Global Comparison of CFD and Wind-Tunnel-Derived Force and Moment Databases for the Space Launch System

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Recently a very large (739 runs) collection of high-fidelity RANS CFD solutions was obtained for Space Launch System ascent aerodynamics for the vehicle to be used for the first exploratory (unmanned) mission (EM-1). The extensive computations, at full-scale conditions, were originally developed to obtain detailed line and protuberance loads and surface pressures for venting analyses. The line loads were eventually integrated for comparison of the resulting forces and moments to the database that was derived from wind-tunnel tests conducted at sub-scale conditions. The comparisons presented herein cover the ranges \(0.5 \leq M_\infty \leq 5\), \(-6^\circ \leq \alpha \leq 6^\circ\), and \(-6^\circ \leq \beta \leq 6^\circ\). For detailed comparisons, slender-body-theory-based component build-up aero models from missile aerodynamics are used. The differences in the model fit coefficients are shown to be relatively small except for the low supersonic Mach number range, \(1.1 \leq M_\infty \leq 2.0\). The analysis is intended to support process improvement and development of uncertainty models.

Nomenclature

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Description</th>
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</thead>
<tbody>
<tr>
<td>BMC</td>
<td>balance moment center (also reference moment center and approximate location of the vehicle center of gravity)</td>
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<tr>
<td>DB</td>
<td>database</td>
</tr>
<tr>
<td>(C_m)</td>
<td>pitching-moment coefficient about the BMC</td>
</tr>
<tr>
<td>(C_{m_0})</td>
<td>zero intercept of the linear fit to the (C_m) versus (C_N) data</td>
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<tr>
<td>(C_N)</td>
<td>normal-force coefficient</td>
</tr>
<tr>
<td>(C_{N_0})</td>
<td>zero intercept of the linear fit to the (C_N) versus (\alpha) data</td>
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<tr>
<td>(C_{n})</td>
<td>yawing-moment coefficient about the BMC</td>
</tr>
<tr>
<td>(C_{n_0})</td>
<td>zero intercept of the linear fit to the (C_n) versus (C_Y) data</td>
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<tr>
<td>CP</td>
<td>center of pressure</td>
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<tr>
<td>(C_Y)</td>
<td>side-force coefficient</td>
</tr>
<tr>
<td>(C_{Y_0})</td>
<td>zero intercept of the linear fit to the (C_Y) versus (\beta) data</td>
</tr>
<tr>
<td>(M_\infty)</td>
<td>Mach number</td>
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<tr>
<td>RANS</td>
<td>Reynolds-Averaged Navier-Stokes</td>
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<tr>
<td>SBT</td>
<td>slender-body theory</td>
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<tr>
<td>SLS</td>
<td>Space Launch System</td>
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<tr>
<td>WT</td>
<td>wind tunnel</td>
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<tr>
<td>WTDB</td>
<td>wind-tunnel-derived database</td>
</tr>
<tr>
<td>(\hat{x}_{C_x})</td>
<td>fit to the (C_x) axial location of the CP, normalized by the reference length (core diameter)</td>
</tr>
<tr>
<td>(\hat{x}_{C_Y})</td>
<td>fit to the (C_Y) axial location of the CP, normalized by the reference length (core diameter)</td>
</tr>
<tr>
<td>(\alpha)</td>
<td>angle of attack in body coordinates</td>
</tr>
<tr>
<td>(\beta)</td>
<td>sideslip angle in body coordinates</td>
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</table>

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Introduction

The NASA Space Launch System (SLS) is intended to be the United States very-heavy launch vehicle for the foreseeable future (see Fig. 1). The first vehicle in the series, the Block 1 SLS-10003, will be used for the launch of the unmanned Orion space capsule around the moon and back (EM-1). For the purposes of various customers, the SLS Aerodynamics Task Team conducted wind tunnel tests in order to derive a comprehensive force and moment (F&M) database for ascent aerodynamics. Those subscale tests and subsequent development of the WT database are discussed in Pinier et al.\(^1\) To complement the force and moment WTDB, a comprehensive CFD campaign was conducted to provide line and protuberance loads and venting data. The CFD campaign obtained very high fidelity Reynolds-Averaged-Navier-Stokes solutions at 739 ascent conditions (full-scale) in the ranges \(0.5 \leq M_\infty \leq 5\), \(-6^\circ \leq \alpha \leq 6^\circ\), and \(-6^\circ \leq \beta \leq 6^\circ\). The overset grid system consisted of 375 million grid points. The details of the CFD campaign are presented in Rogers et al.\(^6\).

It is unusual in the uncertainty community to have such a large set of CFD solutions for comparison with experimental (wind tunnel) data. Of course, the CFD solutions were obtained at full-scale flight conditions while the wind tunnel database was derived from subscale tests, albeit with the model gritted for boundary layer transition. Nevertheless, comparing the CFD with the WTDB for quality control checks and for potential uncertainty modeling seems reasonable and useful. The following four sections present the CFD-to-WTDB comparisons for the normal-force coefficient, \(C_N\), the side-force coefficient, \(C_Y\), the pitching-moment coefficient, \(C_m\), and the yawing-moment coefficient, \(C_n\), respectively. The axial-force and rolling-moment coefficients are not considered in the paper.

The CFD and WTDB results can, of course, be compared by simple subtraction and those results are given below. However, it is possible to do more detailed comparisons using modeling appropriate for the given conditions. For that purpose, slender-body-theory-based component build-up aero models\(^3,5\) from missile aerodynamics are selected. For simplicity, the models will be designated by the acronym, SBT, for the rest of the paper. It should be noted from the references that SBT can be applied to any airframe provided that the airflow is sufficiently "slender".

Normal-Force Coefficient

Simply subtracting the WTDB \(C_N\) results from the CFD \(C_N\) results for all of the conditions in the CFD 739-run matrix gives the plots shown in Fig. 2. Fig. 2(a) shows the differences between the two types of angle-of-attack polars at the breakpoint conditions as a function of \(M_\infty\), revealing that the lower supersonic conditions have almost double the scatter of the rest of the differences. Fig. 2(b) shows the effect of angle of attack on the scatter of the differences. The Mach number breakpoints of the CFD database are visible as the abscissa locations of the columns of symbols in Fig. 2(a). Note that it is difficult to discern much of a pattern in Fig. 2. To be able to discern any patterns associated with the three independent variables of \(M_\infty\), \(\alpha\), and \(\beta\), it is necessary to model the dependent variables as functions of the independent variables. The rest of the paper will describe the modeling choices and the results.

It has long been known that first-order SBT\(^3,4,6,7\) predicts \(C_N\) to be a linear function of \(\alpha\) only, independent of \(\beta\) and \(M_\infty\). Wind tunnel testing of slender airframes over many years has shown, however, that \(C_N\) is a weak function of \(M_\infty\). Given that information, the following linear equation is used to model the \(C_N\) dependence with each \(M_\infty\) treated separately. Least squares is used to choose the coefficients:

\[
C_N = C_{N_0} + C_{N_\alpha} \alpha
\]  

Typical examples of the fits are given in Fig. 3 for \(M_\infty = 0.5, 1.4, 3.0\). Since the data are sensitive, it was necessary to omit the ordinate scale for general publication. The results of Fig. 3(a) are typical of the subsonic cases, while the results of Fig. 3(b) are typical of the low supersonic cases. All of the subsonic, transonic and low supersonic cases have very little deviation from linearity for the angle of attack range of \(-4^\circ \leq \alpha \leq 4^\circ\), the range for which the ascent trajectories are designed. For the higher supersonic cases, Fig. 3(c) is a more typical example. It shows some nonlinearity but not enough to make it unusable for comparison purposes. Note that many of the conditions in Fig. 3...
have combined angles of attack and sideslip, thus making it clear that there is no discernable dependence of \( C_N \) on \( \beta \) as predicted by SBT.

The \( C_f \) fit coefficients are compared in Fig. 4. Fig. 4(a) gives the ratio of the slopes, \( C_{N_f} \), showing remarkable agreement for \( M_\infty \leq 1.1 \) and \( M_\infty \geq 2.5 \), but rather disappointing agreement for the Mach numbers in between. The possibility of wall interference in that Mach region is discussed in Reference 1. But the cause of the lack of agreement has not yet been definitively determined.

Fig. 4(b) gives the differences in the intercept values, showing good agreement for the subsonic and transonic Mach numbers and rather remarkable agreement for \( M_\infty \geq 1.3 \).

### Side-Force Coefficient

The CFD minus WTDB results for \( C_f \) are given in Fig. 5 to the same scale used in Fig. 2. The scatter of the differences is about 1/3 of that observed in Fig. 2 for \( C_N \). This is not too surprising since the side force derivatives with respect to sideslip are roughly 1/3 of those for the normal force with respect to angle of attack. Note in Fig. 1 that the sideslip flow sees a considerably smaller vehicle planform than that of the pitch flow, which sees the core and both SRB planforms. Following the approach used for \( C_N \), the side force coefficient dependence on \( \beta \) is modeled as

\[
C_f = C_{f_0} + C_{f_\beta} \beta \tag{2}
\]

Typical examples of the fits are given in Fig. 6 for \( M_\infty = 0.5, 1.4, 3.0 \), using the same (hidden) scale as Fig. 3. Again, all of the subsonic, transonic and low supersonic cases have very little deviation from linearity for the angle of attack range of \(-4^\circ \leq \alpha \leq 4^\circ \). For the higher supersonic cases, Fig. 5(c) shows some very slight nonlinearity. Also, note again that many of the conditions in Fig. 5 have combined angles of attack and sideslip, thus making it clear that there is no discernable dependence of \( C_f \) on \( \alpha \), as predicted by SBT.

The \( C_f \) fit coefficients are compared in Fig. 7 using the same scales as Fig. 4. Fig. 7(a) gives the ratio of the slopes, \( C_{t_f} \), showing remarkable agreement except for the transonic region, \( 0.9 \leq M_\infty \leq 1.1 \). The possibility of wall interference affecting the WTDB results in this Mach range was discussed in Reference 1. Fig. 7(b) shows that the differences in the fit intercepts are small.

### Pitching-Moment Coefficient

The CFD minus WTDB results for \( C_m \) are given in Fig. 8. The differences are largest for the subsonic and low supersonic Mach numbers. However, in addition to predicting that \( C_m \) is linear in \( \alpha \), SBT predicts that the center of pressure for \( C_m \) is independent of all three parameters, \( \alpha, \beta, M_\infty \). Hence, SBT would predict \( C_m \) to be linear with respect to \( C_N \). Consequently, the following model is used to fit the dependence of \( C_m \) on \( C_N \):

\[
C_m = C_{m_0} + \left( \hat{x}_{m_0} - \hat{x}_{m_{BMC}} \right) C_N \tag{3}
\]

It should be noted that the reference center for the moment coefficients is the balance moment center (BMC) used for the wind tunnel tests \(^1\). The BMC is shown in Figure 1 and is located roughly at the location of the force centers of pressure (CP) throughout the ascent trajectory. Hence, the moments are not expected to be large and any movement of the CP, however slight, can contribute to any nonlinearity observed.

Typical examples of the resulting linear fits are given in Figure 9. The abscissa and ordinate values are omitted to enable publication of the data. The fits are certainly reasonable for comparing the CFD to the WTDB, but nonlinearities are clearly present as well as much more scatter than was seen in the fits for the force coefficients. This is most likely due to the center of pressure being close to the reference center. The fit coefficients are shown in

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Figure 10(a) for the discrepancy in \( \hat{x}_{\gamma} \) and in Figure 10(b) for the discrepancy in \( C_{n_0} \). Except for \( 0.7 \leq M_{\infty} \leq 0.9 \), the \( \hat{x}_{\gamma} \) discrepancies are quite small.

Yawing-Moment Coefficient

The CFD minus WTDB results for \( C_{\gamma} \) are given in Fig. 11 using the scales of Fig. 8. The differences are larger for the higher values of sideslip. Since SBT predicts results for \( C_{\gamma} \) that are similar to those for \( C_{n_0} \), the following model is used to fit the dependence of \( C_{\gamma} \) on \( C_{\gamma} \):

\[
C_{\gamma} = C_{n_0} + \left( \hat{x}_{\gamma} - \hat{x}_{\gamma,\text{MC}} \right) C_{\gamma}
\]

Typical examples of the resulting linear fits are given in Fig. 12. The fits are certainly reasonable for comparing the CFD to the WTDB. Nonlinearities are clearly present but with less scatter than was seen in the fits for the \( C_{n_0} \) coefficients. This is most likely due to the center of pressure being somewhat farther forward of the reference center compared to the \( C_{n_0} \) behavior, as would be expected for the different planforms for each flow. The fit coefficients are shown in Fig. 13(a) for the discrepancy in \( \hat{x}_{\gamma} \) and in Fig. 13(b) for the discrepancy in \( C_{n_0} \). Except for subsonic Mach numbers, the \( \hat{x}_{\gamma} \) discrepancy is larger than for \( \hat{x}_{\gamma} \).

Concluding Remarks

In the opinion of the author, the modeling presented above shows that SBT can be useful for generalizing and comparing the computational and experimental aerodynamic behavior of slender vehicles such as the SLS for conditions during ascent where the angles of attack and sideslip are small. Although it is generally realized that the aerodynamic behavior of slender airplane configurations follows the models used above, it seems less well known that slender airframes, including launch vehicles with boosters and/or fins, behave similarly. In fact, tactical rocket aero design prediction methods based on such behavior have been in use for over 60 years.\(^5\)

Using the SBT models, it is shown that the very-high-fidelity RANS results of Reference 2 compare reasonably well with the wind-tunnel-derived database\(^1\) for both the normal and side forces and the locations of their axial centers of pressure even though the Reynolds numbers are quite different (two orders of magnitude). Reference 2 actually shows that the agreement is within the reported WTDB uncertainties.

Finally, it is expected that additional efforts to locate the regions and causes of the larger discrepancies will help improve future database development, including development of uncertainty models.

Acknowledgements

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References

Figure 1. Plan (left) and side (right) views of the SLS 10003 configuration.
Figure 2. $C_N$ from CFD minus $C_N$ from WTDB for all $\alpha, \beta, M_\infty$. 

(a) Effect of Mach number

(b) Effect of Angle of Attack
Figure 3. Examples of least-squares linear fits to $C_n$ as a function of $\alpha$ for all $\beta$.

(a) $M_\infty = 0.5$

(b) $M_\infty = 1.4$
Figure 3. Concluded.

(c) $M_\infty = 3.0$
Figure 4. Comparison of $C_\alpha$ fit coefficients.

(a) Ratio of slopes

(b) Differences in intercept values
Figure 5. $C_Y$ from CFD minus $C_Y$ from WTDB for all $\alpha$ and $\beta$.

(a) Effect of Mach number

(b) Effect of Angle of Sideslip
Figure 6. Examples of least-squares linear fits to $C_y$ as a function of $\alpha$ for all $\beta$.
Figure 6. Concluded.

(c) $M_{\infty} = 3.0$
Figure 7. Comparison of $C_Y$ fit coefficients.

(a) Ratio of slopes

(b) Differences in intercept values
Figure 8. $C_m$ from CFD minus $C_m$ from WTDB for all $\alpha$ and $\beta$.

(a) Effect of Mach number

(b) Effect of Angle of Attack
Figure 9. Examples of least-squares linear fits to $C_m$ as a function of $C_y$ for all $C_y$. 

(a) $M_\infty = 0.5$

(b) $M_\infty = 1.4$
Figure 9. Concluded.

(c) $M_\infty = 3$

Figure 9. Concluded.
Figure 10. Comparison of fit coefficients for $C_m$ versus $C_y$.

(a) Number of core diameters that CFD CP is forward of the WTDB CP

(b) Differences in intercept values
Figure 11. $C_n$ from CFD minus $C_n$ from WTDB for all $\alpha$ and $\beta$.
Figure 12. Examples of least-squares linear fits to $C_n$ as a function of $C_Y$ for all $C_N$. 

(a) $M_\infty = 0.5$

(b) $M_\infty = 1.4$
Figure 12. Concluded.
Figure 13. Comparison of fit coefficients for $C_n$ versus $C_y$.

(a) Number of core diameters that CFD CP is forward of the WTDB CP

(b) Differences in intercept values