VERITAS Observations
of
Six Bright, Hard-Spectrum Fermi-LAT Blazars


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ABSTRACT

We report on VERITAS very-high-energy (VHE; $E \geq 100$ GeV) observations of six blazars selected from the Fermi Large Area Telescope First Source Catalog (1FGL). The gamma-ray emission from 1FGL sources was extrapolated up to the VHE band, taking gamma-ray absorption by the extragalactic background light into account. This allowed the selection of six bright, hard-spectrum blazars that were good candidate TeV emitters. Spectroscopic redshift measurements were attempted with the Keck Telescope for the targets without Sloan Digital Sky Survey (SDSS) spectroscopic data. No VHE emission is detected during the observations of the six sources described here. Corresponding TeV upper limits are presented, along with contemporaneous Fermi observations and non-concurrent Swift UVOT and XRT data. The blazar broadband spectral energy distributions (SEDs) are assembled and modeled with a single-zone synchrotron self-Compton model. The SED built for each of the six blazars show a synchrotron peak bordering between the intermediate- and high-spectrum-peak classifications, with four of the six resulting in particle-dominated emission regions.

Subject headings: active galactic nuclei: general — gamma rays: individual (RGB J0136+361, RGB J0316+090, RGB J0909+231, RGB J1058+564, RGB J1243+364, RX J1436.9+5639)
1. Introduction

Blazars are active galactic nuclei (AGN) with a relativistic jet pointed close to the Earth line of sight. AGN are thought to be powered by accretion onto a supermassive black hole (∼ 10⁹M☉) at the center of the host galaxy and are characterized by a double-peaked spectral energy distribution (SED) in the νFν representation.

The lower energy peak of the broadband SED is attributed to synchrotron emission of highly relativistic electrons and positrons in the presence of a tangled magnetic field. In leptonic models, the higher energy peak is produced via the inverse-Compton (IC) up-scattering by the relativistic leptons of the synchrotron photons (synchrotron self-Compton, SSC) or a photon field external to the jet (external Compton, EC). More details regarding leptonic non-thermal emission of blazars can be found in Dermer et al. (1992); Maraschi et al. (1992); Sikora et al. (1994) and the references therein. Alternative models associate the higher-energy peak to interaction of relativistic protons with an ambient photon field (Aharonian et al. 2002; Bednarek 1993; Dar et al. 1997; Mannheim 2000; Mücke & Protheroe 2000; Pohl & Schlickeiser 2000), or a hybrid population comprised of both leptons and hadrons (Böttcher 2007).

The blazar population is divided into two subclasses: flat spectrum radio quasars (FSRQs) and BL Lacs. FSRQs are, to first order, more distant, more luminous, and have stronger emission lines than BL Lacs. Historically, BL Lacs have been sub-classified based on their radio and X-ray flux ratios as low-, intermediate- or high-frequency-peaked BL Lac objects (LBL, IBL and HBLs, respectively) as described in Padovani & Giommi (1995) and Bondi et al. (2001). More recently, a classification based on the location of the synchrotron peak in frequency space (νsynch) has been proposed by Abdo et al. (2010a), with low-spectrum-peaked (LSP) BL Lac objects having νsynch below 10¹⁴ Hz, intermediate-spectrum-peaked (ISP) BL Lac objects peaking between 10¹⁴ and 10¹⁵ Hz and high-spectrum-peaked (HSP) BL Lac objects showing a peak above 10¹⁵ Hz. HSP BL Lacs are the most common extragalactic object to be detected at very-high-energies (VHE; E ≥ 100 GeV), comprising 33 of 41 VHE BL Lac objects detected as of June 2012. There have also been 4 ISPs (Acciari et al. 2009a,b, 2010; Ong 2010a) and 4 LSPs (Mazin et al. 2009; Mariotti 2011; Hoffman 2010; Albert et al. 2007) detected since the advent of VHE gamma-ray astronomy in the late 1960s. In addition, three FSRQs, three radio galaxies, and two starburst galaxies complete the catalog of associated extragalactic objects (Horan & Wakely 2008)¹.

The potential scientific impact of increasing the catalog of VHE emitting extragalactic objects is substantial. A significant fraction of the power released from these objects is within

¹http://tevcat.uchicago.edu/
the VHE band. This makes the measurement of VHE blazar spectra an important component of the overall understanding of these objects. With a better sample of well-measured VHE blazar spectra available for study, a population-based investigation of gamma-ray production in these objects through broadband SED modeling will be possible, providing means to answer the long-standing question of whether VHE gamma-ray emission results from leptonic or hadronic processes in AGN jets. We can also apply the model-inferred properties of these new discoveries to understanding how the gamma-ray production may differ among blazar subclasses, exploring the apparent blazar sequence and evolution of the AGN (Fossati et al. 1998; Ghisellini & Tavecchio 2008; Meyer et al. 2011).

VHE blazars can also be used to constrain the optical to near-infrared extragalactic background light (EBL) density and evolution, as well as the nature of cosmic rays. The EBL encodes the integrated history of structure formation and the evolution of stars and galaxies in the Universe. Understanding these characteristics requires detailed theoretical modeling of all the processes that contribute, e.g. structure formation and stellar evolution (Dominguez et al. 2011; Gilmore et al. 2009; Primack et al. 2005; Finke et al. 2010; Franceschini et al. 2008; Hauser & Dwek 2001; Stecker et al. 2006). According to some cosmic ray models, e.g. Essey & Warren (2012), interactions of cosmic rays along the blazar line-of-sight can produce relatively hard VHE gamma-ray spectra as compared with the high-energy gamma-ray spectra, depending on the distance to the blazar.

The current catalog of extragalactic VHE objects largely contains relatively nearby AGN; only three have a confirmed redshift above $z = 0.3$. The highest redshift blazar detected at VHE thus far is the FSRQ 3C279, at a redshift of 0.536 (Albert et al. 2008b). The proximity of these VHE blazars is partly a result of gamma-ray absorption by the EBL. VHE gamma rays that propagate through the intergalactic medium are absorbed by low energy EBL photons via pair production, $\gamma + \gamma \rightarrow e^+ + e^-$ (Nikishov 1962; Gould & Shr´eder 1967; Stecker et al. 1992). The absorption process deforms the VHE gamma-ray spectra emitted by extragalactic objects in a redshift-dependent manner and can be translated to upper limits on the local density of the far-infrared EBL (Stecker et al. 1993; Dwek & Krennrich 2005; Aharonian et al. 2006; Primack et al. 2011; Orr et al. 2011).

The search for new VHE blazars is complicated by the fact that many of these objects do not yet have known redshifts. BL Lac objects, by definition, display very weak or no optical emission or absorption lines used for spectroscopic redshift measurements. Due to the interaction of VHE photons with the EBL, non-detection of blazars with no redshift information can be either attributed to the object being too distant or having an intrinsically low luminosity in the VHE band. This makes VHE discovery observations of blazars with no measured redshift a risky venture, although it has proven successful in the past, as
in the discoveries of VHE emission from 3C 66A, PKS 1424+240, and RXJ0648.7+1516 (Acciari et al. 2009a, 2010; Aliu et al. 2011).

The small field of view of imaging atmospheric Cherenkov telescopes (IACTs, \( \leq 5.0^\circ \)) makes new source discovery from a large-scale sky survey difficult and therefore the hunt for VHE emitting objects has historically involved targeted observations of source candidates selected from surveys at lower frequencies, such as the selection of hard X-ray candidates presented in Costamante & Ghisellini (2002). In this way, VHE blazar candidate selection has relied on experiments such as EGRET onboard the Compton Gamma Ray Observatory (Thompson et al. 1993), Swift (Gehrels et al. 2004) and ROSAT (Turriziani et al. 2007). The launch of Fermi in June 2008 has enhanced VHE blazar discovery programs, leading to new blazar discoveries including RBS0413, RXJ0648.7+1516, 1ES0033+595 and 1RXS J101015.9-311909 (Aliu et al. 2012, 2011; Mariotti 2011; Abramowski et al. 2012).

The Fermi Large Area Telescope (LAT) observes the entire sky in the energy range from 20 MeV to > 300 GeV every three hours and has better sensitivity than its predecessor EGRET. Within the first 11 months of operation, the 1FGL catalog reported the detection of 1451 sources at a significance greater than 5\(\sigma\) (Abdo et al. 2010b). A majority of these sources are, or are expected to be, associated with AGN (Abdo et al. 2010c). It is known that a large majority of these blazars has not yet been detected by VHE instruments, as can be seen by the current TeV catalog, which only contains 51 associated extragalactic objects. However, the proximity of the Fermi-LAT energy band to that of the IACTs makes the 1FGL catalog a good place to search for candidate VHE blazars. We report on VERITAS observations above 100 GeV of six candidate blazars. For the first time, a multiwavelength description of their SED including radio, UV, X-ray and gamma-ray frequencies is assembled and modeled using a SSC model.

2. Target Selection

The energy coverage of IACTs overlaps with that of Fermi-LAT above 100 GeV and extends to tens of TeV. Sources that are most likely to be detected by both Fermi-LAT and IACTs have high fluxes and hard photon indices in the Fermi-LAT energy band. A selection process was established to identify the best VHE candidates within the first Fermi-LAT catalog. Sources at low Galactic latitude (\(|b| < 10^\circ\)) were excluded, with the intent of removing the majority of Galactic sources from the selection. Additionally, sources with low integrated flux above 100 MeV (\(F_{100\text{MeV}} < 2 \times 10^{-9}\text{ph cm}^{-2}\text{s}^{-1}\)), with soft photon indices (\(\Gamma > 2.0\)), or with a low number of associated photons (\(N_{\text{pred}} < 20\)), were excluded from the selection. Just over 200 sources passed this initial set of cuts.
The remaining sources were re-analyzed using the same data set as that used for the 1FGL catalog to search for the presence of curvature in their spectra. A log-parabola parameterization for the spectra was chosen as the alternative to the power-law model (null hypothesis). In addition, the data were analyzed only using data above 1 GeV to confirm that the results from the power-law fit found for the entire Fermi-LAT energy band agreed well with the power-law fit above 1 GeV. Sources showing significant spectral curvature or softening of the spectrum above 1 GeV with an improvement in likelihood value corresponding to $3\sigma$ were excluded from the final selection as these sources are not expected to exhibit bright TeV flux levels. For sources with spectra best fit by a power law that also matched the power-law fit above 1 GeV, the Fermi-LAT spectrum was extrapolated up to the VHE regime (150 GeV - 1 TeV). When a reliable redshift measurement was available, the EBL model of Franceschini et al. (2008) was used to estimate the extrapolated VHE gamma-ray flux. Otherwise an assumed redshift of $z = 0.2$ was used, a conservative value considering that most known TeV sources have redshifts less than $z = 0.2$. Finally, the sources were ranked based on their extrapolated flux in the VHE regime. Six VHE candidates, all BL Lac objects, were selected for observations with the VERITAS telescope between September 2009 and June 2010 for 10 hours each, corresponding to the VERITAS 3\% Crab flux sensitivity exposure timescale.

Table 1 shows the VHE extrapolated integral flux from the 1FGL catalog power-law fits used for selection of the six candidate VHE emitting blazars. These relatively high integral flux values above 150 GeV, shown in percentages of the Crab Nebula flux,$^2$ are shown in comparison to the updated power-law fits from 29 months of Fermi-LAT data for the steady sources, and data from the VERITAS coincidental window for variable sources. The analysis of the extended Fermi-LAT dataset (2008 August 4 to 2011 January 4) is detailed in section 4.2. These updated extrapolations show much lower expected integral flux values for each of the six candidates, reflecting the fact that as more LAT data were collected, better high-energy statistics showed the candidates to be softer and/or dimmer than found with the eleven month data set used in the 1FGL. Additionally, we compare these extrapolated values to the upper limits derived from the VERITAS observations, where the analysis leading to these upper limits is detailed in Section 4.1.

$^2$Flux calculated according to the curved power law presented in Albert et al. (2008a)
3. The Targets

**RGB J0136+391** (1FGL J0136.5+3905) was discovered in the third Bologna sky survey of 408 MHz radio objects (Ficarra et al. 1985). It was later detected as a X-ray bright active galaxy in the Northern ROSAT all-sky survey (Brinkmann et al. 1997), and identified spectroscopically as a BL Lac object, with an IBL sub-classification from Laurent-Muehleisen et al. (1998) and Laurent-Muehleisen et al. (1999), respectively. Based on the optical and X-ray spectral properties, this blazar was proposed as a VHE candidate blazar by Costamante & Ghisellini (2002), under the assumption that the redshift was less than 0.2. Bright gamma-ray emission above 1 GeV was detected from this source by *Fermi* after three months of operation. Only an upper limit below 1 GeV was reported (Abdo et al. 2009). This blazar also showed constant emission in the first eleven months of LAT operation (Abdo et al. 2010b). This blazar remains without a known redshift, with no previous spectroscopic redshift measurements found in the available literature. We attempted a spectroscopic redshift measurement on 2009 September 17 (MJD 55091) using the Keck Low Resolution Imaging Spectrometer (LRIS) instrument, but measured only a featureless power-law spectrum that is characteristic of BL Lac objects (see Figure 1 upper-left panel). Detailed spectroscopic analysis of this data can be found in Kaplan et al. (2012).

**RGB J0316+090** (1FGL J0316.1+0904) was first detected by the NRAO Green Bank 91 meter radio telescope (Becker et al. 1991). It was later optically identified as a BL Lac object (Fischer et al. 1998) and sub-classified as an IBL by Laurent-Muehleisen et al. (1999). No spectroscopic redshift measurements had been made prior to this work. We attempted a spectroscopic redshift measurement on 2011 March 5 using the Keck Echelette Spectrograph and Imager (ESI) instrument, resulting in a featureless power-law spectrum except for two unidentified absorption features (see Figure 1 upper-right panel). The detailed spectroscopic analysis of the Keck ESI data can be found in Kaplan et al. (2012).

**RGB J0909+231** (1FGL J0909.2+2310) was first detected by the NRAO Green Bank radio telescope (Becker et al. 1991) and was later classified as a radio-loud active galaxy (Brinkmann et al. 1997). The BL Lac optical counterpart was identified nearly a decade later (Mickaelian et al. 2006). The redshift reported by the NASA Extragalactic Database (NED) could not be found within the corresponding reference ($z = 0.231$, Brinkmann et al. 2000)) and was therefore taken as unknown. Inspection of publicly available Sloan Digital Sky Survey (SDSS) data revealed two Mg II absorption lines in the otherwise featureless optical spectrum. Assuming that these lines could be intrinsic, or due to the absorption by an intervening cloud, a lower limit on the redshift of $z \geq 0.43$ is derived (see Figure 1 middle-
left panel). This is the only source out of the six selected which has a neighboring Fermi-LAT detected blazar within the VERITAS 3° field of view. 2FGLJ0910.9+2246, associated with TXS0907+230, is located 0.61° away from RGBJ0909+231. The high redshift blazar TXS0907+230 \((z = 2.66\) according to Healey et al. (2008)) shows no signal in the 10-100 GeV band, with an upper limit of \(9.6 \times 10^{-11}\) photons cm\(^{-2}\) s\(^{-1}\) reported in the 2FGL catalog (Nolan et al. 2012). An exclusion region of radius 0.3° centered on the blazar was nevertheless used in the VERITAS analysis to avoid possible contamination.

**RGB J1058+564** (1FGLJ1058.6+5628) was first detected in the 6C radio survey (Hales et al. 1990). It was identified as a BL Lac object during the association of the ROSAT all-sky survey with the Hamburg Quasar Survey (Nass et al. 1996). Bondi et al. (2001) classified the object as an IBL based on the optical and X-ray characteristics. This blazar was detected within the first three months of Fermi-LAT operation (Abdo et al. 2009) with a broadband SED shown in Abdo et al. (2010a) which does not include any VHE information. The blazar also shows a high level of flux variability (probability of variability: 79%) in the high-energy gamma-ray band, as shown in Abdo et al. (2010b) and is the only one out of the six targets in this paper that has a redshift measurement. The absorption lines corresponding to the redshift 0.143 can be seen in the SDSS spectrum shown in Figure 1 (middle-right panel).

**RGB J1243+364** (1FGLJ1243.1+3627) was first reported in the B2 catalog of radio sources (Colla et al. 1973). This target was also determined to be a radio-loud active galaxy by ROSAT (Brinkmann et al. 1997), and specifically classified as a BL Lac in Appenzeller et al. (1998). NED cited the SDSS data for a spectroscopic redshift of \(z = 1.065\). Inspection of this publicly available SDSS data revealed no lines suggesting this redshift but instead revealed Mg II absorption lines that translate to a lower limit of \(z \geq 0.485\) (see Figure 1 lower-left panel). A recent photometric redshift of \(z= 0.5^{+0.14}_{-0.12}\) from Meisner & Romani (2010) is in agreement with the SDSS lower limit.

**RX J1436.9+5639** (1FGLJ1437.0+5640) was detected in the ROSAT all-sky survey and identified as a BL Lac object by Nass et al. (1996). This HBL, as classified by Nieppola et al. (2006) based on the frequency of the synchrotron peak, remains without a redshift, although the redshift of the galaxy cluster within the same region of the sky is known to be \(z= 0.15\) (Bauer et al. 2000). Inspection of the publicly available SDSS data shows a featureless spectrum, shown in Figure 1.
4. Multiwavelength Observations and Analysis

4.1. VERITAS

The VERITAS observatory is an array of four 12-meter diameter IACTs, located in southern Arizona. VERITAS is sensitive to photons between 100 GeV and several tens of TeV with an energy resolution of better than 20%. The instrument has a $5\sigma$ point source sensitivity of $1\%$ of the Crab Nebula flux in less than 30 hours with an angular resolution of less than 0.1$^\circ$ for a Crab-like source with a spectral index of 2.5. See Weekes et al. (2002) and Holder et al. (2006) for a detailed overview of the instrument.

The VERITAS observations of the six VHE candidate blazars were completed between September 2009 and June 2010 (MJD 55122–55383). These observations were taken in *wobble* mode, with an offset of 0.5$^\circ$ from the source position in each of four cardinal directions to allow simultaneous background measurement, as explained in Fomin et al. (1994) and Berge et al. (2007). The radio location of the counterparts as specified by NED were used for source position.

Air shower events initiated by gamma and cosmic-rays are reconstructed following the procedure outlined in Acciari et al. (2008). The recorded shower images are parameterized by their principal moments, giving an efficient method for suppression of the far more abundant cosmic-ray background. Any events with a total charge less than 50 photoelectrons are removed from the analysis. A set of cuts is then applied to the parameters in order to reject background events (see details of this method in Krawczynski et al. (2006)). These parameters reject cosmic-ray like events having the mean scaled width and length of the event camera image smaller than 1.1 or 1.4, respectively. Additionally, the reconstructed altitude of the maximum Cherenkov emission is required to be higher than 8 km above ground level. Gamma-ray like events are extracted from a signal region with a radius of 0.14$^\circ$, centered at the coordinates of the candidate source.

The quality-selected livetime collected for each target ranges from 4.4 to 14.2 hours and results in no detections, with significances ranging from -1.1 to 0.9$\sigma$, calculated with Equation (17) of Li & Ma (1983). The VERITAS observations and analysis results are detailed in Table 2. Integral upper limits at 99% confidence level are calculated using the Rolke et al. (2005) method, assuming a photon index of $\Gamma = 3.0$ for the differential power-law spectrum $dN/dE = N(E/E_o)^{-\Gamma}$. This index was assumed as a moderately softened index as compared to the *Fermi*-LAT index range (1.74–2.9), expected due to absorption by the EBL and possible intrinsic turnover. Additionally, this index value is representative of a typical TeV blazar. Without detection, the real index in the VHE band remains unknown and any integral upper limit derived is dependent on the index assumed. The differential upper
limit is quoted at the decorrelation energy, the energy where the calculated flux has minimal
dependence on the index. Changing the spectral index by $\pm 0.5$ changes the differential upper
limits by less than 10%.

The results are independently reproduced with two different analysis packages, as de-
scribed in Cogan (2008) and Daniel (2008). The upper limits range from 1 to 3% of the
integral Crab Nebula flux above the threshold energy. The energy threshold for each obser-
vation is defined as the energy at which the differential rate of reconstructed gamma rays
from the postulated source reaches its maximum and is accurate to within the 20% energy
resolution of the instrument, a value that is dependent on the observation zenith angle and
sky brightness.

4.2. Fermi LAT

The Fermi-LAT is a pair-conversion telescope sensitive from 20 MeV to $>300$ GeV,
which operates in survey mode. Further details about the characteristics and performance of
the LAT Instrument on the Fermi Gamma-ray Space Telescope can be found in Atwood et al.
(2009). Presented here is the analysis of the Fermi LAT data for the six candidates described
in Section 3. Although the targets were selected based on only 11 months of data, more data
was available after the completion of VERITAS observations and this larger data set is used
for the modeling. More specifically, the LAT data from the time period of 2008 August 4
to 2011 January 4 (MJD 54682.7-55565.0) were used for the modeling analysis. Except for
variable sources, the analysis procedure was identical for each of the sources and proceeded
as follows.

For each candidate, events were extracted from a region of interest (ROI) of $10^\circ$ radius
centered on the target coordinates. Events from the “diffuse class” with zenith angle $< 100^\circ$
and energy between 300 MeV and 300 GeV were selected. Only data taken during periods
when the rocking angle of the satellite was $< 52^\circ$ were used to reduce contamination from
the Earth limb gamma rays, which are produced by cosmic rays interacting with the upper
atmosphere. The significance and spectral parameters were calculated using an unbinned
maximum-likelihood method implemented in the LAT Science Tool gtlike$^4$ (Cash 1979;
Mattox et al. 1996). A background model was constructed including nearby gamma-ray
sources and diffuse emission. All sources within 12$^\circ$ of the central source in the second
Fermi-LAT catalog (2FGL, Nolan et al. (2012)) were included in the model. The spectra
of known pulsars were modeled by a power laws with exponential cutoffs. As in the 2FGL

$^4$ScienceTools v9r20p0 with the post-launch instrument response function (IRF) P6_V11_DIFFUSE.
catalog, a log-parabola function was used for sources with significant spectral curvature. Otherwise, spectra were described as a power law. The spectral parameters of the sources in the ROI were left free during the fitting procedure. Sources outside the ROI, but within the 12° range had their spectral parameters fixed to the 2FGL catalog values. The Galactic diffuse emission and an isotropic component, which is the sum of the extragalactic diffuse gamma-ray emission and the residual charged particle background, were modeled using the recommended files.

The LAT data were first analyzed to calculate the time-averaged gamma-ray flux and spectral parameters of each candidate. A second analysis was then performed to study the impact of the Sun, a bright gamma-ray source, on the flux of candidates located near the plane of the ecliptic (RGB J0316+090 and RGB J0909+231). Removing time intervals when the Sun was in the ROI of each candidate had a negligible effect on the analysis results.

Spectral points and a light curve were calculated for each candidate, and a temporal analysis was performed to search for flux variability. The timescale of this analysis was adjusted based on the specific candidate flux levels. The flux in each energy or time bin was determined with the spectral indices of all sources fixed to the best-fit values over the full energy and time interval. For an energy or time bin with a test statistic (TS; see Mattox et al. [1996]) less than 9 or fewer than 3 predicted photons ($N_{\text{pred}}$), a 95% confidence level upper limit was calculated.

The light curves were analyzed to search for flux variability with a likelihood method assuming a constant flux for the null hypothesis, following the same procedure as used in the 2FGL catalog. Only two sources showed significant evidence of flux variability: RGB J0316+090 and RGB J1058+564. For these two sources, a refined analysis was done, selecting time periods contemporaneous with the VERITAS observation windows. The duration of this contemporaneous period was chosen such that a significant detection ($>5\sigma$) could be attained, resulting in slightly extended windows of MJD 55055-55145 for RGB J0316+090 and MJD 55160-55185 for RGB J1058+564 with respect to the VERITAS observation window. For these candidates, a butterfly corresponding to the 1σ confidence interval was used to represent the spectral information (Figure 3).

The derived spectral indices for the differential power laws obtained for the candidates are relatively hard, ranging from 1.74 to 2.09, as compared to the 2FGL average spectral index of 2.21 ± 0.01. The integral fluxes above 300 MeV range from 1.55 to 12.4×10⁻⁹ph

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5The files used were gll_iem_v02_p6_v11_DIFFUSE.fit for the Galactic diffuse and isotropic_iem_v02_p6_v11_DIFFUSE.txt for the isotropic diffuse component available at http://fermi.gsfc.nasa.gov/ssc/data/p6v11/access/lat/BackgroundModels.html
cm$^{-2}\text{s}^{-1}$, indicating that these six blazars are bright in the high-energy band. The detailed Fermi-LAT results for each of these hard-spectrum, bright BL Lacs are summarized in Table 3.

4.3. **Swift XRT**

The X-ray Telescope (XRT) onboard the Swift satellite (Gehrels et al. 2004) is a focusing X-ray telescope sensitive to photons with energy between 0.2 and 10 keV. The data used for the broadband SED modeling were analyzed as described in Burrows et al. (2005) with the HEASoft package Version 6.9 and XSPEC$^6$ Version 12.6.0. All data were taken in photon counting mode and pile up effects are accounted for when count rates exceeded 0.5 counts per second through the use of an annular source region, with a 1-2 pixel inner radius and a 20 pixel outer radius. Each observation is binned and fit with an absorbed power law between 0.3 and 10 keV, with the neutral hydrogen density taken from the Leiden/Argentine/Bonn survey of Galactic HI (Kalberla et al. 2005).

X-ray variability is commonplace for both VHE detected and non-detected blazars. If more than one exposure exists for an object and no variability is detected, the de-absorbed power-law fit of the combined data set is used to constrain the SED modeling. If variability is observed between multiple exposures, results from these separate exposures are shown independently on the SED plot, as is the case for RGBJ0136+391 and RGBJ1058+564. These blazars show flux variability factors of $\sim$2 and 3 between exposures, respectively. The multiple exposures taken on RGBJ0909+231 do not provide sufficient statistics for application of an absorbed power-law model and are therefore summed before fitting. Using Cash analysis (Cash 1979) did not improve the fitting of the single low-statistics spectra. The summed exposure fit result is shown on the SED.

Each of the absorbed power-law fits applied to the XRT data resulted in photon indices greater than 2, with 2-10 keV integral flux levels between 0.3 and $21\times10^{-12}\text{ergs cm}^{-2}\text{s}^{-1}$. The index values suggest that the synchrotron component peaks below keV energies, characteristic of ISP blazars. The analysis results for each observation are summarized in Table 4.

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4.4. *Swift* UVOT

The *Swift*-XRT observations were supplemented with simultaneous UVOT exposures taken in the V, B, U, UVW1, UVM2, and UVW2 bands (Poole et al. 2008). The UVOT photometry is performed using the HEASoft program *uvotsource*. The circular source region has a 5′′ radius and the background region consists of several 15′′ radii circles of nearby empty sky. The results are reddening corrected using E(B-V) coefficients (Schlegel et al. 1998). The Galactic extinction coefficients are applied according to Fitzpatrick (1999). The uncertainty in the reddening E(B-V) is the largest source of error, especially in the UV bands for blazars that have a large value of E(B-V). If more than one exposure exists in a specific band for an object, the data from the observation closest to the VERITAS exposure are used, although no significant variability is seen across any band for any blazar. A summary of the UVOT analysis results is presented in Table 5.

5. Broadband SSC Modeling

Leptonic models for blazar jet emission attribute the higher energy peak in the SED to the inverse-Compton scattering of lower energy photons off a population of non-thermal, relativistic electrons. These same electrons are responsible for the lower-energy synchrotron emission that makes up the first peak. The target photon field involved in the Compton upscattering can either be the synchrotron photons themselves, as is the case in SSC models, or a photon field external to the jet in the case for EC models.

The previously described multiwavelength data are matched with archival radio data collected from NED. These data are used to test a steady-state leptonic jet model for the broad-band continuum emission from the blazars. Although it has been found that ISP BL Lacs are sometimes better represented by external-Compton models, e.g. Acciari et al. (2009b), taking into account the lack of redshift information and lack of constraints from the broadband data we prefer not to apply an EC model, which would introduce additional free parameters, compared to the SSC model applied here. The model-predicted flux reflects the absorption of VHE gamma rays by the EBL according to the redshift information summarized in Table 1 using the model from Gilmore et al. (2009). The level of TeV absorption resulting from this model is consistent with the absorption derived from the Finke et al. (2010) and Franceschini et al. (2008) models.

The SSC model applied to the multiwavelength data is the equilibrium version of the Böttcher & Chiang (2002) model, as described in Acciari et al. (2009b). In this model, the emission originates from a spherical blob of relativistic electrons with radius $R$. This blob is
moving down the jet with a Lorentz factor \( \Gamma \), which corresponds to a speed of \( \beta \Gamma c \). The jet axis is aligned toward the line of sight with an angle \( \theta_{\text{obs}} \), which results in Doppler boosting with a Doppler factor \( D = (\Gamma [1 - \beta \Gamma \cos \theta_{\text{obs}}])^{-1} \). In order to minimize the number of free parameters, we assume that \( \theta_{\text{obs}} = 1/\Gamma \), often referred to as the critical angle, for which \( \Gamma = D \).

Within the model, nonthermal electrons are injected and accelerated into a power-law distribution \( Q(\gamma) = Q_0 \gamma^{-q} \) between the low- and high-energy cut-offs, \( \gamma_{\text{min}} \) and \( \gamma_{\text{max}} \). The radiation mechanisms considered lead to an equilibrium between particle injection, radiative cooling and particle escape. This particle escape is characterized with an escape efficiency factor \( \eta \), such that the time scale of escape \( t_{\text{esc}} = \eta R/c \). This results in a particle distribution which streams along the jet with a power \( L_e \). Synchrotron emission results from the presence of a tangled magnetic field \( B \), with a Poynting flux luminosity of \( L_B \).

The two parameters \( L_e \) and \( L_B \) allow the calculation of the equipartition parameter \( \epsilon_{Be} \equiv L_B/L_e \). This equipartition parameter is used as an estimator of the feasibility of the model, where models which result in \( \epsilon_{Be} \sim 1 \) are preferred. If the particle energy density greatly dominates over the magnetic field energy density, namely a particle dominated scenario, then the magnetic field cannot serve to collimate the jet. Following this design, acceptable parameters should result in at least partition conditions with \( L_B \geq L_e \).

The broadband SED for each blazar can be seen in Figure 3, with the SSC model parameters for each representation summarized in Table 6. For each blazar, the archival radio data are taken as upper limits as these measurements are believed to contain a large amount of radiation produced in the radio lobes in addition to the synchrotron emission from the jet. The modeling for each of the six blazars shows synchrotron peak locations \( \sim 10^{15} \) Hz, characteristic of borderline ISP/HSP blazars.

**RGB J0136+391:** This blazar is modeled using the lower of the two variable X-ray states for three different assumed redshifts (\( z = 0.2, 0.3 \) and 0.4). The variability timescale of three months suggested by the factor-of-two variability observed between XRT exposures is not short enough to provide a constraint on the size of the model emission region. Under the assumption that the gamma-ray emission remained steady during the *Fermi* and VERITAS observations, the hard LAT spectrum and low VERITAS upper limit derived for this blazar suggest a steepening of the gamma-ray spectrum at \( E \geq 100 \) GeV, which could be caused by the EBL absorption if a redshift \( z \geq 0.4 \) is assumed. Alternatively, this apparent break could originate from uncorrelated variability in the high and very-high energy bands. More specifically, a low flux state in the VHE band during VERITAS observations could provide a redshift-independent explanation of the apparent spectral softening. Figure 3 shows the model predictions, corrected by EBL absorption, for each of the redshift values. Only the
model at $z = 0.4$ is compatible with the VERITAS upper limit, which also results in a framework with balanced radiation and particle energy.

**RGB J0316+090:** Due to the variability detected in the high-energy band, this blazar is modeled with *Fermi*-LAT data which spans the complete time period sampled by shorter VERITAS observations. The UVOT errors resulting from the E(B-V) reddening correction for this blazar are so large that all exposures except the V band are unconstraining to the SED modeling and therefore not shown. Application of the SSC model for an assumed redshift of $z = 0.2$ results in a particle dominated framework consistent with the VERITAS upper limit.

**RGB J0909+231:** The SSC model parameters used to describe the broadband emission of this blazar are determined for a redshift of $z = 0.5$, based on the lower limit derived from Mg II absorption lines found in the SDSS data. The *Swift* X-ray data from three exposures have been summed in order to provide sufficient statistics for absorbed power-law fitting. The model results in a particle-dominated scenario with a slight discrepancy between the model and the *Fermi* upper limit in the $1 - 3$ GeV energy bin. Moving the Compton peak accommodates this upper limit only results in an even less favored, particle-dominated emission state.

**RGB J1058+564:** Due to the variability observed in the LAT band, the SED modeling for this blazar is done with LAT data from the time period coincident with the VERITAS observation window and for the lower of two X-ray states observed. The variability timescale suggested by the variability factor of three observed between XRT exposures is not short enough to provide any constraint on the size of the emitting region. Both X-ray states are shown on the SED for reference. The SSC model shows agreement with the broadband data, although results in a disfavored particle dominated scenario.

**RGB J1243+364:** The modeling for this blazar is completed for a redshift of $z = 0.5$, based on the new lower limit found from Mg II absorption lines in the public SDSS data. The model agrees with the broadband data and permits parameters at equipartition.

**RX J1436.9+5639:** The modeling for the broadband data of this blazar allows marginal agreement with the *Fermi* data when completed for the redshift of the spatially coincident galactic supercluster, more specifically $z = 0.15$. The resulting model parameters are far below equipartition, suggesting a particle dominated jet.
6. Discussion

Six promising TeV blazar candidates were selected from the 1FGL catalog for observations with VERITAS. These observations resulted in no VHE detections. Initial selection of these candidates from 1FGL power-law fit extrapolation suggested integral flux levels between 1 and 12% of the Crab Nebula flux above 150 GeV after accounting for absorption by the EBL. With additional Fermi-LAT data and more information about the blazar redshifts, the expected fluxes were updated to levels between 0.3% and 3.5% of the Crab Nebula flux above 150 GeV. The VERITAS exposure times were allocated based on the initial extrapolated values, resulting in only one upper limit below the updated VHE extrapolation, namely the upper limit for RGBJ0136+391. The non-detection of this blazar suggests spectral steepening of the high-energy spectrum that can either be explained by intrinsic spectral curvature, redshift-dependent EBL absorption or uncorrelated variability.

Multiband observations are presented and allow for the construction and SSC modeling of the radio through TeV broadband SEDs. The model applied to these blazars is sufficient to represent the broadband data for each of the six, with model parameters roughly comparable to those found for other VHE-detected ISPs using the same model (Abdo et al. 2011; Acciari et al. 2009b). The magnetic fields obtained in the modeling are generally low, resulting in disfavored, particle-dominated jets. This condition could be relaxed by including an external photon field for inverse-Compton scattering, allowing solutions closer to equipartition, as was done in Abdo et al. (2011) and Acciari et al. (2009b). However, this scenario has not been explicitly tested. The quality of the data sampling of the inverse-Compton peak and the fact that the redshifts are not well determined for the majority of the objects do not provide sufficient constraints for a model with the extra degrees of freedom associated with adding an external source of seed photons. Similar parameters were also found for the borderline ISP/HSP TeV detected blazar PKS1424+240, with the exception of the spectral index for the injected electron distribution, which was found to be very soft ($q = 5.1$, Acciari et al. (2010)), while the typical values for this model fall between $q = 2.3$ and 2.7.

Each of these blazars shows a synchrotron peak frequency characteristic of the ISP/HSP divide, namely $\nu_{synch} \sim 10^{15}$. Additionally, each of these BL Lacs exhibits a similar flux level within the high-energy gamma-ray band, showing comparable $L_{synch}$ and $L_{IC}$. This commonality of sub-type and peak luminosities is likely a bias introduced to the selection process through the sensitivity of the Fermi LAT instrument being greater in the 300 MeV - 100 GeV band as compared to the $\geq 100$ GeV band. Nearby blazars releasing a large fraction of power above 100 GeV are known to be good candidates for detection in the VHE regime. This type of power emission is characteristic of HSP BL Lacs, the most commonly
VHE-detected type of blazar, but the least frequently detected BL Lac in the 1FGL and 2FGL catalogs.

Although the selection of these six blazars from high-energy Fermi-LAT data did not lead to any new TeV blazar discoveries, the exercise has been very illuminating. We are reminded that when selecting candidates for observation with TeV instruments, accurate redshift information is crucial. Without this information it is difficult to decipher if the lack of TeV emission is due to the intrinsic emission mechanism or the absorption of gamma rays by the EBL. The redshift lower limits that were found in the SDSS for RGBJ0909+231 and RGBJ1243+364 indicate distances where future TeV detection from a deeper exposure is unlikely. Continued TeV observations of RGB J0136+361, RGB J0316+090 and RXJ1436.9+5639, the three sources that remain without solid redshift information, could provide insight into the distance to these sources, while further observations of RBG J1058+564 can be directly applied to further investigate the emission mechanism at work within the blazar jet.

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Facilities: VERITAS, Fermi-LAT, Swift XRT and UVOT, SDSS, Keck, Lick

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Fig. 1.— Optical spectra for the six BL Lacs selected from the 1FGL catalog and observed with VERITAS. The black shows the object spectrum, while the red shows the instrumental noise. Only one BL Lac had a confirmed redshift upon selection (RGB J1058+564, $z=0.143$), confirmed with the SDSS spectrum shown in the middle-right panel. Redshift lower limits for RGB J0909+231 ($z \geq 0.4305$; middle left) and RGB J1243+364 ($z \geq 0.485$; lower left) are found from Mg II absorption lines in the SDSS spectra. A featureless SDSS spectrum is found for RX J1436.9+5639. A redshift measurement attempt for RGB J0136+391 using the Keck LRIS instrument shows a featureless power-law spectrum (upper left). The normalized ESI spectrum of RGB J0316+090 exhibits three unidentified absorption features (upper right). The spectral analysis of the Keck LRIS and ESI spectra are detailed in Kaplan et al. (2012).
Fig. 2.— *Fermi*-LAT light curves, with units of $10^{-9}$ ph cm$^{-2}$s$^{-1}$, are shown for the six candidate VHE-emitting BL Lacs for 29 months of LAT data (MJD 54682-55565; 2008 August 4-2011 January 4). The beginning and end of the VERITAS observations are denoted by vertical grey lines in each panel. The short VERITAS observation periods for RGB J0136+391 and RGB J0316+090 can be seen, representing only 7 days each. Upper limits at 95% confidence level are shown for bins resulting in TS of less than 9, denoted by downward pointing black triangles. Only RGB J0316+090 and RGB J1058+564 show any significant variability. For these two sources, LAT data only for the time periods within the window of VERITAS observations are used to constrain the modeling.
Fig. 3.— Non-contemporaneous broadband SED data for each BL Lac shown with corresponding SSC modeling using the model of Böttcher & Chiang (2002). The modeling results, corrected for the EBL absorption, are shown with grey lines. The archival radio data points are taken from NED and used only as upper limits, accounting for the fact that much of the radio emission may result from diffuse synchrotron emission in the radio lobes of the jet. See the text for a more detailed description of the model parameters, with values listed in Table 5.
Table 1. Summary of the high-energy *Fermi*-LAT power-law extrapolation beyond 150 GeV resulting from power-law fits from the 11 months of data used to select the candidates as well as the extrapolation based on the extended data set spanning 29 months. These extrapolated flux values factor in the gamma-ray absorption resulting from interaction with the EBL and are reported in % Crab Nebula flux units above the same energy threshold in order to allow direct comparison with the upper limits derived from the VERITAS observations.

<table>
<thead>
<tr>
<th>Counterpart Name</th>
<th>Original z Used for Selection</th>
<th>1FGL Index†</th>
<th>1FGL Integral Flux† (\geq 300 \text{ MeV} \times 10^{-9} \text{ ph cm}^{-2}\text{s}^{-1})</th>
<th>1FGL Extrapolated Flux (\geq 150 \text{ GeV} \times \text{% Crab})</th>
<th>Updated z</th>
<th>Updated Extrapolated Flux*</th>
<th>VERITAS UL</th>
</tr>
</thead>
<tbody>
<tr>
<td>RGB J0136+391</td>
<td>0.2</td>
<td>1.73</td>
<td>9.5</td>
<td>4.1</td>
<td>0.2</td>
<td>3.5</td>
<td>1.7</td>
</tr>
<tr>
<td>RGB J0316+090</td>
<td>0.2</td>
<td>1.72</td>
<td>19.6</td>
<td>8.1</td>
<td>0.2</td>
<td>0.3</td>
<td>2.0</td>
</tr>
<tr>
<td>RGB J0909+231</td>
<td>0.2</td>
<td>1.46</td>
<td>6.6</td>
<td>12.1</td>
<td>(\geq 0.4305)</td>
<td>0.5</td>
<td>1.5</td>
</tr>
<tr>
<td>RGB J1058+564</td>
<td>0.143</td>
<td>1.97</td>
<td>15.7</td>
<td>1.5</td>
<td>0.143</td>
<td>1.1</td>
<td>1.1</td>
</tr>
<tr>
<td>RGB J1243+364</td>
<td>0.2</td>
<td>1.74</td>
<td>5.8</td>
<td>2.2</td>
<td>(\geq 0.485)</td>
<td>0.8</td>
<td>3.1</td>
</tr>
<tr>
<td>RX J1436.9+5639</td>
<td>0.15</td>
<td>1.45</td>
<td>3.6</td>
<td>11.7</td>
<td>0.15</td>
<td>0.7</td>
<td>2.4</td>
</tr>
</tbody>
</table>

†Taken from Abdo et al. (2010b). Index and flux extrapolated without error for target selection.

*Computed with power-law fit from 29 months of data for steady sources and VERITAS coincident time window for variable sources (RGB J0316+090 and RGB J1058+564).
Table 2. Summary of VERITAS observations and analysis results. The significances are computed from counts extracted from source and background regions (ON and OFF, respectively) according to Equation 17 of Li & Ma (1983). The VHE 99% confidence level integral upper limits (ULs) are used for the modeling and are derived with an assumed photon index of 3.0. The percentage Crab Nebula flux values are given above the corresponding energy thresholds for each observation.

<table>
<thead>
<tr>
<th>Counterpart Name</th>
<th>$z$</th>
<th>VERITAS Livetime [hr]</th>
<th>Observation Window [MJD]</th>
<th>ON Events</th>
<th>OFF Events</th>
<th>$\alpha$</th>
<th>Significance $\sigma$</th>
<th>Threshold Energy [GeV]</th>
<th>Integral UL Above Threshold $\times 10^{-8}$ ph m$^{-2}$s$^{-1}$ (% Crab)</th>
<th>Decorrelation Energy [GeV]</th>
<th>Differential UL Energy at Decorrelation $\times 10^{-7}$ m$^{-2}$s$^{-1}$TeV$^{-1}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>RGB J0136+391</td>
<td></td>
<td>9.9</td>
<td>55122-55129</td>
<td>1422</td>
<td>11224</td>
<td>0.1277</td>
<td>-0.2</td>
<td>165</td>
<td>4.9 (1.7%)</td>
<td>260</td>
<td>1.6</td>
</tr>
<tr>
<td>RGB J0316+090</td>
<td></td>
<td>4.4</td>
<td>55126-55133</td>
<td>698</td>
<td>5560</td>
<td>0.1282</td>
<td>-0.5</td>
<td>165</td>
<td>5.7 (2.0%)</td>
<td>275</td>
<td>1.5</td>
</tr>
<tr>
<td>RGB J0909+231</td>
<td>≥0.4305</td>
<td>14.2</td>
<td>55150-55204</td>
<td>2141</td>
<td>14310</td>
<td>0.1499</td>
<td>-0.1</td>
<td>165</td>
<td>4.3 (1.5%)</td>
<td>260</td>
<td>1.3</td>
</tr>
<tr>
<td>RGB J1058+564</td>
<td>0.143</td>
<td>9.8</td>
<td>55160-55185</td>
<td>1415</td>
<td>10405</td>
<td>0.1417</td>
<td>-1.1</td>
<td>220</td>
<td>2.1 (1.1%)</td>
<td>350</td>
<td>4.7</td>
</tr>
<tr>
<td>RGB J1243+364</td>
<td>≥0.485</td>
<td>11.5</td>
<td>55189-55291</td>
<td>1617</td>
<td>11133</td>
<td>0.1399</td>
<td>0.9</td>
<td>150</td>
<td>9.9 (3.1%)</td>
<td>250</td>
<td>2.9</td>
</tr>
<tr>
<td>RX J1436.9+5639</td>
<td>0.15*</td>
<td>13.0</td>
<td>55275-55383</td>
<td>1563</td>
<td>12314</td>
<td>0.125</td>
<td>0.6</td>
<td>240</td>
<td>4.1 (2.4%)</td>
<td>375</td>
<td>8.9</td>
</tr>
</tbody>
</table>

1Defined as the energy at which the differential rate of reconstructed gamma rays from the postulated source reaches its maximum, accurate to within the 20% energy resolution of VERITAS.
2Normalization factor between size of the ON and OFF regions.
3The energy at which the upper limit is minimally dependent on the index assumed for the source.

*Redshift of spatially coincident galactic cluster, taken from Bauer et al. (2000).
Table 3. Summary of *Fermi*-LAT observations and analysis results. A variability study was completed using 29 months of data for all sources. For steady sources, the spectral analysis is completed for 29 months of data. The two variable sources RGB J0316+090 and RGB J1058+564 show spectral analysis results for LAT data coincident with a slightly expanded window as compared to the VERITAS observation window so as to allow a 5σ detection.

<table>
<thead>
<tr>
<th>Counterpart Name</th>
<th>1FGL Name</th>
<th>Variability Detected?</th>
<th>Prob*</th>
<th>MJD Fit Window</th>
<th>TS</th>
<th>Index</th>
<th>Integral Flux ≥ 300 MeV [×10⁻⁹ ph cm⁻² s⁻¹]</th>
</tr>
</thead>
<tbody>
<tr>
<td>RGB J0136+391</td>
<td>J0136.5+3905</td>
<td>no</td>
<td>0.39</td>
<td>54682.7-55565.0</td>
<td>1430</td>
<td>1.78±0.04</td>
<td>12.0±0.7</td>
</tr>
<tr>
<td>RGB J0316+090</td>
<td>J0316.1+0904</td>
<td>yes</td>
<td>0.00</td>
<td>55055.0-55145.0</td>
<td>45.7</td>
<td>2.09±0.26</td>
<td>9.5±3.3</td>
</tr>
<tr>
<td>RGB J0909+231</td>
<td>J0909.2+2310</td>
<td>no</td>
<td>0.06</td>
<td>54682.7-55565.0</td>
<td>123</td>
<td>1.68±0.13</td>
<td>1.8±0.5</td>
</tr>
<tr>
<td>RGB J1058+564</td>
<td>J1058.6+5628</td>
<td>yes</td>
<td>0.00</td>
<td>55160.0-55185.0</td>
<td>37.2</td>
<td>1.98±0.24</td>
<td>12.4±3.8</td>
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<tr>
<td>RGB J1243+364</td>
<td>J1243.1+3627</td>
<td>no</td>
<td>0.60</td>
<td>54682.7-55565.0</td>
<td>627</td>
<td>1.76±0.06</td>
<td>5.6±0.5</td>
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<tr>
<td>RX J1436.9+5630</td>
<td>J1437.0+5640</td>
<td>no</td>
<td>0.24</td>
<td>54682.7-55565.0</td>
<td>132</td>
<td>1.74±0.13</td>
<td>1.6±0.4</td>
</tr>
</tbody>
</table>

*Prob is the probability of steady emission as measured from ΔTS per degree of freedom.
Table 4. Summary of Swift-XRT observations and analysis results. Photon counting mode data were fit with absorbed power laws using HI column densities from Kalberla et al. (2005). Unabsorbed data were used for SED modeling. For multiple observations showing no significant variability, exposures were combined to improve statistics. If variability was detected between observations, both results are shown on the SED, although the model is only shown for the XRT observation falling closest to the window of VERITAS observations.

<table>
<thead>
<tr>
<th>Counterpart Name</th>
<th>Observation ID</th>
<th>MJD</th>
<th>Exposure Time [ks]</th>
<th>HI Flux [×10^{20} cm^{-2}]</th>
<th>Flux (2-10 keV) [×10^{-12} ergs cm^{-2} s^{-1}]</th>
<th>Photon Index</th>
<th>( \chi^2 ) (DOF)</th>
<th>Used in SED</th>
</tr>
</thead>
<tbody>
<tr>
<td>RGBJ0136+391</td>
<td>00039107001</td>
<td>55134</td>
<td>1.6</td>
<td>6.17</td>
<td>12.1±1.0</td>
<td>2.24±0.05</td>
<td>38.0 (40)</td>
<td>yes</td>
</tr>
<tr>
<td></td>
<td>00039107002</td>
<td>55223</td>
<td>2.9</td>
<td>6.17</td>
<td>21.4±1.0</td>
<td>2.16±0.03</td>
<td>107.2 (84)</td>
<td>yes</td>
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<tr>
<td>RGBJ0316+090</td>
<td>00038370001</td>
<td>54899</td>
<td>2.6</td>
<td>12.6</td>
<td>7.0±0.6</td>
<td>2.13±0.06</td>
<td>34.7 (35)</td>
<td>yes</td>
</tr>
<tr>
<td>RGBJ0909+231*</td>
<td>00040540001</td>
<td>55338</td>
<td>1.3</td>
<td>--</td>
<td>--</td>
<td>--</td>
<td>--</td>
<td>no</td>
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<td></td>
<td>00040540002</td>
<td>55338</td>
<td>1.9</td>
<td>--</td>
<td>--</td>
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<td>--</td>
<td>no</td>
</tr>
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<td>55342</td>
<td>1.6</td>
<td>--</td>
<td>--</td>
<td>--</td>
<td>--</td>
<td>no</td>
</tr>
<tr>
<td></td>
<td>combined</td>
<td>--</td>
<td>--</td>
<td>4.63</td>
<td>0.3±0.06</td>
<td>2.7±0.2</td>
<td>5.1 (5)</td>
<td>yes</td>
</tr>
<tr>
<td>RGBJ1058+564</td>
<td>00038215001</td>
<td>54852</td>
<td>3.8</td>
<td>0.78</td>
<td>2.2±0.2</td>
<td>2.60±0.05</td>
<td>42.4 (39)</td>
<td>yes</td>
</tr>
<tr>
<td></td>
<td>00038350001</td>
<td>54888</td>
<td>1.0</td>
<td>0.78</td>
<td>6.3±0.6</td>
<td>2.48±0.07</td>
<td>25.9 (27)</td>
<td>yes</td>
</tr>
<tr>
<td>RGBJ1243+364</td>
<td>00038445001</td>
<td>54875</td>
<td>2.1</td>
<td>1.4</td>
<td>8.2±0.5</td>
<td>2.33±0.05</td>
<td>50.4 (46)</td>
<td>yes</td>
</tr>
<tr>
<td>RXJ1436.9+5639</td>
<td>00038435001</td>
<td>54918</td>
<td>5.1</td>
<td>1.55</td>
<td>2.4±0.2</td>
<td>2.31±0.05</td>
<td>40.6 (31)</td>
<td>no</td>
</tr>
<tr>
<td></td>
<td>00038289001</td>
<td>54918</td>
<td>4.9</td>
<td>1.55</td>
<td>2.4±0.2</td>
<td>2.27±0.05</td>
<td>43.3 (31)</td>
<td>no</td>
</tr>
<tr>
<td></td>
<td>combined</td>
<td>--</td>
<td>--</td>
<td>1.55</td>
<td>2.4±0.2</td>
<td>2.29±0.04</td>
<td>72.5 (78)</td>
<td>yes</td>
</tr>
</tbody>
</table>

*Statistics extracted from each single observation are too low to fit an absorbed power law. The combination of the three exposures is used.
Table 5. Summary of Swift-UVOT observations and analysis results.

<table>
<thead>
<tr>
<th>Target ID</th>
<th>Observation ID</th>
<th>Date [MJD]</th>
<th>Band</th>
<th>Frequency [Hz]</th>
<th>$\nu F_\nu$ [Jy Hz]</th>
<th>Used in SED?</th>
</tr>
</thead>
<tbody>
<tr>
<td>RGB J0136+391</td>
<td>00039107001</td>
<td>55134</td>
<td>UVW1</td>
<td>$1.14 \times 10^{15}$</td>
<td>$(1.16 \pm 0.15) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00039107001</td>
<td>55134</td>
<td>UVM2</td>
<td>$1.34 \times 10^{15}$</td>
<td>$(1.44 \pm 0.30) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td>RGB J0316+090</td>
<td>00038370001</td>
<td>54899</td>
<td>V</td>
<td>$5.55 \times 10^{14}$</td>
<td>$(2.24 \pm 1.04) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038370001</td>
<td>54899</td>
<td>U</td>
<td>$8.57 \times 10^{14}$</td>
<td>$(2.92 \pm 3.87) \times 10^{12}$</td>
<td>n*</td>
</tr>
<tr>
<td></td>
<td>00038370001</td>
<td>54899</td>
<td>UVW1</td>
<td>$1.14 \times 10^{15}$</td>
<td>$(2.44 \pm 7.64) \times 10^{12}$</td>
<td>n*</td>
</tr>
<tr>
<td></td>
<td>00038370001</td>
<td>54899</td>
<td>UWW2</td>
<td>$1.48 \times 10^{15}$</td>
<td>$(2.8 \pm 18.5) \times 10^{12}$</td>
<td>n*</td>
</tr>
<tr>
<td>RGB J0909+231</td>
<td>00040540001</td>
<td>55338</td>
<td>V</td>
<td>$5.55 \times 10^{14}$</td>
<td>$(2.35 \pm 0.11) \times 10^{11}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00040540001</td>
<td>55338</td>
<td>B</td>
<td>$6.93 \times 10^{14}$</td>
<td>$(1.77 \pm 0.23) \times 10^{11}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00040540001</td>
<td>55338</td>
<td>U</td>
<td>$8.57 \times 10^{14}$</td>
<td>$(2.26 \pm 0.21) \times 10^{11}$</td>
<td>y</td>
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<tr>
<td></td>
<td>00040540001</td>
<td>55338</td>
<td>UVW1</td>
<td>$1.14 \times 10^{15}$</td>
<td>$(1.78 \pm 0.18) \times 10^{11}$</td>
<td>y</td>
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<tr>
<td></td>
<td>00040540001</td>
<td>55338</td>
<td>UVM2</td>
<td>$1.34 \times 10^{15}$</td>
<td>$(1.63 \pm 0.22) \times 10^{11}$</td>
<td>y</td>
</tr>
<tr>
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<td>55338</td>
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<td>55338</td>
<td>UVM2</td>
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<td>$(1.90 \pm 0.22) \times 10^{11}$</td>
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<td>55338</td>
<td>UWW2</td>
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<tr>
<td>RGB J1058+564</td>
<td>00038215001</td>
<td>54852</td>
<td>V</td>
<td>$5.55 \times 10^{14}$</td>
<td>$(1.42 \pm 0.04) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038215001</td>
<td>54852</td>
<td>U</td>
<td>$8.57 \times 10^{14}$</td>
<td>$(1.45 \pm 0.04) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038215001</td>
<td>54852</td>
<td>UVW2</td>
<td>$1.48 \times 10^{15}$</td>
<td>$(1.49 \pm 0.05) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038453001</td>
<td>54888</td>
<td>B</td>
<td>$6.93 \times 10^{14}$</td>
<td>$(1.64 \pm 0.07) \times 10^{12}$</td>
<td>n</td>
</tr>
<tr>
<td></td>
<td>00038453001</td>
<td>54888</td>
<td>U</td>
<td>$8.57 \times 10^{14}$</td>
<td>$(1.70 \pm 0.06) \times 10^{12}$</td>
<td>n</td>
</tr>
<tr>
<td></td>
<td>00038453001</td>
<td>54888</td>
<td>UVW1</td>
<td>$1.14 \times 10^{15}$</td>
<td>$(1.66 \pm 0.06) \times 10^{12}$</td>
<td>n</td>
</tr>
<tr>
<td>RGB J1243+364</td>
<td>00038445001</td>
<td>54875</td>
<td>V</td>
<td>$5.55 \times 10^{14}$</td>
<td>$(1.28 \pm 0.05) \times 10^{12}$</td>
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</tr>
<tr>
<td></td>
<td>00038445001</td>
<td>54875</td>
<td>B</td>
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<td>$(1.30 \pm 0.04) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038445001</td>
<td>54875</td>
<td>U</td>
<td>$8.57 \times 10^{14}$</td>
<td>$(1.42 \pm 0.05) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038445001</td>
<td>54875</td>
<td>UVW1</td>
<td>$1.14 \times 10^{15}$</td>
<td>$(1.34 \pm 0.05) \times 10^{12}$</td>
<td>y</td>
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<tr>
<td></td>
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<td>54875</td>
<td>UVM2</td>
<td>$1.34 \times 10^{15}$</td>
<td>$(1.49 \pm 0.05) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038445001</td>
<td>54875</td>
<td>UWW2</td>
<td>$1.48 \times 10^{15}$</td>
<td>$(1.47 \pm 0.05) \times 10^{12}$</td>
<td>y</td>
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<tr>
<td>RX J1436.9+5639</td>
<td>00038435001</td>
<td>54918</td>
<td>V</td>
<td>$5.55 \times 10^{14}$</td>
<td>$(9.03 \pm 0.09) \times 10^{11}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
<td>00038435001</td>
<td>54918</td>
<td>B</td>
<td>$6.93 \times 10^{14}$</td>
<td>$(1.14 \pm 0.07) \times 10^{12}$</td>
<td>y</td>
</tr>
<tr>
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<td>54918</td>
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<td>$(9.71 \pm 0.05) \times 10^{11}$</td>
<td>y</td>
</tr>
<tr>
<td></td>
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<td>54918</td>
<td>UVW1</td>
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<td>$(1.06 \pm 0.06) \times 10^{12}$</td>
<td>y</td>
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<td>54918</td>
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<td>$(1.21 \pm 0.07) \times 10^{12}$</td>
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<td>54918</td>
<td>UWW2</td>
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<td>$(1.12 \pm 0.06) \times 10^{12}$</td>
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<td>54918</td>
<td>V</td>
<td>$5.55 \times 10^{14}$</td>
<td>$(8.58 \pm 0.05) \times 10^{11}$</td>
<td>n</td>
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<td></td>
<td>00038289001</td>
<td>54918</td>
<td>U</td>
<td>$8.57 \times 10^{14}$</td>
<td>$(1.06 \pm 0.04) \times 10^{12}$</td>
<td>n</td>
</tr>
<tr>
<td></td>
<td>00038289001</td>
<td>54918</td>
<td>UVW2</td>
<td>$1.48 \times 10^{15}$</td>
<td>$(1.26 \pm 0.06) \times 10^{12}$</td>
<td>n</td>
</tr>
</tbody>
</table>

*Galactic reddening dominates the uncertainty with $E(B-V) = 0.356$ for RGB J0316+090; U, UVW1, UVW2 band flux measurements do not constrain the SED.
Table 6. Summary of the broadband SED SSC modeling parameters. See text for parameter descriptions.

<table>
<thead>
<tr>
<th>Counterpart Name</th>
<th>z</th>
<th>$\gamma_{\text{min}}$</th>
<th>$\gamma_{\text{max}}$</th>
<th>$q$</th>
<th>$t_{\text{esc}}$ [hr]</th>
<th>B [Gauss]</th>
<th>$\Gamma$</th>
<th>Radius [cm]</th>
<th>$L_e$ [ergs/sec]</th>
<th>$L_B$ [ergs/sec]</th>
<th>$L_B/L_e$</th>
<th>$t_{\text{min}}$ [hr]</th>
</tr>
</thead>
<tbody>
<tr>
<td>RGB J0136+391</td>
<td>0.2</td>
<td>$4.0 \times 10^4$</td>
<td>$6.0 \times 10^5$</td>
<td>2.6</td>
<td>120</td>
<td>0.15</td>
<td>15</td>
<td>$3.4 \times 10^{16}$</td>
<td>$1.1 \times 10^{44}$</td>
<td>$2.3 \times 10^{43}$</td>
<td>$2.1 \times 10^{-1}$</td>
<td>25.9</td>
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<td>$5.0 \times 10^5$</td>
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<td>50</td>
<td>0.4</td>
<td>15</td>
<td>$2.4 \times 10^{16}$</td>
<td>$1.0 \times 10^{44}$</td>
<td>$7.8 \times 10^{43}$</td>
<td>$7.5 \times 10^{-1}$</td>
<td>19.2</td>
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<td>15</td>
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<td>$1.3 \times 10^{44}$</td>
<td>$1.5 \times 10^{44}$</td>
<td>1.1</td>
<td>25.9</td>
</tr>
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<td>$7.0 \times 10^5$</td>
<td>2.5</td>
<td>15</td>
<td>0.11</td>
<td>15</td>
<td>$8.0 \times 10^{16}$</td>
<td>$9.8 \times 10^{43}$</td>
<td>$6.5 \times 10^{43}$</td>
<td>$6.6 \times 10^{-1}$</td>
<td>59.2</td>
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<tr>
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<td>2.8</td>
<td>10</td>
<td>0.1</td>
<td>20</td>
<td>$2.5 \times 10^{16}$</td>
<td>$1.5 \times 10^{44}$</td>
<td>$9.4 \times 10^{42}$</td>
<td>$6.4 \times 10^{-2}$</td>
<td>17.4</td>
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<tr>
<td>RGB J1058+564</td>
<td>0.1433</td>
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<td>$3.0 \times 10^5$</td>
<td>2.9</td>
<td>10</td>
<td>0.25</td>
<td>20</td>
<td>$9.0 \times 10^{15}$</td>
<td>$3.9 \times 10^{43}$</td>
<td>$7.6 \times 10^{42}$</td>
<td>$2.0 \times 10^{-1}$</td>
<td>4.8</td>
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<tr>
<td>RGB J1243+364</td>
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<td>$3.0 \times 10^4$</td>
<td>$1.0 \times 10^6$</td>
<td>2.5</td>
<td>30</td>
<td>0.1</td>
<td>20</td>
<td>$1.5 \times 10^{17}$</td>
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<td>$3.4 \times 10^{44}$</td>
<td>1.1</td>
<td>104.1</td>
</tr>
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<td>RX J1436.9+5639</td>
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<td>2.6</td>
<td>100</td>
<td>0.01</td>
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<td>$1.1 \times 10^{43}$</td>
<td>$7.5 \times 10^{-2}$</td>
<td>143.7</td>
</tr>
</tbody>
</table>

* Assumed redshift values due to lack of any redshift constraining measurement.

$^a$ Parameters adjusted to match the model to the data.

$^b$ Quantities derived from the modeling result.