In situ detection of strong Langmuir turbulence processes in solar type III radio bursts

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Received 7 March 2012; revised 29 May 2012; accepted 1 July 2012; published 29 August 2012.

[1] The high time resolution observations obtained by the WAVES experiment of the STEREO spacecraft in solar type III radio bursts show that Langmuir waves often occur as intense localized wave packets. These wave packets are characterized by short durations of only a few ms and peak intensities, which well exceed the supersonic modulational instability (MI) thresholds. These timescales and peak intensities satisfy the criterion of the solitons collapsed to spatial scales of a few hundred Debye lengths. The spectra of these wave packets consist of primary spectral peaks corresponding to beam-resonant Langmuir waves, two or more sidebands corresponding to down-shifted and up-shifted daughter Langmuir waves, and low frequency enhancements below a few hundred Hz corresponding to daughter ion sound waves. The frequencies and wave numbers of these spectral components satisfy the resonance conditions of the modulational instability (MI).

Moreover, the tricoherences, computed using trispectral analysis techniques show that these spectral components are coupled to each other with a high degree of coherency as expected of the MI type of four wave interactions. The high intensities, short scale lengths, sideband spectral structures and low frequency spectral enhancements and, high levels of tricoherences amongst the spectral components of these wave packets provide unambiguous evidence for the supersonic MI and related strong turbulence processes in type III radio bursts. The implication of these observations include: (1) the MI and related strong turbulence processes often occur in type III source regions, (2) the strong turbulence processes probably play very important roles in beam stabilization as well as conversion of Langmuir waves into escaping radiation at the fundamental and second harmonic of the electron plasma frequency, f\textsubscript{pe}, and (3) the Langmuir collapse probably follows the route of MI in type III radio bursts.


1. Introduction

[2] This paper describes the strong Langmuir turbulence processes, namely, the supersonic modulational instability (MI) and collapsing Langmuir solitons observed in the source regions of solar type III radio bursts. These observations are from the WAVES experiment of the STEREO spacecraft [Bougeret et al., 2008], which is designed to detect in situ waves as well as remote radio emissions over a wide frequency range. Ginzburg and Zheleznyakov [1958] were the first to suggest that (1) the electron beams accelerated during solar flares generate Langmuir waves at the local electron plasma frequency, f\textsubscript{pe}[Hz] = 9n\textsubscript{e}/C\textsubscript{0} by the bump-on-tail instability [Bohm and Gross, 1949] (n\textsubscript{e} is the electron density in m\textsuperscript{-3}), and (2) these Langmuir waves are subsequently converted by some nonlinear processes into the observed type III radio emissions at the fundamental and second harmonic of the local electron plasma frequency, f\textsubscript{pe}. Although this hypothesis has been confirmed by the in situ detection of type III burst associated electron beams [Lin, 1970; Lin et al., 1973, 1981; Ergun et al., 1998], and Langmuir waves [Gurnett and Anderson, 1976, 1977; Lin et al., 1986; Gurnett et al., 1993; Thejappa et al., 1993a, 1993b, 1999; Hospodarsky and Gurnett, 1995; Thejappa and MacDowall, 1998; Henri et al., 2009], several issues remain unresolved.

[3] Sturrock [1964] was the first to pose the following question: How does the electron beam preserve the bump-on-tail distribution over 1 AU and beyond against the “quasi-linear relaxation” or “plateauing”, that is known to deplete all the streaming energy within 100 km or less? Here we should note that Sturrock considered a special case of an infinitely extended homogeneous plasma, and neglected the effects of re-acceleration of beam particles by Langmuir waves. Several authors [Zaitsev et al., 1972, 1974; Takakura
and Shibahashi, 1976; Magelssen and Smith, 1977; Escande and de Genouillac, 1978] have treated the wave particle interactions self-consistently all along the path of the inhomogeneous beam, and claimed that the continuous emission and reabsorption of waves enable the beam to survive over the distances of order of ~1 AU and more. However, the Langmuir wave intensities predicted by the inhomogeneous quasi-linear theories are also subject to wave-wave interactions.

[4] In the context of homogeneous beam-plasma interaction, Kaplan and Tsytovich [1968] suggested that the induced scattering of Langmuir waves off ion clouds, which acts as the electrostatic decay of Langmuir waves into backscattered Langmuir waves and forward-going ion sound waves when \( T_e > T_i \) (\( T_e \) and \( T_i \) are the electron and ion temperatures, respectively) can stabilize the beam; these authors considered only the mono energetic beam and isotropic distribution of waves. Although some signatures of the electrostatic decay are detected in type III source regions [Lin et al., 1986; Gurnett et al., 1993; Thejappa et al., 1995, 2003; Thejappa and MacDowall, 1998; Hospodarsky and Gurnett, 1995; Henri et al., 2009], if the assumption of mono energetic beam and isotropic distribution of waves is relaxed, the net scattering of waves out of resonance with the beam by the induced scattering becomes too slow to prevent quasi-linear relaxation [Zheleznyakov and Zaitsev, 1970; Heyvaerts and de Genouillac, 1974].

[5] The supersonic MI, which is popularly known as the oscillating two stream instability (OTSI) [Papadopoulos et al., 1974; Smith et al., 1979; Goldstein et al., 1979] and, soliton formation and Langmuir collapse [Zakharov, 1972; Nicholson et al., 1978; Goldman, 1983] are suggested as the most effective beam stabilization mechanisms. These strong Langmuir turbulence processes are the focus of several theoretical studies [Zakharov, 1972; Goldman, 1984; Robinson, 1997], computer simulations [Bardwell and Goldman, 1976; Nicholson et al., 1978; Degtyarev et al., 1980; Doolen et al., 1985; Russell et al., 1988] and experimental investigations [Whelan and Stenzel, 1981; Michel et al., 1982; Wong and Cheung, 1984; Cheung and Wong, 1985a, 1985b; McFarland and Wong, 1997, 2001; Vyacheslavov et al., 2002] because of the fundamental nature of the problem [Zakharov, 1972; Krasnoselskikh and Sotnikov, 1977], and practical importance in fields, such as the laser heating of pellets [Montgomery et al., 1999], ionospheric modification experiments [DuBois et al., 1993], and several astrophysical applications [Galeev and Krasnoselskikh, 1976, 1978; Gurnett et al., 1981; Pelletier et al., 1988; Baltikhin et al., 1989; Kellogg et al., 1999].

[6] As shown in Figure 1, MI excites a low frequency ion density perturbation of frequency \( \Omega \), the wave number \( q \), which can be either a quasi-mode or a freely propagating ion sound wave. This ion sound wave can beat with two of the beam-resonant Langmuir waves with frequencies and wave numbers \( f_L, k_L \) and produce high frequency sidebands in the non resonant region with frequencies and wave numbers \( f_L - \Omega, k_L - q \) (Stokes mode) and \( f_L + \Omega, k_L + q \) (anti-Stokes mode). The nonlinear effect responsible for MI is the density cavity dug by the action of the ponder motive force of the Langmuir wave electric field, which in turn traps the waves thus deepening it still further. Usually, the Langmuir wave bunch localized in the corresponding density cavity is referred to as the Langmuir soliton. The spatial collapse, on the other hand, occurs due to intensification of the Langmuir wave packet in the self generated density cavity [Zakharov, 1972].

[7] Gurnett et al. [1981] were the first to report the Voyager observations of the possible evidence for MI and spatial collapse of Langmuir waves in the Jovian foreshock region. However, in this case, the daughter ion sound waves were not detected, and the intensities of Langmuir waves were well below the strong turbulence thresholds. The Ulysses observations of Langmuir waves in the form of intense millisecond spikes were interpreted in terms of Langmuir collapse [Kellogg et al., 1992; Thejappa et al., 1993b]. However, in this case, there were some instrumental uncertainties regarding millisecond spikes. We have reported the observations of the coexistence of strong and weak Langmuir turbulence processes [Thejappa and MacDowall, 1998], envelope solitons formed as a result of subsonic MI [Thejappa et al., 1999] and ion-sound waves radiated by the burnt-out cavities left behind by the collapse events.
[Thejappa and MacDowall, 2004] in type III sources. However, the envelope soliton does not correspond to the concept of physical collapse of Langmuir waves because of its relatively large spatial scale $S \gg \lambda$ and low intensity ($\lambda$ is the wavelength of the Langmuir wave). Of much greater interest from the point of collapse is the large-amplitude Langmuir soliton formed in the regime $\frac{m_e}{\lambda_{De}} > (k_e \lambda_{De})^2$ and $\frac{m_i}{\lambda_{De}} > \frac{m_e}{\lambda_{De}}$, which corresponds to the supersonic MI [Zakharov et al., 1985]. Here $\frac{m_e}{\lambda_{De}}$ is the normalized energy density of Langmuir waves, $m_e$ and $m_i$ are the electron and ion masses, $\lambda_{De}$ is the electron Debye length and $k_e$ is the wave number of the Langmuir waves. Even after years of comprehensive in situ high time resolution measurements of the wave and beam parameters [Lin et al., 1986; Kellogg et al., 1992; Gurnett et al., 1993; Thejappa et al., 1993a, 1993b, 1996, 1999, 2003; Hospodarsky and Gurnett, 1995; Thejappa and MacDowall, 1998, 2004; Ergun et al., 2008; Henri et al., 2009], the evidence for the Langmuir collapse has proved to be extremely elusive.

Recently, Thejappa et al. [2012a] reported the high time resolution observations of a very intense localized Langmuir wave packet in the source region of a type III radio burst from the improved Time Domain Sampler (TDS) of the STEREO/WAVES experiment (improved over that of all similar high time resolution receivers flown in earlier spacecraft in precision, linearity, sample length and rate [Kellogg et al., 2009]). The characteristics of this wave packet include: (1) the normalized peak energy density $\frac{m_e}{\lambda_{De}}$, which is well above the threshold for the supersonic MI, (2) timescale of only ~3 ms, and (3) spectrum containing the resonant peak at the local electron plasma frequency, $f_{pe}$, sidebands at frequencies slightly lower than $f_{pe}$ (Stokes mode) as well as slightly higher than $f_{pe}$ (anti-Stokes mode), and low frequency enhancement below a few hundred Hz, which satisfy the frequency and wave number matching rules of MI. The spatial scale determined using the observed timescale of the wave packet is less than the spatial scale derived using the normalized peak intensity, which indicates that the wave packet is probably the Langmuir soliton collapsed to a few hundred Debye lengths. These observations were interpreted in terms of the oscillating two stream instability (OTSI), which is nothing but the supersonic MI and collapsing Langmuir soliton. In a separate study, Thejappa et al. [2012b] applied the trispectral analysis technique and showed that the spectral components of this wave packet are coupled to each other with a high degree of phase coherency (high tricoherence).

In this study, we present a detailed analysis of high time resolution observations of Langmuir wave packets associated with four other type III events. These TDS events represent some of the most intense localized Langmuir wave packets ever detected in type III bursts. The characteristics of these wave packets indicate that MI and associated strong turbulence processes occur in type III radio bursts. The implications of this study include (1) the strong turbulence processes probably play important roles in beam stabilization as well as conversion of Langmuir waves into escaping fundamental and harmonic radio emissions, and (2) the Langmuir collapse in type III sources probably occurs via MI as suggested by Zakharov [1972]. In section 2, we present the observations, in section 3, we present the discussion, and in section 4, we present the conclusions.

2. Observations

[10] The STEREO/WAVES experiment [Bougeret et al., 2008] is uniquely suited to study the nonlinear processes associated with solar type III radio bursts. It can detect a variety of in situ waves as well as remote radio emissions over a wide frequency range. Data from three of the SWAVES instruments are pertinent to this study: the high frequency receiver (HFR), low frequency receiver (LFR) and time domain sampler (TDS).

[11] 1. High frequency receiver (HFR)- The HFR measures the remote electric field signals over the frequency range 125 kHz - 16.025 MHz with 50 kHz increments. The total dynamic range of the HFR is 80 dB and its sensitivity at the preamplifier input is 6 nVHz$^{-1/2}$.

[12] 2. Low frequency receiver (LFR)- The LFR is a direct conversion receiver for spectral processing from 2.5 kHz up to 160 kHz. The frequency range is divided into three 2-octave bands A (2.5–10) kHz, B (10–40) kHz and C (40–160) kHz. The LFR covers a broad dynamic range with a constant amplitude resolution over the whole range, i.e., it adjusts its gain according to the input level. Automatic Gain Control (AGC) determines the receiver gain on a logarithmic scale as a function of the input level. The LFR analog front end normalizes the signal amplitude with an allowable gain variation of 80 dB. Combined with 12-bit analog to digital converters, the LFR provides a total dynamic range of 120 dB. The LFR sensitivity is 6 nVHz$^{-1/2}$ at the preamplifier input.

[13] 3. Time domain sampler (TDS)- The TDS is intended primarily for the study of Langmuir waves at the local electron plasma frequency, $f_{pe}$. It gathers events on three channels simultaneously connected to three orthogonal antennas and to a fourth pseudo-dipole channel obtained by taking the difference of any two monopoles. The TDS has programmable sampling rates with events of maximum length of 16 kilo samples, with possible rates of 250,000, 125,000, 31,250 and 7,812.5 samples per second. The maximum event duration can be as long as 2 s. The use of a 16 bit linear A/D converter provides sufficient dynamic range with better fidelity. The long durations of the TDS events enable us to determine the frequency components of the wave packets accurately. The gain is linear in the frequency range of interest, from typical electron cyclotron frequency, $f_{ec} \sim 100$ Hz to typical electron plasma frequency $f_{pe} \sim 30$ kHz.

[14] In the data, we have identified four type III events, for which the radio emissions drift from high frequencies to low frequencies. The Langmuir wave packets captured by the TDS during these events form the observational basis for the current study. These type III events occurred on 2011 December 19 (event I), 2012 February 12 (event II), 2009 July 18 (event III) and 2006 December 5 (event IV). We assign a typical value of $\sim 10^5$ K for the electron temperature $T_e$ for all four events. Here we note that owing to an unfortunate malfunction of the electron experiment, STEREO does not measure the electron temperature, $T_e$. For the solar wind speed ($v_{sw}$), we use the
values of 450 km s\(^{-1}\), 500 km s\(^{-1}\) and 500 km s\(^{-1}\) for events I, II and III, respectively. These values are from the STEREO/PLASTIC experiment [Galvin et al., 2008]. Since the data of solar wind are not available from the STEREO spacecraft for 2006 December 5 event (event IV), we use a value \(v_{\text{sw}} = 315\) km s\(^{-1}\) obtained by the WIND/SWE experiment [Ogilvie et al., 1995]; the STEREO A, B and WIND were very close to each other during this event. We have deduced the electron densities, \(n_e\), of the solar wind as \(\sim 4.5 \times 10^6\) m\(^{-3}\), \(\sim 3.2 \times 10^6\) m\(^{-3}\), \(\sim 0.8 \times 10^6\) m\(^{-3}\) and \(\sim 1.5 \times 10^5\) m\(^{-3}\) from the frequencies of the dominant spectral peaks of the waveforms of \(\sim 19.12\) kHz, \(\sim 14\) kHz, \(\sim 20.1\) kHz and \(\sim 31.6\) kHz (see, Figures 5, 9, 13, and 17); these frequencies are assumed to correspond to local electron plasma frequencies, \(f_{\text{pe}}\), \(\approx n_e^{1/2}\), \(n_e\) is in m\(^{-3}\)) during events I to IV, respectively.

[15] The solarsoft solar activity reports indicate that events I to IV are likely associated with active regions of N18W28, N29E85, N22W81 and S06E71, respectively. For the electron densities, \(n_e\), of solar wind are not available from the STEREO spacecraft at that frequency for almost 35 min, and, then jump suddenly to \(\sim 25\) kHz at \(\sim 14:20\) UT. These abrupt changes in the frequencies of Langmuir waves are probably due to some density structure in the solar wind. As seen from this spectrogram, at any given instant, the Langmuir waves show substantial frequency spreading of \(\sim 8\) kHz, which is indicative of nonlinear frequency broadening. The Langmuir wave burst corresponding to the TDS event of the present study is shown by an arrow in this figure.

[19] Time Domain Sampler (TDS) has resolved these Langmuir waves as highly localized wave packets. Each of these wave packets contains 16384 samples with an acquisition rate of 250,000 samples per second (a time step of 4 \(\mu\)s for a total duration of 65 ms). In Figure 4, we present one of the most intense wave packets detected during this type III event. This event is selected because of its sideband spectral structure, even though it is slightly weaker than the TDS events, which occur later at \(\sim 13:20\) UT. The peak electric field strength \(E_L\) of this event is \(34.6\) mV m\(^{-1}\) and its \(\frac{1}{c}\) power duration \(\tau\) is \(\sim 5.7\) ms. These values yield (1) the normalized peak energy density \(w_{\text{pe}} = \frac{E^2}{2n_e}\) \(\approx 8.5 \times 10^{-4}\) for \(E_L = 34.6\) mV m\(^{-1}\), (2) the spatial scale of the field structure \(S \approx 250\) m; by using the relation \(S \sim v_{\text{sw}}\tau\) for \(\tau = 5.7\) ms, \(v_{\text{sw}} = 450\) km s\(^{-1}\) and \(\lambda_{\text{De}} \approx 10.3\) m. Since the \(E_x\), \(E_y\) and \(E_{x-y}\) signals are weaker and show the same general features as the \(E_z\) signal, we analyze only the \(E_z\) signal. We justify one-dimensional treatment by assuming that these Langmuir wavefields are probably aligned along the ambient magnetic field.

[20] In Figure 5, we show the spectral characteristics of this wave packet. Figure 5a shows the total spectrum of the waveform from 0 to 65 kHz. The primary spectral peak at \(\sim 19.12\) kHz probably corresponds to Langmuir waves excited by the beam at the local electron plasma frequency, \(f_{\text{pe}}\). Figure 5b shows the narrow logarithmic spectrum around this primary spectral peak. This spectrum contains an intense peak (L) at \(f_{\text{pe}} \approx 19.12\) kHz, corresponding to the beam excited Langmuir wave, a down-shifted sideband (D) at \(\sim 19\) kHz which is slightly less than \(f_{\text{pe}}\) and an up-shifted sideband (U) at \(\sim 19.3\) kHz, which is slightly higher than \(f_{\text{pe}}\). Figure 5c shows linear spectrum below 1.5 kHz. A clear enhancement below 200 Hz seen in this low frequency spectrum probably corresponds to ion-sound waves.
2.2. February 11, 2012 (Event II)

In Figure 6, we present the dynamic spectrum of the second type III event and associated Langmuir waves, observed on February 11, 2012 by the STEREO B spacecraft. The fast drifting type III burst and the non-drifting Langmuir wave emissions are clearly seen in this dynamic spectrum. In Figure 7, we present the frequency-time spectrogram in a narrow frequency range, which shows the clumpy Langmuir waves with substantial frequency spreading. In this figure, the Langmuir wave burst corresponding to the TDS event of the present study is shown by an arrow. In Figure 8, we present one of the most intense wave packets detected by the TDS during this type III event. This TDS event is characterized by the peak electric field strength \( E_L \) of 44.73 mV\( \text{m}^{-1} \) and \( \frac{1}{n} \)-power duration \( \tau \) of \( \sim 7.9 \text{ ms} \). These values yield:

1. \( \frac{n}{\tau} \approx 2.6 \times 10^{-3} \) for \( E_L = 44.7 \text{ mV m}^{-1} \), \( n_e = 2.4 \times 10^6 \text{ m}^{-3} \) and \( T_e = 10^7 \text{ K} \), and
2. \( S \sim 282 \lambda_{\text{De}} \) for the measured values of \( \tau = 7.9 \text{ ms} \) and \( v_{\text{sw}} = 500 \text{ km s}^{-1} \). In Figure 9, we show the spectral characteristics of this wave packet. The total spectrum of the waveform from 0 to 65 kHz is shown in Figure 9a, where the primary spectral peak at \( \sim 14 \text{ kHz} \) corresponds probably to beam excited Langmuir wave at \( f_{pe} \).

The narrow logarithmic spectrum around this primary peak (L) in Figure 9b shows two sidebands, Stokes peak (D) at \( \sim 13.78 \text{ kHz} \) (which is slightly less than \( f_{pe} \)), and an anti-Stokes peak (U) at \( \sim 14.25 \text{ kHz} \) (which is slightly higher than \( f_{pe} \)). In Figure 9c, we present the low frequency spectrum, which shows a clear enhancement \( \leq 250 \text{ Hz} \) corresponding to ion-sound waves.

2.3. July 18, 2009 (Event III)

In Figure 10, we present the dynamic spectrum of the third local type III event and associated Langmuir waves. This event was observed by the STEREO A spacecraft on July 18, 2009. In Figure 11, we present the frequency-time spectrogram of the Langmuir waves, where the arrow shows the location of the TDS event of the present study. Here we note that even though this TDS event (Figure 12) is slightly weaker than some of the later events, it is included in the present analysis because of its sideband spectral structure. This TDS event is characterized by the peak electric field \( E_L \) of 70.9 mV\( \text{m}^{-1} \) and \( \frac{1}{n} \)-power duration \( \tau \) of \( \sim 10 \text{ ms} \). These values yield:

1. \( \frac{n}{\tau} \approx 3.2 \times 10^{-3} \) for \( E_L = 70.9 \text{ mV m}^{-1} \), \( n_e = 5 \times 10^6 \text{ m}^{-3} \) and \( T_e = 10^5 \text{ K} \), and
2. \( S \sim 512 \lambda_{\text{De}} \) for the measured values of \( \tau = 10 \text{ ms} \) and \( v_{\text{sw}} = 500 \text{ km s}^{-1} \). Figure 13 presents the spectral characteristics of this wave packet. Figure 13a shows the total spectrum of the waveform from 0 to 65 kHz, where the dominant spectral peak at \( \sim 20.1 \text{ kHz} \) corresponds to Langmuir waves excited by the electron beam at \( f_{pe} \). Figure 13b shows the narrow logarithmic spectrum around this intense peak (L), which contains two sidebands: a Stokes peak (D) at \( \sim 19.85 \text{ kHz} \) (which is slightly less than \( f_{pe} \)), and an anti-Stokes peak (U) at \( \sim 20.32 \text{ kHz} \) (which is slightly higher than \( f_{pe} \)). In Figure 13c, we present the low frequency spectrum of this
wave packet, which shows a clear enhancement $\leq$250 Hz corresponding to ion-sound waves.

2.4. December 5, 2006 (Event IV)

[23] In Figure 14, we present the dynamic spectrum of the fourth local type III event and associated Langmuir waves, observed by the STEREO B spacecraft. Here the fast drifting emission from $\sim$10:00 to $\sim$12:30 UT is the type III burst, and the non-drifting emissions in the 28–32 kHz frequency range are the Langmuir waves. In Figure 15, we present the frequency-time spectrogram in a narrow frequency band, which clearly shows that the Langmuir waves are very bursty and show substantial frequency spreading $\sim$10 kHz. The arrow shows the Langmuir wave burst corresponding to the TDS event of the present study.

[24] In Figure 16, we present one of the most intense waveforms detected during this type III event, which contains several modulated structures. The most intense structure is characterized by $E_L = 69.7$ mV m$^{-1}$ and $\tau \sim 7$ ms. These measured parameters yield: (1) $\frac{W_L}{n_eT_e} \approx 1.3 \times 10^{-3}$ for $E_L = 69.7$ mV m$^{-1}$, $n_e = 1.23 \times 10^7$ m$^{-3}$ and $T_e = 10^5$ K, and (2) $S \sim 355\lambda_{De}$ for the measured values of $\tau = 7$ ms and $v_{sw} = 315$ km s$^{-1}$. In Figure 17, we show the spectral characteristics of this wave packet, where Figure 17a shows the total spectrum from 0 to 65 kHz. The dominant spectral peak at $\sim$31.6 kHz in Figure 17a corresponds to Langmuir waves excited by the electron beam at the local electron plasma frequency, $f_{pe}$, which corresponds to $n_e$ of $1.23 \times 10^7$ m$^{-3}$. The narrow logarithmic spectrum around this dominant spectral peak (L) as shown in Figure 17b contains several secondary peaks. The peaks at frequencies slightly less than $f_{pe}$ (Stokes modes) are denoted as $U_1$, $U_2$ and $U_3$. The low frequency spectrum as seen in Figure 17c clearly shows the ion acoustic wave associated enhancement $\leq$200 Hz.

3. Discussion

[25] In Table 1, we present the summary of observations. The normalized peak energy density $\frac{W_L}{n_eT_e} = \frac{e^{2E_L}}{2\pi^2v_{sw}}$ controls the non-linear beam-plasma interactions. If

$$\frac{W_L}{n_eT_e} \geq (k_i\lambda_{De})^2,$$

then the dominant nonlinear beam-plasma interactions are the modulational instability [Papadopoulos et al., 1974] and the fully developed soliton formation and collapse.
Nicholson et al., 1978. In the present case, this condition is easily satisfied, since the observed \( \frac{n_e}{n_i} \) of 8.5 \( \times \) 10\(^{-4} \), 2.6 \( \times \) 10\(^{-3} \), 3.2 \( \times \) 10\(^{-3} \) and 1.3 \( \times \) 10\(^{-3} \) are well above the estimated \((k_L \lambda_{De})^2\) values of 1.9 \( \times \) 10\(^{-4} \), 2.2 \( \times \) 10\(^{-4} \), 1.2 \( \times \) 10\(^{-4} \) and 1.4 \( \times \) 10\(^{-4} \), for the events I to IV, respectively. This suggests that the strong turbulence processes can occur, and, therefore, in the context of present observations, we examine MI [Zakharov, 1972; Papadopoulos et al., 1974; Smith et al., 1979; Goldstein et al., 1979], and related soliton formation and Langmuir collapse [Zakharov, 1972; Nicholson et al., 1978; Goldman, 1984].

For the events of the present study, the inequality

\[
k_L \leq \left( \frac{m_e}{m_i} \right)^{1/2} k_D
\]

is easily satisfied, since \( \left( \frac{m_e}{m_i} \right)^{1/2} k_D \) are of 2.3 \( \times \) 10\(^{-3} \) m\(^{-1} \), 1.7 \( \times \) 10\(^{-3} \) m\(^{-1} \), 2.4 \( \times \) 10\(^{-3} \) m\(^{-1} \) and 3.8 \( \times \) 10\(^{-3} \) m\(^{-1} \), whereas, the wave numbers of the beam excited Langmuir waves, \( k_L \), are \~1.3 \( \times \) 10\(^{-3} \) m\(^{-1} \), \~1 \( \times \) 10\(^{-3} \) m\(^{-1} \), \~1.1 \( \times \) 10\(^{-3} \) m\(^{-1} \) and \~1.9 \( \times \) 10\(^{-3} \) m\(^{-1} \) for the events I to IV, respectively. We have estimated the beam speeds for a smooth spiral. The pitch angle scattering is known to increase the path length of electrons by a factor of \( \alpha = 1.3 \) to 1.7 [Alvarez et al., 1975; Lin et al., 1973]. This implies that the beam speeds, corrected for the pitch angle scattering will be much more favorable for MI and spatial collapse.

As far as the inequality \( \frac{n_e}{n_i} > \frac{m_e}{m_i} \) is concerned, it is also easily satisfied for all four events, since the observed \( \frac{n_e}{n_i} \) are 8.5 \( \times \) 10\(^{-4} \), 2.7 \( \times \) 10\(^{-3} \), 3.2 \( \times \) 10\(^{-3} \) and 1.3 \( \times \) 10\(^{-3} \), whereas, \( \frac{m_e}{m_i} \) is \~5.5 \( \times \) 10\(^{-4} \). This indicates that the observed wave packets are in the regime of supersonic modulational instability (MI) as shown schematically in Figure 18. The different regimes of Langmuir wave saturation mechanisms are discussed by Zakharov et al. [1985].

The growth rate of supersonic MI can be estimated as \~4 \times 10\(^{-4} \), \~7 \times 10\(^{-4} \), \~7.7 \times 10\(^{-4} \) and \~4.8 \times 10\(^{-4} \) for the observed \( \frac{n_e}{n_i} \) of 8.5 \( \times \) 10\(^{-4} \), 2.7 \( \times \) 10\(^{-3} \), 3.2 \( \times \) 10\(^{-3} \) and 1.3 \( \times \) 10\(^{-3} \) corresponding to the events of Figures 4, 8, 12, and 16, respectively. On the other hand, the bandwidth of the beam-excited Langmuir waves

\[
\frac{\Delta \omega}{\omega_{pe}} = 3(k_L \lambda_{De})^2 \frac{\Delta k_L}{k_L}
\]
Figure 5. (a) The complete spectrum of the wave packet of Figure 4 from 0 to 65 kHz, (b) the narrow spectrum around $f \sim f_{pe} \sim 19.12$ kHz, where L, D, and U correspond to the beam excited Langmuir wave, down-shifted sideband at $\sim 19$ kHz, and up-shifted sideband at $\sim 19.3$ kHz, respectively, and (c) the low frequency spectrum: the enhancement below 200 Hz corresponds to ion-sound waves. The sideband emissions are probably due to nonlinear interactions between the beam excited Langmuir waves and ion sound waves.
Here $\Delta v_b$ is the velocity dispersion of the beam, $N \approx \ln p$ is the number of linear growth times before the onset of MI, and $p$ is the ratio of the peak electric field amplitude to the thermal background. We estimate that $p \approx 10^4$ for all four events. It is interesting to note that higher the $N$, the longer the time available for the linear beam-plasma instability to grow. For the events of the present study, we estimate the number of linear growth times $N$ as $\geq 9$ for $p \approx 10^4$. Thus, for $N \sim 9$ and $\Delta v_b \approx 0.1 v_p$, we obtain $\frac{\Delta k}{k_{L}} \approx 3.9 \times 10^{-3}$. This value yields $\frac{2\omega}{\omega_p}$ as $2.2 \times 10^{-6}$, $2.5 \times 10^{-6}$, $1.4 \times 10^{-6}$, and $1.6 \times 10^{-6}$ for $k_L$ of $1.3 \times 10^{-3}$ m$^{-1}$, $1 \times 10^{-3}$ m$^{-1}$, $1.1 \times 10^{-3}$ m$^{-1}$, and $1.9 \times 10^{-3}$ m$^{-1}$, and $\lambda_{De}$ of $10.3$ m, $14.0$ m, $14.1$ m, and $9.8$ m, respectively. Thus the growth rates $\frac{1}{\omega_p}$ are much larger than the bandwidths of beam-excited Langmuir wave packets $\frac{2\omega}{\omega_p}$. This suggests that the pump waves are monochromatic enough for excitation of supersonic MI. It is important to note that for the bandwidths of initial pump waves, one should not use the spectral widths from Figures 5b, 9b, 13b, and 17b since these are severely affected by the nonlinear effects.

[29] The peak intensity also decides whether the wave packet is a collapsing Langmuir soliton of the type described by Zakharov [1972] and Nicholson et al. [1978] or not. For the wave packet to be a collapsing soliton, it should satisfy the condition [Thornhill and ter Haar, 1978; Gurnett et al., 1981]

$$\frac{W_L}{n_e T_e} \geq (\Delta k \lambda_{De})^2,$$

where $\Delta k = \frac{2\pi}{\lambda}$ is the wave number characteristic of the envelope. Assuming that the wave packet is convected in the solar wind, we can convert the measured timescales $\tau$ of $\sim$5.7 ms, $\sim$7.9 ms, $\sim$10 ms and $\sim$7 ms into spatial scales $S$ of $\sim$250$\lambda_{De}$, $\sim$282$\lambda_{De}$, $\sim$512$\lambda_{De}$, and $\sim$355$\lambda_{De}$ using the relation $S \sim \tau v_{sw}$ for $v_{sw}$ of 450 km s$^{-1}$, 500 km s$^{-1}$, 500 km s$^{-1}$, and 315 km s$^{-1}$, and $\lambda_{De}$ of $\sim$10.3 m, $\sim$14.1 m, $\sim$9.8 m and $\sim$6.2 m, for the events I to IV, respectively. Thus, the observed $\frac{W_L}{n_e T_e}$ of $\sim$8.5$ \times 10^{-4}$, $\sim$2.7$ \times 10^{-3}$, $\sim$3.2$ \times 10^{-3}$ and $\sim$1.3$ \times 10^{-3}$ are greater than $(\Delta k \lambda_{De})^2$ of $\sim$6.3$ \times 10^{-4}$, $\sim$5$ \times 10^{-4}$, $\sim$1.5$ \times 10^{-4}$ and $\sim$3.1$ \times 10^{-4}$ obtained for the spatial scale $S$ of $\sim$250$\lambda_{De}$, $\sim$282$\lambda_{De}$, $\sim$512$\lambda_{De}$, and $\sim$355$\lambda_{De}$ for events of Figures 4, 8, 12, and 16, respectively. This suggests that the observed wave packets are probably the Langmuir solitons collapsed to the spatial scales of a few hundred Debye lengths.

### 3.1. Spectral Characteristics

[30] As shown in Figure 1, MI, which is a four-wave interaction process excites a low frequency ion density perturbation of frequency and wave number ($\Omega$, $q$). This perturbation can be either a freely propagating ion sound wave or a strongly damped quasi-mode. It can beat with two of the

**Figure 6.** Dynamic spectrum of the second local type III radio burst (fast drifting emission from $\sim$5 MHz down to $\sim$20 kHz) and associated Langmuir waves (non-drifting emissions in the frequency interval 10–14 kHz).
initial pump waves with wave numbers, $k_L$ and frequencies $f_L \sim f_{pe}$ (beam-excited Langmuir waves), and produce the down-shifted (Stokes) and up-shifted (anti-Stokes) side-bands with wave numbers and frequencies of $(k_L - q)$ and $(f_L - \Omega)$, and $(k_L + q)$ and $(f_L + \Omega)$, respectively. Thus, the Fourier spectrum resulting from such a four wave MI should contain the main emission peak at $f_L \sim f_{pe}$ together with two sidebands at $(f_L - \Omega)$ and $(f_L + \Omega)$, and a low frequency component. In the present case, the spectral peaks seen in Figures 5b, 9b, 13b, and 17b probably correspond to the down-shifted sideband (Stokes (D)), beam-excited Langmuir (L) and up-shifted sideband (anti-Stokes (U)) modes, respectively, and the low frequency spectral enhancements seen in Figures 5c, 9c, 13c, and 17c correspond to ion acoustic waves. This suggests that the observed spectral structures of the wave packets show one-to-one correspondence with the characteristic spectral signatures of the MI.

The requirement for the observed spectral components to be associated with MI is that they should satisfy the frequency, wave number and phase matching conditions:

$$2f_L = f_D + f_U$$

$$2k_L = k_D + k_U$$

$$2\phi_L = \phi_D + \phi_U,$$

where, the subscripts L, D and U correspond to the beam-excited Langmuir, down-shifted sideband (Stokes) and up-shifted sideband (anti-Stokes) modes, respectively.

In order for the frequency matching condition to be satisfied, the down-shifted (Stokes) and up-shifted (anti-Stokes) sidebands should be symmetric with respect to the beam-excited Langmuir wave, and their frequency shifts $\Delta f = f - f_{pe}$ should be in the frequency ranges of the respective low frequency enhancements.

For December 19, 2011 event (Figure 5b), the frequency shifts of the down-shifted and up-shifted sidebands are reasonably symmetric with respect to the Langmuir wave pump, being $\sim 120$ Hz and $\sim 180$ Hz, respectively. These frequency shifts are in good agreement with the observed frequencies of ion sound waves of $<200$ Hz (Figure 5c). For February 11, 2012 event (Figure 9b), the frequency shifts of the down-shifted and up-shifted sidebands with respect to Langmuir waves are $\sim 220$ Hz and $\sim 250$ Hz, respectively, which are in agreement with the observed frequencies of ion sound waves of $<250$ Hz. For July 18, 2009 event (Figure 13), the frequency shifts of the D and U modes (Figure 13b) are $\sim 250$ Hz and $\sim 220$ Hz, which are again in good agreement with the observed frequencies $<250$ Hz of the low frequency enhancement as seen in Figure 13c. The U mode in this case is weaker than the D mode, probably because it suffers higher damping rate in comparison with that of the D mode.

Figure 7. Frequency-time spectrogram of the Langmuir waves associated with the type III burst of February 11, 2012. The Langmuir wave emissions, which are very bursty show substantial frequency spreading. The arrow shows the Langmuir wave burst corresponding to the current TDS event.
In the case of TDS event of December 5, 2006, the spectral peak at \( \omega_{\text{C24}} \) at 31.6 kHz in Figure 17b corresponds to the beam-excited Langmuir wave, the spectral peaks below this frequency correspond to Stokes modes propagating anti-parallel to the beam, and the spectral peaks at frequencies higher than 31.6 kHz correspond to anti-Stokes modes propagating parallel to the beam. The connection between the structure in the spectrum of this event and the MI is as follows. The MI excites low frequency ion density perturbations of frequency and wave number \((\omega, q)\), which can beat with two of the initial beam excited Langmuir waves, producing high frequency sidebands with wave numbers \(\bar{k}_L - \bar{q}\) (Stokes mode) and \(\bar{k}_L + \bar{q}\) (anti-Stokes mode). These Stokes and anti-Stokes modes can again couple to the pump wave and regenerate the initial low frequency density perturbation at \(\bar{q}\); they can also couple to each other to generate another low frequency density perturbation at \(2\bar{q}\). This new density perturbation can beat with the pump to produce another pair of Stokes and anti-Stokes waves at \(\bar{k}_L \pm 2\bar{q}\) which beat together to produce another density perturbation at \(4\bar{q}\), which, in turn, will also generate new Stokes and anti-Stokes waves. A hierarchy of low frequency density perturbations together with higher order sidebands are thus generated. This process of producing higher order sidebands can proceed as long as the level of ion density fluctuations are above the thermal level. The frequency shifts of all the Stokes and anti-Stokes modes are in the frequency range of the low frequency enhancement seen in Figure 17c.

We can estimate the wave numbers of the downshifted and up-shifted sidebands and examine whether they satisfy the required matching condition as follows. The wave numbers \(k\) of the sidebands can be estimated using the expression for the Doppler shift:

\[
\Delta \omega = \frac{\nu_{\text{sw}}}{2\pi\lambda_{\text{De}}} (k\lambda_{\text{De}}) \cos \theta,
\]

where \(\theta\) is the angle between \(\bar{k}\) and \(\nu_{\text{sw}}\), i.e., \(\theta = 0\) and \(\theta = \pi\) correspond to the anti-Stokes and Stokes modes propagating away from and toward the Sun, respectively. The pump Langmuir waves, the sidebands and ion sound waves should satisfy the conservation relation

\[
\bar{k} = \bar{k}_L \pm \bar{q},
\]

where the plus and minus signs correspond to anti-Stokes and Stokes modes, respectively. Since the wave numbers of the pump waves \(k_L\) are usually much smaller than \(q\), the \(|\bar{q}|\) should be \(\approx |\bar{k}|\).

In the case of TDS event of Figure 4, the frequency shifts of the D and U modes of \(\sim 120\) Hz and \(\sim 180\) Hz as seen in Figure 5b are <200 Hz, which is the frequency range of the low frequency enhancement seen in Figure 5c.
Figure 9. (a) The complete spectrum of the wave packet of Figure 8 from 0 to 65 kHz, where the primary peak at ~14 kHz corresponds to beam excited Langmuir wave, (b) the narrow logarithmic spectrum around $f \sim f_{pe} \sim 14$ kHz, where the spectral peaks $L$, $D$, and $U$ correspond to the beam excited Langmuir wave, down-shifted sideband $\sim 13.78$ kHz, and up-shifted sideband $\sim 14.25$ kHz, respectively, and (c) the low frequency linear spectrum: the enhancement below 250 Hz corresponds to ion-sound waves. The non-linear interactions between the beam-excited Langmuir waves and the ion sound waves observed are probably responsible for the sideband emissions.
Figure 10. Dynamic spectrum of the third local type III radio burst (fast drifting emission from ~5 MHz down to ~20 kHz) and associated Langmuir waves (non-drifting emissions in the frequency interval 19–22 kHz).

Figure 11. Frequency-time spectrogram of the Langmuir waves associated with type III burst of July 18, 2009. The arrow shows the Langmuir wave burst corresponding to the current TDS event.
Substituting these measured values $\Delta f \sim 120$ Hz and $\Delta f = 180$ Hz in equation (12), we estimate $k\lambda_{De}$ as $\sim -0.017$ and $\sim 0.026$, for the Stokes and anti-Stokes modes with $\theta = 180^\circ$ and $\theta = 0^\circ$, respectively. These values indicate that the pump Langmuir waves with $k_L\lambda_{De} \sim 0.014$ are very efficiently converted into those of forward and backward propagating daughter Langmuir waves with large wave numbers. As far as the ion sound waves seen in Figure 5c are concerned, since their phase velocities are usually well below $v_{sw}$, their frequency spectrum in the spacecraft frame of reference is solely determined by the Doppler shifts. Thus, the upper limit of the wave numbers of these waves can be estimated as $q = 2\pi \Omega \nu_{sw} / v_{sw} \sim 3.1 \times 10^{-3}$ m$^{-1}$ and $q \lambda_{De} \approx 0.029$ for $\Omega = 200$ Hz and $v_{sw} = 450$ km s$^{-1}$. Thus, the wave numbers of the observed low frequency waves are comparable to those of the sideband emissions.

[37] In the case of TDS event of Figure 8, the frequency shifts of the D and U modes $\sim 220$ Hz and $\sim 250$ Hz are in agreement with the observed frequencies of ion sound waves of $\sim 250$ Hz. Substituting these measured values $\Delta f \sim 220$ Hz and $\Delta f = 250$ Hz in equation (12), we can estimate $k\lambda_{De}$ as $\sim -0.039$ and $\sim 0.044$, for the Stokes and anti-Stokes modes, respectively, and compare them with $k_L\lambda_{De} \sim 0.015$. The upper limit of the wave numbers of the ion sound waves of Figure 9c can be estimated as $q = 2\pi \Omega / v_{sw} \sim 3.1 \times 10^{-3}$ m$^{-1}$ and $q \lambda_{De} \approx 0.044$ for $\Omega = 250$ Hz and $v_{sw} = 500$ km s$^{-1}$. Thus, the wave numbers of the observed low frequency waves are comparable to those of the sideband emissions.

[38] In the case of July 18, 2009 event (Figure 12), for the Stokes and anti-Stokes modes with $\theta = 180^\circ$ and $\theta = 0^\circ$, we can plug the measured values of $\Delta f \sim 250$ Hz and $220$ Hz in equation (12) and estimate $k\lambda_{De}$ as $\sim -0.031$ and $\sim 0.027$, respectively. These values of $k\lambda_{De}$ are much higher than $k_L\lambda_{De} \sim 0.015$. As far as the ion sound waves seen in Figure 13c are concerned, the upper limit of their wave numbers are estimated as $q = 2\pi \Omega / v_{sw} \sim 3.1 \times 10^{-3}$ m$^{-1}$ and $q \lambda_{De} \approx 0.031$ for $\Omega = 250$ Hz and $v_{sw} = 500$ km s$^{-1}$. Thus, the wave numbers of the observed low frequency waves are comparable to those of the sideband emissions, which indicates that the observed low frequency waves are probably involved in parametric interaction with the Langmuir waves.

[39] In the case of Dec 5, 2006 event (Figure 16), we can estimate $k_L\lambda_{De} \sim 1.2 \times 10^{-2}$ for $k_L$ is $\sim 1.9 \times 10^{-3}$ m$^{-1}$ and $\lambda_{De} \sim 6.2$ m. Plugging the maximum measured $\Delta f$ of $200$ Hz in equation (12), we estimate that the wave numbers $k$ of the anti-Stokes and Stokes modes should be as large as $\sim 2k_L$. Thus, the spectral observations of Langmuir wave packets show that the beam excited Langmuir waves presented in Figure 16 with $k_L\lambda_{De}$ of $\sim 0.012$ are very efficiently converted into those of the forward and backward propagating daughter waves with large wave numbers up to $k\lambda_{De} \sim 0.025$. As far as the wave numbers of low frequency waves (Figure 17c) are concerned, we can estimate $q = 2\pi \Omega / v_{sw} \sim$...
Figure 13. (a) The complete spectrum of the wave packet of Figure 12 from 0 to 65 kHz, where the dominant spectral peak at ~20.1 kHz corresponds to the Langmuir wave excited by the electron beam at local electron plasma frequency, $f_{pe}$, (b) the narrow spectrum around the main Langmuir wave peak at $f \sim f_{pe} \sim 20.1$ kHz, where $L$, $D$, and $U$ correspond to the beam excited Langmuir wave, down-shifted sideband at ~19.85 kHz, and up-shifted sideband peak at ~20.32 kHz, respectively, and (c) the low frequency spectrum: the enhancement below 250 Hz probably corresponds to ion-sound waves.
Figure 14. Dynamic spectrum of the fourth local type III radio burst (fast drifting emission from \(\sim 5\) MHz down to \(\sim 32\) kHz) and associated Langmuir waves (non-drifting emissions in the frequency interval 28–32 kHz).

Figure 15. Frequency-time spectrogram of the Langmuir waves associated with the type III burst of December 5, 2006. The Langmuir wave emissions are very clumpy and show substantial frequency spreading. The arrow shows the location of the current TDS event.
4 \times 10^{-3} \text{ m}^{-1} \) and \( q \lambda_{De} \approx 0.025 \) for \( \Omega = 200 \text{ Hz} \) and \( v_{sw} = 315 \text{ km s}^{-1} \). These wave numbers of low frequency waves are comparable to those of sideband waves.

[40] Thus, the sidebands and low frequency enhancements contained in the spectra of TDS events of Figures 4, 8, 12, and 16 probably correspond to the daughter waves excited by the MI.

3.2. Trispectral Analysis

[41] The conventional spectral analysis does not yield complete information about the nonlinear processes from the waveform data. To extract the phase coherence amongst spectral components of the wave packets, the higher order spectral analysis techniques, namely, the bispectral and trispectral analysis techniques should be used. The bi-spectral analysis can extract the nonlinear characteristics between three spectral components with frequency sum and difference [Kim and Powers, 1979]. Since bi-spectral analysis is related to the skewness of a signal, it can detect only the asymmetric non-linearities, such as the electrostatic decay [Henri et al., 2009; Balikhin et al., 2001; Walker et al., 2002, 2003] and harmonic generation [Bale et al., 1996]. The amplitude of the bispectrum at the bifrequency \( (F_k, F_l) \) measures the amount of coupling between the frequencies, \( F_k, F_l \) and \( F_{k+l} \). The trispectral analysis, on the other hand, is used to demonstrate the occurrence of cubic nonlinear interactions in the system. The tricoherence spectra can isolate the individual cubic interactions between quartets of spectral components, especially those associated with oscillations at the power spectral primary peak frequency and its sidebands, as well as their combination tones.

[42] A cumulant based trispectral method has been developed and applied to synthetic [Kravtchenko-Berejnoi et al., 1995a, 1995b; Lefevre et al., 1995] and simulated data [Soucek et al., 2003]. The expression for the trispectrum can be written as [Kravtchenko-Berejnoi et al., 1995a]

\[
T(k, l, p) = E\left[X_k X_l X_p^* X_q^*\right] - N(k, l, p, q),
\]

where \( X_k, X_l, X_p \) and \( X_q \) are the complex Fourier components of the signal at frequencies \( F_k, F_l, F_p \) and \( F_q = F_k + F_l - F_p \), and \( N(k, l, p, q) = E[X_k X_l X_p X_q] + E[X_k X_p X_q] + E[X_l X_p X_q] + E[X_k X_l X_q] + E[X_k X_l X_p X_q] \). Here \( E[] \) is the expectation operator. In the present case, the subscripts \( k \) and \( l \) refer to beam excited Langmuir waves, and \( p \) and \( q \) refer to down-shifted (Stokes) and up-shifted (anti-Stokes) sidebands, respectively.

[43] The normalized trispectrum is usually referred to as the tricoherence. It is more useful because it eliminates the dependence of trispectrum on the amplitudes of the signals. The tricoherence quantifies the phase coherence amongst the spectral components of the wave packet, i.e., it quantifies the fraction of the total product of powers at the frequency quartet, \( (F_k, F_l, F_p, F_q) \), that is owing to cubicly phase-coupled modes. The expression for the square
Figure 17. (a) The complete spectrum of the wave packet of Figure 16 from 0 to 65 kHz, where the dominant spectral peak at $\sim 31.6$ kHz corresponds to the beam excited Langmuir wave, (b) the narrow spectrum around $f \sim f_{pe} \sim 31.6$ kHz, where the main peak is denoted as $L$, the down-shifted (Stokes) sidebands are denoted as $D_1$ and $D_2$, and up-shifted (anti-Stokes) sidebands are denoted as $U_1$, $U_2$ and $U_3$. (c) The low frequency spectrum, which shows a clear enhancement below 200 Hz corresponding to ion-sound waves.
tricoherence can be written as \cite{Kravchenko-Berejnoi1995}:  
\[
\hat{t}^2(k,l,p) = \frac{|T(k,l,p)|^2}{(E[X_kX_lX_p] + |N(k,l,p,q)|)^2},
\]  
(15)

A unit value of the tricoherence indicates perfect coupling, a zero value indicates no coupling, and any value between zero and one indicates partial coupling.

The tricoherence is zero for a Gaussian process. Even for a Gaussian process, due to statistical fluctuations, the estimate of the tricoherence from a finite data record will not be zero. Therefore, the method of periodograms is usually used to estimate the trispectrum and tricoherence. This method involves the division of the data record into M segments; an appropriate window is applied to each segment to reduce leakage; the trispectrum as well as tricoherence are computed for each segment by using the DFT; finally, the trispectrum and tricoherence are averaged across segments to reduce the variance of the estimator.

**Table 1.** Relevant Parameters

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**Figure 18.** Different regimes of Langmuir turbulence.

**Figure 19.** The three dimensional representation of the tricoherence spectrum \(\hat{t}^2(F_k, F_p, F_p)\) of the TDS event of 2011 December 19 (Figure 4). The tricoherence \(\hat{t}^2\) exhibits peak value of \(\sim 0.35\) at \((\sim 19.25, \sim 19, \sim 18.75)\) kHz, which quantifies the phase relation \(2\phi_D = \phi_D + \phi_L\), where \(\phi_L\) and \(\phi_D\) and \(\phi_U\) are the phases of the beam-excited Langmuir wave at \(\sim 19.12\) kHz, Stokes \((\sim 19)\) kHz and anti-Stokes \((\sim 19.3)\) kHz modes, respectively.
As seen from equation (14), the trispectrum estimator is symmetric with respect to permutations of its arguments $F_k$, $F_l$, and $F_p$. The principal domain for the interaction of the type $F_k + F_l = F_p + F_q$ is determined as $[Kravtchenko-Berejnoi et al., 1995a] 0 \leq F_k \leq f_N, 0 \leq F_l \leq F_k, 0 \leq F_p \leq F_l$, and $F_p \leq F_k + F_l - F_p \leq f_N$, where $f_N$ is the Nyquist frequency.

As seen from equation (14), the trispectrum estimator is symmetric with respect to permutations of its arguments $F_k$, $F_l$, and $F_p$. The principal domain for the interaction of the type $F_k + F_l = F_p + F_q$ is determined as $[Kravtchenko-Berejnoi et al., 1995a] 0 \leq F_k \leq f_N, 0 \leq F_l \leq F_k, 0 \leq F_p \leq F_l$, and $F_p \leq F_k + F_l - F_p \leq f_N$, where $f_N$ is the Nyquist frequency.

Thejappa et al. [2012a] reported the observations of a highly localized Langmuir wave packet captured by the TDS during one of the local type III events. This wave packet was characterized by a short duration of $\sim 3$ ms, peak intensity exceeding the threshold for the OTSI as well as spatial collapse, and the spectrum, containing the signatures of OTSI, namely, an intense peak corresponding to beam-resonant Langmuir wave, two sidebands corresponding to down-shifted and up-shifted daughter Langmuir waves, and a low frequency enhancement corresponding to daughter ion sound waves; the frequencies and wave numbers of these spectral components satisfied the resonance conditions of OTSI. Since the spectral methods are phase blind, Thejappa et al. [2012b] had developed cumulant based trispectral analysis tools to extract the phase information. Using those tools, these authors have clearly demonstrated that the spectral components of the observed wave packet are coupled to each other with a high degree of phase coherency as expected of OTSI.

In this study, we will use the trispectral analysis tools developed by Thejappa et al. [2012b] to extract the phase information amongst the spectral components of the wave packets of Figures 4, 8, 12, and 16. For this purpose, as expected of the periodogram method, we divide the time series into $M = 16$ segments with segment length $N = 1000 (0.004 s)$. We apply the Hamming window for each segment, and calculate the tricoherence spectrum as a function of three frequencies $F_k$, $F_l$, and $F_p$.

The true values of tricoherence are known to be zero for Gaussian noise [Hasselman et al., 1963; Nikias and Raghuveer, 1987; Dalle Molle and Hinich, 1989]. The 90% significance level for estimate of tricoherence for Gaussian noise is given by [Chandran et al., 1994]

\[
\hat{t}^2_{G k l p} = \frac{4.6}{2M}.\]

In the present case, for $M = 16$, we obtain $\hat{t}^2_{G} \sim 0.144$. Therefore, the tricoherence $\hat{t}^2$ values above 0.144 are significant at the 90% significance level for $M$ blocks of the data, i.e., for 2 $M$ degrees of freedom. Although any tricoherence value between zero and one indicates a partial coupling of four waves, we consider only those values which are greater than 0.144 as significant.

As seen from equation (14), the trispectrum estimator is symmetric with respect to permutations of its arguments $F_k$, $F_l$, and $F_p$. The principal domain for the interaction of the type $F_k + F_l = F_p + F_q$ is determined as [Kravtchenko-Berejnoi et al., 1995a] $0 \leq F_k \leq f_N, 0 \leq F_l \leq F_k, 0 \leq F_p \leq F_l$, and $F_p \leq F_k + F_l - F_p \leq f_N$, where $f_N$ is the Nyquist frequency.

Thejappa et al. [2012a] reported the observations of a highly localized Langmuir wave packet captured by the TDS during one of the local type III events. This wave packet was characterized by a short duration of $\sim 3$ ms, peak intensity exceeding the threshold for the OTSI as well as spatial collapse, and the spectrum, containing the signatures of OTSI, namely, an intense peak corresponding to beam-resonant Langmuir wave, two sidebands corresponding to down-shifted and up-shifted daughter Langmuir waves, and a low frequency enhancement corresponding to daughter ion sound waves; the frequencies and wave numbers of these spectral components satisfied the resonance conditions of OTSI. Since the spectral methods are phase blind, Thejappa et al. [2012b] had developed cumulant based trispectral analysis tools to extract the phase information. Using those tools, these authors have clearly demonstrated that the spectral components of the observed wave packet are coupled to each other with a high degree of phase coherency as expected of OTSI.

The true values of tricoherence are known to be zero for Gaussian noise [Hasselman et al., 1963; Nikias and Raghuveer, 1987; Dalle Molle and Hinich, 1989]. The 90% significance level for estimate of tricoherence for Gaussian noise is given by [Chandran et al., 1994]

\[
\hat{t}^2_{G k l p} = \frac{4.6}{2M}.\]

In the present case, for $M = 16$, we obtain $\hat{t}^2_{G} \sim 0.144$. Therefore, the tricoherence $\hat{t}^2$ values above 0.144 are significant at the 90% significance level for $M$ blocks of the data, i.e., for 2 $M$ degrees of freedom. Although any tricoherence value between zero and one indicates a partial coupling of four waves, we consider only those values which are greater than 0.144 as significant.
using the tricoherence technique of the TDS event of 2011 December 19, constructed with the coherence spectrum displayed as a function of two frequencies and one of the anti-Stokes modes, respectively.

This value of tricoherence, which is a statistically significant measure for 2012 February 11 event, we present the 3-D tricoherence spectrum in Figure 20. This 3-D spectrum is constructed using the computed $r^2$ values in the frequency range $[14.25 \text{ kHz} \leq F_k \leq 14.75 \text{ kHz}, 14 \text{ kHz} \leq F_l \leq 14.25 \text{ kHz}, 13.75 \text{ kHz} \leq F_p \leq 14.25 \text{ kHz}]$. Although, the tricoherence spectral feature has an extended shape in 3-D frequency space, it exhibits a peak of $\sim 0.58$ at $(\sim 14.5, \sim 14.25, \sim 13.75)$ kHz. This peak quantifies the phase coherence between two beam excited Langmuir waves at $\sim 14$ kHz and a down-shifted sideband at $\sim 13.75$ kHz and an up-shifted sideband at $\sim 14.25$ kHz.

[52] For the 2009 July 18 event, the 3-D tricoherence spectrum is presented in Figure 21. The tricoherence used in this figure are computed in the frequency ranges $[20 \text{ kHz} \leq F_k \leq 20.75 \text{ kHz}, 20 \text{ kHz} \leq F_l \leq 20.25 \text{ kHz}, 19.75 \text{ kHz} \leq F_p \leq 20.25 \text{ kHz}]$. In this case, the $r^2$ attains a peak of $\sim 0.32$ at $(\sim 20.25, \sim 20, \sim 19.75)$ kHz. Again, since it is statistically significant, it quantifies the phase relation $2\phi_L = \phi_D + \phi_U$.

[53] In Figure 22, we present the 3-D tricoherence $r^2$ spectrum of the 2006 December 5 event. In this case, the tricoherence $r^2$ attains a peak of $\sim 0.32$. This is computed in the frequency range $[31.625 \text{ kHz} \leq F_k \leq 31.875 \text{ kHz}, 31.5 \text{ kHz} \leq F_l \leq 31.875 \text{ kHz}, 31.25 \text{ kHz} \leq F_p \leq 31.75 \text{ kHz}]$. The corresponding frequency quartet is $(\sim 31.5, \sim 31.5, \sim 31.25, \sim 31.75)$ kHz. This tricoherence quantifies the phase relation $2\phi_L = \phi_D + \phi_U$.

[54] The results of tricoherence estimates can also be displayed as 2-D cross-sections at different frequencies. For example, for 2006 December 5 event, we present the cross sections of the tricoherence spectrum in Figure 23. The top panel shows the cross-section at $F_p = f_D = 31.25$ kHz, which shows a clear tricoherence spectral peak at $(\sim 31.5, \sim 31.5, \sim 31.25, \sim 31.75)$ kHz. The peak value of $r^2$ in this case is $\sim 0.32$. This value agrees with that of 3-D analysis. This peak tricoherence quantifies the phase relation $2\phi_L = \phi_D + \phi_U$, where $\phi_L$, $\phi_D$ and $\phi_U$ are the phases of the beam-excited Langmuir wave $(\sim 31.5)$ kHz, one of the Stokes $(\sim 31.25)$ kHz and one of the anti-Stokes $(\sim 31.75)$ kHz modes, respectively. The maximum tricoherence $r^2 \sim 0.32$ in this case is statistically significant.

[55] In the second panel, we present the cross-section of the tricoherence spectrum at $F_p = f_D = 31.50$ kHz. It shows two comparatively weaker spectral features, with peak values of $\sim 0.21$ at $(\sim 31.75$ kHz, $\sim 31.625$ kHz, $\sim 31.50$ kHz) as well as at $(\sim 32, \sim 31.875, \sim 31.50)$ kHz. These relatively smaller tricoherence values reflect the weaker four wave interactions involving the sidebands themselves. In the bottom panel, we present the tricoherence cross-section at $F_p = f_U = 30.75$ kHz, which also shows a weak peak of $\sim 0.21$ at $(\sim 31.75, \sim 31.875, \sim 32.125)$ kHz. This again reflects the four wave interactions involving the sidebands themselves. However, the tricoherence spectral peak $r^2 \sim 0.32$ (top panel) is more significant in the statistical sense. Therefore, the dominant four wave interaction in this case is $2\phi_L \rightarrow f_D + f_U$. One should note that, the structure in the spectrum of this TDS event as seen in Figure 17b with multiple sidebands and the tricoherence peaks in the second and third panels indicate that the four wave interactions between the sidebands also can be possible.
Figure 23. The cross-sections of the tricoherence spectrum of the TDS event of Figure 16, at (top) $F_p = f_D = 31.25$ kHz, (middle) $F_p = f_L = 31.50$ kHz, and (bottom) $F_p = f_U = 31.75$ kHz, where $f_D$, $f_L$, and $f_U$ are the frequencies of the down-shifted sideband, beam-excited Langmuir and up-shifted sideband modes, respectively.
supersonic modulational instability as suggested by lapse in the case of type III bursts probably follows the of the electron plasma frequency, and (3) the Langmuir col-
spatial collapse probably play important roles in beam radio bursts, (2) the supersonic modulational instability (MI) presented in this study are that (1) the strong Langmuir tur-
ments strongly suggest that the wave packets are collapsing high levels of tricoherence amongst their spectral compo-
is satisfy the criterion of solitons collapsed to spatial scales of a few hundred Debye lengths. The spectra of these wave packets consist of intense peaks corresponding to beam-resonant Langmuir waves, two or more sidebands corresponding to Stokes and anti-Stokes modes, and low frequency enhancements below a few hundred Hz corresponding to ion sound waves. The frequencies and wave numbers of these spectral components satisfy the resonance conditions of the modulational instability. The tricoherence values are computed using the trispectral analy-
s techniques, which show that the spectral components of the wave packets are coupled to each other with a high degree of coherency, as expected of the MI. The observed high intensities, short scale lengths, sideband spectral structures and low frequency enhancements of the wave packets, and high levels of tricoherence amongst their spectral compo-
nents strongly suggest that the wave packets are collapsing Langmuir solitons, and they are in the supersonic modula-
tional instability regime. The implications of the observations presented in this study are that (1) the strong Langmuir tur-
bulence processes occur in the source regions of solar type III radio bursts, (2) the supersonic modulational instability (MI) and spatial collapse probably play important roles in beam stabilization as well as conversion of Langmuir waves into escaping radiation at the fundamental and second harmonic of the electron plasma frequency, and (3) the Langmuir collapse in the case of type III bursts probably follows the supersonic modulational instability as suggested by Zakharov [1972].

Acknowledgments. The research of T. G. is supported by the NASA Grants NNX08AO02G, NNX09AB19G and NNX12AH47G. The SWAVES instruments include contributions from the Observatoire de Paris, University of Minnesota, University of California, Berkeley, and NASA/ GSFC. We thank the referees for constructive comments and helpful suggestions.

Philippa Browning thanks the reviewers for their assistance in evaluating this paper.

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