

DISENTANGLING X-RAY EMISSION PROCESSES IN VELA-LIKE PULSARS

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This grant was to support analysis of *XMM-Newton* observations of two young neutron stars, PSRs B1823–13 and B1046–58. The X-ray emission from these sources was expected to be a complicated combination of emission from an associated supernova remnant, from a wind-powered synchrotron nebula, from magnetospheric pulsations, and from the surface of the neutron star itself. It is only with *XMM*'s unique combination of spectral, temporal and angular resolution that all these different processes can be separated and studied.

Observations for both targets were carried out successfully. For PSR B1823–13, we detected two components to the X-ray emission: an elongated core of extent $30''$ immediately surrounding the pulsar, embedded in a fainter, diffuse component of emission $5'$ in extent, seen only on the southern side of the pulsar. The pulsar itself was not detected, either as a point source or through its pulsations. Both components of the X-ray emission were well fit by a power law spectrum, with photon index $\Gamma \approx 1.6$ and X-ray luminosity (0.5–10 keV) $L_X \approx 9 \times 10^{32}$ erg s $^{-1}$ for the core, and $\Gamma \approx 2.3$ and $L_X \approx 3 \times 10^{33}$ erg s $^{-1}$ for the diffuse emission, for a distance of 4 kpc. We interpreted both components of emission as corresponding to a pulsar wind nebula — we argued that the core region represented the wind termination shock of this nebula, while the diffuse component indicated the shocked downstream wind. We proposed that the asymmetric morphology of the diffuse emission with respect to the pulsar is the result of a reverse shock from an associated supernova remnant, which has compressed and distorted the pulsar-powered nebula. Such an interaction might be typical for pulsars at this stage in their evolution. The associated supernova remnant was not detected directly, most likely being too faint to be seen in existing X-ray and radio observations. This work was presented as a poster at the meeting “High Energy Studies of Supernova Remnants and Neutron Stars” in Houston TX in Oct 2002, and was published in May 2003 by Gaensler et al, *The Astrophysical Journal*, **588**, pp441–451 (2003). A copy of this publication is appended.

For PSR B1046–58, our previous *ASCA* observations showed a complex region of X-ray emission, possibly corresponding to several sources. Our *XMM* data confirmed this analysis, and allowed us to identify non-pulsed, non-thermal and extended X-ray emission at the position of PSR B1046–58, which we interpreted as an associated pulsar wind nebula. The surprisingly low X-ray luminosity of this nebula implies a conversion efficiency almost two orders of magnitude below the empirical relationship determined by Possenti et al (2002). Equally intriguing, the spectrum is much harder than that predicted by the empirical relationship of Gotthelf (2003). Taken by itself, we could simply dismiss PSR B1046–58 as an outlier. But when considered with PSR B1823–13 above and also PSR J1119–6127 (Gonzales & Safi-Harb 2003), both other pulsars that show no detected radio nebula, a clear trend emerges. In all three cases, the X-ray nebulae are under-luminous and their spectra steeper than expected. This suggests that other factors, e.g. low ambient densities, must be influencing the underlying processes that power these nebulae. This work was presented as a poster at the IAU Symposium “Young Neutron Stars and Their Environments” in Sydney Australia in Jul 2003. We are actively working on finalizing analysis and drafting a manuscript, that we hope to submit to *The Astrophysical Journal* by the end of 2003.

XMM-NEWTON OBSERVATIONS OF PSR B1823–13: AN ASYMMETRIC SYNCHROTRON NEBULA
AROUND A VELA-LIKE PULSAR

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ABSTRACT

We present a deep observation with the *X-Ray Multimirror Mission* of PSR B1823–13, a young pulsar with similar properties to the Vela pulsar. We detect two components to the X-ray emission associated with PSR B1823–13: an elongated core of extent $30''$ immediately surrounding the pulsar embedded in a fainter, diffuse component of emission $5'$ in extent, seen only on the southern side of the pulsar. The pulsar itself is not detected, either as a point source or through its pulsations. Both components of the X-ray emission are well fitted by a power-law spectrum, with photon index $\Gamma \approx 1.6$ and X-ray luminosity (0.5–10 keV) $L_X \approx 9 \times 10^{32}$ ergs s^{-1} for the core and $\Gamma \approx 2.3$ and $L_X \approx 3 \times 10^{33}$ ergs s^{-1} for the diffuse emission, for a distance of 4 kpc. We interpret both components of emission as corresponding to a pulsar wind nebula, which we designate G18.0–0.7. We argue that the core region represents the wind termination shock of this nebula, while the diffuse component indicates the shocked downstream wind. We propose that the asymmetric morphology of the diffuse emission with respect to the pulsar is the result of a reverse shock from an associated supernova remnant, which has compressed and distorted the pulsar-powered nebula. Such an interaction might be typical for pulsars at this stage in their evolution. The associated supernova remnant is not detected directly, most likely being too faint to be seen in existing X-ray and radio observations.

Subject headings: ISM: individual (G18.0–0.7) — pulsars: individual (PSR B1823–13) — stars: neutron — supernova remnants — X-rays: ISM

1. INTRODUCTION

Many of the energetic processes that take place in the aftermath of a supernova explosion are significant sources of X-ray emission. The supernova remnant (SNR) that results as the expanding debris collides with the interstellar medium (ISM) can generate both thermal X-rays from shock-heated gas and ejecta as well as nonthermal X-ray synchrotron emission from the shock acceleration of electrons. Meanwhile, the core collapse that triggers such explosions often forms a rapidly rotating central pulsar. Pulsars also produce X-rays in the form of pulsed non-thermal emission from the pulsar magnetosphere, quasi-blackbody emission from the neutron star surface, or an extended pulsar wind nebula (PWN) powered by a relativistic particle outflow from the central source.

Each of these processes dominates at different stages in the system's evolution. For example, the youngest systems are generally well suited for studying SNR blast waves and magnetospheric emission from pulsars, while in middle-aged systems one can better study the neutron star surface. Potentially the most challenging group of objects to understand are the “Vela-like” pulsars (Becker & Trümper 1997), which are at an evolutionary stage for which *all* the above processes may be taking place. Such pulsars, typified by the Vela pulsar, have spin periods $P \sim 100$ ms, spin-down luminosities $\dot{E} \equiv 4\pi^2 I \dot{P} / P^3 \sim 10^{36} - 10^{37}$ ergs

s^{-1} , and characteristic ages $\tau \equiv P/2\dot{P} \sim 10 - 20$ kyr (where \dot{P} is the pulsar's spin-period derivative and I is the star's moment of inertia). Until recently, about 10 Vela-like pulsars were known; the highly successful Parkes Multibeam Pulsar Survey has approximately doubled this number (Manchester et al. 2002).

The Vela pulsar itself is sufficiently nearby (~ 250 pc) that its different emission processes can be clearly distinguished: X-rays from the neutron star surface, the pulsar magnetosphere, the surrounding PWN, and the associated SNR have all now been identified (Strickman, Harding, & de Jager 1999; Lu & Aschenbach 2000; Helfand, Gotthelf, & Halpern 2001; Pavlov et al. 2001). However, the rest of the known Vela-like pulsars are at much greater distances, resulting in fainter, smaller, and more heavily absorbed X-ray sources. It has thus been difficult to detect and interpret the emission from these objects.

Radio observations of the pulsar B1823–13 have measured a spin period $P = 101$ ms and a period derivative $\dot{P} = 7.5 \times 10^{-14}$ $s s^{-1}$, implying a spin-down luminosity $\dot{E} = 2.8 \times 10^{36}$ ergs s^{-1} and a characteristic age $\tau = 21.4$ kyr (Clifton et al. 1992), all properties typical of a Vela-like pulsar. The distance to PSR B1823–13, as inferred from the dispersion of its radio pulses, is 3.9 ± 0.4 kpc (Cordes & Lazio 2002). Previous X-ray studies of PSR B1823–13 have been carried out with both *ROSAT* and *ASCA*. The *ROSAT* data showed a complicated source at the pulsar's position, with a morphology best modeled as a point source embedded in a compact nebula of radius $20''$ and surrounded by an elongated diffuse region $\sim 5'$ across (Finley, Srinivasan, & Park 1996). However, because of the low number of photons detected, it was not possible to determine whether the extended component of the source was an SNR or PWN, while the nature of emission from the pulsar itself also remained unclear. *ASCA* data also indicated the presence of one or more compact components embedded in

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a more extended region of emission but lacked the spatial resolution needed to properly separate and interpret these sources (Sakurai et al. 2001).

As the first part of an effort to understand the X-ray emission from Vela-like pulsars, here we present observations of pulsar B1823–13 with the *X-Ray Multimirror Mission* (*XMM-Newton*), the first observatory to provide sufficient angular, spectral, and temporal resolution to separately identify all the processes described above. *XMM-Newton* also has a much higher sensitivity (effective area 4650 cm² at 1.5 keV) than previous missions, making it well suited for studying faint sources such as seen here.

2. OBSERVATIONS AND ANALYSIS

Observations of PSR B1823–13 were carried out with *XMM-Newton* on 2001 October 16 and 18, in two observations each of length ~ 30 ks, as summarized in Table 1. The data described here correspond to the three X-ray imaging instruments on board *XMM-Newton*. The European Photon Imaging Camera (EPIC) first and second metal oxide semiconductor detectors (MOS1 and MOS2) were operated in the standard “full-frame” mode, in which seven CCDs in each detector are used to produce an approximately circular field of view of diameter 30', with a time resolution for each CCD frame of 2.6 s. The EPIC PN detector was operated in “small-window” mode, in which only a central 4' \times 4' square region is active but for which the time resolution is 5.7 ms. In this mode a significant fraction of each 5.7 ms frame is used to read out the CCD, resulting in a dead-time fraction of 30% (Kuster et al. 1999). To avoid optical contamination of the field, the medium and thin blocking filters were used for the MOS and PN detectors, respectively.

Initial processing of the data was carried out at the *XMM-Newton* Science Operations Centre (SOC). We analyzed the resulting event files using the *XMM-Newton* Software Analysis System (SAS), version 5.3.0. The data were first filtered to remove hot pixels and other bad data and to allow only standard event grades (patterns 0–12 for MOS1 and MOS2, patterns 0–4 for PN). Each of these data sets was examined for periods of high background by considering only events with energies in the range 10–15 keV. Times at which flares or high levels were seen in the X-ray count rate for this bandpass were excluded. This latter filtering excluded a significant fraction of each observation; the resulting useful exposure times are listed in Table 1.

2.1. Imaging

For each observation and detector listed in Table 1, the data set was corrected for vignetting losses using the SAS

TABLE 1
XMM-Newton OBSERVATIONS OF PSR B1823–13

DATE OF OBSERVATION	LENGTH OF OBSERVATION (ks)	EFFECTIVE EXPOSURE (ks)		
		MOS1	MOS2	PN
2001 Oct 16	30.5	17.8	18.3	10.2
2001 Oct 18	29.8	23.0	23.2	12.6
Total	60.3	40.8	41.5	22.8

NOTE.—Note that the effective exposures have been corrected for dead-time effects.

task EVIGWEIGHT. Next, an energy filter was applied to include only events falling in the energy range 0.5–10 keV, other energies being dominated by background. Finally, all the MOS and PN data were combined to form a single image for each type of detector. For the MOS CCDs, the flat-fielded images have a plate scale of 1".1 pixel⁻¹; for the PN CCD, the flat-fielded image has a plate scale of 4".4 pixel⁻¹.

2.2. Spectroscopy

The high background count rate for *XMM-Newton* requires careful analysis in order to extract accurate source spectra. The background contribution to the data consists of two main components (see Lumb et al. 2002): a diffuse X-ray background, assumed to have a smooth spatial distribution over the observed field, and a particle background, which shows spatial variations across the detector.

For a source not much larger than the point-spread function (PSF), we determine the spectrum by extracting data from both a small region enclosing the source and a reference spectrum immediately adjacent. The spectrum of interest is then obtained by subtracting the reference spectrum from the source spectrum, scaling appropriately to account for the differing areas of the extraction regions for the two fields. This approach assumes that because of the proximity of the two extraction regions, the differences in the background components are negligible between source and reference spectra.

Analysis of extended sources is complicated by the spatially varying nature of the background. Both the diffuse X-ray background and the source photons are vignettted by the *XMM-Newton* mirrors, resulting in a background component that varies dramatically with detector position. The particle background, however, is unaffected by the mirrors, resulting in a background component that is largely independent of detector position. Given these complications, we follow the prescription given in Appendix A of Arnaud et al. (2002), in which reference spectra from adjacent regions are used to correct for the X-ray background, and spectra from blank-field observations supplied by the SOC⁶ are used to account for the particle background.

Once these corrections had been applied, spectra were regrouped so that there were at least 50 counts per spectral bin for compact sources and 100 counts per bin for extended sources. For analysis of EPIC MOS data, we used the responses supplied by the SOC⁷ to provide information on the redistribution matrix and effective area of each CCD: specifically, for the MOS1 CCD we used the response file “m1_medv9q19t5r5_all_15.rsp,” while for MOS2 we used “m2_medv9q19t5r5_all_15.rsp.” For analysis of EPIC PN data, we generated our own response files using the SAS tasks RMFGEN and ARFGEN. Subsequent spectral fitting and analysis were carried out using XSPEC, version 11.1.0.

2.3. Timing

In the “small-window” mode used here, the EPIC PN CCD has sufficient time resolution to search for X-ray pulsations from PSR B1823–13. Because of improvements in the time tagging of events since the time when pipeline

⁶ See <http://xmm.vilspa.esa.es/ccf/epic/#background>.

⁷ See <http://xmm.vilspa.esa.es/ccf/epic/#responses>.

processing was carried out on the data by the SOC, we completely reprocessed the PN data from the raw observation data files, using the SAS task EPCHAIN. After filtering the data as described in § 2 above, we then corrected the arrival times of each event to correspond to a reference frame at the solar system barycenter. For a given extraction region and energy range, a pulsation analysis could then be carried out on the resulting arrival times.

3. RESULTS

3.1. Imaging

The EPIC MOS image of PSR B1823–13 is shown in Figures 1 and 2. Because a significant contribution to the background emission is from high-energy particles, the vignetting correction that has been applied to the data gives the false impression that the background emission is stronger around the edges of the fields. The O star HD 169727 (outside the extent of the region shown in Fig. 1) is detected as a prominent X-ray source; comparison of its position in our data with that in the Tycho Reference Catalog (Høg et al. 1998) suggests that the absolute astrometry of our observation is correct to within $0''.5$ in each coordinate.

The data clearly demonstrate that the X-ray emission associated with PSR B1823–13 consists of at least two components. The brightest X-ray emission is seen immediately surrounding the position of the pulsar, in a compact region shown in Figure 2 that we subsequently refer to as the

“core.” However, the morphology of the core suggests that it is not an unresolved source but rather is elongated along an extent of $\sim 30''$ in the east-west direction. This extension can be seen independently but at lower resolution in the EPIC PN image (not shown here). We quantify this in Figure 3, in which we show the intensity profile of X-ray emission around the pulsar along both east-west and north-south axes. We compare these profiles with the *XMM-Newton* PSF for MOS1 at 1.5 keV, which we generated using the King profile parameters included in the *XMM-Newton* calibration file “xrt1_xpsf_0005.ccf.”⁸ Above levels of $\sim 25\%$ of the peak flux in the image, the intensity profile in the north-south direction (i.e., in the bottom panel of Fig. 3) appears to be largely consistent with this PSF. However, the profile shown in the top panel of Figure 3 clearly demonstrates that along its east-west axis, the core region is significantly broader than the PSF, only falling to the apparent background level at offsets of $\approx 15''$ on the eastern and western sides of the pulsar. Also present is a secondary clump of emission $20''$ east of the pulsar, seen at coordinates R.A. = $18^{\text{h}}26^{\text{m}}14^{\text{s}}.5$, decl. = $-13^{\circ}34'47''$ (J2000.0) in Figure 2 and visible as a flux enhancement at an offset of $-20''$ in the top panel of Figure 3.

At larger angular scales, Figure 1 demonstrates the existence of a faint, amorphous component of emission

⁸ See <http://xmm.vilspa.esa.es/docs/documents/CAL-SRN-0100-0-0.ps.gz> for more information.

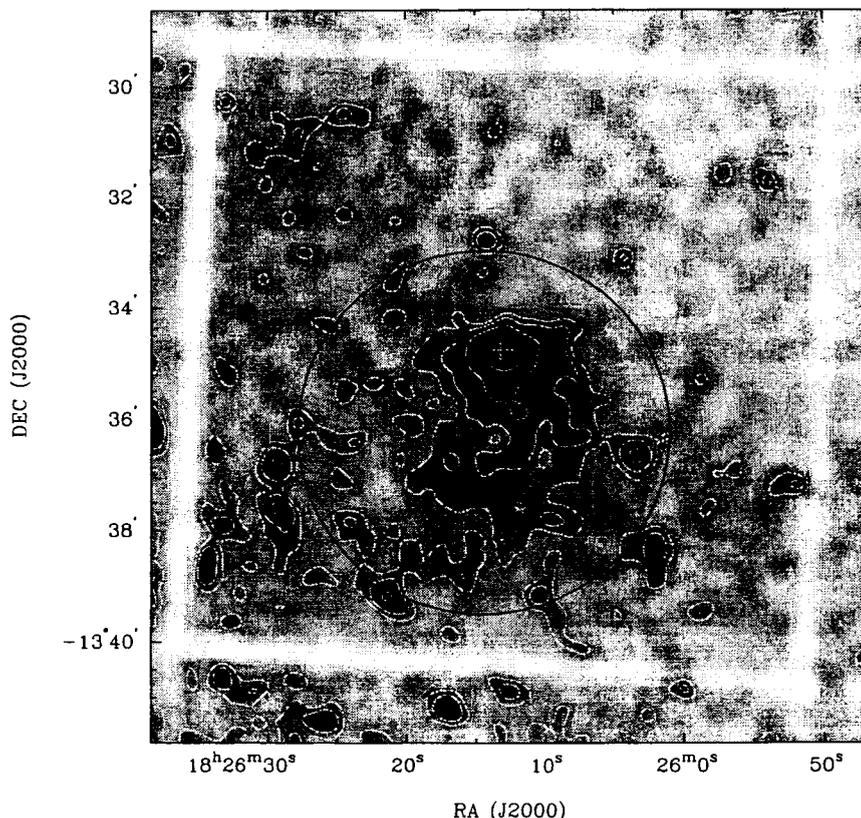


FIG. 1.—EPIC MOS image of the field surrounding PSR B1823–13, produced by combining all MOS1 and MOS2 data listed in Table 1 in the energy range 0.5–10 keV. The image has been corrected for off-axis vignetting and then convolved with a Gaussian of FWHM $20''$. The image is displayed using a linear transfer function ranging between 0% and 30% of the peak intensity in the field; the white contours are plotted at levels of 13%, 15%, 20%, 35%, 75%, and 90% of the peak. The white cross marks the position of the pulsar as listed in Table 3, while the circle shows the outer edge of the extraction region used for the spectrum of the diffuse component of the source. Linear structures in the image correspond to gaps between individual CCDs.

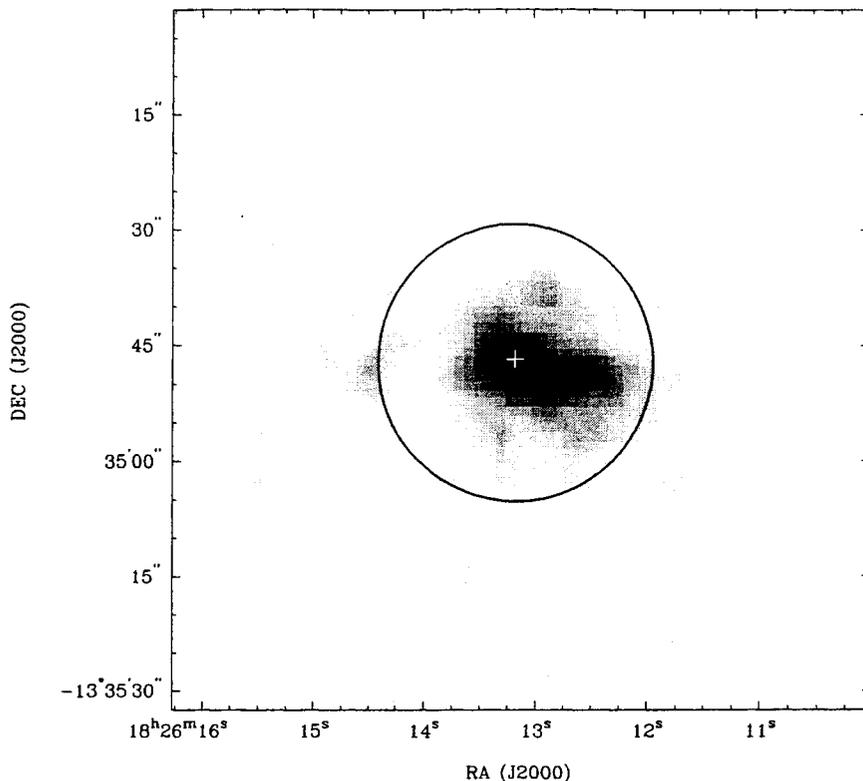


FIG. 2.—As in Fig. 1, but showing emission from the central $90'' \times 90''$ region surrounding the pulsar and smoothed with a Gaussian of FWHM $3''$. The image shows a linear transfer function ranging between 0% and 100% of the peak intensity. The white cross marks the position of the pulsar as listed in Table 3, while the circle shows the extraction region for the spectrum of the core.

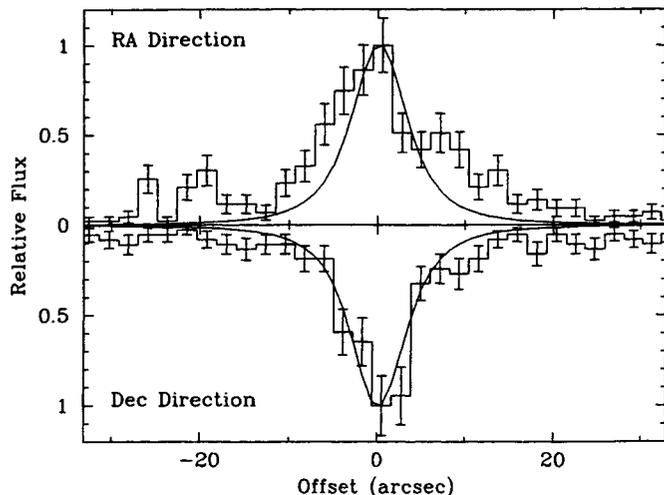


FIG. 3.—Intensity profiles through the core region of G18.0-0.7. The histograms plot the summed flux of EPIC MOS data in successive binned regions along each of the right ascension and declination axes. The size of each region is 2 MOS pixels ($2'' \times 2''$) along the axis of offset and 3 MOS pixels ($3'' \times 3''$) along the axis for which the data are summed. Positive offsets along the right ascension axis (*top*) correspond to more westerly points, while positive offsets along the declination axis (*bottom*) correspond to more northerly data; in both cases the pulsar position is at zero offset. The smooth curve on each plot is the profile of the EPIC MOS PSF at an energy of 1.5 keV, after similar summing as applied to the image.

extending to much larger distances from the pulsar than does the central core; we subsequently refer to this region as the “diffuse” component. This diffuse emission is located predominantly to the south of the core region. The morphology of the diffuse component is approximately elliptical, centered $\sim 1.5'$ south of the pulsar, with a major axis of $\sim 5'$ aligned north-south and a minor axis of $\sim 4'$. In all directions, the diffuse component decreases in brightness with increasing distance from the pulsar. To the south, east, and west of the pulsar, the diffuse component fades into the background over several arcminutes; to the north of the pulsar the brightness profile decreases much more rapidly. We demonstrate this in Figure 4, in which we show the brightness profile of the X-ray emission along the north-south axis of the source. As a comparison, we have extracted the same profile from the blank-field observations supplied by the SOC and have normalized the profile for the blank-field data to that of the source by comparing count rates in an identical source-free region. As shown in Figure 4, a comparison of profiles between the source field and the blank field clearly shows that X-ray emission from the source extends at least $300''$ to the south of PSR B1823-13 but extends no more than $\sim 40''$ to the north of the pulsar. Note that the ranges of offset shown in Figure 4 correspond to the extent of the central CCD of the detector. At larger offsets the background dominates, and in these regions the differences between the background levels in the source field and in the blank field make it impossible to determine whether there is any significant signal above the noise. For this reason we cannot place useful constraints on the presence of a

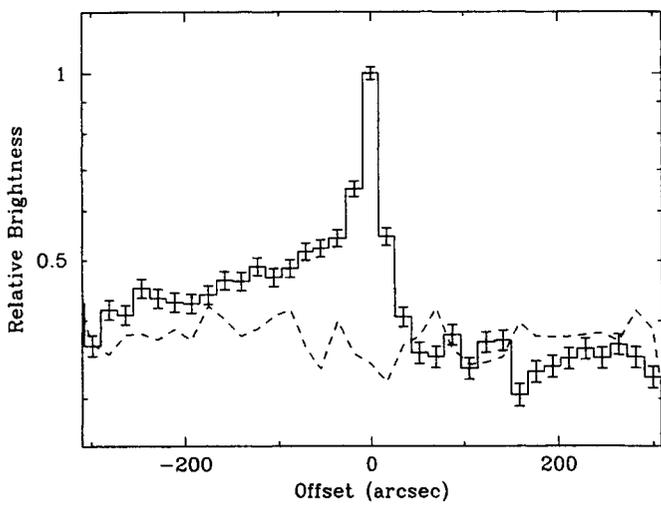


FIG. 4.—Intensity profiles through G18.0–0.7 along the north-south direction. The solid histogram plots the summed flux of EPIC MOS data in successive binned regions along the declination axis. The size of each region is 16 pixels (17"6) along the offset axis and 208 pixels (3'8) along the axis for which the data are summed. Increasing offsets correspond to more northerly data; the pulsar position is at zero offset. The broken line corresponds to an identical analysis for blank-field observations, after appropriate normalization (see text for details).

faint limb-brightened shell of emission at larger radii, as might be produced by an associated SNR (see further discussion in § 4.3 below).

3.2. Spectroscopy

We have extracted spectra for the core from both the MOS and PN data, using a circular extraction region centered on the pulsar and with a radius of 18" (as shown in Fig. 2). Although only ~75% of the energy from an unresolved or slightly extended source is encircled at this radius,

this extraction region minimizes contamination from the surrounding diffuse component. The core region is sufficiently compact that we can correct for background using the simple process of subtracting a reference region from an adjacent source-free region. In this case, the reference region was located ~3' distant from the core and was placed so as to avoid any contribution from the diffuse component of the emission.

Six spectra resulted, one for each of MOS1, MOS2, and PN for each of the two observations listed in Table 1. We then fitted power-law and blackbody models to these six spectra simultaneously, modified by photoelectric absorption; the best-fit results are summarized in Table 2. We find that the core region is well fitted by a power law, with an absorbing column $N_H = (1.0 \pm 0.2) \times 10^{22} \text{ cm}^{-2}$ and a photon index $\Gamma = 1.6^{+0.1}_{-0.2}$. At a distance to the system of 4d₄ kpc, the unabsorbed 0.5–10 keV luminosity implied by this fit is $L_X \approx 9d_4^2 \times 10^{32} \text{ ergs s}^{-1}$, where we have corrected for the ~25% of the flux from this source that presumably falls outside our extraction region. A blackbody model gives a much poorer fit to the spectrum.

The diffuse component occupies a much larger area on the detectors than does the core region and indeed fills the entire field of view of the EPIC PN image. Because there is no source-free region on the detector, it is impossible to accurately account for the background of the diffuse emission for the EPIC PN. Instead, we only considered MOS data in this analysis. We extracted a spectrum from a circular region centered on R.A. = 18^h26^m14^s.4, decl. = –13°36'14"8 (J2000.0), with radius 200", as shown in Figure 1. We excluded from this region a 30" radius region centered on the pulsar and 20" radius regions around the point sources seen at R.A. = 18^h26^m03^s.5, decl. = –13°36'41"7 (J2000.0) and R.A. = 18^h26^m10^s.6, decl. = –13°39'08"4 (J2000.0). As a reference spectrum, we considered all emission outside the main extraction region but falling within a square of side 9'5 centered on the pulsar. Both source and

TABLE 2
SPECTRAL FITS TO X-RAY EMISSION FROM G18.0–0.7

REGION	TOTAL COUNTS			MODEL	N_H (10^{22} cm^{-2})	$\Gamma/kT/E_0$ (.../keV/keV)	f_X^a ($10^{-13} \text{ ergs cm}^{-2} \text{ s}^{-1}$)	χ^2_ν
	MOS1	MOS2	PN					
Core ^b	455 ± 23	458 ± 22	815 ± 33	PL	1.0 ± 0.2	1.6 ^{+0.1} _{-0.2}	3.5 ^{+1.0} _{-0.8}	42/38 = 1.1
				PL	1.0 (fixed)	1.6 ± 0.1	2.9 ± 0.3	42/39 = 1.1
				PL	1.4 (fixed)	1.8 ± 0.1	3.0 ± 0.3	49/39 = 1.3
				BB	<0.2	1.19 ± 0.08	2.1 ± 0.5	60/38 = 1.6
Diffuse component	1843 ± 135	2360 ± 138	...	PL	1.4 ^{+0.5} _{-0.2}	2.3 ^{+0.4} _{-0.3}	18 ⁺¹² ₋₇	130/133 = 1.0
				PL	1.0 (fixed)	1.9 ± 0.1	13 ± 2	135/134 = 1.0
				PL	1.4 (fixed)	2.3 ^{+0.1} _{-0.2}	17 ± 3	130/134 = 1.0
				Brem	1.1 ^{+0.4} _{-0.2}	3.9 ^{+2.2} _{-1.1}	12 ⁺⁴ ₋₃	129/133 = 1.0
				SRCUT	1.4 ± 0.2	~1	15 ± 1	130/133 = 1.0
				RS	1.3 ^{+0.2} _{-0.3}	3.4 ^{+1.3} _{-0.8}	13 ⁺³ ₋₂	133/133 = 1.0
				MEKAL	1.3 ^{+0.2} _{-0.3}	3.4 ^{+1.2} _{-0.8}	13 ⁺³ ₋₂	136/133 = 1.0
				NEI	1.2 ^{+0.3} _{-0.2}	3.8 ^{+1.3} _{-1.1}	38 ⁺⁵ ₋₉	126/132 = 1.0

NOTES.—Uncertainties are all at 90% confidence. All models assume interstellar absorption using the cross sections of Bałucińska-Church & McCammon 1992, assuming solar abundances. Models used: "PL" indicates a power law of the form $f_\epsilon \propto \epsilon^{-\Gamma}$ where Γ is the photon index; "BB" indicates a blackbody emitting at a temperature T ; "Brem" indicates thermal bremsstrahlung; "SRCUT" indicates a synchrotron spectrum from a power-law distribution of electrons in a homogeneous magnetic field, with an exponential cutoff at a characteristic energy E_0 (Reynolds & Keohane 1999); "RS" indicates a Raymond-Smith spectrum of temperature T (Raymond & Smith 1977); "MEKAL" indicates a hot plasma model at temperature T ; "NEI" indicates a nonequilibrium ionization plasma at temperature T with solar abundances.

^a Flux densities are for the energy range 0.5–10 keV and have been corrected for interstellar absorption.

^b Counts and flux densities listed here for the core correspond to only ~75% of the total for this region (see text for details).

TABLE 3

RADIO TIMING EPHEMERIS FOR PSR B1823-13

Parameter	Value
R.A. (J2000.0).....	18 26 13.175(2)
Decl. (J2000.0).....	-13 34 46.7(2)
Frequency (Hz).....	9.85564535524(1)
Frequency derivative (Hz s ⁻¹).....	-7.2930848(2) × 10 ⁻¹²
Frequency second derivative (Hz s ⁻²).....	1.51530(6) × 10 ⁻²²
Epoch of ephemeris (MJD).....	50,750.000000
Range of validity (MJD).....	49,156-52,359

NOTES.—Units of right ascension are hours, minutes, and seconds, and units of declination are degrees, arcminutes, and arcseconds. These data have been supplied by M. Kramer from timing observations with Jodrell Bank Observatory. Numbers in parentheses indicate the uncertainty in the last significant figure of each value.

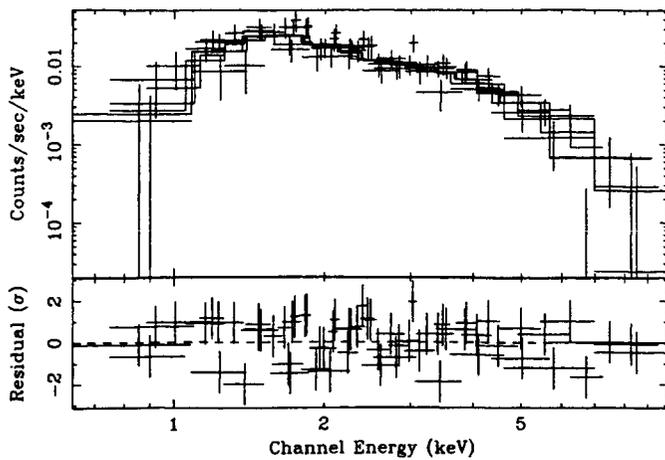


FIG. 5.—*XMM-Newton* spectrum of the diffuse component of G18.0-0.7. The data points in the top panel correspond to the data from each of the MOS1, MOS2 detectors for each of the two observations, while the solid lines show the corresponding best-fit absorbed power-law model. The bottom panel shows the number of standard deviations by which the model and the data differ in each bin. The data have been plotted so as to give a signal-to-noise ratio of at least 3 in each bin.

reference spectra were corrected for the particle background by subtraction of blank-field data, as explained in § 2.2. All four MOS spectra were then fitted simultaneously to a variety of absorbed models (see Table 2 for details): a power law, a thermal bremsstrahlung, a synchrotron spectrum with an exponential cutoff (model “SRCUT”), and three thermal plasma models (“RAYMOND,” “MEKAL,” and “NEI”). In all six cases, we find a good match to the data with an absorbing column of $N_H \sim 1.2 \times 10^{22} \text{ cm}^{-2}$. For the power-law fit shown in Figure 5, the photon index is $\Gamma \approx 2.3$, while the thermal models all imply temperatures $kT \sim 3.5 \text{ keV}$. The cutoff synchrotron spectrum can be extrapolated to radio wavelengths, at which the corresponding 1 GHz flux density implied for the source is $15 \pm 1 \text{ mJy}$.⁹

3.3. Limits on a Point Source

While it is clear that the core component of X-ray emission surrounding PSR B1823-13 is extended, an important consideration is how much of the X-ray emission detected could be from an unresolved source potentially corresponding to the pulsar itself. We here consider the upper limits we can place on the temporal, morphological, and spectral contributions of a central point source.

3.3.1. Timing Analysis

In principle, PSR B1823-13 may be directly detectable in these data through its pulsations. Because any emission from the pulsar is embedded in the surrounding extended component, pulsations are inevitably contaminated by this extended emission, and the signal-to-noise ratio of the pulses thus depends on the extraction radius and energy range in which data are considered. We therefore searched for pulsations by considering events in five possible circular extraction regions of radii 2'', 5'', 10'', 15'', and 20'', each centered on the pulsar position. For each extraction radius, we considered photons in the energy ranges 0.5-2, 2-5, 5-10, and 0.5-10 keV. We then determined the phase of each

photon by folding the data at the ephemeris listed in Table 3, which was kindly supplied to us by M. Kramer from pulsar timing observations at the Jodrell Bank Observatory. Because our two observations were only separated by 2 days, we combined data from both runs when searching for pulsations.

With the phases corresponding to the known pulsar period in hand, we applied to the data the Z_n^2 test (Bucherer et al. 1983), a binning-independent method of looking for the n th harmonic of a pulsating signal in sparsely sampled data. In the absence of pulsations, the quantity Z_n^2 will be distributed like χ^2 with $2n$ degrees of freedom. For each of the extraction radii and energy ranges listed above, we have calculated Z_n^2 for our observations for harmonics in the range $1 \leq n \leq 5$. Among these data, the most significant statistic is $Z_2^2 = 9.5$, which occurs for an extraction radius of 15'' and an energy range 0.5-10 keV. The probability of such a signal emerging by chance is 4.9%, which is consistent with chance occurrence given the number of different extraction radii and energy ranges searched. We therefore conclude that there is no significant pulsed signal in the data.

We can estimate an upper limit on the amplitude of any pulsations as follows: For a sinusoidal signal of pulsed fraction f embedded in an event list containing N photons, we expect $Z_1^2 = 0.5f^2N$ (Leahy, Elsner, & Weisskopf 1983). Since we do not know the optimal extraction radius in which pulses are most likely to dominate over surrounding extended emission, we conservatively consider a 20'' extraction radius. This corresponds to $\sim 75\%$ of the energy of a central point source¹⁰ but also most likely suffers from significant contamination from surrounding emission. The resulting values of Z_1^2 can be used to determine upper limits on f for the entire core region. For 0.5-2, 2-5, 5-10, and 0.5-10 keV, we determine upper limits on f for a sinusoidal pulsed signal from PSR B1823-13 of 11%, 4%, 14%, and 2%, respectively. The upper limits on the pulsed fraction for an unresolved pulse are a factor of 2 lower in each case.

3.3.2. Imaging Analysis

The intensity profiles in Figure 3 show that while cuts through the pulsar along a north-south axis are consistent with a point source, those in an east-west direction demonstrate the presence of extended emission immediately

⁹ A millijansky is $10^{-29} \text{ W m}^{-2} \text{ Hz}^{-1}$.

¹⁰ See http://heasarc.gsfc.nasa.gov/docs/xmm/uhb/XMM_UHB/node17.html.

surrounding the pulsar. However, unlike other cases in which an unresolved source is clearly superposed on a smooth nebula (e.g., Murray et al. 2002), there is no obvious decomposition of the profile seen in the top panel of Figure 3 into two such separate components. The profile shown could be interpreted as a centrally peaked nebula, with no central point source, but it could equally be regarded as an unresolved source embedded in lower level nebular emission. Thus, we must adopt a conservative upper limit on a point source, in which we assume that all the flux in the central bin corresponds to unresolved emission. Integrating under the appropriately normalized PSF, we find that in the energy range 0.5–10 keV, ~ 1010 counts in the combined MOS1+MOS2 data set could be from an unresolved source at the pulsar position. (Note that only $\sim 75\%$ of these counts would fall within the $18''$ extraction radius shown in Fig. 2, owing to the large wings of the *XMM-Newton* PSF.) The corresponding upper limit on the EPIC MOS count rate from the pulsar is 0.0123 counts s^{-1} .

3.3.3. Spectral Analysis

The spectrum of the core region is well fitted by a power law. If the emission from the pulsar has a blackbody spectrum, then it should be spectrally distinct from the rest of the core. Thus, we can constrain the temperature of a central source by determining what contribution a thermal source could make to the core's spectrum (Slane, Helfand, & Murray 2002). Specifically, we adopt a hard upper limit on the foreground column density of $N_H < 2 \times 10^{22}$ cm^{-2} , above all values of N_H listed in Table 2. We assume that the neutron star's radius as viewed by a distant observer is $R_\infty = 12$ km and that the distance to the source is 4 kpc. We can then fit a two-component spectral model to the core's spectrum (as measured by the combination of MOS1, MOS2, and PN), in which we allow the photon index and flux of the power-law component to be free parameters and find the maximum temperature of the blackbody component that still yields an acceptable fit to the spectrum. This results in an upper limit on the surface temperature of the blackbody (as viewed at infinity) $T_\infty < 147$ eV and an unabsorbed bolometric luminosity $L_\infty^{bol} < 8.7 \times 10^{33}$ ergs s^{-1} . We note that a blackbody model will overestimate a neutron star's temperature when compared to more realistic atmosphere models (Pavlov & Zavlin 2000), strengthening the upper limits made from blackbody fits here.

4. DISCUSSION

4.1. Comparison with Earlier Results

ROSAT observations of this source suggested the presence of three components of emission: an unresolved source centered on the pulsar embedded in a compact region of radius $20''$ and further surrounded by a diffuse nebula of extent $\sim 5'$ (Finley et al. 1996). From the *ROSAT* data, there was the suggestion that this latter diffuse component had a cometary morphology, with a fanlike tail extending out behind the pulsar at a position angle $\sim 210^\circ$ (north through east). In the data presented here, we confirm the presence of both a compact region surrounding the pulsar and a diffuse nebula at larger extents. However, our results differ from those of Finley et al. (1996) in two important respects.

First, based on the analyses in § 3.3 above, we can only put upper limits on the contribution of a central unresolved

source. Given the many more counts detected here, we ascribe the apparent detection of such a source by *ROSAT* to statistical fluctuations in the presence of a low signal-to-noise ratio. We can estimate the expected contribution from the pulsar as follows: The X-ray spectrum of the Vela pulsar can be modeled as a blackbody with $R_\infty = 2.1$ km, $T_\infty = 128$ eV, and $L_\infty^{bol} = 1.5 \times 10^{32}$ ergs s^{-1} , combined with a power-law component of photon index $\Gamma = 2.7$ and an unabsorbed 0.2–8 keV luminosity $L_X = 4.2 \times 10^{31}$ ergs s^{-1} (Pavlov et al. 2001). The pulsed fraction in the 0.1–10 keV band is 7% (Helfand et al. 2001). At a distance of 4 kpc, and adopting $N_H = 1 \times 10^{22}$ cm^{-2} as representative of the minimum of the range of possible absorbing columns derived in Table 2, we would expect to detect ~ 35 EPIC PN counts in our observation from such a source. When embedded in the surrounding core region, the effective pulsed fraction would be less than 1%, well below the limits determined in § 3.3.1. The resulting MOS count rate for such a source would be $\sim 4 \times 10^{-4}$ counts s^{-1} , also well below the upper limit on the point-source count rate determined in § 3.3.2. Finally, the expected temperature and luminosity of the blackbody component are both much less than the spectral limits on such a source determined in § 3.3.3. Thus, the failure to detect PSR B1823–13 in our data is entirely consistent with the emission properties expected of a Vela-like pulsar. Higher resolution observations with *Chandra* will be needed to separate out any unresolved source from the surrounding extended component and to thus directly detect the pulsar.

Second, while we confirm the presence of diffuse emission surrounding the core region, we find no evidence for the cometary morphology proposed by Finley et al. (1996). Rather, it seems likely this appearance resulted from the much lower sensitivity of their data (~ 280 total source counts for the *ROSAT* Position Sensitive Proportional Counter [PSPC] compared to ~ 4200 for the *XMM-Newton* EPIC MOS) and simply corresponded to the brightest regions of the larger source seen here.

Our measurements also differ somewhat from the *ASCA* results presented by Sakurai et al. (2001). While Sakurai et al. (2001) concluded that the X-ray emission consists of a core region embedded in a more diffuse nebula, the diffuse component in their data seems symmetrically distributed around the pulsar, as opposed to the distinctly one-sided morphology seen here. Sakurai et al. (2001) find that the core and diffuse components can both be fitted with power-law spectra, with photon indexes of $\Gamma = 1.9 \pm 0.2$ and 2.0 ± 0.1 , respectively. While these spectra are broadly consistent with our measurements, the unabsorbed flux densities (2–10 keV) for the core and diffuse regions inferred from *ASCA* are $\sim 2 \times 10^{-11}$ and $\sim 2 \times 10^{-12}$ ergs cm^{-2} s^{-1} , respectively, an order of magnitude higher than those inferred here. Since Sakurai et al. (2001) provide no information on their analysis procedure, fitted absorbing column, or assumed distance, it is difficult to assess whether there is a genuine discrepancy between our results and theirs. Pending a more detailed analysis of the *ASCA* data, we ascribe these differences to the much poorer angular resolution and sensitivity of *ASCA* when compared to *XMM-Newton*.

4.2. Morphology of the X-Ray Emission

Spectral fits to the diffuse component of the X-ray emission cannot distinguish between thermal and nonthermal

models. In the case of a thermal spectrum, this emission could be interpreted as shock-heated gas from an SNR associated with the pulsar. However, the temperatures inferred by the corresponding fits listed in Table 2 are higher than seen even in very young (<2000 yr) SNRs (e.g., Blanton & Helfand 1996; Gotthelf & Vasisht 1997). Fits with enhanced abundances can lower the inferred temperature (e.g., Pivovarov et al. 2001); we have attempted to model the data in this way here, but the quality of the fit then becomes unacceptable. We thus conclude that the spectrum of the diffuse component is most likely nonthermal and is best described by either a power-law or synchrotron-cutoff spectral model.

The synchrotron-cutoff spectrum (model “SRCUT” in Table 2) corresponds to synchrotron emission from shock-accelerated electrons in SNRs (Reynolds & Keohane 1999). It is not unreasonable to suppose that an SNR associated with PSR B1823–13 is accelerating such electrons and generating the diffuse component of the emission. Indeed, the photon index of a power-law fit to the spectrum, $\Gamma \approx 2.3$, is similar to that seen from other SNRs emitting synchrotron X-rays when fitted with such a model (Slane et al. 1999, 2001). Furthermore, the inferred 1 GHz flux density from the SRCUT fit, 15 ± 1 mJy, is well below the upper limit on 1.4 GHz radio emission from this source (~ 50 mJy, 3σ) as determined by Braun, Goss, & Lyne (1989) from their radio observations of this field. However, in this interpretation the morphology of the X-ray emission should take the form of an approximately circular limb-brightened region associated with the expanding SNR shock (e.g., Koyama et al. 1995). The amorphous, centrally brightened emission seen here does not meet this expectation.

We therefore interpret the diffuse component of the X-rays seen surrounding PSR B1823–13 as synchrotron emission produced by a PWN, which we designate G18.0–0.7 based on its Galactic coordinates. The fact that the diffuse component steadily fades with increasing distance from the pulsar supports the interpretation that PSR B1823–13 is the central energy source for this emission. The inferred photon index for the diffuse component, $\Gamma \sim 2.3$, is typical of PWNe (Becker & Trümper 1997; Possenti et al. 2002). The unabsorbed 0.5–10 keV luminosity for this source is $L_X \sim 3d_4^2 \times 10^{33}$ ergs s⁻¹; the corresponding efficiency with which the pulsar’s spin-down is converted into nebular X-rays is $L_X/\dot{E} = 1.1 \times 10^{-3}$, a value also typical for young pulsars (Becker & Trümper 1997).

Contained completely within the diffuse component, the core region is a compact region with a hard power-law spectrum centered on the pulsar. Bright central structures are seen around many other young pulsars that power PWNe (e.g., Helfand et al. 2001; Gaensler et al. 2002a), and similarly we argue here that the core region is a central region of synchrotron emission powered by the pulsar. We thus interpret both the core and diffuse regions of X-ray emission as parts of a single PWN of total unabsorbed luminosity (0.5–10 keV) $L_X \approx 4d_4^2 \times 10^{33}$ ergs s⁻¹ and foreground column density $N_H \approx 1.2 \times 10^{22}$ cm⁻². Spectral fits with the “SRCUT” model demonstrate that a continuous synchrotron spectrum for G18.0–0.7 extrapolated to radio frequencies results in a 1 GHz flux density well below the sensitivity of radio searches for a nebula of this size, accounting for the nondetection of any PWN surrounding PSR B1823–13 at radio wavelengths (Braun et al. 1989; Gaensler et al. 2000).

In such a source, we expect that the pulsar wind flows freely out to the point at which the wind pressure is balanced by that from the nebula. In this region the wind decelerates, a termination shock is formed, and synchrotron emission is generated. Recent high-resolution observations with the *Chandra X-Ray Observatory* have provided beautiful images of termination shocks around young pulsars. In some cases these images demonstrate that the termination region is not a spherical shock but rather has a toroidal morphology presumably corresponding to a wind focused into the pulsar’s equatorial plane (Gaensler et al. 2002a; Lu et al. 2002).

While we generally expect any X-ray PWN to be brightest near the pulsar, Figure 3 demonstrates that the core is a distinct, well-defined region superposed on the much fainter diffuse background. We thus argue that the core component seen here represents the pulsar termination shock, its elongated morphology possibly indicative of anisotropic outflow from the pulsar. The radius of the termination shock is then $r_s \approx 15'' = 0.3d_4$ pc, an extent similar to that seen around the Crab pulsar and PSR B1509–58 (Weisskopf et al. 2000; Gaensler et al. 2002a) but about an order of magnitude larger than that inferred for the Vela pulsar (Helfand et al. 2001). We expect particles emitting in this region to have undergone minimal synchrotron losses and to thus have a harder spectrum than the rest of the nebula (e.g., Slane et al. 2000; Gaensler et al. 2002a). We can directly compare the spectra of the core and diffuse components by fitting each to a power law with fixed N_H . The results of such fits are given in Table 2 for $N_H = 1.0 \times 10^{22}$ cm⁻² (corresponding to the best-fit absorbing column for the core) and $N_H = 1.4 \times 10^{22}$ cm⁻² (corresponding to the best fit for the diffuse region). From these fits we see that the core is indeed harder than the rest of the nebula at a statistically significant level, with a difference in photon index $\Delta\Gamma \approx 0.3$ –0.5.

The X-ray luminosity of the core region is about 25% of that from the total nebula. This is somewhat higher than the $\sim 10\%$ contribution made by the central termination shock region in other such sources (Slane et al. 2000; Gaensler et al. 2002a; Lu et al. 2002), potentially indicating a difference in the efficiency or length scale over which electrons and positrons are accelerated in the shock seen here compared to other sources. However, we caution that a significant contribution to the flux in the core region could be from the pulsar itself (see § 3.3.2 above).

If the core indeed represents the termination shock, the extent of this region then provides a constraint on the magnetic field in G18.0–0.7. Specifically, we expect pressure balance between the pulsar wind and the surrounding nebula,

$$P_{\text{neb}} \approx \frac{\dot{E}}{4\pi\Omega r_s^2 c}, \quad (1)$$

where P_{neb} is the pressure in the PWN and Ω is the fraction of a sphere through which the wind flows. The limited angular resolution of *XMM-Newton* makes it difficult to derive a specific geometry for outflow from the core morphology; we here leave Ω as a free parameter with a hard upper limit $\Omega \leq 1$. The corresponding nebular pressure needed to confine the pulsar wind is $P_{\text{neb}} = 9.2\Omega^{-1}d_4^{-2} \times 10^{-12}$ ergs cm⁻³. At distances far downstream of the shock, we expect approximate equipartition between particles and magnetic fields (Kennei & Coroniti 1984), so that $B^2/8\pi \approx 0.5P_{\text{neb}}$,

where B is the mean nebular magnetic field. In this case we can infer $B \approx 10\Omega^{-1/2}d_4^{-1} \mu\text{G}$. For electron/positron pairs emitting at 1.5 keV (the peak of the X-ray flux seen in Fig. 5), the corresponding synchrotron lifetime is $t_{\text{synch}} \sim 1000\Omega^{3/4}d_4^{3/2} \text{ yr}$.

The simplest expectation for a PWN is that, beyond the termination shock, the emitting particles should flow out radially in all directions until they either radiate all their energy or reach the outer edge of the PWN. However, G18.0–0.7 differs from this in two important respects. Most fundamentally, rather than forming an approximate ellipsoid with the pulsar and termination shock at its center (e.g., Slane et al. 2000, 2002), G18.0–0.7 is distinctly one-sided, showing little emission to the north of the pulsar but a broad diffuse region to the south. Also problematic is the fact that the southern region of the PWN extends to a distance $L \sim 5' = 6d_4 \text{ pc}$ from the pulsar. If this extent represents the outer edge of the expanding nebula in this direction, then the expansion speed of the nebula is $V \approx L/\tau \approx 300d_4 \text{ km s}^{-1}$ (e.g., van der Swaluw et al. 2001). Alternatively, if the true nebular extent is much larger than that seen in X-rays, then we expect a flow speed at the edge of the X-ray nebula of $V \approx (r_s/L)^2(c/3) \approx 250 \text{ km s}^{-1}$ (Kennel & Coroniti 1984). Both these calculations are at odds with the speed needed to transport particles to the southern end of the PWN within their synchrotron lifetimes, $L/t_{\text{synch}} \approx 6000\Omega^{-3/4}d_4^{-1/2} \text{ km s}^{-1}$. Thus, the size of the PWN is much larger than can be easily accounted for if the nebula has expanded unimpeded on its southern side.

One possible explanation for the PWN's properties is if the pulsar's space velocity is supersonic with respect to the sound speed of its environment. In this case the PWN is confined through ram pressure, and a bow-shock PWN results (e.g., Gaensler, Jones, & Stappers 2002b). In such PWNs, the pulsar is at one end of an elongated X-ray nebula; the PWN has a bright head at or near the pulsar position, with a progressively fainter tail extending opposite to the direction of motion (e.g., Kaspi et al. 2001a). While there is no proper motion information for PSR B1823–13 from which we might estimate its space velocity, the gross morphology of G18.0–0.7 could be interpreted as a tail extending to one side of the pulsar, suggestive of rapid motion of the central source in a northerly direction. However, there are at least two difficulties with this interpretation.

First, in other bow-shock PWNs, the shocked pulsar wind interior to the bow shock is tightly confined by the flow of surrounding material, and the X-ray tail correspondingly consists of a very narrow column of emission, marginally resolved or unresolved even with *Chandra* (Bucciantini 2002; Stappers et al. 2003). On the other hand, the diffuse emission seen here forms a very broad region behind the pulsar, showing no evidence for containment or collimation. Second, the length of the proposed tail region seen here, $L \approx 6d_4 \text{ pc}$, is too long to be a trail of emission left behind by a fast pulsar. Even if the pulsar had an extreme space velocity $V_{\text{PSR}} = 2000 \text{ km s}^{-1}$, the lifetime computed above for X-ray-emitting pairs limits the length of such a synchrotron “wake” to $V_{\text{PSR}}t_{\text{synch}} \sim 2\Omega^{3/4}d_4^{3/2} \text{ pc}$, noticeably smaller than the measured extent of G18.0–0.7. Alternatively, if the emitting particles have not been left behind by the pulsar but are part of the wind flow driven by the pulsar, one is left with the uncomfortably high-flow speed discussed above, $L/t_{\text{synch}} \approx 6000 \text{ km s}^{-1}$. Wang, Li, & Begelman (1993) argue that such high-speed backflows behind the

pulsar can be attained when the wind is forced to travel through a narrow nozzle, but no such collimated region behind the pulsar is seen here.

A second possibility to explain the morphology and extent of this PWN is that the nebula is intrinsically symmetric around the pulsar but that there is a relativistic bulk flow of particles within the nebula that results in Doppler boosting of the approaching component over receding material. Such a situation has been proposed to explain the one-sided PWN seen around PSR B1509–58 (Gaensler et al. 2002a), in which a flow velocity $\beta = v/c \gtrsim 0.2$ has been inferred. However, in PWNs for which such high-speed flows have been claimed, this outflow manifests itself in collimated jets of material directed along the pulsar spin axis (Gaensler et al. 2002a; Lu et al. 2002); the lack of any such collimated structures here argues against this possibility. The large extent of the PWN can be explained in this case, since the flow time from the pulsar to the edge of the nebula is only $t_{\text{flow}} \approx 20\beta^{-1} \text{ yr}$ (ignoring projection effects). However, at the relativistic speeds $\beta \gtrsim 0.1$ needed for Doppler boosting to be effective, we find that $t_{\text{flow}} \ll t_{\text{synch}}$. We thus would expect minimal synchrotron cooling of the outer nebula (i.e., a photon index $\Gamma \approx 1.6$ as is seen in the core), in contradiction with the steeper photon index ($\Gamma \approx 2.3$) seen for the diffuse component of G18.0–0.7.

A final alternative is that the system is in an evolved state, in which the reverse shock from a surrounding SNR has propagated back inward, colliding with the PWN (Reynolds & Chevalier 1984). If different components of this reverse shock collide with the PWN at different times, emitting particles can be rapidly convected away from the pulsar, generating a nebula that can be significantly offset to one side of the pulsar (Chevalier 1998). Such an effect has been specifically proposed to account for the one-sided morphology of the PWN powered by the Vela pulsar (Blondin, Chevalier, & Frierson 2001), and we propose that this process may similarly be occurring here for the Vela-like pulsar B1823–13. Such an asymmetric PWN will result only if the PWN crushing timescale is of the order of or is less than the lifetime of emitting pairs, and if the velocity of convective motions is sufficient to produce the observed PWN extent. Simulations of this process suggest a crushing timescale of $\sim 2000 \text{ yr}$ and convective motions at speeds of $\sim 1000 \text{ km s}^{-1}$ (van der Swaluw et al. 2001). Assuming that the PWN had a radius of a few parsecs before interacting with the reverse shock, these values are roughly consistent with the observed nebular extent and the lifetimes of emitting electrons inferred above, and we thus think an asymmetric reverse-shock interaction the most likely explanation for the observed nebular morphology. One important consideration in this interpretation is that we do not directly detect the SNR required to generate this effect; we discuss this issue in § 4.3 below.

Recent evidence has suggested that characteristic ages can overestimate (Kaspi et al. 2001b) or underestimate (Gaensler & Frail 2000) a pulsar's true age by up to an order of magnitude, depending on assumptions made about the initial period of and braking torque acting on the neutron star. The PWN/reverse-shock interaction that we have proposed here is expected to take place $\sim 10^4 \text{ yr}$ after the supernova explosion (Reynolds & Chevalier 1984), in approximate agreement with the characteristic age $\tau \equiv P/2\dot{P} = 21.4 \text{ kyr}$ measured for PSR B1823–13. Thus, within the reverse-shock scenario, the available evidence

suggests that the pulsar's characteristic age and true age are broadly consistent.

4.3. The Lack of an Associated Supernova Remnant

The interpretation that we have adopted in § 4.2 requires the presence of a surrounding SNR in order to generate the reverse shock that consequently interacts with the PWN. However, no such SNR has been detected at radio wavelengths (Braun et al. 1989), nor do we see any surrounding shell of X-ray emission from such a source in the data presented here.

The lack of detectable SNRs around pulsars with small characteristic ages has long been problematic (e.g., Reynolds & Aller 1985; Frail et al. 1995). It is thought that such SNRs must exist, since otherwise the PWN would expand freely into the ambient medium and produce a highly extended and unobservable “ghost nebula” (Blandford et al. 1973). In the specific case of Vela-like pulsars such as PSR B1823–13, Braun et al. (1989) interpret the absence of a detectable radio SNR as resulting from the rapid expansion of the SNR blast wave into a low-density cavity, followed by a collision with a dense surrounding shell evacuated by the progenitor wind. The SNR's radio and X-ray emissions then fade rapidly as the SNR's energy is dissipated through radiative shocks. Simulations of this process clearly show that the SNR/shell interaction drives a reverse shock back toward the center of the system (e.g., Tenorio-Tagle et al. 1990). Furthermore, either an offset of the progenitor from the center of its wind bubble or pressure variations within the turbulent SNR/bubble interaction region can cause different parts of the reverse shock to collide with the PWN at different times (Różyczka et al. 1993; Dwarkadas 2001), compressing the PWN in a highly asymmetric fashion, as is required here.

We can consider whether such an SNR would be detected by scaling the emission seen from the Vela SNR. Lu & Aschenbach (2000) find that X-rays from the Vela SNR can be modeled as a two-component Raymond-Smith plasma, with temperatures for the two components of $kT_1 \approx 0.15$ keV and $kT_2 \approx 1$ keV. The volume-integrated emission measure for the cooler component is ~ 10 times that of the hotter component, and the total unabsorbed luminosity (0.1–2.5 keV) is 2.2×10^{35} ergs s^{-1} for a distance of 250 pc. Scaling to a distance of 4 kpc, the expected unabsorbed 0.1–2.5 keV flux densities for the two components are $f_1 = 1.0 \times 10^{-10}$ ergs $cm^{-2} s^{-1}$ and $f_2 = 9.5 \times 10^{-12}$ ergs $cm^{-2} s^{-1}$.

Using the W3PIMMS tool,¹¹ we can determine that the expected EPIC MOS count rate from such a source is 0.18–0.29 counts s^{-1} for absorbing columns in the range $(1.0\text{--}1.4) \times 10^{22}$ cm^{-2} . Again scaling from the Vela SNR, we expect this emission to be distributed around a shell of diameter $\sim 20'$. Assuming a shell thickness 20% of its radius, the expected surface brightness for EPIC MOS would be $(4\text{--}7) \times 10^{-7}$ counts $s^{-1} arcsec^{-2}$, several times below the observed background level $10'$ off-axis of $\sim 3 \times 10^{-6}$ counts $s^{-1} arcsec^{-2}$ (after correction for vignetting). We can make similar calculations for archival observations of this source with *ROSAT* and *ASCA*: for the *ROSAT* PSPC we expect an SNR surface brightness of $(5\text{--}9) \times 10^{-8}$ counts $s^{-1} arcsec^{-2}$ against an observed background in these data of $\sim 2 \times 10^{-7}$ counts $s^{-1} arcsec^{-2}$, while for the *ASCA* Gas

Imaging Spectrometer we expect an SNR surface brightness of $(4\text{--}6) \times 10^{-8}$ counts $s^{-1} arcsec^{-2}$ against the observed background of $\sim 2 \times 10^{-7}$ counts $s^{-1} arcsec^{-2}$. Thus, we expect that any emission from a Vela-like SNR associated with PSR B1823–13 would not be detectable in any of these observations. This conclusion is strengthened by the fact that the characteristic age of PSR B1823–13 is approximately twice that of the Vela pulsar; if this indicates a similar difference in these pulsars' true ages, any associated SNR would most likely be even fainter than considered here.

For the radio observations carried out in this direction by Braun et al. (1989), we estimate a 1 GHz surface brightness sensitivity for an SNR of this extent of $\Sigma \gtrsim 1.0 \times 10^{-21}$ W $m^{-2} Hz^{-1} sr^{-1}$ (3σ), where we have assumed a typical radio photon index $\Gamma = 1.5$ and have accounted for primary beam attenuation. There are certainly many SNRs fainter than this, including relatively young SNRs such as G266.2–1.2 (Duncan & Green 2000). Thus, the failure to detect an SNR associated with PSR B1823–13 is not inconsistent with the possibility that a rapidly expanding, rapidly fading SNR has produced the pressure that has confined and distorted the PWN.

5. CONCLUSIONS

We have carried out a deep X-ray observation of the Vela-like pulsar B1823–13 with *XMM-Newton*. We find that all the detected X-ray emission results from the interaction between the relativistic pulsar wind and its environment: a bright elongated core immediately surrounding the pulsar most likely represents the wind termination shock, while a surrounding faint region with a softer spectrum corresponds to a pulsar wind nebula with a mean magnetic field of ~ 10 μ G. Emission from the pulsar itself, either in the form of a central point source or as pulsations at the 101 ms pulsar period, is not detected.

The nebula, which we designate G18.0–0.7, is distinctly one-sided: most of its emission is distributed in a region of extent $5'$ restricted to the southern side of the pulsar. We interpret this morphology as resulting from the compression and distortion of the pulsar wind by the asymmetric reverse shock produced by an unseen surrounding supernova remnant. Deeper observations of PSR B1823–13 and other young pulsars at meter wavelengths, carried out with the Very Large Array and the Low Frequency Array, are probably the most promising way to directly detect these associated SNRs. The reverse-shock interaction between a pulsar nebula and its surrounding SNR is expected to occur 10–20 kyr after the supernova explosion, suggesting that many Vela-like pulsars might be in this phase of evolution. Future observations of other such pulsars and their nebulae with *XMM-Newton* can explore this possibility.

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¹¹ See <http://asc.harvard.edu/toolkit/pimms.jsp>.

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