Title: Development of a Fabry-Perot Interferometer for Ultra-Precise Measurements of Column CO₂

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Abstract

A passive Fabry-Perot based instrument is described for detecting column CO$_2$ through absorption measurements at 1.58$\mu$m. In this design, solar flux reaches the instrument platform and is directed through two channels. In the first channel, transmittance fringes from a Fabry-Perot interferometer are aligned with CO$_2$ absorption lines so that absorption due to CO$_2$ is primarily detected. The second channel encompasses the same frequency region as the first, but is comparatively more sensitive to changes in the solar flux than absorption due to CO$_2$. The ratio of these channels is sensitive to changes in the total CO$_2$ column, but not to changes in solar flux. This inexpensive instrument will offer high precision measurements (error $<$1%) in a compact package. Design of this instrument and preliminary ground-based measurements of column CO$_2$ are presented here as well as strategies for deployment on aircraft and satellite platforms.
1. Introduction

Prior to the industrial revolution of the last 250 years, the concentration of carbon dioxide (CO₂) in the atmosphere was 280 ppm, a value that had remained approximately constant (within ± 10 ppm) for several thousand years (Houghton et al., 2001). When Charles Keeling (Keeling, 1997: Keeling et al., 2004) started measuring atmospheric carbon dioxide from Mauna Loa observatory in 1958, the concentration was on the order 315 ppm (Keeling, 2005). This value has steadily increased to its current value of 380 ppm (Keeling, 2005). This recent increase, the rate of which has been unparalleled for the last 20,000 years, is largely due to the anthropogenic emissions of the modern industrial revolution (Houghton et al., 2001).

The majority of CO₂ variability occurs in the lower atmosphere (~1000 to 800 mbar). Ground-based absorption measurements of the CO₂ column are commonly carried out by Fourier Transform Infrared spectrometry (FT-IR) coupled with a sun tracking system. Typically, these exploited the broad spectral range capable by FT-IR to also include species such as O₂, O₃, CH₄, HCl, N₂, N₂O, HNO₃, HF, and H₂O. Analysis generally involves using an algorithm to fit individual spectral features to reduce noise. Yang, et al. (Yang et al., 2002) measured CO₂ at 6300 cm⁻¹ and O₂ at 7900 cm⁻¹ from the McMath telescope in Kitt Peak, Arizona. The ratio of CO₂ to O₂ is then used to scale to the mean O₂ volume mixing ratio. Precision is estimated at 0.5% after a HITRAN (Rothman, 1998) based spectral fitting routine (JPL) is used to reduce the noise. Goldman, et al. (Goldman et al., 1999) compared several FT-IR based instruments (JPL’s MKIV, NCAR’s Bruker 120 M, and NPL’s Bruker 120 M) measuring a variety of trace species including CO₂ from the Jet Propulsion Laboratory’s Table Mountain Facility in
Wrightwood, CA and found standard deviations from the mean column to range from 0.5 to 2% for CO2.

While these FT-IR techniques offer good sensitivities, cost and size limit the viability for this method fill in the sparse spatial network of surface measurements and ultimately resolve major source and sink regions. Conversely, while space-based measurements of this greenhouse gas offer good spatial coverage as well as temporal resolution, historically, they have not offered adequate precision for resolving key sources and sinks. (O'Brien and Rayner, 2002; Sarmiento and Gruber, 2002) Transport inversion model experiments indicate global column measurements with an uncertainty of 2.5 ppm or better (<1%) would greatly improve surface flux estimates (Pak and Prather, 2001; Rayner and O'Brien, 2001)

In this article, we present an application of a Fabry-Perot interferometer for measurements of spectral absorption of sunlight by CO2. Since its inception in 1899, Fabry-Perot interferometers have found wide applications as background filters in atmospheric research. Mack, et al. (Mack et al., 1963) described the PEPSIOSTM (polyetalon pressure-scanned interferometric optical spectrometer) that used several Fabry-Perot etalons to isolate a spectral bandpass. With each etalon in its own pressure chamber, scanning the etalon was accomplished by varying the pressure of the inert gas between the etalon plates, and thus varying the refractive index of the inert gas. Larar et al. (Larar et al., 1998) proposed a double-etalon Fabry-Perot interferometer as a nadir-viewing satellite instrument for global tropospheric and total ozone monitoring at 9.6 μm. In the accompanying feasibility study, Larar and Drayson (Larar and Drayson, 1998) predict precisions of 12% and 3.6% respectively. In addition, a retrieval method using a
nonlinear least-squares iterative technique to achieve convergence of the measured and calculated forward model signals is described. Skinner, et al. (Skinner et al., 1987) investigated a triple etalon Fabry-Perot interferometer for measuring Doppler shifts of rotational lines of O₂ in the atmosphere from the Upper Atmosphere Research Satellite as a means for determining wind vectors in the stratosphere and mesosphere. Finally, McKay (McKay, 1999) compares the application of single, double and triple Fabry-Perot etalon pairs as solar background filters in lidar. Tunable air-gap etalons are aligned to lidar signal spectrum to reduce signal interference caused by the scattering of sunlight into the field of view. Reflective coupling between etalons was found to limit potential contributions of second and third etalons.

The instrument presented here offers a smaller, less expensive, and mechanically less-complex alternative to FT-IR techniques with equal sensitivity. Additionally, this technique can be packaged to shoe-box size for use in either ground-based applications or deployment on an aircraft or satellite platform. We have assembled a laboratory prototype system and have conducted validation experiments in the lab as well as measured CO₂ in the atmosphere in ground measurements using both direct sunlight, and sunlight reflected off of a Lambertian surface. Preliminary results indicate that the method will be able to achieve the sensitivity and signal-to-noise detection required to measure column CO₂ at the target specification. In addition, we have begun development of a similar instrument for measuring O₂ to correct for temperature fluctuations as well as changes in surface pressure. Methods for increasing sensitivity and stability of the instrument will also be addressed.
2. Instrument Design

The instrument described here employs a Fabry-Perot Interferometer to measure the carbon dioxide column through absorption of sunlight at 1.58μm. In this design, sunlight passes through the atmosphere where it undergoes some absorption by atmospheric carbon dioxide. The light is then introduced into the instrument and split into two channels. The first channel simply measures the total light intensity. This channel responds to changes in overall intensity arising from changes in atmospheric transmission, changes in atmospheric CO$_2$, and when viewing surface reflected light, changes in ground albedo. The second channel is identical to the first except that it includes a solid etalon Fabry-Perot interferometer tuned to pass light at the wavelengths that CO$_2$ absorbs and reflect light at other wavelengths. This channel responds to albedo and atmospheric transmission changes just as the first channel but proportionately much more strongly to changes caused by changing CO$_2$. The ratio of the signals of the two channels can be used to infer the atmospheric CO$_2$ abundance. A schematic of this instrument is shown in Figure 1.

In this design, an off-axis parabolic mirror (OAP) focuses incoming light onto a 2 mm diameter aperture to restrict the field of view of the instrument. Following the aperture, the light is modulated at 400 Hz with a chopper (Stanford Research Systems, SR540). Chopped light is then re-collimated with a second OAP and passes through a 1567-1574 nm bandpass filter (Barr Associates). Approximately 10% of of this light is split with a (fused Si) beamsplitter (CVI), and focused onto the first of two thermoelectric cooled InGaAs photodiode detectors (Oriel Instruments). This first (reference) channel monitors changes in solar flux.
Light passing through the beam splitter continues through a solid Fabry-Perot interferometer (CVI), and is focused onto the second InGaAs detector forming the Fabry-Perot channel. Both the prefilter and Fabry-Perot interferometer optics are mounted in ovens (Andover) and temperature controlled to 1/100th degree. Both Fabry-Perot and Reference signals are processed through lock-in amplifiers (Stanford Research Systems, SR830 DSP) referenced to the frequency of the chopped signals. Custom LabVIEW software (National Instruments, version 6.1) has been developed for controlling the measurement system as well as analyzing the detected signals from both channels. The ratio of these signals can then be related to changes in the CO₂ column.

A. Fabry-Perot Characteristics

The Fabry-Perot in this instrument is solid fused silica with a free spectral range (FSR) of 0.306 nm, a refractive index of 1.443 at λ=1571 nm, and a clear aperture of 50 mm. The interference produced by a Fabry-Perot Interferometer (Figure 2) can be described by the Airy function (Born and Wolf, 1999: Ingle and Crouch, 1988):

\[
\frac{\Phi}{\Phi_m} = \frac{1}{1 + C_F \sin^2 \left( \frac{\delta}{2} \right)}
\]

where \(\Phi/\Phi_m\) is the fractional radiant power, \(C_F\) is the coefficient of finesse, and \(\delta\) is the phase difference between the interfering waves. The finesse (F) of the etalon is related to the reflectance (R) of its surfaces as shown in the following equation. Also detailed are the free spectral range (FSR) and full width at half maximum (FWHM) of the interference fringes.
The free spectral range is the distance between these interference fringes, determined by the wavelength ($\lambda$) of the interfering light, the thickness of the etalon ($d$) and the refractive index ($n$) of the etalon material where the FSR, $\lambda$, and $d$ are in units of nm:

$$F = \frac{FSR}{FWHM} = \frac{\pi \sqrt{R}}{1 - R} = \frac{\pi \sqrt{C_F}}{2}$$

(2)

Thus, the spacing between fringes can be approximately matched to the spacing of CO$_2$ lines through etalon thickness. Similarly, the width of fringes can be approximately matched to CO$_2$ line width by appropriate choice of reflectance of the etalon surfaces.

B. Etalon Temperature Tuning

Alignment between fringes and CO$_2$ absorption transitions can be further adjusted by temperature tuning the Fabry-Perot. As the temperature changes, the refractive index of a Fabry-Perot optical medium changes as:

$$\mu = \mu_0 + \beta(T - T_0)$$

(4)

where $\mu_0$ is the refractive index at temperature $T_0$, and $\beta$ is the refractive index change coefficient. With increasing temperature, etalon fringes shift to longer wavelengths by typically 0.02 nm per °C for fused silica.(Slyusarev, 1984)

Alignment between fringes and CO$_2$ lines at two different temperatures is demonstrated in Figure 3(a), where two laser scans of carbon dioxide (1 atm) in an absorption cell are obtained with the instrument etalon at temperatures of 42° and 54°C. Fringes from the Fabry-Perot channel shown in light and dark gray respectively are
compared with CO₂ lines from the reference channel in black. Fringes and rotational features are shown with the bandpass prefilter defining the wavelength range of the instrument. The 54°C scan clearly shows better overlap between fringes and lines than the 42°C scan. To validate in the laboratory that improved overlap corresponds with greater instrument sensitivity, changes in the ratio of Fabry-Perot to Reference signals were monitored for CO₂ pressures ranging from 0 to 1000 torr in a 1.4 m absorption cell at both Fabry-Perot temperatures. Laboratory measurements using a 1.4 meter absorption cell at 1000 Torr are comparable to a column absorption of approximately ¾ that of the true atmospheric column. Results for this experiment are shown in Figure 3(b) where the ratio of Fabry-Perot/Reference signals decrease as the pressure in the gas cell increase. While some change in the ratio is seen with poor overlap (42°C), at Fabry-Perot temperatures where there is good overlap (54°C), the ratio is considerably more sensitive to changes in CO₂ pressure.

3. Ground Testing

The Fabry-Perot instrument has two potential modes of operation; one as a ground based sensor, and one as an aircraft or satellite deployable instrument. In the first configuration, the CO₂ column is measured through absorption of light by CO₂ in the atmosphere directly between the sun and the ground instrument. We accomplished this by collecting light with a small telescope fixed to an equatorial mount, aligned to track the sun throughout the day. An optical fiber coupled at the rear of the collimator brings light into the instrument. In the second configuration, light passing through the atmosphere reflects on the Earth’s surface before entering the instrument platform. This is the way the
instrument would be used for remote sensing of the CO$_2$ column from an airborne or spaceborne platform. We have tested the instrument’s performance in this mode simply by introducing light reflected off the ground. The flux entering the instrument is independent of its height above the ground because the area included in its field of view increases at the same rate as its overall collection efficiency decreases. While this approach should be more sensitive to changes in CO$_2$ because the light passes through the atmosphere twice before entering the instrument, the retrieval of the actual CO$_2$ column from the measurement is complicated by scattering effects in the atmosphere.

A. Sun tracker results

In general, the ratio arising from direct illumination should depend upon the path length through the atmosphere (airmass) regardless of changes in cloud and aerosol distribution.

The ratio then should follow Beer’s law in the form:

$$ R_{calc} = R_o + I_o e^{-k(Airmass)} $$

(5)

for modest values of airmass, where airmass is estimated by

$$ Airmass = \sec \left( \frac{2}{\pi} - ALT \right) $$

(6)

and ALT is the solar altitude in radians. (U.S. Naval Observatory, 2003) It should however be noted that this formula assumes that there is no saturation by the absorber.

Ground measurement of the CO$_2$ column made from Edwards, CA on May 13, 2004 are shown in Figure 4. The ratio of Fabry-Perot to Reference signals (filled circles) clearly follows changes in the airmass throughout the day.

To validate that changes in the ratio signal were indeed due to changes in the CO$_2$ column, outdoor experiments were conducted in which a 1.4 m absorption cell was
aligned at the front of the instrument, and CO$_2$ pressures in the cell were varied from 0 to 1000 torr. Experiments were conducted at midday when changes in the airmass were minimal. The results of these trials are shown in Figure 5, where the ratio of Fabry-Perot to reference signals are shown as a function of total CO$_2$ measured in the cell and column. In this Figure, it was assumed that 1 torr of CO$_2$ in the cell is equivalent to 0.263 ppm of CO$_2$ in the column. This experiment also serves to calibrate the observed ratio to the total CO$_2$ measured. However, this is only useful in the linear region of this figure extending to approximately 600 ppm. Beyond this, the affects of self-broadening of CO$_2$ in the absorption cell become evident.

B. Reflected flux measurements

Preliminary tests of the instrument in the reflected light viewing mode were also carried out on the ground by having the instrument view a Lambertian surface (Spectralon) at ground level with a downward viewing mirror approximately a meter above the surface. Radiation incident on a surface is partially direct and partially diffuse. The direct component of this radiation is a result of the direct solar beam passing through the atmosphere while the diffuse component is a result of atmospheric scattering. For a clear sky condition, the direct component is about 80% of total incident solar radiation in visible and infrared regions. (Qin et al., 2001) We estimate the direct component of sunlight by measuring total intensity and then blocking the sunlight over the region observed by the instrument to infer the indirect illumination.

Figure 6 shows ground data taken in this method in July 2005. Data for both days are shown with 15 second averaging. The general form of the ratio signal reflects
changes in the airmass throughout the day. Noisier regions show increased cloud cover in the morning of July 20, and after 3:00 p.m. (15) on July 21, as well as low signal levels in morning and evening. The signal spikes (seen at 30 minute intervals) show the change in ratio when the direct light is blocked and only indirect (scattered) light reaches the instrument. The direction that spikes point depends on whether the scattered light travels a longer path (decreases ratio) or shorter path (increases ratio) through the column than the direct light.

An illustration of this is shown in Figure 7. For the high sun example at left, the scattered light has a longer absorption pathlength through the column than the direct light. With low sun (at right), the absorption pathlength of the scattered light is shorter than that of the direct light because there is very little CO$_2$ along the high altitude path through the atmosphere before the scatter occurs. While blocking the direct light results in an increase in the ratio for data taken on this date, data taken on other dates have shown the ratio to decrease with blocking when the sun is low in the sky. This effect is complicated by varying cloud heights, aerosols, altitude variability of the CO$_2$, etc. and poses a significant retrieval challenge.

C. Sensitivity and Noise Estimate

In both sun tracker and reflected flux measurements, changes in the CO$_2$ ratio were found to track changes in the airmass throughout the day. This temporal correlation was exploited to produce the calibration curve shown in Figure 8(a) where changes in the CO$_2$ ratio are plotted as a function of airmass. For this plot, ground data acquired with the sun tracker were used (See Figure 4). It can be seen that the ratio changes by ~0.10
for an approximate airmass change of 4 (about 0.025 for an airmass change of 1). One airmass is approximately equal to 360 ppm of CO$_2$. Thus, the sensitivity of the ratio to a 1 ppm change in CO$_2$ is $0.025/360 = 0.000069$.

An expansion of the noise in this calibration curve taken at 1/10 second intervals is shown in Figure 8(b), with the ten point average shown as a solid line. The standard deviation in noise for data taken at 1/10 second intervals is $\pm 0.00025$. In one second, the mean error is equal to $0.00025/\sqrt{10} = 0.000079$. If the detection limit is defined as a 2:1 SNR, twice the noise is $0.000079 \times 2 = 0.00016$, then the current instrument detection limit is $0.00016 \times 1 \text{ ppm}/0.000069 = 2.3$ ppm in a one second average. This implies the ground based instrument performance would be better than the desired sensitivity of 1 ppm in less than ten seconds of averaging.

4. Conclusions and Future Directions

Presented here is a prototype design for a passive Fabry-Perot based instrument to make absorption measurements of total column carbon dioxide using sunlight at 1.58$\mu$m. This instrument has been demonstrated as a ground based sensor when used in conjunction with a positioning system to follow sun movement. Signals from these data trials were found to consistently track changes in airmass. Additionally, laboratory experiments validate the instrument’s sensitivity to changes in CO$_2$. An estimate of the system performance indicates that this design has the capability to resolve changes in the CO$_2$ column as small as 2.3 ppm with a one second average and better than 1 ppm in less than 10 seconds averaging.
Preliminary ground data have also been presented using surface reflectance to measure the CO₂ column. While this second configuration is a clear lead to deployment of the Fabry-Perot instrument on an airborne platform, it also presents some additional retrieval challenges due to contributions of aerosol scattering in the reflected signal. While contributions from aerosol scattering can be measured and easily corrected during ground measurements using surface reflectance, deconvolution of scattering in airborne measurements poses a far greater challenge. This clearly will also be an issue for other passive instruments measuring absorption of the CO₂ column as well as other trace gases.
Acknowledgments

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Figure Captions

Figure 1. Optical and electronic schematic for Fabry-Perot instrument. Incoming light is divided into Fabry-Perot (detector 2), and reference (detector 1) channels, measuring changes in CO₂ absorption and solar flux, respectively. The ratio of these signals can then be related to changes in the CO₂ column.

Figure 2. Interference in a Fabry-Perot Interferometer. The full width at half maximum (FWHM) and free spectral range (FSR) are indicated by arrows.

Figure 3. In (a), fringe – CO₂ line overlap is compared at two Fabry-Perot temperatures. In (b), the change in ratio as a function of CO₂ pressure in an absorption cell is compared for good (54°C) and poor (42°C) fringe-line overlap.

Figure 4. Ground measurement of the CO₂ column on May 13, 2004. In this configuration, light enters the instrument via a fiber-coupled collimator on an equatorial mount, which tracks sun throughout the day. The ratio (Fabry-Perot/Reference) signal (filled circles) clearly follows changes in the airmass (dashed line).

Figure 5. Field testing instrument validation at mid-day. CO₂ pressure in a 1.4 m absorption cell placed in front of the instrument ranged from 0 to 1000 torr. The ratio of Fabry-Perot to Reference signals is plotted as a function of total CO₂ observed.

Figure 6. Ground data taken with sun light reflected on a Lambertian surface. Direct light hitting this surface is blocked every 30 minutes to illustrate the contribution of scattered light to the signal.

Figure 7. Contribution of scattered light to the ratio. With the sun in a higher position (left), the optical pathlength of the scattered light may be longer than that for the direct
sunlight’s path. If the sun is lower in the sky (right), the scattered light may travel a shorter path through the column.

**Figure 8.** In (a), the CO₂ ratio from ground measurements on May 13, 2004 is calibrated to the airmass. Enlargements of noise for data taken at 1/10 second is shown in (b) with a ten point average shown as a solid line.
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