Characterization of Meta-Materials Using Computational Electromagnetic Methods

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Abstract

An efficient and powerful computational method is presented to synthesize a meta-material to specified electromagnetic properties. Using the periodicity of meta-materials, the Finite Element Methodology (FEM) is developed to estimate the reflection and transmission through the meta-material structure for a normal plane wave incidence. For efficient computations of the reflection and transmission over a wide band frequency range through a meta-material a Finite Difference Time Domain (FDTD) approach is also developed. Using the Nicholson-Ross method and the Genetic Algorithms, a robust procedure to extract electromagnetic properties of meta-material from the knowledge of its reflection and transmission coefficients is described. Few numerical examples are also presented to validate the present approach.

Introduction

In the recent years there has been a great deal interest in studying the material (called meta-material [1-3]) having simultaneously negative permittivity and permeability. These meta-materials do not exist in nature and are artificially formed by embedding properly designed metallic/non-metallic inclusions in binding mediums. In such a material, the wave vector $\hat{k}$ and the Poynting vector $\hat{P}$ are anti-parallel causing reversal of some basic electromagnetic wave propagation characteristics such as Snell’s law and Doppler effects. The reversal of these basic propagation properties has lead to some potential applications of these materials such as high resolution lenses, sub wavelength resonators and enhancement of antenna gains. One of the challenges faced by the researchers in this field is lack of a robust method to characterize these materials. The characterization of the meta-materials is carried out in two steps. First, the reflection and transmission properties of these materials are estimated/measured using a computational/experimental method. In the second step, a proper inversion process is used to extract effective electromagnetic properties from the knowledge of reflection and transmission coefficients. Experimental procedures are too expensive and time consuming to explore all design space, and are is usually not preferred in early stages of design cycle. In the present paper, an efficient and powerful computational method is presented to characterize/estimate effective electromagnetic properties of meta-materials. In this method, first, the metallic inclusions of interest are modeled using a commercially available Computer Aided Design (CAD) package. Using the Finite Difference Time Domain method the reflection and transmission through these inclusions are then estimated. In the final step using an inversion procedure based on two approaches, Nicolson-Ross method [4] and Genetic Algorithm, effective properties of the selected inclusions are extracted. Few examples with known effective properties are presented to validate the present method.

Analysis of Meta-Materials

In this section, the development of computational methods based on the Finite Element Methodology and also Finite Difference Time Domain (FDTD) method are presented. Figure 1(a) shows a metallic inclusion in a split ring form and Figure 1(b) shows a dielectric medium embedded with a periodic arrangement of split ring shown in Figure 1(a). It is assumed that the structure is infinite in the x-y plane and has finite width in the z-direction. It also assumed that the structure is illuminated by a plane wave with normal incidence as shown in Figure 1(b).
**Finite Element Method:** To estimate the reflection and transmission coefficients of such a structure, a unit periodic cell as shown in Figure 1(b) is considered. Using the electric and magnetic symmetries, the unit cell can be terminated by electric and magnetic conducting boundaries at $y = \pm a/2$ and $x = \pm b/2$, respectively. The electromagnetic field incident at the input port of the unit cell can be written as

$$\hat{E}_{\text{in}} = x E_0 e^{-j k_0 z}, \quad \hat{H}_{\text{in}} = -x E_0/\eta_0 e^{-j k_0 z}$$

(1)

The electromagnetic field scattered in the regions $z \leq 0$ and $z \geq d$ can be expressed in terms of the Floquet modes. Using the FEM formulation, the weak form of wave equation to be solved for the electric field in the unit cell takes following form [5]

$$\int_V \left( \nabla \times \hat{\tau} \cdot \left( \frac{1}{\mu_r} \nabla \times \hat{E}^{\text{II}} \right) - \left( \frac{k_0^2 \varepsilon_r}{k_0} \hat{E}^{\text{II}} \right) \right) dv = 2j \omega \mu_0 \gamma_0 \int_{S_1} \hat{\tau} \cdot \hat{E}_{\text{in}} \; ds \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \quad \ Quad
results confirms validity of the FEM formulation and the also the modeling procedure followed while using the CST Microwave Studio.

Using these techniques the reflection and transmission characteristics of metallic inclusions of various shapes and sizes are computed and results will be shared with the audience at the time of presentation.

**Extraction of Electrical Properties:**

Following the Nicholson-Ross [4] method to estimate electric and magnetic properties of meta-materials, it is assumed that the dielectric slab with the metallic inclusions is replaced by a homogeneous material slab having effective permittivity \( \varepsilon_r^{\text{eff}} \) and permeability \( \mu_r^{\text{eff}} \). The reflection and transmission coefficients of this effective slab can be written as

\[
S_{11} = \frac{r}{1 - r^2}, \quad S_{21} = \frac{(1 - r^2)t}{1 - r^2 t^2}
\]

where \( r = e^{-jk_0d\sqrt{\varepsilon_r^{\text{eff}}}} \) and \( t = \left( \sqrt{\frac{\mu_r^{\text{eff}}}{\varepsilon_r^{\text{eff}}} - 1} \right) / \left( \sqrt{\frac{\mu_r^{\text{eff}}}{\varepsilon_r^{\text{eff}}} + 1} \right) \). The problem of extraction of electric properties of an equivalent slab can now be defined as determination of \( \varepsilon_r^{\text{eff}}, \mu_r^{\text{eff}} \) using (3) from the knowledge of \( S_{11}, S_{21}, d \). Defining new variables \( V_1 = S_{11} + S_{21}, V_2 = S_{21} - S_{11} \) and \( X = (1 - V_1 V_2)/(V_1 - V_2) \), it can be shown that

\[
r = X \pm \sqrt{X^2 - 1} \quad \text{and} \quad t = (V_1 - r)/(1 - V_1 r)
\]

The steps involved in extraction process are:

(a) determine \((r, t)\) using (4).

(b) express \( \sqrt{\mu_r^{\text{eff}}/\varepsilon_r^{\text{eff}}} = k_r = \frac{1}{k_0}\Phi_t-2\pi n + j\ln(|r|) \), where \( n = ..., -3, -2, -1, 0, 1, 2, 3, ... \)

(c) with proper choice of \( n \), parameters are given by \( \mu_r^{\text{eff}}/(1 + r), \varepsilon_r^{\text{eff}}/(1 + r)k_r \).

Due to ambiguity in selecting proper value of \( n \), earlier researchers [7] assuming \( d \ll \lambda \) had simplified the calculations in the step (b). However, the assumption \( d \ll \lambda \) limits application to very thin slab. In this work we have develop...
oped a procedure to track correct value of n without making the assumption. The proper values of effective electric properties are also estimated using the Genetic Algorithm.

**Numerical Results**

In this section numerical data for the effective parameters of a slab embedded with split rings arranged in a periodic fashion, shown in Figure 3(a), are presented. First, the unit cell including a single layer of split ring is modeled using a CAD package. The transmission and reflection coefficients as a function of frequency are then computed using the CST Microwave Studio software. Using the extraction procedure described above the effective parameters are estimated and are shown in Figure 3(b). Using the extracted parameters the reflection and transmission coefficients are calculated and compared with the numerical data obtained using the CST Microwave Studio.

**Conclusions**

A computational electromagnetic method using the FEM and FDTD approaches have been presented to estimate the reflection and transmission through a meta-material slab. Nicholson-Ross technique[4] and also Genetic Algorithm procedure in conjunction with the knowledge of reflection and transmission coefficients have been used to extract effective permittivity and permeability of a meta-material slab. The reflection and transmission coefficient of actual meta-material slab have compared well with the numerical data obtained using the extracted effective permittivity and permeability. Due to page limitation the numerical results for a metallic inclusion in the shape of split ring only has been shown here. The metallic inclusions of other shapes will be discussed at the time of presentation.

**References**