oceans as large as 20–25  $Wm^{-2}$ . Figures 1b and 1c show  $\delta CLR(\delta T_g)$  and  $\delta CLR(\delta q)$ , respectively. The pattern and magnitude of the  $\delta CLR$  distribution over land is very close to  $\delta CLR(\delta T_g)$ , while over

Clear Sky OLR Dif ( $Wm^{-2}$ ) January 1998  
(Contour intervals ( $5Wm^{-2}$ , starting at  $-5$ ) are for negative values)

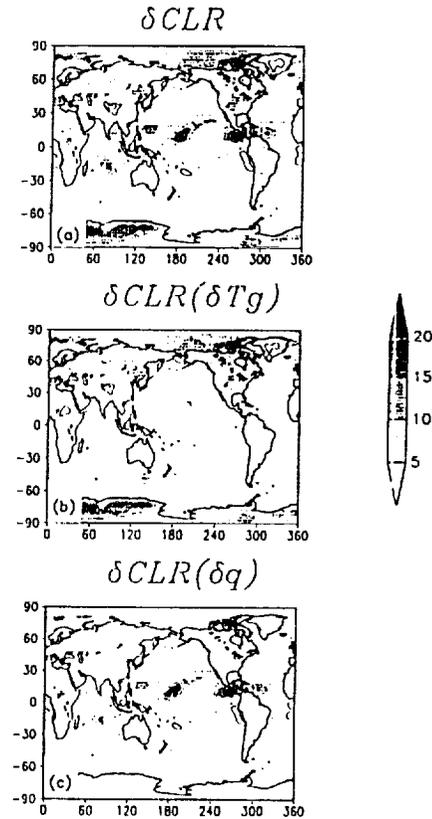


Fig. 1

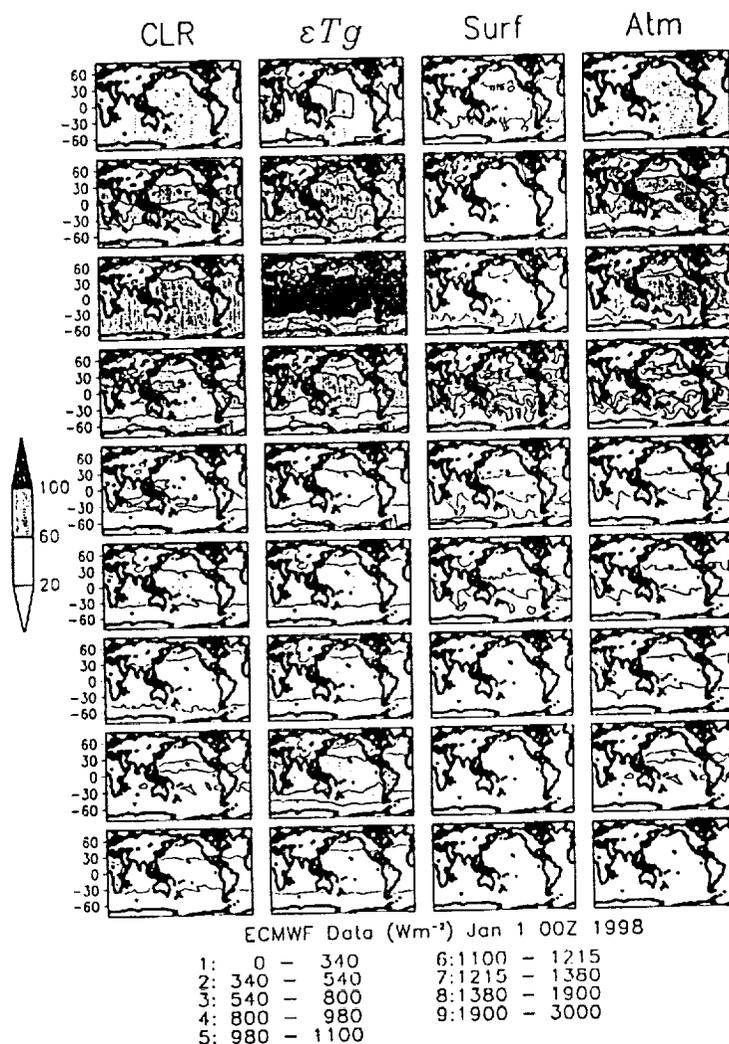
Figure 1. Monthly mean  $\delta CLR$  ( $Wm^{-2}$ ) based on the GEOS2 data and the ECMWF data: (a) differences based on  $\delta T_g$ ,  $\delta q$ , and  $\delta T$ ; (b) differences induced by  $\delta T_g$ ; and (c) differences induced by  $\delta q$ . Contours are for negative values; the intervals are the same as the shaded values.

oceans,  $\delta CLR$  is very similar to  $\delta CLR(\delta q)$ . This indicates that for these two assimilated data sets the largest discrepancies in CLR over the land poleward of about  $40^\circ$  latitude are due to ground temperature differences, while the largest discrepancies in CLR over the oceans are due to differences in the moisture profiles.

To better understand the sources of the CLR from the atmosphere and land, we look in more detail at the various band contributions to the CLR (Figures 2 and 3). The gaseous absorption included in the radiative transfer code [Chou and Suarez, 1994] are  $H_2O$ ,  $CO_2$ ,  $O_3$ , and trace gases (see section 2.2). In the following we will focus on the water vapor effect. Column 1 in Figure 2 shows the CLR computed from the ECMWF data in each of the nine bands. The first three bands are the contributions from the water vapor rotational bands. The next three are the contributions from the 800 to 1215  $cm^{-1}$  window region, and the last three represent contributions from water vapor vibration and rotation bands. Column 2 shows the emission from the surface ( $\epsilon_s T_g$ ). Column 3 is the contribution to the CLR from the first two terms on the RHS

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**Figure 2.** Contributions to CLR ( $\text{Wm}^{-2}$ ) from different parts of the atmosphere and different bands. From top down the band ranges from band 1 to band 9, the wave number for each band is shown at the bottom of the figure. Value  $\epsilon T$  is surface emission, Surf (surface contribution) is the first two terms of the RHS of (1) and Atm (contribution from the atmosphere) is the third term of the RHS of (1). Contour intervals are  $10 \text{ Wm}^{-2}$  except the last one in a column, which is  $5 \text{ Wm}^{-2}$ .

of (1), and column 4 shows the contribution to the CLR from the atmosphere (third term on the RHS of (1)).

The strength of surface emission follows the Planck function which peaks in the third band ( $540\text{--}800 \text{ cm}^{-1}$ ). The intensity from the other bands decreases gradually toward both sides. Note that the band widths are not uniform. The atmospheric transmission function strongly modifies the surface emission, especially in bands with strong gaseous absorption. The degree of the atmospheric effect depends on the opacity of the atmosphere. Comparing, in Figure 2, the surface emission with the amount reaching the top of the atmosphere, we find a significant greenhouse effect due to gaseous absorption. The greenhouse effect is very strong in the water vapor rotational band ranging from zero to  $800 \text{ cm}^{-1}$  (which include the  $15 \mu\text{m}$   $\text{CO}_2$  band) and the vibrational and rotational band ranging from

$1215$  to  $3000 \text{ cm}^{-1}$ . The largest contributions to CLR are from bands 2 to 4, ranging from  $340$  to  $980 \text{ cm}^{-1}$ , and most of the contributions are from the atmosphere because of the opacity of the water vapor rotational bands and  $\text{CO}_2$  bands. For example, the surface has the largest emission in band 3, but the surface contribution to CLR is relatively small due to the opacity of the atmosphere. In particular, over the tropical oceans the surface contribution to the CLR from this band is about 5%, while the rest is from the atmosphere (see Figure 3). Figure 3 includes the effective pressure level of the contributions from the atmosphere averaged over the month. This is determined by finding the peak of the weighting function computed from the mean transmission functions. In bands 2 and 3, most contributions over land are from layers close to surface (lowest 300 hPa), while most contributions over the tropical

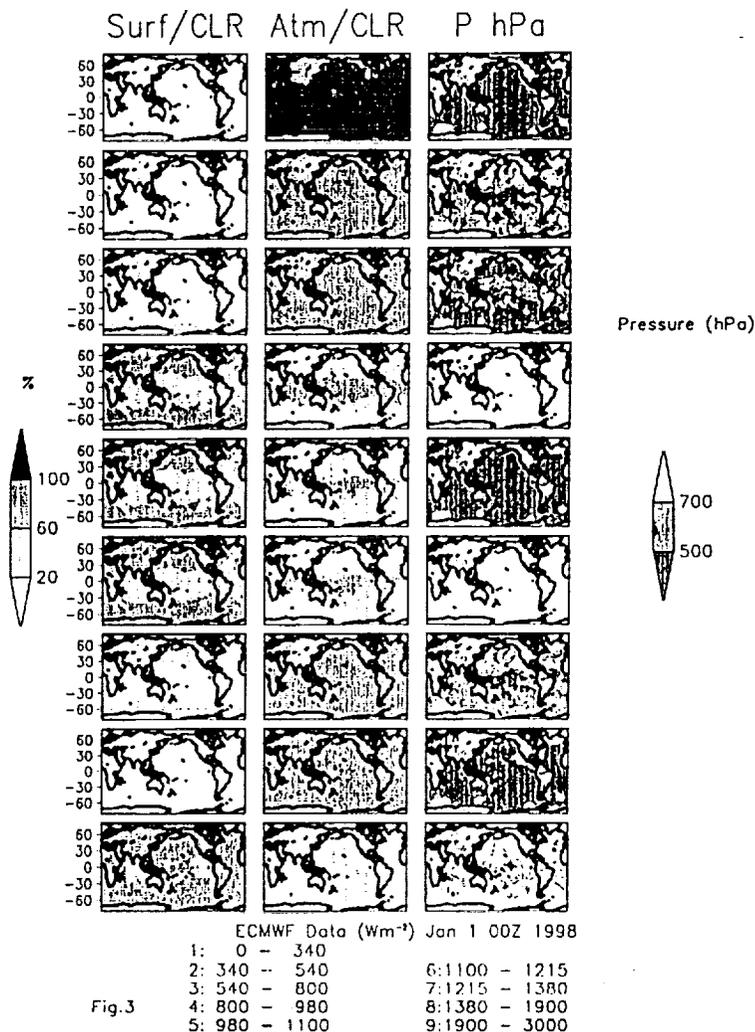


Fig. 3

Figure 3. Percentage contributions to CLR (%) from Surf and Atm and mean peak of the weighting function (hPa).

oceans are from layers ranging from 200 to 500 hPa, where water vapor concentration is high. Away from the tropics the opacity decreases, and the layer of largest emission is lowered to between 500 and 700 hPa. In the tropics the atmosphere contributes more than the surface in all bands. The largest contributions are from between 200 to 500 hPa and some are from 500 to 700 hPa depending on the opacity of the atmosphere.

In summary, the major contribution to CLR over the tropical oceans (30°S to 10°N) is from the atmosphere between 200 and 500 hPa. In the subtropics the contribution tends to be from lower in the atmosphere between 500 and 700 hPa. Over land, for the Northern Hemisphere winter, most of the contribution to CLR is from the surface (from window bands) and from the lowest (300 hPa) layer of the atmosphere.

F4

Figure 4 shows the regression results for  $T_g$  (land only) for January 1998 based on 6 hourly data. We note that sea surface temperature (SST) is specified from observations in both anal-

yses, and we do not consider any differences in the SST data. The regression is carried out on the CLR differences resulting from perturbing only the ground temperature using the radiative transfer equation (1). The extent to which the regression explains the results from the full radiative transfer calculation is thus a measure of the adequacy of the linear approximation for the effect of  $T_g$ . The estimated sensitivity parameter  $\partial CLR / \partial T_g = \alpha_T$  (Figure 4a) is highest in the middle and high latitudes (ranging from 1.2 to 1.6  $Wm^{-2}$  per 1°C change in  $T_g$ ) and lowest over the tropical land masses. This is consistent with our previous analysis of the various contributions to the CLR (Figure 2) and reflects the latitudinal dependence of the atmospheric moisture content. The explained variance is shown in Figure 4b. Remarkably, the regression line (5) explains more than 90% of the variance over most areas (Figure 4b). These results suggest that the sensitivity of CLR to changes in  $T_g$  can be reasonably estimated by the simple linear approximation (5) of the radiative transfer equation.

ECMWF and GEOS2 January 1998

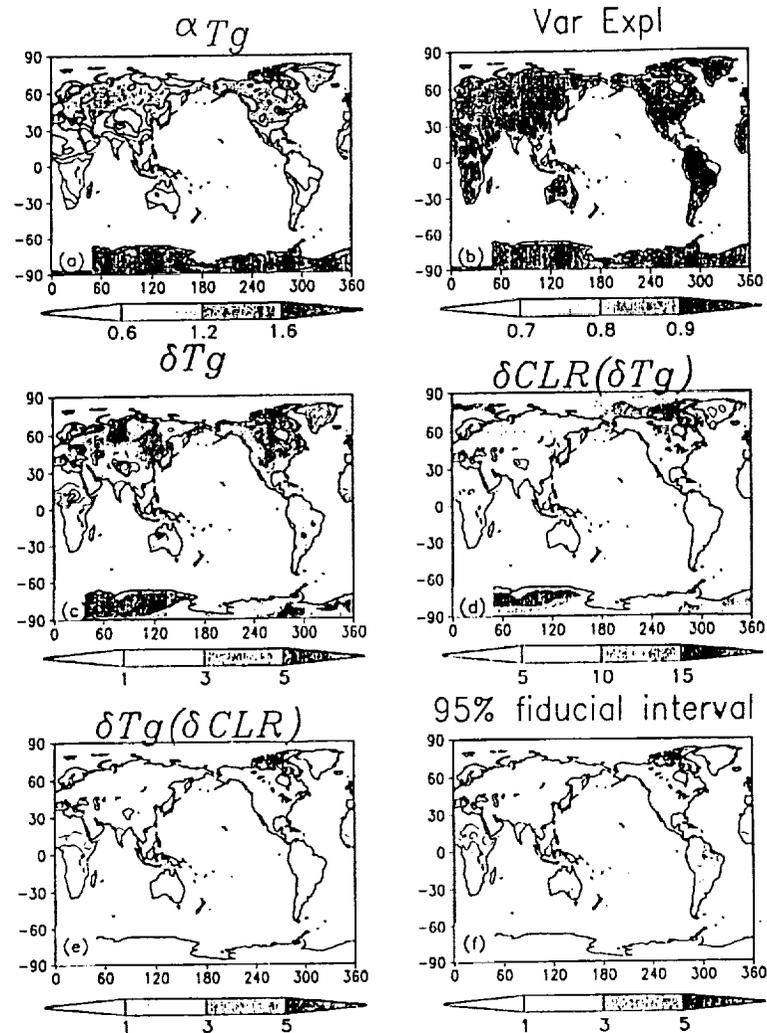


Figure 4. Estimation of  $\delta T_g$  based on linear fit: (a)  $\alpha T_g$  ( $\text{Wm}^{-2} \text{C}^{-1}$ ), contour intervals are  $2 \text{ Wm}^{-2} \text{C}^{-1}$ ; (b) total variance explained; (c) monthly mean  $\delta T_g$  ( $^{\circ}\text{C}$ ); (d) estimated  $\delta \text{CLR}(\delta T_g)$  ( $\text{Wm}^{-2}$ ) based on Figures 4a and 4c; and (e) errors ( $^{\circ}\text{C}$ ) in estimation of  $\delta T_g$  based on Figure 4d and (f) 95% fiducial interval. Contours are for negative values in Figures 4c-4f; the intervals are the same as the shaded values.

We are interested in the ability of the regression to represent the systematic difference in the CLR fields. Figure 4c shows the mean ground temperature difference  $\delta T_g$  between the two data sets for January 1998. The mean temperature differences are quite large over the cold continents accounting for most of the differences in CLR over those regions. Figure 4d shows the mean CLR difference from the regression equation evaluated for the time mean  $\delta T_g$ . Comparison with Figure 1b shows that the linear approximation (the regression line) with just  $\delta T_g$  provides a good approximation to the time mean differences in CLR over land. The reader is reminded that the results in Figure 1b are computed with the radiative transfer equation and include the effects of  $q$ ,  $T$ , and  $T_g$ . We shall see in the next

section that the differences in  $T_g$  for the most part reflect a cold bias in this version of GEOS2.

We next obtain an estimate of  $\delta T_g$  by inverting the regression equation (5): i.e., by solving for  $\delta T_g$  using the estimate of  $\alpha T_g$ . For brevity we refer to this estimate of  $\delta T_g$  as the "retrieved" value. Figure 4e shows the difference between our retrieved  $\delta T_g$  and the actual values shown in Figure 4c. The magnitude of the errors in  $\delta T_g$  are in most regions less than  $1^{\circ}\text{C}$ . Figure 4f shows that the 95% fiducial intervals for the inverse regression [Drafer and Smith, 1981] or retrieved values are typically  $\pm 3^{\circ}\text{C}$ . Note that here we only show one side of the two-sided confidence interval corresponding to the sign of the actual error in Figure 4e.

ECMWF and GEOS2 January 1998

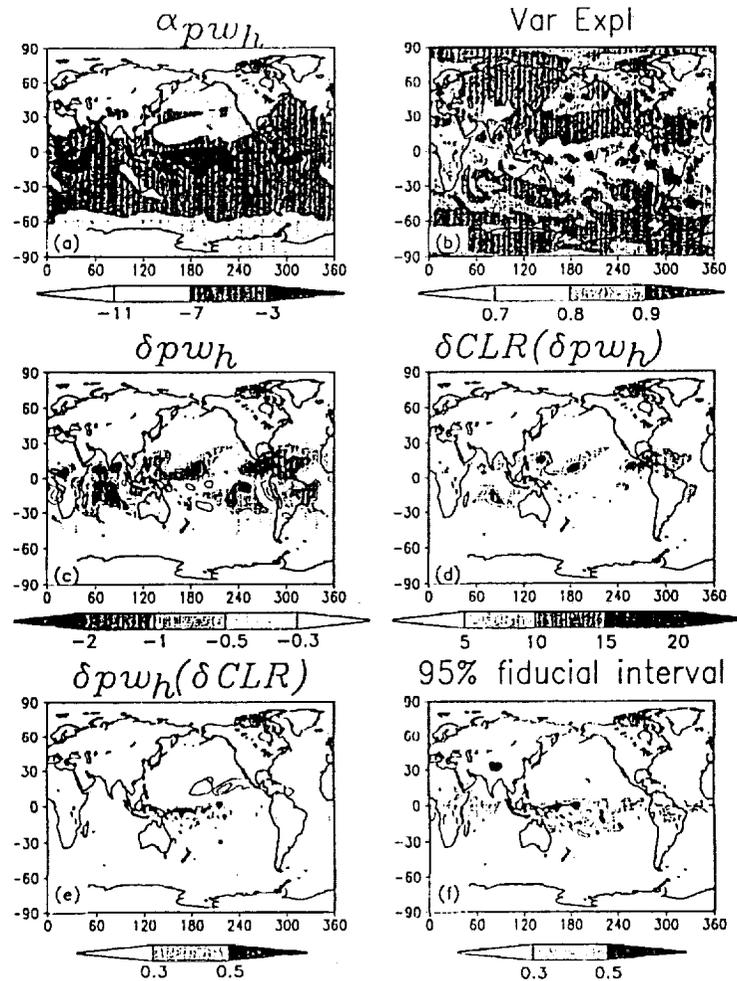


Figure 5. Same as in Figure 4 except for  $\delta p w_h$ ; units are  $(Wm^{-2} mm^{-1})$ . Contour intervals are for values opposite to the shaded values.

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Figure 5 is the same as Figure 4 except for the mean moisture difference in the layer between 200 and 500 hPa ( $\delta p w_h$ ). Figure 5a shows the sensitivity parameter ( $\alpha_{p w_h}$ ) determined from the regression

$$\delta CLR(\delta p w_h) = \alpha_{p w_h} \delta p w_h + \epsilon. \quad (6)$$

The spatial pattern of the sensitivity to upper level moisture is largely zonally symmetric, with the lowest sensitivity ( $<3 Wm^{-2}/mm$  precipitable water in the 200–500 hPa layer) in the convective regions of the tropics, and increasing sensitivity away from the tropics. A region of enhanced sensitivity ( $>9 Wm^{-2}/mm$  precipitable water in the 200–500 hPa layer) occurs over the subtropical regions of the North Pacific. We note that bands 2, 7, and 8 (Figure 2) show a tendency for enhanced contributions to the CLR from the subtropics especially over the North Pacific, suggesting the increased sensitivity is coming largely from these bands. Figure 5b shows that the regression line explains more than 90% of the variance in the CLR over

most of the globe. Scattered regions of the tropics and subtropics, the extratropical storm tracks, and the Himalayas, explained variance less than 80%. With those exceptions the linear approximation to the full radiative transfer calculations for the impact of the upper level moisture appears to be quite good. Figure 5c shows the mean difference in the upper level moisture ( $\delta p w_h$ ). The differences tend to be negative (GEOS2 is wetter) and largest just outside the tropics. The results of the regression for the time mean  $\delta CLR(\delta p w_h)$  (Figure 5d) again show a good approximation to the full calculation (Figure 1c), suggesting that much of the systematic difference between GEOS2 and ECMWF CLR over the subtropics are a result of differences in the upper level moisture. The results of the inverse calculation (Figure 5e) show that the magnitude of the errors in the retrieval of  $\delta p w_h$  is in most regions less than 0.5 mm. Figure 5f displays the 95% fiducial intervals on the retrieved values. The typical values range from 0.3 to 0.5 mm.

Figure 6 is the same as Figure 5 except for the mean mois- F6

ECMWF and GEOS2 January 1998

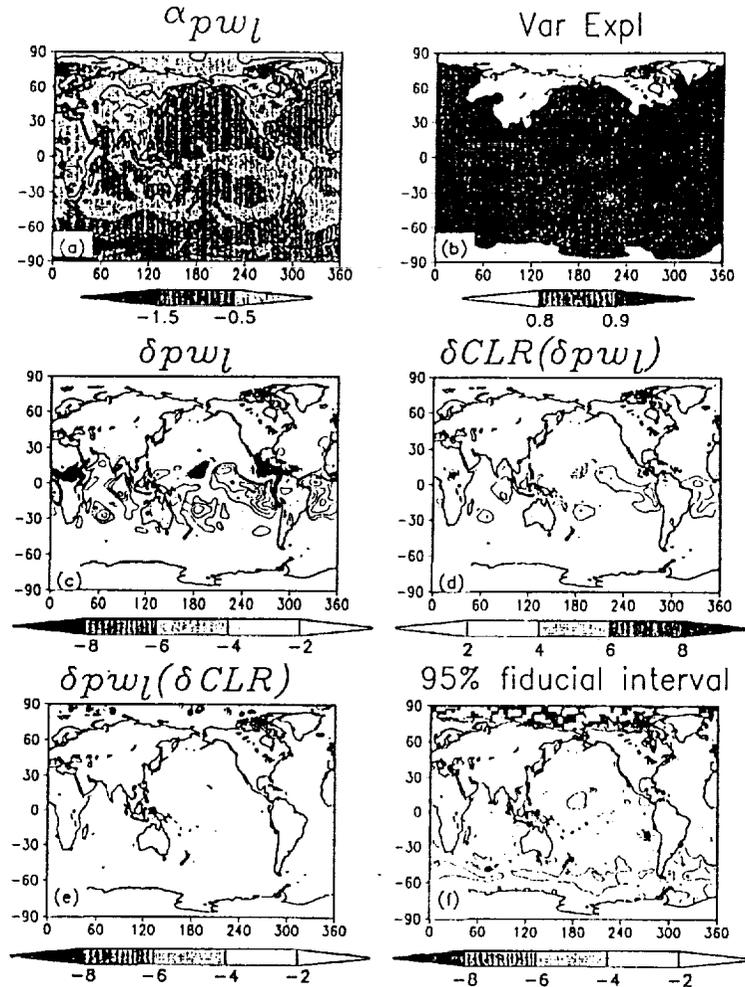


Figure 6. Same as in Figure 4 except for  $\delta p w_l$ ; units are ( $\text{Wm}^{-2} \text{mm}^{-1}$ ). Contour intervals are for values opposite to the shaded values.

ture difference in the layer between 500 hPa and the surface ( $\delta p w_l$ ). Figure 6a shows that the sensitivity parameter ( $\alpha_{p w_l}$ ) is about an order of magnitude smaller than the upper level sensitivity parameter ( $\alpha_{p w_u}$ ). The greatest sensitivity (magnitude  $>1$  to  $2 \text{ Wm}^{-2}/\text{mm}$  precipitable water) occurs at high latitudes in regions that are probably ice covered. Figure 6b shows the variance explained by the regression line which explains more than 90% of the variance in the CLR over most of the subtropics. Generally, less than 80% of the variance is explained over the cold continents and the polar regions. Figure 6c shows the mean difference in the lower-level moisture ( $\delta p w_l$ ). The differences tend to be positive (GEOS2 is drier) throughout much of the tropics and subtropics, with the largest values occurring off the west coasts of Africa and South America in regions characterized by the presence of low-level stratus clouds. There are some regions with negative values in the Northern Hemisphere subtropics that coincide with large negative values at upper levels (compare Figure 5c). The results of

the regression for the time mean  $\delta \text{CLR}(\delta p w_l)$  is shown in Figure 6d. This shows that some of the negative values along the west coasts of Africa, South America, and Africa in the full calculation (Figures 1a and 1c) are due to the differences in the low-level moisture. Figure 6e shows that the magnitude of the errors in the retrieval of  $\delta p w_l$  is in most regions less than 2 mm. The 95% fiducial intervals (Figure 6f) are in the range  $\pm 2$ –4 mm with the largest uncertainties in the Southern Hemisphere and the north polar region.

We next look more closely at the vertical dependence of the sensitivity of CLR to moisture. The panels on the left-hand side of Figure 7 show the differences between the ECMWF and the GEOS2 time mean moisture. The sensitivity parameters are the slopes of the regression lines through the scatterplots of  $\delta \text{CLR}$  versus  $\delta p w$  (panels on the right-hand side of Figure 7). These are computed, in this case, from the monthly mean fields, and each point in the scatterplots represents a different grid point (ocean only). This clearly shows that  $\delta \text{CLR}$

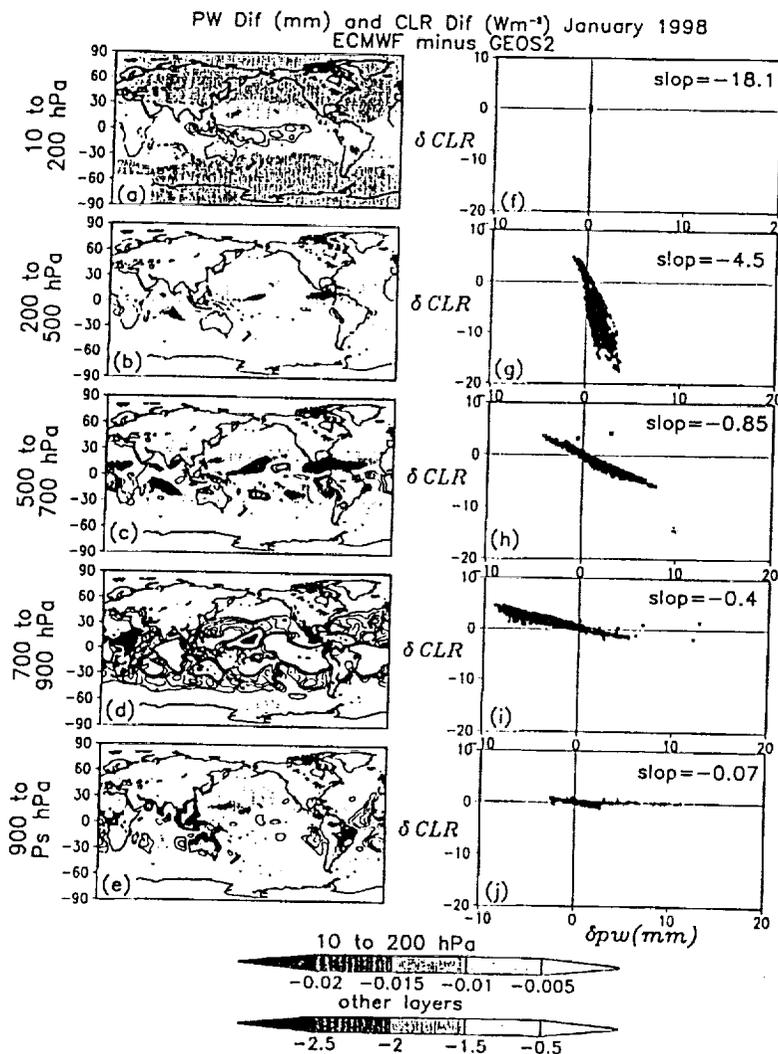


Figure 7. Sensitivity of  $\delta CLR$  ( $Wm^{-2}$ ) to  $\delta pw$  (mm). The left-hand side are precipitable water ( $pw$ ) differences between ECMWF and GEOS2; the right-hand side are scatterplots of  $\delta CLR$  and  $\delta pw$ . Contour intervals are for values opposite to the shaded values.

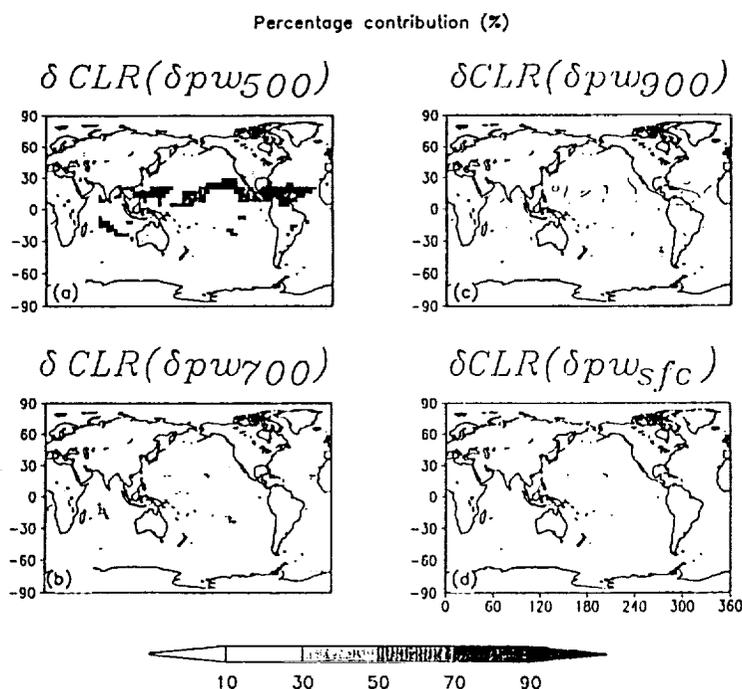
is most sensitive to  $\delta pw$  in the top layers. The sensitivity decreases the farther the layer is from the top of the atmosphere. Thus the moisture in the layer close to the surface tends to have little effect on CLR, consistent with our previous analysis of the contributions to the total CLR (Figure 2). Also, above 200 hPa the amount of moisture is too small to have an appreciable impact on the CLR despite the strong sensitivity. This is illustrated in Figure 8, which shows the relative contributions by layer to the time mean difference in the CLR due to moisture differences. Between 50 and 90% of the major CLR differences are due to  $pw$  differences between 200 and 500 hPa and another 10–30% comes from the layer between 500 and 700 hPa.

We next extend the results of the previous analysis to consider the case where, in addition to the total CLR measurements, we also have available CLR measurements for the win-

dow region. Figure 9 shows the sensitivity parameters  $\rho$  computed separately for the window region (RadWn, left-hand side panels) and for the total CLR minus RadWn (right-hand side panels). The latter consists of the rotation and vibration and rotation (RVR) bands or nonwindow bands; see discussion of Figure 2). A comparison of Figures 9a and 9d shows that the  $T_g$  sensitivity in the window region tends to peak in middle latitudes, while in the nonwindow or RVR, the sensitivity shows a general increase with latitude. As such, the greatest sensitivity to  $T_g$  over the United States, Europe, and parts of China occurs in the window region. A comparison of Figures 9b and 9e with Figure 5a shows that the  $pw_s$  sensitivity is dominated by the RVR bands or nonwindow bands. Over the tropical oceans the sensitivity to the low-level moisture is greatest in the window region (compare Figures 9c and 9f). The large sensitivity over the eastern North Pacific and Atlan-

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**Figure 8.** Relative contribution (%) of layer  $\delta p w$  to  $\delta \text{CLR}$ . Contour intervals are for values opposite to the shaded values.

tic seen in the total CLR calculation (Figure 6a) is due to the nonwindow bands or RVR bands (Figure 9f).

By considering the linearization (2) for both the RadWn and the RVR bands, we obtain two equations relating the CLR differences in each band to changes in the geophysical parameters. If we further limit the approximation of the CLR change to be primarily the result of just two geophysical parameters, we can estimate or "retrieve" these parameters. The previous results suggest that over land the CLR differences can be approximated by

$$\delta \text{RadWn} = \alpha_{T_g(\text{RadWn})} \delta T_g + \alpha_{p_w(\text{RadWn})} \delta p_w, \quad (7)$$

$$\delta \text{RVR} = \alpha_{T_g(\text{RVR})} \delta T_g + \alpha_{p_w(\text{RVR})} \delta p_w, \quad (8)$$

while over ocean, the approximation is

$$\delta \text{RadWn} = \alpha_{p_w(\text{RadWn})} \delta p_w + \alpha_{p_w(\text{RadWn})} \delta p_w, \quad (9)$$

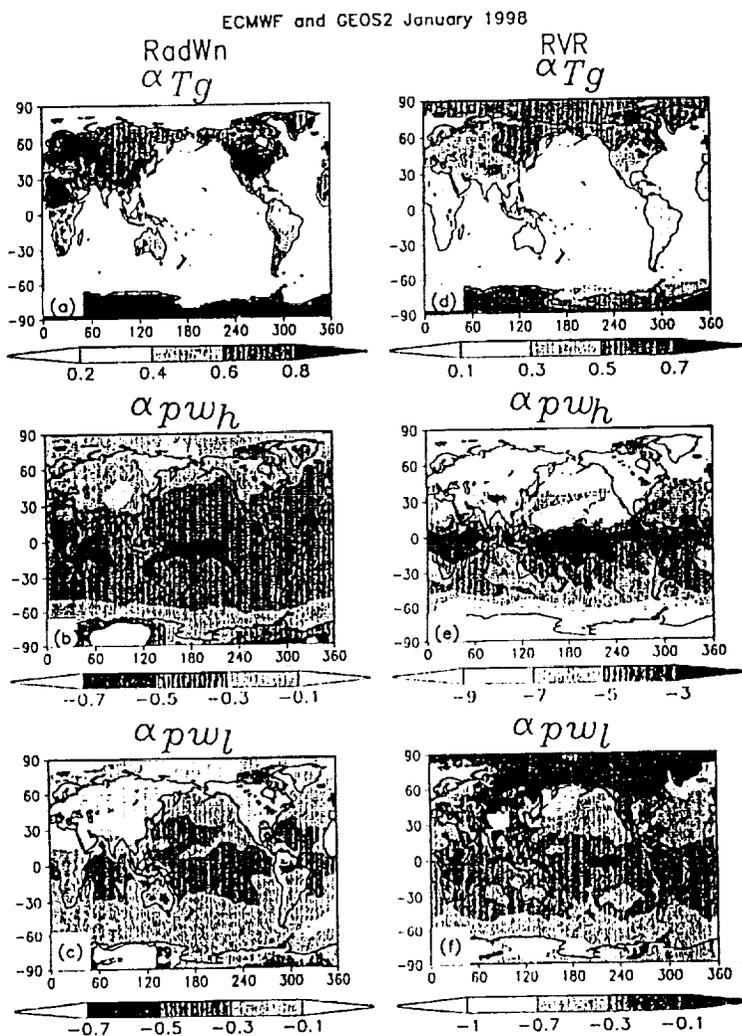
$$\delta \text{RVR} = \alpha_{p_w(\text{RVR})} \delta p_w + \alpha_{p_w(\text{RVR})} \delta p_w. \quad (10)$$

Note that in the above relationships the sensitivity parameters (the  $\alpha_x$  in Figure 9) are obtained from the 6 hourly data

according to (5). By inverting (7)–(8) and (9)–(10), we obtain the retrieved estimates of the differences in the geophysical parameters assuming the  $\alpha_x$  and  $\delta \text{RadWn}$  and  $\delta \text{RVR}$  are known. For example, the retrieved monthly mean differences in the geophysical parameters ( $T_g$ ,  $p_w$ , and  $p_w$ ) are shown in Figure 10 (left panels). The actual differences in  $T_g$ ,  $p_w$ , and  $p_w$  from the full radiative transfer equation shown earlier are presented again in the right-hand panels to facilitate the comparison with the estimated values. The global mean bias in the retrieved  $\delta T_g$  is  $-0.04^\circ\text{C}$ . The global mean root-mean-square error is  $1.77^\circ\text{C}$ , which corresponds to about 7% of the standard deviation of the true values. The global mean bias in the retrieval of  $\delta p_w$  is  $-0.1 \text{ mm}$ . The rms error is  $0.26 \text{ mm}$ , which corresponds to about 8% of the standard deviation of the true values. Similarly, the global mean bias in the retrieval of  $\delta p_w$  is  $0.4 \text{ mm}$ . The rms error is  $2.15 \text{ mm}$ , which corresponds to about 20% of the standard deviation of the true values.

The above results show that (7) through (10) give a close approximation to the mean differences in the ground temper-

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**Figure 9.** Coefficients of linear fit from Figures 9a to 9c are based on  $\delta\text{RadWn}$  and from Figures 9d to 9f are based on the rotation and vibration-rotational bands ( $\delta\text{RVR}$ ): (a)  $\alpha T_g$  ( $\text{Wm}^{-2} \text{ } ^\circ\text{C}^{-1}$ ), (b)  $\alpha p w_h$  ( $\text{Wm}^{-2} \text{ mm}^{-1}$ ), (c)  $\alpha p w_l$  ( $\text{Wm}^{-2} \text{ mm}^{-1}$ ), (d)  $\alpha T_g$  ( $\text{Wm}^{-2} \text{ } ^\circ\text{C}^{-1}$ ), (e)  $\alpha p w_h$  ( $\text{Wm}^{-2} \text{ mm}^{-1}$ ), and (f)  $\alpha p w_l$  ( $\text{Wm}^{-2} \text{ mm}^{-1}$ ).

ature and moisture and suggest that these relationships can be used to infer errors in the geophysical parameters. In the following section we examine the differences between the synthetic GEOS2 CLR and the CLR from a preliminary CERES/TRMM data set.

### 5. A Comparison With CERES/TRMM CLR and RadWn

In the previous section we compared synthetic clear-sky fluxes (CLR and RadWn) computed from two 4-DDA data sets to develop some simple but quantitative measures of the sensitivity of clear-sky fluxes to changes in ground temperature and moisture. Insight into the sensitivity of the clear-sky fluxes to changes in these geophysical quantities was obtained by

computing separately the various band contributions to the clear-sky fluxes. Reasonable approximations to the sensitivity of clear-sky fluxes to changes in the geophysical parameters were obtained with simple linear relationships using ground temperature and moisture in two broad layers of the atmosphere (the upper and lower troposphere) to predict the clear-sky flux changes in the RadWn and RVR bands (or non-window bands): the single predictor relationships (5) and (6) as well as the two parameter relationships (7) through (10).

We found that compared with ECMWF, GEOS2 has substantially colder ground temperature during January 1998, with differences ranging from  $3^\circ\text{C}$  to more than  $7^\circ\text{C}$  over high latitudes. Over these regions the sensitivity of CLR to  $T_g$  is between  $1.2$  and  $1.8 \text{ Wm}^{-2}$  per  $1^\circ\text{C}$ , so differences in the CLR range from about  $5$  to  $15 \text{ Wm}^{-2}$  in these regions, while parts of

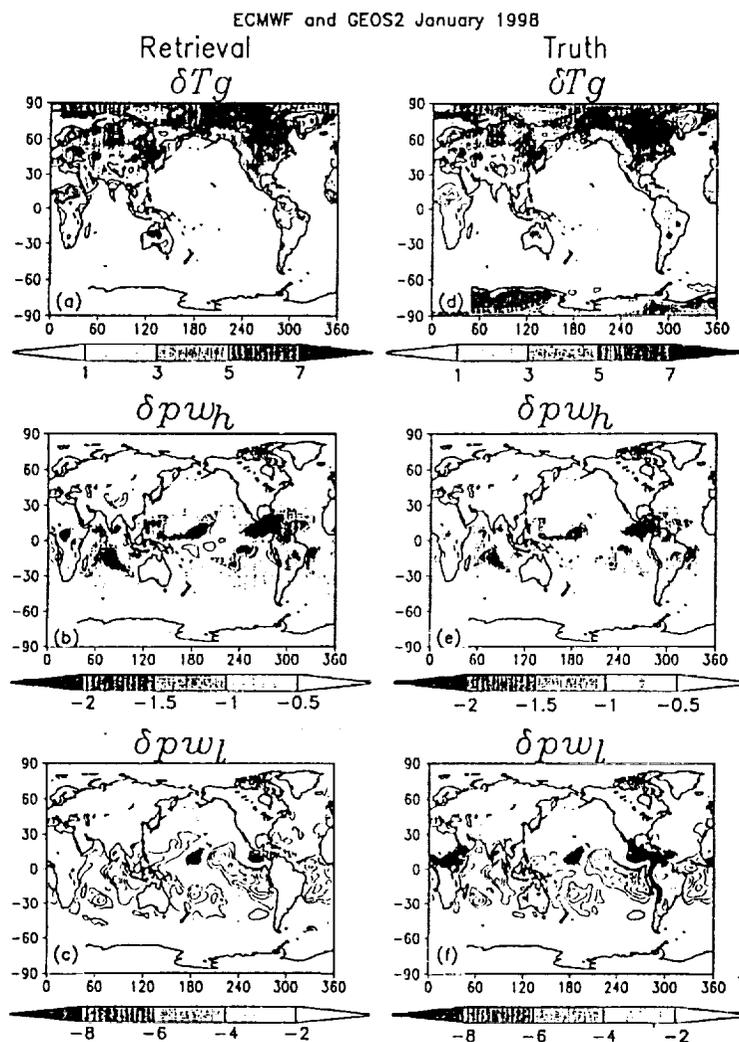


Figure 10. Retrieval of (a)  $\delta T_g$  ( $^{\circ}\text{C}$ ); contour intervals are  $1^{\circ}\text{C}$ ; (b)  $\delta p w_4$  (mm); contour intervals are 1 mm; and (c)  $\delta p w_1$  (mm); contour intervals are 2 mm. True values of (d)  $\delta T_g$  ( $^{\circ}\text{C}$ ); (e)  $\delta p w_4$  (mm); and (f)  $\delta p w_1$  (mm).

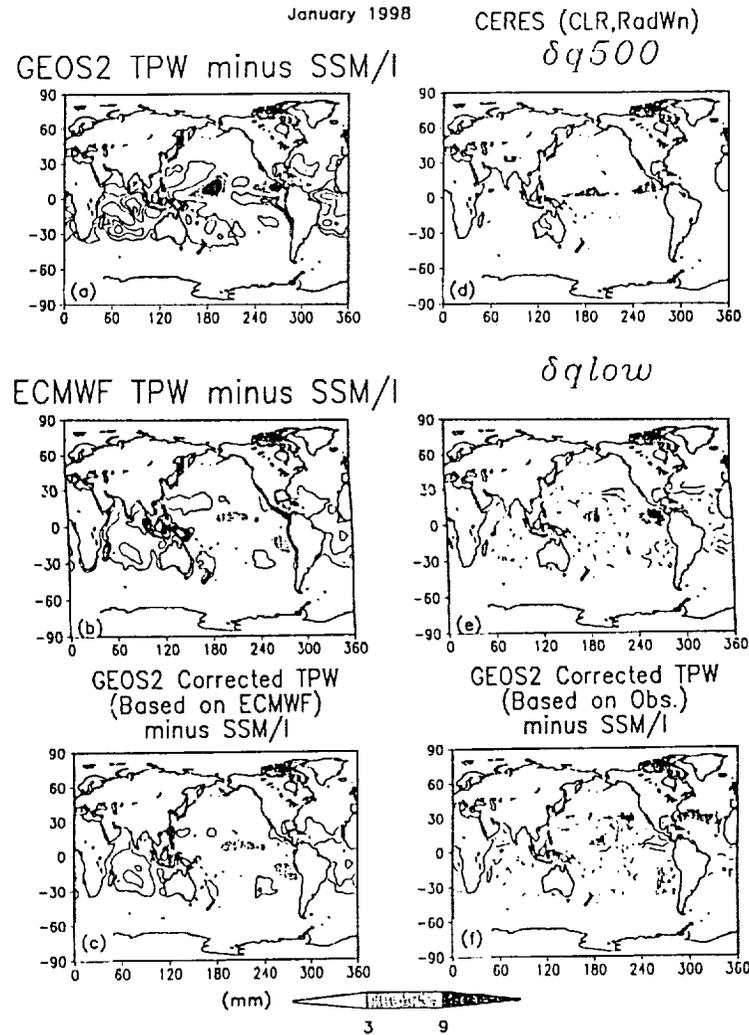
the Antarctic show differences exceeding  $20 \text{ Wm}^{-2}$ . We also showed (Figures 10e and 10f) that compared with ECMWF, GEOS2 is wetter in the upper troposphere during January 1998 with differences as large as 2.5 mm over the tropical oceans. In the lower troposphere, GEOS2 is drier over most of the tropical oceans, with differences as large as 10 mm.

The above results give us confidence in the methodology, though they only provide us with difference (or "error") fields relative to the ECMWF products. We will now use some preliminary CERES/TRMM total CLR and RadWn data for January 1998 (in place of the ECMWF analysis data) to evaluate the quality of the GEOS2 ground temperature over land and moisture over the oceans.

The results for the moisture are shown in Figure 11. We will validate our results against the SSM/I total precipitable water (TPW). For purposes of illustration and to demonstrate the reliability of our algorithm, we first apply our methodology to

the synthetic  $\delta\text{CLR}$  and  $\delta\text{RadWn}$  data. Here we assume that the ECMWF data are the ground truth. The results Figure 11c show that indeed, we have adjusted the GEOS2 TPW field to be close to that of ECMWF.

We next demonstrate our algorithm by applying it to the satellite data (CLR and RadWn) and assume that the satellite data are the ground truth. Figures 11d–11f show the results of applying our algorithm using the January 1998 CERES observations of  $\delta\text{CLR}$  and  $\delta\text{RadWn}$ . Figure 11d indicates that our moisture between 200 and 500 hPa is excessive over the tropical regions. Figure 11e indicates that at low levels GEOS2 is too dry in subtropical regions around  $\pm 20^{\circ}$ – $30^{\circ}$  in both hemispheres. The corrected GEOS2 TPW (based on Figures 11d and 11e) is shown in Figure 11f. Comparing with Figure 11a, we see that the corrections have removed the major bias in the GEOS2 TPW fields. We note some amount of overcorrection, especially in the Northern Hemisphere subtropics.



**Figure 11.** Diagnose of  $q$ : synthetic studies of TPW differences (mm): (a) GEOS2 minus SSM/I, (b) ECMWF minus SSM/I, and (c) GEOS2 corrected TPW based on ECMWF data minus SSM/I, and satellite-based TPW and layer (d) retrieved  $\delta p_{w_0}$ , based on  $\delta CLR$  and  $\delta RadWn$  from CERES, (e) the same as Figure 11d except for  $\delta p_{w_0}$ , and (f) GEOS2-corrected TPW based on Figures 11d and 11e. Contour intervals are for negative values, which are opposite to the shaded values.

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Finally, we diagnose  $T_g$  for January 1998. Since we do not have global ground temperature observations, we will again compare with the ECMWF analyses. Figure 12a is the difference map of the retrieved  $\delta T_g$  minus actual (Figure 10a minus Figure 10d) based on the ECMWF data. As noted previously, the global mean bias is  $-0.04^\circ\text{C}$ , and standard deviation is  $1.25^\circ\text{C}$ . This again gives us confidence that our algorithm can correct the ground temperature to within those error limits. Figure 12b is the mean difference in the CLR between GEOS2 and the CERES observations over the land areas between  $40^\circ\text{S}$  and  $40^\circ\text{N}$  (the fraction of the globe covered by the CERES instrument). Differences are negative over the Sudan, Saudi Arabia, and India. Positive values occur over southeastern Australia, southern Africa, the Andes, Central America, northern South America, and the Tibetan highlands. Figure 12c

shows the implied GEOS2  $T_g$  errors using the algorithm described above. The errors are quite large in some areas. For example, over the Sudan GEOS2 is more than  $10^\circ\text{C}$  too cold, while over southeast Australia, GEOS2 is more than  $12^\circ\text{C}$  too warm. Determining the errors in these estimates of the GEOS2 ground temperature bias will require, among other things, reliable estimates of the bias in the CERES measurements.

## 6. Discussion and Conclusions

We have introduced a simple but quantitative method for relating errors in model-based estimates of clear-sky longwave fluxes (CLR and RadWn) to errors in geophysical parameters. The primary motivation for this work is the underlying assumption that by linking the radiation errors to errors in the geo-

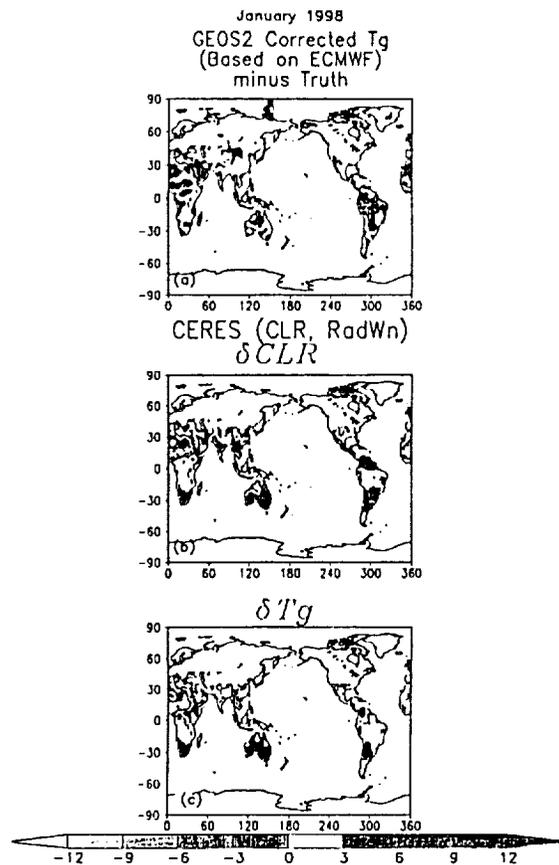


Figure 12. Diagnose of  $T_g$ : synthetic studies of  $\delta T_g$  ( $^{\circ}\text{C}$ ) based on ECMWF: (a) GEOS2-retrieved  $\delta T_g$  minus  $\delta T_g$  between ECMWF and GEOS2. Satellite diagnose of  $\delta T_g$ : (b)  $\delta\text{CLR}$  between CERES and GEOS2, and (c) retrieved  $\delta T_g$ . Contour intervals are for negative values, which are opposite to the shaded values.

physical parameters, we can provide greater insight to the model developer on potential model errors. We show for a test case (January 1998) that the method can potentially be used to obtain quantitative estimates of errors in ground temperature ( $\delta T_g$ ) and moisture ( $\delta\rho_w$ ) from satellite observations. In particular, our analysis of synthetic total and window region clear-sky flux differences (computed from two different assimilated data sets) shows that a simple linear regression employing  $\delta T_g$  and broad layer  $\delta\rho_w$  provides a good approximation to the full radiative transfer calculations, typically explaining more than 90% of the 6 hourly variance in the flux differences. These simple regression relations can be inverted to "retrieve" the errors in the geophysical parameters. Uncertainties (normalized by standard deviation) in the monthly mean retrieved parameters range from 7% for  $\delta T_g$  to about 20% for the lower tropospheric moisture ( $\delta\rho_w$ ).

Our initial application of the methodology employed an early CERES/TRMM data set (total and window region clear-sky longwave fluxes) to assess the quality of the GEOS2 data. The results showed that over the tropical and subtropical oceans, GEOS2 is, in general, too wet in the upper tropo-

sphere (mean bias of 0.99 mm) and too dry in the lower troposphere (mean bias of  $-4.7$  mm). We note that these errors, as well as a cold bias in the ground temperature, have largely been corrected in the current version of GEOS2 with the introduction of a land surface model, a moist turbulence scheme and the assimilation of SSM/I total precipitable water.

The methodology described in this paper was developed as a result of our efforts to validate monthly mean fields from the GEOS2 global data assimilation system, but the methodology should also prove useful for validating the climatological fields of global atmospheric models. The accuracy of the methodology depends on the accuracy of the radiation code, surface emissivity, the ozone profile, as well as the accuracy of the satellite estimates of total and window region clear-sky longwave fluxes. While the initial results are promising, further work is required to fully assess the sensitivity to errors in the input parameters and to account for the mismatch between satellite observations and grid-scale model fields.

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