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Fundamental Study into Rotor Outwash and Dust Kick-up under Mars-like Conditions

TESIS

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FÍSICA

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Abstract

Humans have reached a limit to how much information can be collected from a purely ground-based rover exploration mission on Mars [1]. Therefore, an airborne mission would broaden our knowledge of this planet and at the same time it would allow for a different approach to scientific research. Thus, the Mars Helicopter Scout (MHS) Project was founded. However, as this is only a technology demonstrator, its first and short-term goal is to fly in the vicinity of the Mars Rover and assess where the rover can go. This thesis project attempts to take a first step towards understanding how and if the MHS’ coaxial rotor system may produce dust kick-up and how the in and out of ground effect contribute to a possible “brownout”. Brownout here on Earth means the complete or partial blockage of the pilot’s vision caused by the lifting of dust particles found on the ground. In our case, we refer to blockage or miscalculations of the sensors, such as the altimeter of the MHS, caused by the kick-up of Martian sand. This thesis further consists on the study of rotor downwash/outwash.
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Chapter 1

Introduction

This thesis project is to be thought of as a first approximation to a very complex problem. The dust cloud generation during take-off and landing of a rotorcraft, that may interfere with navigation systems (brownout). The landing of a helicopter here on Earth can become dangerous to the people on-board when trying to land over a granulated surface, i.e. snow or sand. Thus, this problem caught the attention of NASA AMES Aeromechanics branch, as shown in the paper “Rotorcraft Downwash Flow Field Study to Understand the Aerodynamics of Helicopter Brownout” written by Alan J Wadcock and co-workers [2]. Likewise, designing a vehicle capable of multiple take-off and landings in a different planet has also become one of NASA’s main goals. The Mars Helicopter Scout project was thus proposed to be the first vehicle to accomplish this. However, given it is still a rotorcraft vehicle, the threat of a brownout during take-off and landing must be addressed. It is unknown to what extent this phenomenon occur and how it can endanger the Mars helicopter.

The main objective of the project is conducting small-scale hover testing of a coaxial rotor system in and out of ground effect to study both rotor downwash/outwash and to examine the conditions under which dust kick-up and “brownout” does or does not occur under Mars-like conditions, with potential application to Mars Helicopter development. Given the exploratory nature of this project, it is limited to a proof of concept of testing techniques that may eventually be used on the full scale MHS. These future tests will help determine whether or not dust kick-up can pose a problem on the performance of the MHS’s altimeter.

The main limitation to this project will be the pressure at which both
the tuft testing and the dust testing will take place. Given this thesis is just a first approximation, it is believed to be sufficient to do testing at room pressure. Since it is harder to produce saltation at low air densities \[3\], it is believed that at Mars-like conditions, there will be minimum dust saltation. Therefore, if tests are done in conditions that are more prone to produce saltation and yet there is none, it can then be concluded that at lower densities, or conditions closer related to those in Mars, the phenomenon will not appear. However, in the case that it is proven that saltation is possible, further testing shall be necessary. The experimental set-up was designed for future testing inside a vacuum chamber where Mars’ low air densities can be reached \[4\].

1.1 The Mars Helicopter Scout

Mars is a very difficult terrain to move around in. Its dunes, rocks and valleys make it a challenge for our vehicles to move about without danger. That is why the current Mars Rover has many difficulties when it comes to making a decision on when and towards where it should move. It relays on the images taken by itself and previous missions that have photographed The Red Planet in order to allow the control station back on Earth to decide whether it is safe for it to keep moving. This is very limiting in nature, given the amount of time lost in communicating between Earth and Mars. This is the seed that lead to the inspiration of the vertical lift vehicle that will be the subject of this project.

The Mars Helicopter Scout is, up until now, a proof-of-concept technology demonstration of what later hopes to become an auxiliary vehicle for the Mars Rover. The short term purpose of this helicopter would be to explore the terrain ahead of the rover, for target selection, path selection, and geologic context. Although, for the safety of the rover, the helicopter will not land nor fly near it. Individual daily flights would be limited to 2–3 minutes maximum due to power consumption, but it should be able to fly approximately 600 m at an altitude of 100 m from the surface.

Besides there not being a set design for this rover, there are set physical features it shall include. These features are: a total mass of 1 kg and a 1.1 m rotor diameter, powered by solar charged batteries; a high resolution camera and a communication system to relay data to the rover. The current design consist of a co-axial main rotor and a cube-shaped 140 mm body with four
1.1. THE MARS HELICOPTER SCOUT

legs as landing gear [1, 5].

The differences between the Martian atmosphere and that of the Earth is what makes Mars such a harsh environment for flight. Since it is composed almost entirely out of CO$_2$, a density instability, or abrupt change, is caused by this molecule’s polar states, gaseous or solid, at pressures and temperatures commonly found in Mars (table 1.1). This poses as an extra difficulty when developing rotary-wing designs [6].

Moreover, as can be seen in Table 1.1, despite Mars’ gravity being a little bit over a third of Earth’s gravity, the extreme low temperatures and pressure make it very difficult to produce the necessary lift to carry the helicopter given its dependency on atmospheric density (equation 3.1).

Therefore, the autonomous Vertical Take-off and Landing (VTOL) aircraft will have to be extremely light weight and with large lifting surfaces. Thus, facing us with an extremely complex problem given the high Reynolds number and possible compressible flow aerodynamics (see table 1.2). The Reynolds number is defined as

$$Re = \frac{\rho u_{tip} R}{\mu},$$

(1.1)

where, for the purpose of this thesis, $\rho$ was the atmospheric density, $u_{tip}$ the tip velocity of the rotor, $R$ was the rotor radius and lastly, $\mu$ was the dynamic viscosity. On the other hand, if $u_s$ is the speed of sound, the Mach number is

$$Ma = \frac{u_{tip}}{u_s}.$$

(1.2)

Different designs have proposed rotor such as a tiltrotor, a rotor that is able to change its position from perpendicular to the ground to parallel;

<table>
<thead>
<tr>
<th></th>
<th>Mean Radius (km)</th>
<th>Gravity (m/s$^2$)</th>
<th>Mean Surface Atmos. Temp. (°K)</th>
<th>Mean Surface Atmos. Pressure (Pa)</th>
<th>Mean Surface Atmos. Density (kg/m$^3$)</th>
<th>Atmos. Gases</th>
</tr>
</thead>
<tbody>
<tr>
<td>Earth</td>
<td>6371</td>
<td>9.82</td>
<td>288.2</td>
<td>101,300</td>
<td>1.23</td>
<td>N$_2$ 78%</td>
</tr>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>O$_2$ 21%</td>
</tr>
<tr>
<td>Mars</td>
<td>3390</td>
<td>3.71</td>
<td>214</td>
<td>636</td>
<td>1.55 x 10$^{-2}$</td>
<td>CO$_2$ 95%</td>
</tr>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>N$_2$ 2.7%</td>
</tr>
<tr>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>Ar 1.6%</td>
</tr>
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<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
<td>O$_2$ 0.1 %</td>
</tr>
</tbody>
</table>

Table 1.1: Planetary Description [7].
quad-rotor, four rotors with different rotor centers; and coaxial rotors, two rotors that share the same rotor-center. The latter being the main focus of study of this project\cite{[6]}.

<table>
<thead>
<tr>
<th></th>
<th>Earth</th>
<th>Mars</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>Full-scale</td>
<td>Model</td>
</tr>
<tr>
<td>Re</td>
<td>$51.18 \times 10^5$</td>
<td>$7.07 \times 10^5$</td>
</tr>
<tr>
<td>Ma</td>
<td>0.4066</td>
<td>0.1106</td>
</tr>
</tbody>
</table>

Table 1.2: Calculations of Reynolds number and Mach number for both full-scale MHS and the model used in this thesis.
Chapter 2

Theoretical Background

Given the complexity of the proposed problem, initial background research had to be done in order to fully understand the problem. Basic knowledge of the Martian atmospheric conditions and the fundamentals of fluid dynamics, helicopter theory and computational fluid dynamics, were essential to the development of this thesis.

2.1 Fluid Dynamics Theory

Fluid mechanics governs the flight of helicopters and their interaction with its surroundings, like a dusty landing spot. Throughout this thesis, the fluid in question shall be the air surrounding the helicopter. Air is considered a Newtonian fluid, and so the flow rate of strain (the symmetrised velocity gradient) is linearly proportional to the stress tensor within the fluid. Let's use the Eulerian description of a fluid, so for example the fluid velocity field \( \mathbf{u}(\mathbf{x}, t) \) is interpreted as the velocity at a specific time \( t \) of an element of fluid that coincidently passed through the point in space \( \mathbf{x} \).

In the Eulerian description, mass conservation is represented by the continuity equation

\[
\frac{\partial \rho}{\partial t} = \nabla \cdot (\rho \mathbf{u}),
\]

where \( \rho \) is the density of air, that relates the local density change with the mass flux. The rate of change of the volume of the fluid elements is given by \( \nabla \cdot \mathbf{u} \). Therefore, a flow is considered incompressible when \( \nabla \cdot \mathbf{u} = 0 \). This is equivalent, by mass conservation (2.1), to considering the density of the
fluid elements to remain constant, expressed in the material derivative of the

density

\[ \frac{D\rho}{Dt} = \frac{\partial \rho}{\partial t} + \mathbf{u} \cdot \nabla \rho \]  

(2.2)

where in the case of an incompressible flow, \( \frac{D\rho}{Dt} = 0 \). If the flow of an initially
homogeneous fluid is incompressible, the density field remains constant in
time and space [9]. Therefore, despite air being a compressible fluid, the
air flow caused by the rotor will be considered an incompressible flow. This
is justified by the condition of incompressible flow that relates the speed at
which a solid moves through the fluid in hand, and the speed of sound in the
same fluid, i.e. the Mach number. In this specific case, the Mach number of
the blades is less than 0.3 thus it is safe to assume incompressibility [8].

On the other hand, the air is a Newtonian fluid and the stress tensor, for
an incompressible flow, can be expressed as

\[ \sigma(x, t) = -pI + 2\mu e, \]  

(2.3)

where \( e \) is the rate of strain tensor \( e = \frac{1}{2} \left( \nabla \mathbf{u} + (\nabla \mathbf{u})^T \right) \), \( \mu \) is the shear
viscosity of the fluid and \( p(x, t) \) is the mechanical pressure, defined as the
average normal stress acting on a fluid element.

The stress tensor is related to the total surface force \( F_S \) that is applied
to a solid by the fluid surrounding it by

\[ F_S = \int_{\partial\Omega} \sigma \cdot n dS, \]  

(2.4)

where \( \partial\Omega \) is the solid boundary in contact with the fluid. On the other hand,
the total thrust \( T \) of a helicopter in hover, i.e. a helicopter in stationary
flight, is equal to the weight of the helicopter \( W \). In other words, it can be
express as the following

\[ T = -W. \]  

(2.5)

Thus since the thrust, in this case, is equal to the total force exerted by
the fluid, we can combine the previous two equations. The resulting equation is
the relation between the stress tensor and the thrust:

\[ T = \int_{\partial\Omega} \sigma \cdot n dS. \]  

(2.6)

Now, in order to obtain the equation of motion for a volume \( \Omega \) that is
fixed in space, with a boundary \( \partial\Omega \), we recur to Newton. Newton proposes
that the rate of change of momentum is equal to the sum of all the forces acting on the fluid plus the momentum flux through $\partial \Omega$. Therefore, if the momentum of the fluid inside $\Omega$ can be written as

$$\int_{\Omega} u \rho dV \quad (2.7)$$

Thus, the rate of change can be expressed as

$$\frac{d}{dt} \int_{\Omega} u \rho dV = - \int_{\partial \Omega} (\rho u) u \cdot n dS + F, \quad (2.8)$$

If it is now assumed that the total force acting on a part of the fluid is the sum of both the volumetric and superficial forces. Then, the total force acting on the volume would be

$$F = \int_{\Omega} f dV + \int_{\partial \Omega} \sigma \cdot n dS, \quad (2.9)$$

where $f$ is the body force density and $\sigma \cdot n$ is the contact force per unit area or stress vector. Using the Divergence Theorem, we can rewrite the superficial forces and momentum fluxes to obtain

$$\int_{\Omega} \frac{\partial}{\partial t} (\rho u) dV = \int_{\Omega} (\nabla \cdot \sigma + f) dV - \int_{\Omega} \nabla \cdot (\rho uu) dV. \quad (2.10)$$

Given that the selected volume $\Omega$, is an arbitrary region in space, we can say that the relation exists between the integrands

$$\frac{\partial}{\partial t} (\rho u) = (\nabla \cdot \sigma + f) - \nabla \cdot (\rho uu). \quad (2.11)$$

If the derivatives are now opened and the different parts rearrange, the equation is now:

$$u \left[ \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho u) \right] + \rho \frac{\partial u}{\partial t} = \nabla \cdot \sigma + f - \rho u \cdot \nabla u. \quad (2.12)$$

However, thanks to equation $2.1$, it is known that the parenthesis in the left side of the equation, since the fluid is considered incompressible, is equal to zero

$$\rho \left[ u \cdot \nabla u + \frac{\partial u}{\partial t} \right] = \nabla \cdot \sigma + f. \quad (2.13)$$
Now, that which is inside the parenthesis in the left side of the equation, can be seen as the material derivative of $u$

$$\rho \frac{Du}{Dt} = f + \nabla \cdot \sigma. \quad (2.14)$$

If we now substitute the expression found for the stress tensor in equation 2.3 in the previous equation we obtain what is commonly known as Navier-Stokes equation of motion:

$$\rho \frac{Du}{Dt} = f - \nabla p + \mu \nabla^2 u. \quad (2.15)$$

With equations 2.1 and the condition of incompressibility $\nabla \cdot u = 0$ and the equations 2.3, 2.15 there are now 4 equations to solve for the velocity $u$ and mechanical pressure $p$. The initial conditions, such as a no-slip boundary conditions of the fluid around a solid, give further information to solve the problem. The no-slip condition is when the fluid at the boundary has the same velocity as the solid boundary itself.

### 2.2 Helicopter Theory

In this section both the momentum theory and the blade element theory shall be discussed. These two theories are used to both explain and simulate the flight of a rotorcraft. For the purpose of this thesis, they are important because the software used to make the simulations, RotCFD, combines these two in order to simulate rotorcraft downwash. Furthermore, both theories produce variables of extreme importance to the description of the phenomena addressed by this thesis. Example of such variables are $v_H$, the wake velocity, and $C_T$, the thrust coefficient. The wake velocity is the maximum velocity the fluid in the wake will reach and thus is taken to be the one responsible for saltation (the kick-up of dust). Additionally, $C_T$ was used in this research to assure that the scaling of the problem was done correctly and that we are solving a dynamically similar problem to that in Mars.

A helicopter is an aircraft that uses rotating wings to provide lift, propulsion and control. In contrast to an fixed-wing aircraft, a helicopter is capable of producing aerodynamic forces even when the velocity of the vehicle itself is zero. Thus, allowing it to take-off and land vertically [10].

A helicopter is in hover when it has no velocity relative to the ground, let that be vertical or horizontal. While vertical flight, as the name suggest,
is a vertical movement. It implies axial symmetry of the rotor, causing the velocities and loads on the rotor blades to be independent of the azimuth position.

In order to do a general analysis of the problem, momentum theory can be used. Momentum theory applies the conservation laws of fluid mechanics (mass, momentum and energy) to the rotor and flow as a whole to estimate the rotor performance. In this theory, the rotor is modelled as an actuator disk. That is, a circular surface of zero thickness that can support a pressure difference and thus accelerate the air through the disk. Moreover, the actuator disk can also support torque, which can thus impart angular momentum to the fluid that passes through the disk [10]. Figure 2.1 shows a diagram of

![Diagram exemplifying the Momentum Theory and actuator disk approximation.](image)

Figure 2.1: Diagram exemplifying the Momentum Theory and actuator disk approximation.

the idealized flow across the actuator disc in hover. The flow is considered to be irrational and incompressible above and below the actuator disc. The idealized flow across the actuator disc consists of a streamtube, and Ω represents a section of the streamtube split by the actuator disc somewhere in the middle. In Momentum Theory, a relation for the thrust on the rotor can be found by first considering the pressure difference in both sides of the disc and using the Bernoulli equation. Take a stream line starting way above the disc, at the point 0 of figure 2.1 and ending right above the disc, at point 1. Similarly a second streamline can be used below the disc, say from point 2 below the disc to point 3 far away. Using Bernoulli in both streamlines we obtain
Where $p_0$ to $p_3$ are the pressure at points 0 to 3 respectively, $\rho$ is the uniform fluid density, $u_H$ is the fluid speed across the disc in points 2 and 3, $u_0$ and $w$ are the fluid speed in points 0 and 3 respectively. Pressures $p_0$ and $p_3$ are approximately the atmospheric pressure, and the pressure difference across the disc represent the magnitude of the thrust $T$ distributed over the disc area $A$, giving $T/A = p_2 - p_1$. Adding the relations (2.16) we obtain

$$T/A = \frac{1}{2} \rho w^2.$$  

Now consider the momentum balance in $\Omega$ given by equation (2.9). Since the velocity field is steady and density is constant the momentum in $\Omega$ remains constant and the left hand side of equation (2.9) is zero. The actuator disc imposes a force $-T\hat{e}_z$ on the flow, that appears as a source term in the momentum balance 2.9, leading to the following expression:

$$0 = -\int_{\partial\Omega} \rho u(u \cdot \hat{n})ds + \int_{\partial\Omega} \sigma \cdot \hat{n}ds + \int_{\Omega} f dV - T\hat{e}_z.$$  

The first integral, the flux of momentum, is roughly $-A_0 \rho u_0^2 \hat{e}_z + A_3 \rho w^2 \hat{e}_z$, where $u_0$ is so small it is negligible compared to $w$. Meanwhile, the stress integral is approximately $A_3 p_3 \hat{e}_z - A_0 p_0 \hat{e}_z$, neglecting the contribution coming from the side boundaries. The force applied by the actuator disc, equivalent to the weight of the helicopter, must be considerably larger than the weight of the air in $\Omega$, the integral of $f$. Also, as $p_0$ and $p_3$ are basically the atmospheric pressure, it is reasonable to consider the effect of stresses on $\partial\Omega$ as well as the momentum flux across $A_0$ to be much smaller than $T$, reducing equation (2.18) to $-T\hat{e}_z + A_3 \rho w^2 \hat{e}_z = 0$. Since mass is conserved, mass flux $A_3 \rho w$ is equal to $A \rho u_H$, and thrust can be expressed as

$$T = A \rho u_H \omega.$$  

Using the relation (2.17) an estimate to the velocity below the rotor is obtained

$$u_H = \sqrt{\frac{T}{2 \rho A}}.$$  

$$p_0 + \frac{1}{2} \rho u_0^2 = p_1 + \frac{1}{2} \rho u_H^2,$$

$$p_2 + \frac{1}{2} \rho u_H^2 = p_3 + \frac{1}{2} \rho w^2.$$  

(2.16)
2.2. HELICOPTER THEORY

On the other hand, in order to explain Blade element theory, lets first introduce some of the terms related to the rotor blade (see Figure 6.4). In this context an airfoil is a 2D representation of a rotor blade which is used to study the aerodynamic forces present when interacting with an air stream. The imaginary straight line that exist between the leading edge and the trailing edge is known as the chord line. The leading edge is the first point of contact between the airfoil and the air stream. While the trailing edge is the opposite end of the airfoil. The blade twist refers to a change (a twist) of the chord line. The angle of attack is the angle between the blade chord line and the direction of the free air stream. Flapping is when the blade is allowed to increase and decrease the angle of attack of each blade during rotation. What is known as the hinge offset is the movement of the flapping hinge from the center of the mast. The pitch angle is the angle between the chord line and the reference plane, which is the same one as that of the rotor hub. If we look at figure 2.3, this reference plane would be the horizontal axis. The collective pitch refers to the pitch angle at which all blades are from the reference plane. Meanwhile, the cyclic pitch refers to the change in angle of each blade depending in its azimuthal position. Lastly, when the helicopter is making a vertical take-off, there are two main forces acting on the rotor blades, the lift that allows it to carry the weight of the helicopter and the centrifugal force on the rotating blades. The result of these two forces on the blades is the coning of their path, as opposed to staying in their original flat position. The angle at which they cone is known as the cone angle.

Figure 2.2: Airfoil section diagram where the geometric properties of an airfoil are pointed out.

The blade element theory is based on the lifting-line assumption. The Lift-lining assumption considers only rigid flap motion and no lag or pitch
degrees of freedom. Blade Lag is the forward and backwards movement of the blade along the plane of rotation. Lift-lining also assumes small angles and neglects effects of stall, compressibility and radial flow. Stall is when the lift produced by the airfoil is reduced, this occurs when the angle of attack is increased passed the critical angle of attack. Moreover, the blade must have a constant Chord and linear twist. A uniform induced velocity, or the velocity of the fluid surrounding the blades caused by their movement, is also considered [10].

In figure 2.3 the velocities and forces acting on the blade are shown. The free air stream has two velocity components in this figure, $u_T$ air velocity component tangential to the disk plane, and $u_P$ is the air velocity perpendicular to the disk plane. Therefore, the air stream has a speed

\[ U = \sqrt{u_T^2 + u_P^2}, \]

and makes an angle with respect to the disk plane given by

\[ \phi = \tan^{-1} \frac{u_P}{u_T} \]

Given that the air stream can be blowing in any direction, since it is defined by the weather conditions, this angle can be different to zero in hover. Thus, the angle of attack can be defined mathematically as

\[ \alpha = \theta - \phi, \]
2.2. HELICOPTER THEORY

where $\theta$ is the angle between the chord line and the plane of rotation. This angle can also take values different to zero since it is the pitch angle of the blades. Additionally, the forces that are being produced by the air flow around the blade are lift, $L$, and drag, $D$. These forces, as shown on the diagram (figure 2.3), are parallel to the related wind ($U$), in the case of the drag, and perpendicular in the case of the lift. The lift coefficient, $C_L$, and the drag coefficient, $C_D$, can be defined as

$$L = \frac{1}{2} \rho U^2 CC_L$$

(2.24)

$$D = \frac{1}{2} \rho U^2 CC_D$$

(2.25)

where $\rho$ is the air density, $C$ the chord.

However, the resultant force can be split into two components $F_x$ and $F_z$ that are parallel and perpendicular to the disk plane, respectively. They are calculated as follows:

$$F_x = L \sin \phi + D \cos \phi$$

(2.26)

$$F_z = L \cos \phi - D \sin \phi$$

(2.27)

From these, important characteristics of the rotor can be approximated such as the thrust

$$T = N \int_0^R F_z dr,$$

(2.28)

where $N$ is the number of blades. These are the values that the blade element provides that are useful when designing a rotor.

The way this theory is implemented into the simulations is by including what is known as the airfoil tables. Given that most blades have smooth surfaces, it is possible to approximate how these 3D objects interact with the fluid, the integral in eqn. (2.28) by solving a 2D boundary layer problem with some of its cross-sections. The airfoil tables contain information about the forces on the cross-sections for different angles of attack.

Another concept in helicopter theory that is essential for this thesis is Ground effect. Ground effect is the name given to the phenomenon that occurs when the rotor dish is in proximity to the ground. What happens is that the induced velocity of the airflow through the rotor disk as well as the rotor wake are constricted and modified by the surface friction. When
the helicopter is out of ground effect, there are two “inner vortices” that produce stall. These vortices appear at the blades tip and close to the rotor mast, reducing lift in these portions of the blade. Interaction with the ground restricts the generation of blade tip vortices making a larger portion of the blade produce lift (Figure 2.4). All these factors contribute to the increase in thrust. This effect is limited to height of about a diameter of the rotor above the ground, and reaches its maximum over a firm smooth surface [11].

![Figure 2.4: Air circulation patterns change when hovering out of ground effect and when hovering in ground effect [11].](image)

The Rotorcraft brownout is the result of the entrainment of dust and/or sand particles due to helicopter downwash. It can also be referred to as whiteout when landing over powdered snow. This effect while landing poses a problem since the dust can reduce or completely block the pilot’s visibility. Brownout, can be divided into 3 parts: helicopter downwash, particle entrainment and visual obscuration. Particles become airborne outside the rotor disk at the same distance that the maximum outwash velocity is reached [2]. The three main transport modes for particles on a surface are: surface creep, saltation and suspension. The first one occurs for low-speed flows, since the force being exerted on them only have energy to roll them on the surface. At higher-speed flows, there is saltation, a process in which the particle becomes airborne [13]. Suspension happens once the particles are off the ground and the flow has a high enough speed to beat the gravitational pull.
2.3 Saltation

Saltation is one of the many ways the wind can transport particles on the surface. The word “saltation” comes from the Latin word “saltare”, which means to leap, making the word seem appropriate since the grains being moved by saltation are actually bouncing. Bagnold’s “The Physics of Blown Sand and Desert Dunes” [13] talks about two different threshold flow velocities to produce saltation, the static threshold and the dynamic threshold. For the purpose of this thesis, we shall focus purely on the static threshold velocity. That is the necessary wind conditions that will allow the particles of a given size and density to initiate motion from rest. The threshold surface stress ($\tau_t$) necessary to start particle motion is directly proportional to the threshold friction velocity

$$\tau_t = \rho u_*,$$

(2.29)

where $\rho$ is the fluid density. According to Bagnold, the threshold friction velocity on Earth can be determine by

$$u_* = A \sqrt{\frac{(\rho_p - \rho)gD_p}{\rho}},$$

(2.30)

where $\rho_p$ is the particle density, $g$ is the gravitational acceleration, $D_p$ is the mean particle diameter and $A$ is an experimental coefficient. Bagnold believed $A$ to be a unique function of the particles friction Reynolds number

$$Re = \frac{u_*D_p}{\nu},$$

(2.31)

where $\nu$ is the kinematic viscosity. Bagnold and Sagan tried to find a relationship between these two coefficients such that $A = A(Re)$ [13]. This would mean that if we knew Earth’s $A$ and $Re$ we could extrapolate this to treat saltation under Mars conditions. However, the results obtained by Iversen et al. [3] combined with those by Weinberger and Adlon [14], suggest that the universality of $A(Re)$ for small particles at low pressures is invalid.

The effects of winds upon a loose particulate surface at low pressure, as is the case of the Martian surface, is still unknown with complete certainty [13]. Thus, causing there to be a wide range of results for the predicted threshold wind speed, or minimum speed required to cause saltation, at such a low pressure. This project will take the results found by R Greeley et al. on “Dust Storms on Mars: Considerations and Simulations” to select the size
and material of the particles for the experiments [3]. Despite not satisfying all parameters involved in simulating an aeolian motion in Mars when doing an experiment inside a wind-tunnel conducted on Earth, this paper found an optimum grain size for Mars. That is, the size at which saltation occurs by the lowest winds.

Since Gravity on Mars is 0.38 that of Earth, less force is required to initiate particle lift and movement in Mars. For threshold simulations conducted on Earth, the particles must be 0.38 the density of their equivalent in Mars. Ground walnut shells are the right density and shape to that of the natural windblown materials in Mars [3].

Moreover, in the research done by Greeley, White et al. [3], it was revealed that the particle size with the lowest friction threshold speed at pressures within $10^5 \text{Pa}$ to $400 \text{Pa}$, were the samples of 64 $\mu m$ and 330 $\mu m$ average diameter of ground walnut shells [3]. It was further determined that the minimum threshold friction speed to cause saltation on Mars at its average surface pressure of 500$\text{Pa}$ is 2.5 $m/s$. Additionally, it is concluded that the particle size that is easiest to move is that of 160 $\mu m$ in diameter. On the other hand, given that the speed is inversely proportional to the pressure of the test section, the threshold velocity at a pressure of 1000 $mb$ is of 22.61 $cm/s$ ($0.02261 \text{ m/s}$). Therefore, we assume this value to be the minimum speed necessary to move particles.

### 2.4 Computational Fluid Dynamics Tool

Flow simulations used in this work were produced with a program called RotCFD. RotCFD is a computational fluid dynamics (CFD) tool developed at Surka Helitek, Inc. that emphasizes user-friendliness and efficiency that streamline the design process from geometry to CFD solution. RotCFD brings together CFD and Integrated Design Environment (IDE) specifically for rotorcrafts. One of RotCFD’s key components is its geometry module. With it, the user is able to import a body geometry in a STL or Plot3D format, which are standard CAD program formats (Computer-Aided Design). Additionally, the user may alter these designs by applying transformations such as rotations, translations and scaling [15].

Another key component is its semi-automated grid generation module or “UGen”. This module generates a Cartesian Octree grid that allows refinement in specified regions of interest. Octree is a data structure in which a
given number of grid cubes are subdivided into 8 smaller grid cubes, producing a local mesh refinement. RotCFD also uses "Body Conforming Grid", a technique that moves the vertices of the grid to fit a given boundary, like the fuselage of a rotorcraft, distorting the shape of the neighboring cell. The cells that intersect the body or are surrounding it are replaced by tetrahedra. To further illustrate this technique, Figure 2.5 shows two drawings. The left side shows the original grid and the curve represents the outline of an object. The right hand shows the result of this grid generation caused by the presence of the object [16].

The grid is not the only approximation, since a new body surface is being generated form these new edges. This produces distortions on the body which are obvious on the example geometry found in Figure 2.6 [15].

The rotor module is also one of the most important features of RotCFD. This module is based on the concept of rotor blades represented as momentum sources. The rotor momentum sources are primarily a function of the local velocity of the flow, given by the numerical algorithm the solves the Navier-Stokes equations and the two-dimensional airfoil characteristics of the rotor blades. The Navier-Stokes equations and the blade element theory are coupled implicitly in order to obtain the rotor performance and flow field. While the blade element theory provides the forces acting on the rotor blades, the numerical algorithm that solves the Navier-Stokes equations yields the flow field near the rotor blades, including the induced momentum sources.

RotCFD utilizes an incompressible, unsteady Navier-Stokes solver with a
staggered grid, based on finite volume method with either implicit or explicit time integration. The pressure field is computed using a SIMPLE algorithm (Semi-Implicit Method for Pressure Linked Equations) \[15\].

Lastly, a flow visualization and analysis module is included in this program. The idea of having a Graphical User Interface (GUI) is to facilitate the CFD post-processing by allowing the user to visualize the results.
Chapter 3

Design of a 30% Scale Model

The 30% scale model of the MHS was designed specifically for the experiments described in this thesis. The model was built trying to preserve the shape and the rotor configuration of the full size prototype. Previously, a half-scale version was made in order to run flight tests inside a wind tunnel. The tunnel was designed in order to fit inside a vacuum chamber located at the NASA AMES Research Centre. However, this model was not used in the tests for this thesis.

The dimensions of the experimental model were based on data that has been made public by California Institute of Technology’s Jet Propulsion Laboratory, also known as NASA JPL (Figure 3.1). This centre is the main research facility in charge of the Mars Helicopter Scout (MHS) project. Real and accurate dimensions were not made available for the purpose of this thesis except for the real size of the cube box that will be the body of the helicopter.

The electric and mechanical systems used in the design of the model were taken from an off-the-shell RC helicopter. This is a GPTOYS G610 11” (279.4mm) Durant Built-in Gyro Infrared Remote Control Helicopter. This helicopter was selected for the size of the blades since it was the closest we could find off-the-shelf of approximately 30% of the full size of the MHS rotor disk diameter \( D \) (\( D = 1.1m \)).

Moreover, the size of the rotor disk was not the only feature of the MHS that had to be scaled in order to consider this model. This model had to fit inside a fuselage that also represented the 30% of the MHS cubed body. The model will now be referred to as the M3HS. At the same time, there were things that had to remain the same, such as the Thrust coefficient, \( C_T \). The
Thrust Coefficient of a rotor is defined by the formula

\[ C_T = \frac{T}{\rho \pi u_{tip}^2 R^2}, \]  

where \( T \) represents the thrust, \( \rho \) the density of the atmosphere, \( R \) the radius of the rotor disk and \( u_{tip} \) is the rotors tip velocity. Given equation 2.20 and knowing that the area of the disk is \( A = \pi R^2 \) then the scaled wake velocity is given by

\[ \frac{u_H}{u_{tip}} = \sqrt{2C_T}. \]  

The problem can now be considered equivalent if the \( C_T \) is specified. Since we consider \( u_H \) to be the responsible for saltation, we may compare it to that of the full-scale problem.

Therefore, a thrust test was necessary to determine these values. In order to do the thrust test, and further tuft and dust testing, the motors and the rotor had to be removed from the original toy’s fuselage. The motors were then placed in a new fuselage that would hold them steady and, at the same time, would have a similar shape and size of the real Mars Helicopter Scout.
fuselage. This was done using a mechanical design software and 3D CAD software called Inventor. However, given the size of the gear system, the shape of the fuselage could not be fully cubed. Although we could adjust both the height and width of the cube to be the right size, the depth had to be modified to fit the gears. Once the size of the box was design, an adaptor was then added in order to attach the fuselage to the mounts (Figure 3.2). The dimensions of the fuselage were $39 \text{ mm} \times 51 \text{ mm} \times 40 \text{ mm}$. This design includes side holes to keep cool the motors. The motors would heat up so much it melted older designs that had no ventilation. Lastly, the fuselage was printed using a 3D printer with 30% infill. This level of filling guaranteed the stiffness necessary to avoid vibrations and movement of the body and mast.

Once the model was ready, it was placed on a sting mount that was to be used for the thrust experiment. This experiment consisted on placing the model directly above a ground plate of $1.5 \times 1.5 \ D$, with $D$ now being the diameter of the scaled rotor. In our case $D = 55.8 \ cm$ (22 in). The plate was placed on top of a scale in order to measure the thrust by means of the hydrodynamic force of the flow, produced by the rotor right above, pushing down the plate (figure 3.3). Meanwhile, a tachometer was held steadily and perpendicular to the rotor disk. By placing a reflective sticker on one of the
blade tips, the revolutions per minute (RPM) of the model could be measured at all times. The results of this experiment are shown on Figure 3.3. The blue points represent the data points from the experiment. The blue line represents the linear trend line. From the RPM shown on figure 3.3 (top) the tip velocity $u_{tip}$, was calculated using

$$u_{tip} = \frac{RPM \times R \times \pi}{60s},$$

(3.3)

where R is the radius of the rotor disk and the $u_{tip}$ is in m/s. Meanwhile, the red points in figure 3.3 represent the average values taken from these results, and were used as the tip velocity input in the RotCFD simulations.

From these graphs and using equations 2.20 and 3.1, the following comparison was done. It can be seen on table 3.1 that most variables are roughly the same percentage except for the thrust coefficient that should be the same. Furthermore, this helicopter model seems to be good enough for these initial studies since the $u_H$ is bigger than the predicted threshold velocity needed to move the dust particles under Earth conditions [3].

<table>
<thead>
<tr>
<th></th>
<th>Full Size MHS</th>
<th>Scaled Model</th>
<th>Percentage</th>
</tr>
</thead>
<tbody>
<tr>
<td>Volume</td>
<td>$140mm \times 140mm \times 140mm$</td>
<td>$39mm \times 51mm \times 40mm$</td>
<td>28.99%</td>
</tr>
<tr>
<td>$u_{tip}$</td>
<td>138.2301 m/s</td>
<td>37.6005 m/s</td>
<td>29.76%</td>
</tr>
<tr>
<td>$u_H$</td>
<td>8.08128 m/s</td>
<td>2.3107 m/s</td>
<td>28.59%</td>
</tr>
<tr>
<td>$C_T$</td>
<td>0.0068</td>
<td>0.0064</td>
<td>93.68%</td>
</tr>
</tbody>
</table>

Table 3.1: Summary of model’s Characteristics compared to real MHS in Earth conditions.

One of the original restrictions set to this model were costs. This model’s original purpose was to function as an outreach project. Therefore, although the components were selected to be as accurate as possible, they were possibly not the best choice. Moreover, all of the testing done in this thesis is at sea level Earth pressure. Further testing is suggested at Martian pressures, however it is possible that this toy model will not function at such low pressures.
Figure 3.3: Thrust vs. $RPM^2$ graph based on thrust test of M3HS (top). Thrust vs. $u_{tip}^2$ graph based on thrust test of M3HS (bottom).
Chapter 4

Dust Test Cell Experiment

The Dust Test Cell Experiment included the design of an experimental setup that would allow us to visualize the dust cloud produced by the rotor of the scaled MHS model in Mars-like conditions.

4.1 Design of Experiment Test Section

The Dust Test Cell was designed thinking about the size and scale of the Mars Scout Model that would be tested inside and the possible tests that could be done. The main objective of this cell is to enable a controlled environment to simulate the possible dust saltation due to the helicopter’s coaxial rotor in Mars-like conditions. Therefore, it was decided that the Cell had to be an enclosed area, in order to keep the dust from contaminating its surroundings, specially considering future testing will take place inside a vacuum chamber in order to better emulate Mars conditions. The cell was made using a 91.76 cm × 102.55 cm × 120.33 cm cardboard box (inner measurements). The bed section had to be large to reduce wall effects on our experiment. We chose the floor of the box to be more than $2D \times 2D$ where $D$ is the diameter of the rotor being used to ensure the wake of the rotor to be falling completely on the test section. However, since we were unsure at what point or if there would be dust saltation, this area had to be considerably larger to allow us to see the effect. Since both the half-scale and the third-scale models are to be tested, the box’s dimensions were selected for the half-scale rotor. The test of the half-scale model will not be addressed on this thesis.
Moreover, it was important to make both the top of the box and the door of the cell transparent in order to see and record the experiments. These were to be filmed and photographed from different angles in order to later compare with simulations. The see-through sections are acrylic windows that are strong enough to resist impact from a broken rotor.

Lastly, the bottom part of the box had to have a wooden barrier to help keep the dust inside the box when it was open. This barrier further helped in the flattening process of the dust. The bottom of the box was also covered by a particle board, material designed to be very smooth with an even surface in order to ensure the flatness of the sand was not altered by an uneven bottom surface (Figure 4.1).

The material used to simulate the dust on the Martian surface for this experiment was selected based on the research done by R. Greeley et al. [3]. They proposed that the best material to simulate the dust grains are crushed walnut shells due to their density and size. Both of these qualities help compensate for Earth’s gravity and simulate the weight of the particles in Mars, at the same time it keeps the right ratio between the atmospheric density and that of the particles. For purposes of this thesis, only particles of size 60 $\mu m$ were selected since that is the size that proved to be, in their experiments, the most prone to saltation.

Three heights were selected to gather information about the interaction between the model and the dust. These heights are “wheels on ground”, 15 cm drop, and $3R$, where $R$ is the rotor radius. The first height, “wheels on ground”, as the name suggests is the height at which the helicopter rotor when it is about to take off. This height is the closest to the ground and therefore it is believed to be the height at which it is most likely to see dust saltation. The second height, 15 cm, refers to the scaled version of that at which the full scale Mars Helicopter Scout (MHS) turns-off the rotor and lands. This height is of extreme importance given that the design of the helicopter has an altimeter that will determine when it has reached the 15 cm, measured from the “wheels” of the helicopter to the ground, and shut down to land softly. It is believed that if the helicopter is ever going to be in danger of brownout, it would be at this point given that it is both the moment of landing plus it experiences in-ground effects. Lastly, the third height will allow us to study saltation effects from a position outside ground-effects. Since the model that was used for this thesis was the 30% scale model of the full size MHS, these heights were scaled in order to see similar effects. Therefore, the heights used for these experiments, shown in Table 4.1
Figure 4.1: Test cell box. The fuselage inside the box is the half-scale model.

the half point between the two rotor disks of the model on the mast as their reference point.
36

CHAPTER 4. DUST TEST CELL EXPERIMENT

<table>
<thead>
<tr>
<th>Height Type</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>Wheels on ground</td>
<td>99 mm</td>
</tr>
<tr>
<td>15 cm drop</td>
<td>147 mm</td>
</tr>
<tr>
<td>3R</td>
<td>419.1 mm</td>
</tr>
</tbody>
</table>

Table 4.1: Heights used on both Tuft testing and Dust testing.

4.2 Methodology

After the dust test section box was built, it was time to place the dust inside of the box. The depth that was selected for this experiment was the width of the wooden guides on the side of the box, of $3/8$ in or 9.525 mm. The method of “dusting” the inside of the box was of the highest importance. The dusting had to be done with extreme caution to avoid compacting the crush walnut-shells (CWS), while keeping the depth of the dust bed even. This was achieved by first kneeling inside the Test box placing the pail and the screed near you. By placing oneself in a comfortable and strategic location, one was able to place the CWS starting from the far back, forward. Then, the test bed was mentally divided into 4 strips parallel to the far wall. With the help of a plastic scoop, the CWS was generously spread out on the first strip. That is, the one furthest from the entrance (Figure 4.2). The amount recommended was approximately 3 scoops. It is better to put less CWS at a time than a lot. The screed was then used to flatten the CWS by gently pull it towards you, using the flattest side. One must make sure not to apply any downward force since this could alter the thickness of the layer at the same time that one would compact the CWS. It was important to keep the screed at a constant angle while moving it. Furthermore, the direction in which one moves the screed had to remain constant every time. Once the CWS was no longer being pushed, another scoop was grabbed to cover any holes that might exist in the previously flattened strip. This was repeated until there were no more holes on the sand. This process was repeated for the second strip and fourth as well. For the third strip, the procedure changed. Once one is done flattening the strip, there was a “mountain” of CWS in the middle. This was caused by the presence of the pedestal in the centre of the Test Box. In order to level this section the screed was used, always making sure that both ends were on the side guides, and it was rotated around the pedestal. This left an “X” shape pattern on the centre strip area. Therefore, the left and the right sides were flattened separately. One should never try to flatten all in one go.
4.2. METHODOLOGY

Figure 4.2: Placing crushed walnut shells inside dust test section box.

Once the dust bed was set, the MHS 30% scale model was placed on a mount that could be attached to the centre of the pedestal. Given that the sides of the model were exposed, it was decided to cover them with black tape in order to avoid the dust from getting in between the gears. The height at which the model was held was determined by the length of the mount that was attached to the pedestal. This height was measured not from the bottom of the test bed, but from the surface of the dust bed.

4.2.1 White Light Experiment

For lighting, a LED lamp was placed on the top part of the test cell in order to illuminate as much as possible the inside of the box. All experiments were done with the room lights turned on. It was realised that placing a lamp above was better than by the side window. The reason for this was that the reflection of the lamp on the window did not allow the camera to correctly capture the phenomenon. Moreover, the lamp was not placed directly above the model to avoid creating a shadow directly underneath. For this experiment, the camera used was that of an Iphone 6S, and was placed perpendicular to the MHS disk plane. The camera allowed us to film the experiment at a frame rate of 120 frames per second.
4.2.2 Laser Wall Experiment

These experiments were done in similar conditions to the white light testing. The material used in the bed of the box was CWS. However, the main difference between these two experiments falls in the lighting and the recording of the phenomenon. The lighting consisted of a 2 watts green laser, which was opened using a glass stirring rod in order to produce a “laser wall”. This laser wall was aligned with the centre of the rotor and the vertical axis of the helicopter (Figure 4.3). Furthermore, the laser wall was placed perpendicular to the camera axis. The camera used for these experiments was a phantom camera. Given the lens size available and the distance at which the camera could be placed, outside the box, the area on the laser plane that was visualized was roughly of about 30 cm across by 20 cm tall. Moreover, the camera’s frame rate was of 500 frames per second with an exposure time of 2000 µs.

Figure 4.3: Diagram of the view inside the box with the helicopter illuminated with a laser wall.
Chapter 5

Tuft Testing

In order to validate the RotCFD simulations, an experiment that could allow the visualization of the flow field had to be designed. When studying the flow field for normal earth aircrafts, a method commonly referred to as “Tuft testing” is used. The test consists on laying a grid of “tufts” on the most relevant and easily comparable planes. On rotorcraft theory, it is believed that the rotor wake has radial symmetry when there is no external body to alter this assumption. Therefore, the main planes of interest are those radial to the rotor shaft and that which is perpendicular to it, i.e. the ground plane. The testing methodology used to study the UH-60 helicopter brownout by Alan Wadcock et al in “Rotorcraft Downwash Flow Field Study to understand the Aerodynamics of Helicopter Brownout” [2], was the base of the methodology used in this thesis. However, there were major factors that had to be accounted for that were cause of changes to the experimental setup.

The flow field experiments were done using the 30% scale model that was described in the previous section. Given the difference in scale between our 30% scale-MHS and the life scale UH-60, the size and material of the tufts had to be changed. To represent the air-flow movement it is important to make the tufts as small as possible. When studying the UH-60, given its size, normal pieces of string of 16.51 cm long were used. Our model is so small, that a tuft of that size would not align with the airflow. Thus, it was decided to go for the smallest tufts available, micro-fiberglass tufts. These tufts have a micrometric diameter, which is approximately a tenth of a human hair. They are hardly visible to the human eye when seen with visible light. However, when shone with a UV lamp, the florescent property
of the material increases its visibility to a millimetric diameter, allowing us to capture their movements with a camera.

There are several advantages to using a small scale model. One of those being that by using a much larger box compared to the model’s size, it is possible to simulate free-flight without the model actually being outside. Furthermore, by having an enclosed test area, it was possible to reduce the airflow in the room and thus determine that the movement of the tufts was caused strictly by the wake of the rotor.

The tufts used were selected because of their micro-size and material. Their micro-metric diameter allowed them to be easily moved by the wake of the helicopter, at the same time, it made them the least invasive “visualizing particles”. However, it is important to point out that although it is the “least invasive” solution, the tufts are still a lot heavier than the air that moves them. Furthermore, despite all the efforts done to make this the less intrusive as possible, one should always keep in mind that the board itself holding the tufts is very intrusive and thus alters the flow.

An important limitation these tufts had was in fact their size. They were so thin that it made them very fragile and easy to break or bend. Furthermore, they would occasionally get stuck with the ruffles of the cardboard on which they were glued to. Not to mention the orientation they were originally glued to and gravity caused them to have a preferred orientation. Also, the presence of static electricity interfered with their freedom of movement. With the purpose of reducing this problem, the boards were rubbed using dryer sheets.

5.1 Experimental Setup

The experimental setup included black poster-boards of different sizes that were used to glue the tufts on. The poster board thickness, 1.27 cm, was selected by keeping in mind the need for the board to be as less intrusive as possible, at the same time it needed to be stiff in order to avoid the deformation of the board. The colour of the back is important, previous testing showed that the visualization of these tufts was improved by having a dark background as opposed to a light-bright background. Three different sized boards were made, two vertical test boards and one ground board. The two vertical ones where used to include the three selected heights when studying the radial planes. The dimensions of the biggest one of the two,
from now on referred to as the Vertical Big Board (VBB), were 28 cm × 0.4 cm × 40.6 cm, while the Vertical Small Board (VSB) was only 28 cm × 0.4 cm × 7.6 cm.

Both vertical boards had an angle of inclination of 1.75° from the vertical position. The reason for this was to avoid the air flow being affected by the edge of the board while trying to maintain the board as vertical as possible. The VBB had 42 vertical rows and 66 horizontal rows giving a total number of 2772 tufts. Meanwhile, the VSB had a total number of 132 tufts. Both boards had borders and separation between rows of 0.64 cm. The ground board was a 91.44 cm square with 71 rows of tufts. Each tuft row ran all across the board and had a separation of 0.64 cm. The total number of tufts on the ground plane was 5041.

In order to make the boards it was important to indicate accurately where the string of tufts would be placed. This allows to insure that they are all evenly spaced and parallel to each other. Additionally, by also drawing lines perpendicular to the tuft rows, it was clear where one was to put the non-reflecting glue. This was done in order to secure the tuft after cutting the roll of material. The tufts were not to be touched in order to avoid static
electricity or anything that could cause the tuft to stick on the board. All this was done with the help of a hand-held UV-lamp that was not bright enough to cause eye problems, but facilitated the visualization of the tufts.

The design of the mount represented a challenge on itself. Given the ground floor testing, the model had to be held horizontally since it could not be held down to the board. That is why it was decided that for this test, the model had to be held by a “sting-mount”. A sting mount consists on a strong vertical base that has a “sting” held on to it at a 90° angle to which the model is attached to (Figure 5.2). It was of extreme priority to keep the diameter of the tube holding the 3D model as small as possible to be the least invasive as possible, without compromising the stability and stiffness of the mount. For these tests, the base of the mount was magnetic, so to make it a stronger base if placed over a stiff metal surface. The 3D printed version of the model was modified in order to be attached to both the tuft testing sting-mount and the dust testing vertical mount.

Figure 5.2: Example of a Tuft Test Experimental setup.
5.2 Methodology

The Tuft Experiment consisted on short runs of 30 seconds each in which the 30% scale model was turned on and both video and photographs were taken. Three boards were placed in different positions from the model each time, in order to visualize the phenomenon the best as possible. Analogous to the Dust Testing, the same three heights were selected for these experiments. However, this study was taken one step further by including planes with different inclinations with respect to the plane of the rotor. These angles ranged from 0° to 30° with an increase of 5° each time. By combining both the heights and the angles selected, we obtain a total number of 18 possible configurations. Nonetheless, the total number of runs needed are a lot more, approximately 54, since for each configuration we wanted to get at least three different shots. The variety in geometries of the tuft boards allowed for the visualization of the air flow underneath and next to the rotor and fuselage, from views parallel and perpendicular to the rotor plane.

To capture the phenomenon, both pictures and videos were taken. Two cameras were used in this experiment, one for the recollection of data and the other to record a video that would facilitate the comparison. A professional camera was placed as perpendicular to the tuft plane as possible. Yet, this was not always the case since it was not always the most convenient setup. Since a large UV lamp was used to illuminate the entire board, placing the camera directly in front of the board would usually result in an over saturation of several pixels in the image due to the lamp’s reflection on the smooth black board. Furthermore, the body of the model was cover with black tape in order to avoid having a bright reflecting white body. This also helped keep the aerodynamics of the body as close to the original MHS shape as possible. Given the lighting conditions were controlled and kept constant, the camera settings for all pictures were the same. These settings are shown on table 5.1.

<table>
<thead>
<tr>
<th>Exposure time:</th>
<th>2 sec.</th>
</tr>
</thead>
<tbody>
<tr>
<td>f (Aperture):</td>
<td>16</td>
</tr>
<tr>
<td>ISO sensitivity:</td>
<td>1,600</td>
</tr>
<tr>
<td>Frame rate:</td>
<td>1 frame/5 sec.</td>
</tr>
</tbody>
</table>

Table 5.1: Camera settings for all Tuft Testing Experiments.

Notice the long exposure time, the reason for such exposure time is to
capture the motion-cone of each individual tuft [5.3]. Thus, making it possible to see the region in which each tuft moved in the given time lapse. The ISO and aperture were selected by consulting the department’s photography expert and corroborated via trial and error, to optimise the amount of light entering the camera in order to avoid the over saturation of pixels while maximizing the recollection of information.

![Figure 5.3: Close up of a sample tuft board in which tuft coning can be seen.](image)

On the other hand, the camera used to make the videos was the one on an Iphone 6S. The video was used to accurately record when the camera started taking pictures with respect to the moment in which the model was turned on. This is used when comparing experimental results with simulations.

The experiment were done in an enclosed environment, a room, with no wind currents, low humidity and protection from external contamination by limiting the access to this room. Before setting the tuft board in position, it is imperative to wipe the board with anti-static cloth so to decrease the static electricity. This must be done with extreme caution to avoid damaging the tufts.
Chapter 6

CFD Flow Field Simulations

The Computational Fluid Dynamics’ simulations were done using the program RotCFD. However, in order to start the simulations, several steps had to be taken beforehand. These steps will be referred to as the pre-simulations steps. Amongst these are: 3D scanning and processing of the model’s blades; creation of the airfoil tables and their implementation into RotCFD.

6.1 Blade Scanning and Post-Processing

In order to get the correct geometry of the blades into the simulation, they were first 3D scanned. A “Metrascan 3D” handheld laser scanner, produced by “Creaform” was used to do the 3D scanning [17].

The scanner produces a CAD file from which the geometry of the blade can be extracted (Figure 6.1). Two scanning methods were tested. The first one was a target based scan. As the name implies, this scanning method relies on target stickers that are placed around the blade in order to calibrate the scanner. However, a down side to this method is that, in order for it to work properly, the blade cannot move. Given the size and weight of the blade, this made it very difficult since we could not find a non-invasive way of securing it in such a way that it would no be affected by external vibrations. Therefore, a second method was utilized for the scanning. This second method does not require placing targets on the blade, but instead uses the targets that were already on the hand-held scanner to triangulate the position. This was the method of choice, although it introduced some errors since we couldn’t avoid partially blocking the signal with our bodies while scanning.
CHAPTER 6. CFD FLOW FIELD SIMULATIONS

Figure 6.1: Screenshot of the Blade scanned. The blade is the blue smooth figure in the centre. The blue shapes that surround the blade are partial scans of the objects that surround it while it was being scanned. These were all removed during the cleaning process.

Once the laser scanning was done, the CAD file produced was analysed using the 3D CAD modelling software package named “Rhino”. This tool allows for the “cleaning” of the blades scan. This refers to the creation of points that “fill up” for the areas of the blades that were not correctly scanned. This was done visually by the person in charge of the post-processing stage at NASA Ames, Eduardo Solis. Once all gaps are filled and the surface of the blades are smooth a point cloud is produced (see figure 6.2). Out of this point cloud, 20 cross-sections are selected and exported into text files that are then used to create the airfoil tables.
6.2 Airfoil Table Creation

The text-files containing the 20 cross-sections were then input into an already existing Matlab code written by Carl Russell designed to plot the cross-sections from the text-file. With this, each individual cross-section was graphed and placed next to each other (Figure 6.3). It was decided that five cross-sections represented the different geometries, given the simplicity of our toy helicopter blade design.

As was previously explained in the theory chapter of this thesis section 2.2, Airfoil Tables are of extreme importance since they determine how the air interacts with the rotor blades at different angles of attack. Therefore, obtaining accurate airfoil tables was a high priority. That is why it was decided to use the program “Xfoil”, a program designed by MIT to produce airfoil tables, instead of the more commonly used “C81 gen”.

6.3 Blades and Body Geometry

RotCFD allows for the user to introduce the specific geometry of the rotor blades by allowing the specifications of different parameters such as: radius,
CHAPTER 6. CFD FLOW FIELD SIMULATIONS

Figure 6.3: All cross-sections available. The top left corner is the rotor tip, while the bottom right corner is the part of the blade that is attached to the rotor mast. They are all different shapes since the blade design and its twists have to be extrapolated from these cross sections. The flat edges that can be identified in the last couple of cross-sections is the part from which the blade is held to the mast.

chord, twist, number of blades and cutout radius. All this information is extracted from the point cloud and exported into a text file produced by the same Matlab script as that of the airfoil tables. This text file exports the Chord and Twist value at specific $r/R$ where $r$ is the radial position from the center of the rotor and $R$ is the full radius of the blade. It is important to note that RotCFD requires consistency in the units that are used in its simulations. By graphing chord $/R$ vs $r/R$ (Figure 6.5), we create
6.3. BLADES AND BODY GEOMETRY

Figure 6.4: Sample Airfoil table. On top right hand side is the list of variables the table provides for each angle of attack.

a dimensionless plot that allows us to see the smoothness of our blade. The interaction of each blade and the fluid surrounding it is being extrapolated from only 3 airfoil tables. Therefore, the smoother the line is, the better this approximation is. The smoothness would indicate that the blade does not have bizarre changes between the positions of the airfoil tables, that may not be correctly simulated.

On the other hand, the twist of the blade is used to optimise the efficiency in hover. When graphing twist angle vs $r/R$ (Figure 6.6), we see a smooth curve that peaks at roughly 0.3 $r/R$ and then decreases the further away one is to the mast.

Moreover, it is also possible to determine how the rotor blades are positioned and moving by altering the following factors: cone angle, collective pitch, hinge offset, cyclic pitch, and flapping. Given that the model selected for the experiment was built from an off-the-shelf RC helicopter, all of the previously mentioned values are fixed by the manufacturing company and were all zero.
Figure 6.5: Chord/$R$ vs $r/R$ graph used to produce the shape of the blade. The data points had to be manually entered from the text file produced by the Matlab program to RotCFD.

6.4 Methodology

The simulations were done using a program that is currently being developed at Surka Helitek, Inc named “RotCFD”. This program specializes in the fluid dynamics that surround rotating objects such as the rotors of a vertical lift vehicle. The main purpose of these simulations was to obtain the maximum velocity values for each height and inclination. In order to correctly represent the real experiments, we have to input the same initial conditions of the system, which are as follows. The flight conditions shown in the figure 6.1 describe a rotor that is not moving relatively to the airflow.

The flow properties (figure 6.2) describe the type of atmosphere that surrounds the rotor.

Moreover, these are the values necessary to calculate the Reynolds number and set a unit system, in this thesis we shall use metric units. Lastly, the wall properties help determine whether the rotor is completely isolated or if there
Figure 6.6: Twist vs $r/R$ graph used to produce the shape of the blade. The data points had to be manually entered from the text file produced by the Matlab program to RotCFD.

<table>
<thead>
<tr>
<th>Flight Condition Type</th>
<th>General</th>
</tr>
</thead>
<tbody>
<tr>
<td>Free Stream Velocity</td>
<td>$0.0000$ $x$</td>
</tr>
<tr>
<td></td>
<td>$0.0000$ $y$</td>
</tr>
<tr>
<td></td>
<td>$0.0000$ $z$</td>
</tr>
</tbody>
</table>

Table 6.1: Flight conditions for all simulations.

is a "wall" to consider (figure 6.3). For these simulations, a non-slip condition was placed on the floor while keeping the other walls at atmospheric pressure and far away, 5D each side, to make them “invisible” to the rotor.

Additionally, the rotor configuration for the simulations had to match that of the real helicopter used for the experiments. Thus, once the geometry of each blade was scanned and introduced, the rotor properties where as shown on figure 6.7. The two blades and the radius are both part of the full size requirements and the scaling down. Meanwhile, the rotor tip speed was determined by the manufacturing company of the motors used to rotate the
rotors. However, as previously mentioned in the 30% model design chapter, it is consistent with the tip speed required. It is also important to point out that the rotation of one of the rotor disks was opposite to that of the other rotor.

### 6.4.1 Grid Production and Time Step Selection

The Grid selection was designed in order to maximize the number of cells surrounding the area of interest. The number of cells was limited to the memory availability on the graphics processor available. For this thesis all simulations were done using a Tesla C2070 Nvidia Graphics Card with a total memory of 5376 MB. Each simulation had the external boundary, which contained the

<table>
<thead>
<tr>
<th>Static Density:</th>
<th>1.225</th>
</tr>
</thead>
<tbody>
<tr>
<td>Static Temperature:</td>
<td>288.16</td>
</tr>
<tr>
<td>Gas constant:</td>
<td>287.05</td>
</tr>
<tr>
<td>Specific Heat Ratio:</td>
<td>1.4</td>
</tr>
<tr>
<td>Dynamic Viscosity:</td>
<td>$1.75 \times 10^{-5}$</td>
</tr>
<tr>
<td>Static Pressure:</td>
<td>101325.0</td>
</tr>
<tr>
<td>Calculate Viscosity</td>
<td>[Standard Units]</td>
</tr>
<tr>
<td>Calculate Viscosity</td>
<td>[Standard Units]</td>
</tr>
</tbody>
</table>

Table 6.2: Flight properties of all simulations

<table>
<thead>
<tr>
<th>X-Min type</th>
<th>Static Pressure</th>
</tr>
</thead>
<tbody>
<tr>
<td>X-Min Pressure</td>
<td>101325.0</td>
</tr>
<tr>
<td>X-Max type</td>
<td>Static Pressure</td>
</tr>
<tr>
<td>X-Max Pressure</td>
<td>101325.0</td>
</tr>
<tr>
<td>Y-Min type</td>
<td>Static Pressure</td>
</tr>
<tr>
<td>Y-Min Pressure</td>
<td>101325.0</td>
</tr>
<tr>
<td>Y-Max type</td>
<td>Static Pressure</td>
</tr>
<tr>
<td>Y-Max Pressure</td>
<td>101325.0</td>
</tr>
<tr>
<td>Z-Min type</td>
<td>Viscous Wall</td>
</tr>
<tr>
<td>Z-Max type</td>
<td>Static Pressure</td>
</tr>
<tr>
<td>Z-Max Pressure</td>
<td>101325.0</td>
</tr>
</tbody>
</table>

Table 6.3: Wall properties of all simulations
Figure 6.7: Rotor Properties used for all simulations.

information about the helicopters surroundings, and two refinement boxes. The refinement boxes did not input any additional information, they were strictly design to increase the density of cells, and resolution, in the areas of interest.

Figures 6.8a and 6.8b illustrate the grid density at different points of the simulations. It is important to point out that the grid was not “fitted to body” as it generally is. The reason for this was that given the cubed nature of the fuselage, it was not necessary to do so and the grid calculations were less likely to diverge if this was the case. The second refinement box, illustrated in both figures as the white line closest to the body, was selected in order to have a denser grid near the areas of interest, i.e. where the boards were placed on the real experiments.

On the other hand, the Time Step selection was directly related to the rotating speed of the MHS scaled model of 2500 RPM. This was measured experimentally using a hand-held digital Tachometer by placing a reflective
CHAPTER 6. CFD FLOW FIELD SIMULATIONS

Figure 6.8: Grid visualization of both planes where the outer white line represents the first refinement box while the inner white line is the second refinement box.

sticker on the tip of one of the blades and holding the tachometer in a set position, perpendicular to and directly above the reflective sticker. The experiment was done whilst it was being held still to avoid the helicopter from flying and moving. The Courant-Friedrich-Lewy condition (CFL) is a necessary condition in order to assure the numerical scheme used to solve the partial differential equation will be stable [18]. The CFL condition states

\[ \frac{u \Delta t}{\Delta x} \leq C_{\text{max}}, \]  

(6.1)

where \( u \) is the characteristic velocity of the phenomena, and \( \Delta t \) and \( \Delta x \) are the minimum time and distance steps respectively. \( C_{\text{max}} \) is the maximum Courant number which, given we are using a fully implicit scheme, is 1. Therefore, if we make \( u = u_{\text{tip}} \) then the simulations time step is selected to be the time it took for one of the rotor blades to move 1°, which was: \( 6.667 \times 10^{-5} \) s. This time step was selected in order to give enough time for the rotor to move enough in order to have a different result from the previous step, without losing too much information of what happened in between. Furthermore, this time step was the smallest time step that would not increase the total simulation real time past four days. This led to the following time settings shown in figure 6.9.

Figure 6.9 shows the time grid for all simulations. It also points out the selection of the “time scheme”, in this case fully implicit, that is related to the time discretization method used by the program. The total time simulated was of 1 second. This was primarily limited by the computation time available.
Figure 6.9: Time step setting for all simulations.

It is important to keep in mind that thanks to the CFL condition, the minimum grid cube length was now restricted.
Chapter 7

Flow field, Dust and Tuft Testing Results

Interesting results can be extracted from each individual test and simulations. However, there is much to gain from the comparison between the three.

7.1 Dust Testing Results

The results of the white light experiment are qualitative results. By selecting frames from the videos, we can appreciate that some dust particles are being moved. An example of this is shown in figure 7.1 (bottom), the first frame at which CWS saltation is observed (wheels on ground). Figure 7.1 (top) is the pixel difference between two frames, the frame where saltation is first observed minus the initial frame (no saltation).

However, it is important to point out that the dust cloud that was created appeared to have a very low density. This can be seen in figure 7.2 0.34 s after saltation started, where the bottom image is a picture of the dust cloud at its maximum, which is hard to capture with the camera. The top image is the difference between an image of the test bed before the helicopter model was turned on and that of the test bed 0.7 s after it was turned on. This was done in order to visualize the results a lot easier since the white dots are the CWS cloud generated. Given that the test were done in optimal conditions for saltation, we can thus conclude that the helicopter, which will be expose to less than ideal conditions during its flight, will not create a dense enough cloud to pose an obstruction problem to the altimeter (brownout).
Figure 7.1: Difference between a frame with no CWS movement and that of the first frame where there is movement (0.36 s since helicopter was turned on) (top). 1st Still frame of CWS in saltation (bottom).

From the laser wall experiments, we see that the distance at which saltation occurs is less than the distance at which, according to the simulations,
7.1. DUST TESTING RESULTS

Figure 7.2: Difference between a frame with no CWS movement and that of the eleventh frame where there is movement (0.70 s since helicopter was turned on) (top). 11th Still frame of CWS in saltation (bottom).

the maximum wake speed is reached. Thus, it may be concluded that it is not necessary to reach such high speeds in order to induce saltation. This
is concurs with the theory since, as was mentioned in the introduction, the speed necessary to induce saltation is much less. To see all the dust testing images see appendix A.

7.2 Tuft Testing Results

The results from these tests are presented with a single photograph for each configuration, i.e. height and plane inclination. Each picture has an exposure time of 2 seconds, thus showing some degree of tuft coning. Where coning is the production of a cone-like shape caused by the movement of the tuft around a point centre. The pictures selected were all taken in the first 2 seconds of the experiment. This was done to compare the images to the simulations.

It is important to point out that the simulation snapshot showing the horizontal plane (such as Figure 7.3 etc.) have been dewarped in order to have an image as if the camera were directly above it. Placing the camera directly above was not a feasible solution since the body of the helicopter would “block” the area of interest for this study: that directly under the rotor.

Given the lack of time, the only angle of inclination tested for heights 15 cm and 3R, apart from zero, was 30°. To see all the tuft testing results go to appendix B.

7.3 Simulation Results

The results of the simulations are shown with a single still image of when the simulation was at exactly 1 second. The colour scales in the bottom right corner of every picture illustrates the velocity magnitude. The white squares that are next to or under the rotor represent the location at which the boards on the tuft experiments where. This was done in order to facilitate the comparison. Moreover, some of the vectors on the images are white, this means that they are exceeding the maximum value shown on scale. The scale was selected for each individual picture in order to gather the most detailed information of the area of interest.

A main distinction between the images presented for all angles of height “15 cm drop” and 3R, and those of height “wheels on ground” are the lack
7.4 Comparison

When comparing the results from the simulations and the tuft testing we see a good resemblance between the simulations and the tuft experiments. Meaning that the direction suggested by the simulation falls comfortably inside of the tuft conning (Figure 7.7, 7.8 and 7.9). It is also possible to see how the rotor wake becomes thinner, the further away you get from the rotor. This last statement is mostly noticeable for the $3R$ height (Figure 7.9).

Thus, seeing these comparisons, one can begin to analyse the results. When looking at the ground plane, there is an almost symmetrical pattern of tufts/vectors pointing radially outwards from where the rotor wake hits the surface (Figures 7.10, 7.11 and 7.12). Furthermore, in these simulations in all cases where there was no inclination the maximum velocity observed falls outside of the area directly below the rotor disk. Simulations suggest an azimuthal symmetry and a radial dependency of the speed. Therefore
justifying the placement of a single vertical board during the tuft testing to visualize the entire phenomena. Moreover, by comparing the ground plane
7.4. COMPARISON

(a) Vertical Big Board  (b) Vertical Small Board
(c) Ground Board

Figure 7.6: Example of the flow field visualization taken from the simulation with height “wheels on ground” and plane inclination zero.

simulations, it is evident that the inclination does not affect the maximum velocity reached at ground level (see figure 7.13).

Unlike the previous conclusion, the heights appear to influence the maximum speed in an inversely proportion. By comparing figures 7.14a, 7.14b and 7.14c it can be concluded that the maximum wake velocity decreases as we get further away from the floor.

For height $3R$, it can be seen that there are vortices forming at ground level (figure 7.9). This is specially interesting considering that it is the one height “out-of-ground effect”. It is important to point out that the other two heights also present a vortex. Yet these occur sooner, at 0.35 s as opposed to the 1 s of the $3R$ height.

What is more, the maximum velocities reached in all cases of zero incli-
CHAPTER 7. FLOW FIELD, DUST AND TUFT TESTING RESULTS

Figure 7.7: Flow field visualization of vertical plane with height “wheels on ground” and plane inclination zero(right). Tuft pattern visualization of vertical plane, under the rotor, with height “wheels on ground” and plane inclination zero(left).

Figure 7.8: Flow field visualization of vertical plane with height “15cm” and plane inclination zero (left). Tuft pattern visualization of vertical plane, under the rotor, with height “15cm” and plane inclination zero (right).

nation were all not at a same distance from the centre of the rotor (table 7.1).

The maximum velocity ring would shift from directly under the rotor as the inclination of the plane would increase. This ring seemed to increase in diameter as the heights grew taller and would also start distorting slightly into an oval. Despite this not being obvious in the tuft testing results, it is clear in the simulations.

An other important fact to point out from the previous table, (table 7.1) is how the distance at which the CWS begins saltation also increases as the
7.4. COMPARISON

Figure 7.9: Flow field visualization of vertical plane with height $3R$ and plane inclination zero (left). Tuft pattern visualization of vertical plane, under the rotor, with height $3R$ and plane inclination zero (right).

<table>
<thead>
<tr>
<th>Height 1</th>
<th>Height 2</th>
<th>Height 3</th>
</tr>
</thead>
<tbody>
<tr>
<td>$u_{\text{max}}$</td>
<td>149.2 mm ±0.5 mm</td>
<td>183.7 mm ±0.4 mm</td>
</tr>
<tr>
<td>Saltation</td>
<td>116.6±0.2 mm</td>
<td>166.8±0.2 mm</td>
</tr>
</tbody>
</table>

Table 7.1: Comparison between the distances at which the simulations reach their maximum velocity and the distance at which the CWS particles first began their saltation process in the dust laser wall experiment. These last distances were obtained by doing a particle image velocimetry (PIV) analysis.

height of the rotor increases. Furthermore, when comparing the velocities at these points given by the PIV and the simulations we see how in each case the velocity is of 2.6 m/s (figures 7.15, 7.16 and 7.17).

On the other hand, it is important to point out that the wind velocity rapidly decreases after its maximum (Figures 7.7, 7.8 and 7.9). This would suggest that although it may be enough to initiate saltation, the amount of dust lifted may not be significant. Meaning that the dust cloud would not be very dense, which can be easily seen in both dust experiments (for example figures 7.15 and 7.2).
Figure 7.10: Tuft pattern visualization of horizontal plane, under the rotor, with height “wheels on ground” and plane inclination zero (top). Flow field visualization of horizontal plane with height “wheels on ground” and plane inclination zero (bottom).
Figure 7.11: Tuft pattern visualization of horizontal plane, under the rotor, with height “15cm” and plane inclination zero (top). Flow field visualization of horizontal plane with height “15cm” and plane inclination zero (bottom).
Figure 7.12: Tuft pattern visualization of horizontal plane, under the rotor, with height “3R” and plane inclination zero (top). Flow field visualization of horizontal plane with height “3R” and plane inclination zero (bottom).
7.4. COMPARISON

Figure 7.13: Flow field visualization of the horizontal planes with height “wheels on ground” and plane inclination of 5° (left) and plane inclination of 20° (right).

(a) Wheels on ground  (b) 15 cm

(c) 3R

Figure 7.14: Flow field visualization of the vertical plane with no angle of inclination for all three heights.
CHAPTER 7. FLOW FIELD, DUST AND TUFT TESTING RESULTS

Figure 7.15: Comparison between simulation at 0.36 seconds and the PIV analysis done to laser wall experiment at height “wheels on ground”.

Figure 7.16: Comparison between simulation at 0.45 seconds and the PIV analysis done to laser wall experiment at height “15cm”.

Figure 7.17: Comparison between simulation at 0.56 seconds and the PIV analysis done to laser wall experiment at height “3R”.
Chapter 8

Conclusion

The experiments and simulations in this thesis were all done in optimal conditions for dust saltation, while staying as close to the “real problem” as possible. That is, that despite the helicopter model being a close scaled representation of what is planned to be the full-scale MHS, the environment surrounding it different. The density and composition of the atmosphere were the main differences, since gravity was accounted for in the selection of the “Martian dust” emulator, crushed wallnuts. From the simulations done we see that the case in which the air reaches its maximum velocity is when the plane is at inclination zero and height “wheels on ground” (figures 7.14a, 7.6b and 7.6c).

In conclusion, the maximum wind speed the model reached at ground level is: 3.326 m/s (figure 7.14a). The direction of this velocity is parallel to the ground and its magnitude is much larger than the theoretically necessary to induce saltation at Martian atmospheric density (0.5366 m/s) and at Earth’s (0.02261 m/s). However, saltation begins at a distance greater than the rotor radius at all heights. Therefore, given that the cloud density was very low (for example figure 7.2) and begins far from the rotor, it is possible to say with the information at hand, that the dust cloud created by the MHS will not pose a problem when taking off and landing, since it will not be dense enough to block or obstruct the sensors.

Despite these tests suggesting a positive result regarding the lack of danger caused by the cloud, further tests are recommended. The next suggested test would be repeating both the dust test and the tuft test at reduced pressures. This would give us more information regarding saltation at low atmospheric densities. Moreover, dust simulations using RotCFD’s dust unit
are recommended for both the dust testing presented in this thesis and those proposed at lower pressure.
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• NASA I² Program
• Agencia Espacial Mexicana
• Sofia Ixchel Michaelian Martínez
Appendices
Appendix A

Dust Results

A.0.1 White Light Experiment
Figure A.1: Difference between a frame with no CWS movement and that of the first frame where there is movement (0.36 s since helicopter was turned on) (top). 1st Still frame of CWS in saltation (bottom).
Figure A.2: Difference between a frame with no CWS movement and that of the second frame where there is movement (0.40 s since helicopter was turned on) (top). 2nd Still frame of CWS in saltation (bottom).
Figure A.3: Difference between a frame with no CWS movement and that of the third frame where there is movement (0.43 s since helicopter was turned on) (top). 3rd Still frame of CWS in saltation (bottom).
Figure A.4: Difference between a frame with no CWS movement and that of the forth frame where there is movement (0.46 s since helicopter was turned on) (top). 4th Still frame of CWS in saltation (bottom).
Figure A.5: Difference between a frame with no CWS movement and that of the fifth frame where there is movement (0.50 s since helicopter was turned on) (top). 5th Still frame of CWS in saltation (bottom).
Figure A.6: Difference between a frame with no CWS movement and that of the sixth frame where there is movement (0.53 s since helicopter was turned on) (top). 6th Still frame of CWS in saltation (bottom).
Figure A.7: Difference between a frame with no CWS movement and that of the seventh frame where there is movement (0.56 s since helicopter was turned on) (top). 7th Still frame of CWS in saltation (bottom).
Figure A.8: Difference between a frame with no CWS movement and that of the eighth frame where there is movement (0.60 s since helicopter was turned on) (top). 8th Still frame of CWS in saltation (bottom).
Figure A.9: Difference between a frame with no CWS movement and that of the ninth frame where there is movement (0.63 s since helicopter was turned on) (top). 9th Still frame of CWS in saltation (bottom).
Figure A.10: Difference between a frame with no CWS movement and that of the tenth frame where there is movement (0.66 s since helicopter was turned on) (top). 10th Still frame of CWS in saltation (bottom).
Figure A.11: Difference between a frame with no CWS movement and that of the eleventh frame where there is movement (0.70 s since helicopter was turned on) (top). 11th Still frame of CWS in saltation (bottom).
Figure A.12: Difference between a frame with no CWS movement and that of the twelfth frame where there is movement (0.73 s since helicopter was turned on) (top). 12th Still frame of CWS in saltation (bottom).
A.0.2 Laser Wall Experiment

Figure A.13: Vertical Laser plane visualization of the CWS at 0.366 s since helicopter was turned on (Right). It is the first moment at which saltation is seen for height “wheels on ground”. The left image is exactly two frames after, at 0.370 s, and it is the first moment in which the PIV program detects the particle and assigns it a value of 2.6 m/s.

Figure A.14: Vertical Laser plane visualization of the CWS at 0.436 s since helicopter was turned on (Right). It is the first moment at which saltation is seen for height “15 cm”. The left image is exactly one frame after, at 0.438 s, and it is the first moment in which the PIV program detects the particle and assigns it a value of 2.6 m/s.
Figure A.15: Vertical Laser plane visualization of the CWS at 0.556 s since helicopter was turned on (Right). It is the first moment at which saltation is seen for height $3R$. The left image is exactly one frame after, at 0.558 s, and it is the first moment in which the PIV program detects the particle and assigns it a value of 2.6 $m/s$. 
Appendix B

Tuft test Results
Figure B.1: Tuft pattern visualization of the vertical and horizontal boards with height “wheels on ground” and plane inclination zero.
Figure B.2: Tuft pattern visualization of the vertical and horizontal boards with height “wheels on ground” and plane inclination of 5°.
Figure B.3: Tuft pattern visualization of the vertical and horizontal boards with height “wheels on ground” and plane inclination of 10°.
Figure B.4: Tuft pattern visualization of the vertical and horizontal boards with height “wheels on ground” and plane inclination of 15°.
Figure B.5: Tuft pattern visualization of the vertical and horizontal boards with height "wheels on ground" and plane inclination of 20°.
Figure B.6: Tuft pattern visualization of the vertical and horizontal boards with height “wheels on ground” and plane inclination of 25°.
Figure B.7: Tuft pattern visualization of the vertical and horizontal boards with height “wheels on ground” and plane inclination of 30°.
Figure B.8: Tuft pattern visualization of the vertical and horizontal boards with height “15 cm drop” and plane inclination zero.
Figure B.9: Tuft pattern visualization of the vertical and horizontal boards with height “15 cm drop” and plane inclination of 30°.
Figure B.10: Tuft pattern visualization of the vertical and horizontal boards with height $3R$ and plane inclination zero.
Figure B.11: Tuft pattern visualization of the vertical and horizontal boards with height $3R$ and plane inclination of $30^\circ$. 
Appendix C

Simulation Results

Figure C.1: Flow field visualization of the vertical and horizontal planes with height “wheels on ground” and plane inclination zero.
Figure C.2: Flow field visualization of the vertical and horizontal planes with height “wheels on ground” and plane inclination of 5°.
Figure C.3: Flow field visualization of the vertical and horizontal planes with height “wheels on ground” and plane inclination of 10°.
Figure C.4: Flow field visualization of the vertical and horizontal planes with height “wheels on ground” and plane inclination of 15°.
Figure C.5: Flow field visualization of the vertical and horizontal planes with height “wheels on ground” and plane inclination of 20°.
Figure C.6: Flow field visualization of the vertical and horizontal planes with height “wheels on ground” and plane inclination of 25°.
Figure C.7: Flow field visualization of the vertical and horizontal planes with height “wheels on ground” and plane inclination of $30^\circ$.
Figure C.8: Flow field visualization of the vertical and horizontal planes with height “15 cm drop” and plane inclination zero.

Figure C.9: Flow field visualization of the vertical and horizontal planes with height “15 cm drop” and plane inclination of 5°.
Figure C.10: Flow field visualization of the vertical and horizontal planes with height “15 cm drop” and plane inclination of 10°.

Figure C.11: Flow field visualization of the vertical and horizontal planes with height “15 cm drop” and plane inclination of 15°.

Figure C.12: Flow field visualization of the vertical and horizontal planes with height “15 cm drop” and plane inclination of 20°.
APPENDIX C. SIMULATION RESULTS

Figure C.13: Flow field visualization of the vertical and horizontal planes with height “15 cm drop” and plane inclination of 25°.

Figure C.14: Flow field visualization of the vertical and horizontal planes with height “15 cm drop” and plane inclination of 30°.
Figure C.15: Flow field visualization of the vertical and horizontal planes with height $3R$ and plane inclination zero.

Figure C.16: Flow field visualization of the vertical and horizontal planes with height $3R$ and plane inclination of $5^\circ$.

Figure C.17: Flow field visualization of the vertical and horizontal planes with height $3R$ and plane inclination of $10^\circ$. 
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Figure C.18: Flow field visualization of the vertical and horizontal planes with height $3R$ and plane inclination of $15^\circ$.

Figure C.19: Flow field visualization of the vertical and horizontal planes with height $3R$ and plane inclination of $20^\circ$.

Figure C.20: Flow field visualization of the vertical and horizontal planes with height $3R$ and plane inclination of $25^\circ$.
Figure C.21: Flow field visualization of the vertical and horizontal planes with height $3R$ and plane inclination of $30^\circ$. 

(a) Vertical Big Board (Left)  
(b) Vertical Big Board (Right)  
(c) Ground Board
Bibliography


