

PROBING THE INFLOW/OUTFLOW AND ACCRETION DISK OF CYGNUS X-1 IN THE HIGH STATE WITH THE *CHANDRA* HIGH ENERGY TRANSMISSION GRATING

Y. X. FENG,^{1,2} A. F. TENNANT,³ AND S. N. ZHANG^{1,2}

Received 2003 May 2; accepted 2003 July 26

ABSTRACT

Cygnus X-1 was observed in the high state at the conjunction orbital phase (0) with the *Chandra* High Energy Transmission Grating (HETG). Strong and asymmetric absorption lines of highly ionized species were detected, such as Fe xxv, Fe xxiv, Fe xxiii, Si xiv, S xvi, Ne x, etc. In the high state the profile of the absorption lines is composed of an extended red wing and a less extended blue wing. The red wings of higher ionized species are more extended than those of lower ionized species. The detection of these lines provides a way to probe the properties of the flow around the companion and the black hole in Cyg X-1 during the high state. A broad emission feature around 6.5 keV was significantly detected from the spectra of both the *Chandra*/HETG and the *RXTE*/Proportional Counter Array. This feature appears to be symmetric and can be fitted with a Gaussian function rather than the Laor disk line model of the fluorescent Fe K α line from an accretion disk. The implications of these results on the structure of the accretion flow of Cyg X-1 in the high state are discussed.

Subject headings: accretion, accretion disks — binaries: general — black hole physics — stars: individual (Cygnus X-1) — X-rays: binaries

1. INTRODUCTION

Cygnus X-1 is a black hole candidate that consists of a compact object orbiting a $\sim 17 M_{\odot}$ O9.7 I supergiant with a period of 5.6 days (Webster & Murdin 1972; Bolton 1972). The dynamically determined mass of the compact object that significantly exceeds the neutron star upper limit provides evidence of a black hole in Cyg X-1. The companion star almost fills its Roche lobe in Cyg X-1 (Gies & Bolton 1986). Therefore, the configuration of material being transferred from the companion to the black hole may be intermediate between Roche lobe overflow and stellar wind accretion. The X-ray emission of Cyg X-1 is believed to be powered by the accreted stellar wind material from the companion. The X-ray spectrum is composed of soft and hard spectral components (see, e.g., Liang & Nolan 1984; Tanaka & Lewin 1995). Various spectral shapes have been found and classified as the so-called low, intermediate, and high spectral states (see, e.g., Oda 1977; Liang & Nolan 1984; Tanaka & Lewin 1995; Belloni et al. 1996). This classification is primarily based on the power-law index obtained by fitting the soft (2–10 keV) X-ray spectrum with a single power law. In the low state, the power-law index is about 1.6. In the intermediate state, the power-law index is around 2.0, and it is about 2.5 in the high state (see, e.g., Oda 1977; Liang & Nolan 1984; Tanaka & Lewin 1995; Belloni et al. 1996).

Further studies in timing reveal more differences among the states from the extensive observations of Cyg X-1 with *Ginga* and *Rossi X-Ray Timing Explorer (RXTE)*, such as the various shapes of power density spectrum (PDS) (see, e.g., van der Klis 1995; Belloni et al. 1996; Cui et al. 1997)

and various spectral evolutions in both short timescales (~ 1 s) (Feng, Li, & Chen 1999; Li, Feng, & Chen 1999) and long timescales (\sim several hours) (Wen, Cui, & Bradt 2001). Although there is no model that provides a consistent explanation of both the temporal and the spectral properties of the states, the various states have been attributed to different geometric structures and physical properties of the accretion flow (e.g., Esin, McClintock, & Narayan 1997) which may be closely related to the accretion rate and determined by the configuration of the mass transfer (e.g., Gies et al. 2003).

Spectroscopy of Cyg X-1 in the various states can greatly enhance the knowledge of the configuration of the mass transfer in the system. Redshifted absorption lines of highly ionized species have been detected with the *Chandra* High Energy Transmission Grating (HETG) during the low state at orbital phases 0.73 (Schulz et al. 2002), 0.84 (Marshall et al. 2001), and 0.77 (Miller et al. 2002a) and in intermediate states at orbital phase 0.88 (Y. X. Feng et al. 2003, in preparation). These lines show that there is a highly ionized inflow with a variable configuration. This inflow may play an important role in controlling the spectral state and in triggering the state transition. However, the previous observations of Cyg X-1 with *Chandra* covered only a limited range of orbital phase (0.73–0.88) and spectral states (the low and intermediate states), which prevents us from obtaining a complete picture of the stellar wind. Spectroscopy of Cyg X-1 in both the high and the low states around orbital phase 0 is particularly interesting because the column density may be enhanced and produce strong absorption lines. Furthermore, comparison of the properties of absorption lines detected in the different states and orbital phases will show the variations in the configuration of mass density and shed light on the mysteries of the spectral state and its transition.

The fluorescent Fe K α line may be a powerful way to probe the accretion disk in Cyg X-1 (Fabian et al. 1989). Fe K α line emission detected during the low state indicates that there is a cold accretion disk where the iron fluorescent

¹ Department of Physics, University of Alabama in Huntsville, Huntsville, AL 35899.

² National Space Science and Technology Center, NASA Marshall Space Flight Center, SD50, 320 Sparkman Drive, Huntsville, AL 35805.

³ NASA Marshall Space Flight Center, SD50, Huntsville, AL 35812; fengy@email.uah.edu, allyn.tennant@msfc.nasa.gov, zhangsn@email.uah.edu.

emission line is generated (e.g., Barr, White, & Page 1985; Ebisawa et al. 1996; Miller et al. 2002b). The detection of a ~ 6.6 keV broad line emission in 1996 during the intermediate state with *ASCA* (Cui et al. 1998) has been attributed to the Fe $K\alpha$ emission line from the highly ionized inner region of the accretion disk rather than the cold disk of the low state. In addition, in the high state, the inner edge of the accretion disk may extend down to the last stable orbit. The high resolution obtained with *Chandra*/HETG in the detection of the possible Fe $K\alpha$ line emission will constrain the line emission model. In this paper, we present results from the first spectroscopy of Cyg X-1 in the high soft state around orbital phase 0 with *Chandra*/HETG.

2. OBSERVATION AND ANALYSIS

Monitoring of Cyg X-1 with the *RXTE* All Sky Monitor (ASM) shows that from 2001 September to 2002 October the X-ray flux was 2–4 times higher than ~ 20 counts s^{-1} of the low state, as seen in Figure 1, indicating that the source might have entered the soft state (Zhang et al. 1997). Furthermore, the hardness ratio of less than 0.6 and the positive correlation coefficient between hardness ratio and total count rate were found during the observational period, as seen in Figure 1. These are consistent with Cyg X-1 having entered the high soft state (Wen, Cui, & Bradt 2001). A *Chandra* X-ray Center Director's Discretionary Time (DDT) observation (observation ID 3724) proposed to catch Cyg X-1 in the high state occurred from 2002 July 30 17:27:17 to 2002 July 31 01:52:08 with a simultaneous

RXTE observation. An *RXTE* Proportional Counter Array (PCA) spectrum of Cyg X-1 was found to be the typical shape of the high state, and the power density spectrum was in the shape of a power law. All these timing and spectral properties confirmed that Cyg X-1 was in the genuine high soft state during the observation. Also the companion star was closest along our line of sight, which is the so-called orbital phase 0 based on the ephemeris provided by Gies et al. (2003).

The *Chandra*/HETG observation was operated in continuous-clocking (CC) graded mode in order to eliminate pile-up. To reduce telemetry saturation, a 100 pixel region centered on the zeroth-order image was transmitted with a 10% duty cycle. However, because of the very high X-ray flux, the observation still suffered from telemetry saturation and frame dropping.

In CC mode ACIS reads out the CCD in virtual frames. This ensures that all the onboard processing is identical for CC and TE modes. Because the CCD is clocked differently, there is a slight gain difference between these modes. This gain is used only for order sorting, and we have verified that the order sorting is still selecting the correct data. For high count rate sources, data are generated faster than the telemetry can handle. If ACIS has processed a frame of data, which represents about 1.5 s of CC mode data from one CCD and that data will not fit in the telemetry stream, it discards the entire frame. This results in a data gap for that CCD but otherwise has no impact on the data or processing. However, it is important to realize that each CCD will now have a different integration time, which is easily calculated from the "Good Time Intervals" given in the FITS file.

To extract the spectrum of a point source from the grating instrument of *Chandra*/HETG, it is important to precisely locate the source position on the CCD. We located Cyg X-1's position on ACIS-S CCD chip S3 in three steps. First, we located the Cyg X-1 position from its zeroth-order image. Then, we compared it with the location calculated from the Galactic coordinate of Cyg X-1. Finally, we adjusted the source position to make the spectral features in both minus and plus order spectra line up. On the basis of the source position, we used CIAO 2.2.1 and followed the updated threads provided by CXC to extract the spectrum. The *Chandra*/HETG spectrum was also extracted with a software developed by us. The spectral features reported here are consistent in both spectra. For a bright X-ray source, such as Cyg X-1, the background detected with *Chandra*/HETG in the 0.7–4 keV region is negligible. The spectrum in this range is analyzed without background subtraction. When we searched for the broad feature near 6.5 keV, we did subtract a background that contributed about 4% of the counts in this band. The background was estimated at a given location from ACIS-S energies just above and below the dispersed grating energy for the location.

Before looking for the broad Fe $K\alpha$ emission line, we examined the HEG spectrum in the range of 4–5 keV. It showed that the plus and minus orders do not agree. The data that fall near the edge of S3 are about 10% lower than the corresponding data from S2. At a distance of about 250 pixels from the S3 edge, the two rates roughly agree. This ripple can also be seen in the data from the ACIS calibration source (see Fig. 5 of Townsley et al. 2002) and is weak at small CHIPY but increases as CHIPY increases. This ripple is due to a combination of serial and parallel clocking charge transfer inefficiencies.

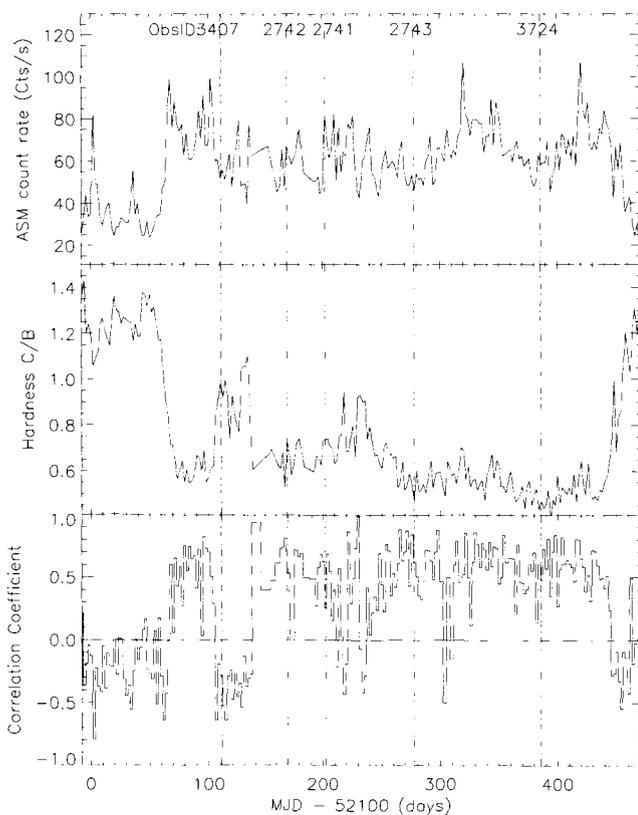


FIG. 1.—*RXTE*/ASM monitoring of Cyg X-1 from 2001 July 1 to 2002 November 1. The vertical lines in the plot show observation time with *Chandra*/HETG.

To correct for this ripple effect, we applied a simple correction factor to our ARF files. From the Mn line calibration source data, at the chip coordinate where the Cyg X-1 data were taken, we know that the efficiency is about 10% lower near the readout and close to full efficiency at the maximum distance (256 pixels). For the model we assumed that the correction was linear between these two points. This correction was applied to all four nodes on S3, resulting in a zig-zag correction. When we used the corrected ARF we now find much better agreement between the plus and minus orders.

When we jointly fit *Chandra* and *RXTE* spectra in the range between 2.5 and 15 keV, we allowed the *RXTE* absolute effective area to vary by a constant (independent of energy) amount. It is worth noting that although the *RXTE* and *Chandra* data overlap in time, the fact that the *Chandra* data were continuous, whereas the *RXTE* data were broken by gaps, means that the data are not from precisely the same time interval. Thus some of the differences could be due to the source variability. Furthermore, we added 1.0% systematic uncertainty to the *RXTE* data.

3. RESULTS

Fitting the spectrum with a single power law in the 2–10 keV energy range gives a power-law index of about 3.2, and although it was not a good fit, it does indicate that Cyg X-1 was in the high state during the observation. A better fit can be obtained by adding either a second power law or a disk blackbody. In the residuals of the fitting, absorption lines of highly ionized species show up significantly. A broad feature near 6.5 keV is also detected. In this paper we report the broad and asymmetric absorption lines and the broad emission feature. Comprehensive studies of the absorption lines will be reported elsewhere.

3.1. Asymmetric Absorption Lines

Strong and asymmetric absorption lines are detected in Cyg X-1 during the high state. The line profile is composed of an extended red wing and a less extended blue wing, as seen in Figure 2. The red wings of higher ionized species, such as Mg XII, Si XIV, S XVI, and Fe XXII, are more extended than that of lower ionized species, such as Mg XI, Si XIII,

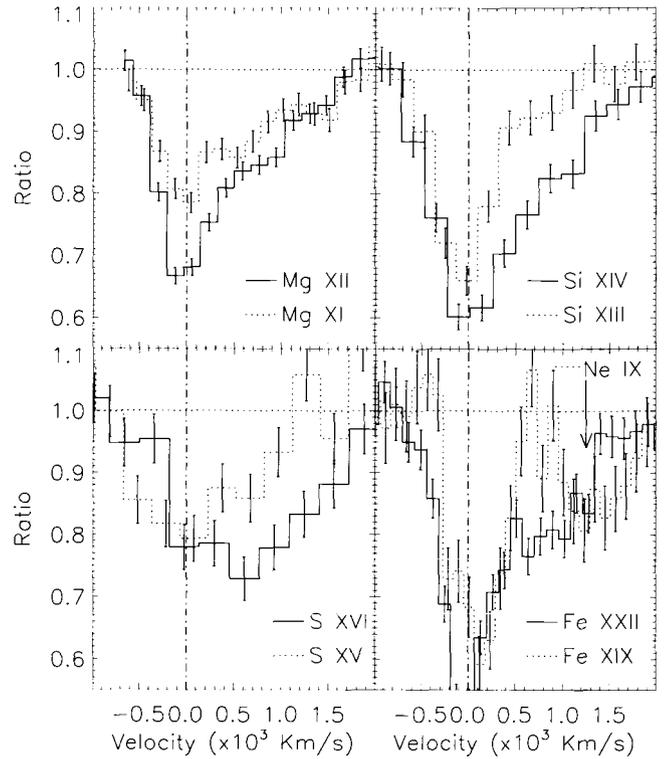


FIG. 2.—Strong absorption-line features were observed in the high state 2002 July 30 around orbital phase 0. The red wing of lower ionized species is less extended than that of higher ionized species.

S XV, and Fe XIX, as shown in Figure 2 and Table 1. Comparison of the absorption line of Mg XII detected in the different orbital phases and states shows that the bottoms of the absorption lines are near zero velocity (see Fig. 3), whereas those of previous observations in the low and intermediate states are redshifted. It shows also that an absorption line of the high state has a more extended blue wing than that of the low and intermediate states. During the high state (phase 0) the red wing is more extended than that in the low state (phase 0.74) and that in the intermediate state (phase 0.88), as shown in Figure 3.

TABLE 1
SELECTED ABSORPTION LINES IDENTIFICATION AND PROPERTIES

Ion	Rest (Å)	Measured ^a (Å)	Shift (km s ⁻¹)	Red Wing ^b (km s ⁻¹)	Blue Wing ^b (km s ⁻¹)	EQW (mÅ)	N_j^c (10 ¹⁵ cm ⁻²)	R^d
Mg XII	8.4192	8.4252 ^{+0.0018} _{-0.0014}	213 ⁺⁴⁸ ₋₆₃	1776	-534	10.0(0.3)	57.7(1.6)	2.6(0.2)
Mg XI	9.1688	9.171 ^{+0.014} _{-0.005}	55 ⁺¹⁷¹ ₋₄₆₆	1794	-655	7.5(0.4)	13.6(0.6)	4.9(1.2)
Si XIV	6.1822	6.1857 ^{+0.0005} _{-0.0005}	167 ⁺²² ₋₂₅	1935	-850	10.8(0.4)	77.0(2.8)	3.9(0.3)
Si XIII	6.6480	6.6469 ^{+0.0004} _{-0.0004}	-48 ⁺¹⁹ ₋₁₉	1350	-903	6.0(0.4)	20.4(1.5)	5.4(0.9)
S XVI	4.7229	4.7323 ^{+0.0018} _{-0.0004}	593 ⁺²⁵ ₋₁₁₅	1264	-688	6.6(0.6)	80.1(7.0)	6.0(1.5)
S XV	5.0387	5.0401 ^{+0.0010} _{-0.0015}	80 ⁺⁹² ₋₅₈	1039	-1043	4.2(0.5)	24.0(3.0)	4.4(0.4)
Fe XXII	11.770	11.7711 ^{+0.0013} _{-0.0011}	28 ⁺³⁴ ₋₃₄	1905	-701	19.3(0.7)	23.2(0.8)	2.8(0.2)
Fe XIX	13.518	13.5134 ^{+0.0006} _{-0.0013}	-102 ⁺²⁹ ₋₁₃	499	-499	10.2(0.8)	16.4(1.3)	3.0(0.5)

^a The wavelengths are calculated by fitting Gaussians to the line features. 1 σ errors are quoted.

^b Width is distance between the edge of wing and the deepest absorption, as seen in Fig. 2.

^c The column densities of the absorption lines are calculated by using, $W_\lambda = (\pi e^2 / m_e c^2) N_j \lambda^2 f_{ij}$, where N_j is the column density of a given species, f_{ij} is the oscillator strength, W_λ is the equivalent width of the line which is calculated by fitting the spectrum in a narrow range to a power law and a negative Gaussian, and λ is the wavelength of the line in centimeters (Spitzer 1978).

^d Ratio between the column densities derived from HEG spectra observed at phase 0 and at phase 0.77.

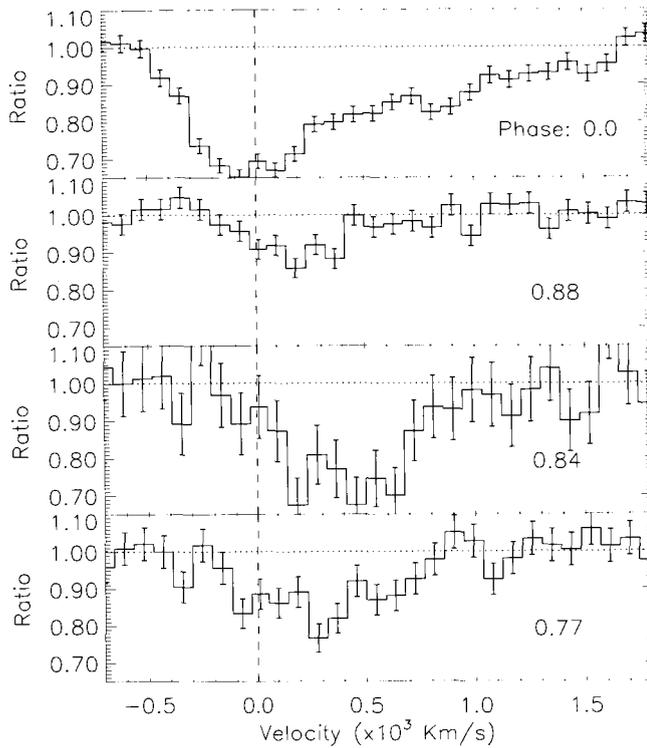


FIG. 3.—Comparison of Mg XII absorption line (8.4192 Å) features observed in the high state on 2002 July 30 around orbital phase 0 with that in the low state, observed on 2001 January 4, at phase 0.77 and 2000 January 12 at phase 0.84, and that in the intermediate state observed on 2001 October 28 at phase 0.88.

3.2. Broad Emission Feature

We first fitted the joint *RXTE* and *Chandra* spectra to a simple continuum model consisting of two power laws. This gave a χ^2 of 1573.6 for 835 degrees of freedom (dof); the residuals are plotted in Figure 4a. We next added three

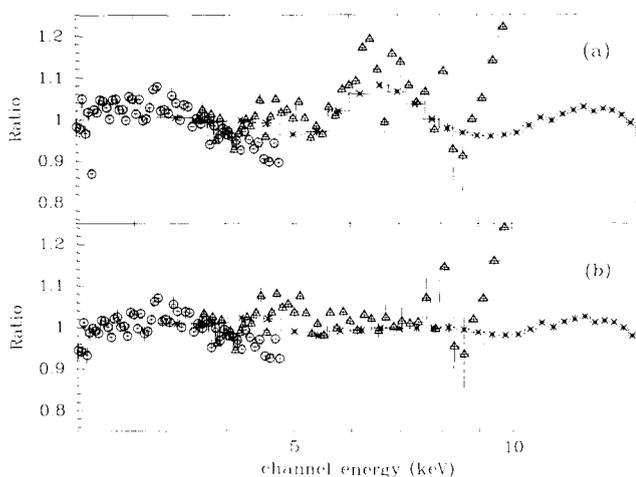


FIG. 4.—Broad emission feature is detected from jointly fitting spectra of Cyg X-1 during the high state observed with *Chandra*/HETG and *RXTE*/PCA. (a) Residuals of the fit with a simple continuum model consisting of two power laws; (b) residuals of the fit with the same continuum model as that of (a) and three Gaussians, two with negative norms (at 2.62 and 6.67 keV) for the hydrogen-like Si and the He-like Fe lines and one broad Gaussian at about 6.50 keV.

Gaussians, two with negative norms (at 2.62 and 6.67 keV) for the hydrogen-like Si and the He-like Fe lines. The third broad Gaussian was added at about 6.5 keV. The new χ^2 was 1060.8 for 826 dof, and the broad line has a central energy of 6.49 ± 0.04 keV and width σ of 0.51 ± 0.04 keV. If we replace the low-energy power law with a disk blackbody component then χ^2 becomes 1028.4, but the line is now centered at 5.44 ± 0.12 keV with width σ of 1.42 ± 0.07 keV. Errors are 68% confidence range for a single parameter.

We tried to replace the broad Gaussian with the more physical Laor (1991) model. At no time did the Laor model generate a better χ^2 than the Gaussian and often χ^2 was significantly worse. Thus with the two power-law continua, the Laor resulted in $\chi^2 = 1061.7$, and for this case the best-fitting parameters for the disk had an inner radius of $25r_g$ ($r_g = GM/c^2$), an outer radius of $100r_g$, and an inclination of 86° . With the disk blackbody plus power-law continuum, the best χ^2 obtained was 1266.4 for a “disk” extended from $23.5r_g$ to $26r_g$ and an inclination of 86° . For both these models we included the two negative Gaussians for the Si and Fe absorption lines. We also note that often the fitting routine would make the negative Fe line broad and positive, effectively discarding the Laor model; formally, this did result in a better fit. Physically reasonable Laor models tend to be asymmetric and the fitting tended to find nonphysical solutions which made the line symmetric. Although an asymmetry could be smeared by the lower resolution of the *RXTE* spectrum, the higher resolution of the HETG makes it impossible to hide such an asymmetry. Part of the problem is due to the disk blackbody and power-law components being roughly equally important in the 5–6 keV range. Since the disk blackbody spectrum is falling rapidly, the precise nature of the fall is sensitive to assumptions made. Merloni, Fabian, & Ross (2000) have shown that self-consistent disk models produce more flux at higher energies than the XSPEC *diskbb* model predicts. Thus the excess seen in 5–7 keV may be partially attributed to the underestimation of the flux at higher energies with the *diskbb* model.

4. DISCUSSION

Strong absorption lines of highly ionized species are detected in Cyg X-1 during the high state around orbital phase 0, as seen in Figure 3. This suggests that there is a high mass density in the region between the companion and the black hole, as shown in the simulation of high-mass binaries (Blondin, Stevens, & Kallman 1991). Similar species of absorption lines are found during observations in the high, intermediate, and low states. The slight change of the ionization state of the stellar wind is consistent with the X-ray flux increase by only a factor of 2–4 in the high state from that in the low state, if the photoionization is the dominated ionization process in Cyg X-1. Strong absorption lines of Fe xxv to Fe xvii have been detected, indicating a large range of the ionization states in the wind. The large variation in the ionization state is not surprising since the density of the wind decreases as it moves away from the companion and the irradiation flux depends on the distance to the compact object. This combination results in a large variation in the ionization parameter along the line of sight.

The absorption lines are broad and asymmetric and consist of an extended red wing and a less extended blue

wing. The extended red wing can be attributed to the inflow with a projected radial velocity larger than that of the outflow. The inflow is the region where the angle between the direction of motion of the material and our line of sight is larger than 90° . According to the wind structure of the radiatively driven wind model of Castor, Abbott, & Klein (1975), this region is closer to the X-ray source than the region where the blue wing is generated. The closer to the source the wind is, the higher photo-ionization parameter and higher ionization degree it should have. Indeed, the red wings of higher ionized species, such as Si XIV, S XVI, and Fe XXIII, are found to be more extended than that of lower ionized species, such as Si XII, S XV, and Fe XIX, as seen in Figure 2 and Table 1, consistent with the above picture of mass transfer. As the wind material approaches the X-ray source, its projected radial velocity increases because of the strong gravity around the black hole; in addition, the ionization parameter should also increase because the wind receives a higher X-ray flux. This picture is consistent with the existence of the lower (or non) ionized stellar wind at high latitude (far from the source), where the blueshifted absorption $H\alpha$ line is produced and has been detected by Gies et al. (2003). The projected radial velocity of the blueshifted absorption line of $H\alpha$ is $\sim 380 \text{ km s}^{-1}$, which is comparable with that of the blue wing seen in Figure 2. The asymmetric distribution of the projected radial velocity implies a focused stellar wind toward the black hole.

The 5.6 day orbital modulation was detected in the X-ray emission of Cyg X-1 during the low state (Holt et al. 1979; Brocksopp et al. 1999; Wen et al. 1999), which is consistent with the distribution of X-ray intensity dips measured with *RXTE*/ASM (Bałucińska-Church et al. 1999) and *RXTE*/PCA (Feng & Cui 2002). It suggests enhanced mass density in the region where the companion faces the black hole during the low state. However, no orbital modulation was detected from X-ray flux (2–12 keV) during the high state (Wen et al. 1999). Since the velocity distribution of the stellar wind with high ionization degree has been found to be broader than those in the other states, as seen in Figure 3, we think that one possible reason for the lack of orbital modulation is that the stellar wind is highly ionized in a larger region during the high state. Therefore, the column density of cold stellar wind could be lower and produce weaker modulation during the high state, which is consistent with the speculation of Wen et al. (1999). Another possibility is that the distribution of the stellar wind in the region where X-ray emission has been attenuated becomes more isotropic from the viewpoint of observer during the high state. If this is true, the difference between the stellar wind in the low and the high states implies that state transitions are triggered by variations in the configuration of the stellar wind. Comparison with a low-state HETG observation near phase 0 will help resolve which model is more correct.

Part of the highly ionized stellar wind in Cyg X-1 flows away from the black hole and the other part is accreted by the black hole. There are two possible ways for the accreted material to fall into the horizon of the black hole, i.e., through the corona or the accretion disk. If it is via the disk, the highly ionized stellar wind needs to condense to form the cold disk, for example, via Comptonization cooling of low-energy photons (Igumenshchev, Illarionov, &

Abramowicz 1999). As the highly ionized plasma is cooled down to form the cold disk, recombination emission lines are expected. However, no recombination line emission has been detected. Therefore, the highly (or fully) ionized stellar wind most likely goes into the horizon mainly via the corona. The existence of the corona during the high state is consistent with the presence of the hard/nonthermal component in Cyg X-1 during the high state, as well as in the low and intermediate states. In studies of the X-ray intensity dips, two types of dips have been found (Feng & Cui 2002). One type (called type A) exhibits strong energy-dependent attenuation of X-ray emission at low energies during a dip, which is characteristic of photoelectric absorption from cold material. The other type (called type B) shows nearly energy-independent attenuation, which is consistent with the partial obstruction of the X-ray emission region by dense streams. The dense stream could be cold and neutral, which may exist mainly around the equatorial plane and form the disk.

Comparison of the absorption lines observed in different orbital phases shows that the blue wing is more extended at phase 0 than those at other orbital phases. The blue wings may be attributed to the stellar wind flowing away from the source. Less extended blue wings and redshifted absorption lines suggest a focused wind structure. However, the variation may also be attributed to the change of the spectral state, because the absorption lines have been detected not only at different orbital phases but also in different states.

On the basis of the radiatively driven wind model (Castor et al. 1975), the column density viewed from phase 0 is expected to be about 2 times higher than that viewed from phase 0.77. The column densities in the lines observed at phase 0 in the high state are more than 2 times those at phase 0.77, as seen in Table 1. This implies a stronger stellar wind in the high state which can cause higher accretion rate. However, this interpretation is dependent upon the wind model which may not suitable to apply to Cyg X-1. Again, comparison with the *Chandra*/HETG observation of Cyg X-1 around orbital phase 0 in the low state will reveal the difference in stellar wind between the low and high states. Spectroscopy of Cyg X-1 with more complete orbital phase coverage in the low and high states is necessary to probe the structure of the mass transfer between the companion and the black hole in Cyg X-1.

There is a very broad and significant feature in the spectrum near 6.5 keV that can be modeled with a broad Gaussian function. Attempts to model the feature with a Laor model have resulted in worse fits. This is most likely due to the Laor model being asymmetric whereas the residuals seen in Figure 4a appear fairly symmetric. The precise shape of this excess depends on the continuum model, with small changes in the continuum resulting in large changes to the excess. Thus we have been unable to use this broad feature to derive physically meaningful parameters. Nevertheless, the appearance of this broad feature in the high state reveals that the spectral evolution in Cyg X-1 is much more complicated than that detected with low-resolution detectors. The spectral evolution suggests that the geometric structure and physical properties of the accretion flow may be associated with the spectral state.

We thank *Chandra* director Harvey Tananbaum and all the other members of the *Chandra* team for their enormous

efforts to conduct the DDT observation. We thank all the members of the *RXTE* team for their enormous efforts to conduct a simultaneous observation with *RXTE*. We thank M. C. Weiffkopf, M. K. Joy, and K. K. Ghosh for helpful discussions. We gratefully acknowledge the support for this work that was provided by NASA through Chandra award number GO2-3061X and DD2-3018X issued by the

Chandra X-Ray Observatory Center, which is operated by the Smithsonian Astrophysical Observatory for and on behalf of NASA under contract NAS 8-39073. This work was also supported in part by NASA Marshall Space Flight Center under contract NCC80200 and by NASA Long Term Space Astrophysics Program under grants NAG 5-7927 and NAG 5-8523.

REFERENCES

- Bałucińska-Church, M., Church, M. J., Charles, P. A., Nagase, F., LaSala, J., & Barnard, R. 2000, *MNRAS*, 311, 861
- Barr, P., White, N. E., & Page, C. G. 1985, *MNRAS*, 216, P65
- Belloni, T., Mendez, M., van der Klis, M., Hasinger, G., Lewin, W. H. G., & van Paradijs, J. 1996, *ApJ*, 472, L107
- Blondin, J. M., Stevens, I. R., & Kallman, T. R. 1991, *ApJ*, 371, 684
- Bolton, C. T. 1972, *Nature*, 235, 271
- Brocksopp, C., Tarasov, A. E., Lyuty, V. M., & Roche, P. 1999, *A&A*, 343, 861
- Castor, J. I., Abbott, D. C., & Klein, R. I. 1975, *ApJ*, 195, 157
- Cui, W., Ebizawa, K., Dotani, T., & Kubota, A. 1998, *ApJ*, 493, L75
- Cui, W., Zhang, S. N., Focke, W., & Swank, J. H. 1997, *ApJ*, 484, 383
- Ebisawa, K., Ueda, Y., Inoue, H., Tanaka, Y., & White, N. E. 1996, *ApJ*, 467, 419
- Esin, A. A., McClintock, J. E., & Narayan, R. 1997, *ApJ*, 489, 865
- Fabian, A. C., Rees, M. J., Stella, L., & White, N. E. 1989, *MNRAS*, 238, 729
- Feng, Y. X., & Cui, W. 2002, *ApJ*, 564, 953
- Feng, Y. X., Li, T. P., & Chen, L. 1999, *ApJ*, 514, 373
- Gies, D. R., & Bolton, C. D. 1986, *ApJ*, 304, 371
- Gies, D. R., et al. 2003, *ApJ*, 583, 424
- Holt, S. S., Kaluzienski, L. J., Boldt, E. A., & Serlemitsos, P. J. 1979, *ApJ*, 233, 344
- Igumenshchev, I. V., Illarionov, A. F., & Abramowicz, M. A. 1999, *ApJ*, 517, L55
- Laor, A. 1991, *ApJ*, 376, 90
- Li, T. P., Feng, Y. X., & Chen, L. 1999, *ApJ*, 521, 789
- Liang, E. P., & Nolan, P. L. 1984, *Space Sci. Rev.*, 38, 353
- Marshall, H. L., Schulz, N. S., Fang, T., Cui, W., Canizares, C. R., Miller, J. M., & Lewin, W. H. G. 2001, in *X-Ray Emission from Accretion onto Black Holes*, ed. T. Yaqoob & J. H. Krolik (Baltimore: Johns Hopkins Univ.), 45
- Merloni, A., Fabian, A. C., & Ross, R. R. 2000, *MNRAS*, 313, 193
- Miller, J. M., et al. 2002a, preprint (astro-ph/0208463)
- . 2002b, *ApJ*, 578, 348
- Oda, M. 1977, *Space Sci. Rev.*, 20, 757
- Schulz, N. S., Cui, W., Canizares, C. R., Marshall, H. L., Lee, J. C., Miller, J. M., & Lewin, W. H. G. 2002, *ApJ*, 565, 1141
- Spitzer, L. 1978, *Physical Processes in the Interstellar Medium* (New York: Wiley-Interscience), 52
- Tanaka, Y., & Lewin, W. H. G. 1995, in *X-Ray Binaries*, ed. W. H. G. Lewin, J. van Paradijs, & E. P. J. van den Heuvel (Cambridge: Cambridge Univ. Press), 126
- Townsley, L. K., Broos, P. S., Nousek, J. A., & Garmire, G. 2002, *Nucl. Instrum. Methods Phys. Res.*, 486, 751
- van der Klis, M. 1995, in *X-Ray Binaries*, ed. W. H. G. Lewin, J. van Paradijs, & E. P. J. van den Heuvel (Cambridge: Cambridge Univ. Press), 252
- Webster, B. L., & Murdin, P. 1972, *Nature*, 235, 37
- Wen, L., Cui, W., & Bradt, H. V. 2001, *ApJ*, 546, L105
- Wen, L., Cui, W., Levine, A. M., & Bradt, H. V. 1999, *ApJ*, 525, 968
- Zhang, S. N., Cui, W., Harmon, B. A., Paciesas, W. S., Remillard, R. E., & van Paradijs, J. 1997, *ApJ*, 477, L95