Computational Study of the Side-by-Side Urban Air Taxi Concept

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High-fidelity computational fluid dynamics simulations of NASA’s Side-by-side air taxi concept have been carried out. The three-dimensional unsteady Navier-Stokes equations are solved on overset grids using high-order accurate schemes, dual-time stepping, and a hybrid turbulence model. The flow solver has been loosely coupled with a helicopter comprehensive analysis code in order to get the trimmed flight solution. The vehicle simulated is a six-passenger side-by-side intermeshing rotor helicopter with hybrid propulsion for air taxi operations, also known as urban air mobility applications. This concept vehicle is intended to focus and guide NASA research activities in support of aircraft development for emerging aviation markets, in particular vertical take-off and landing air taxi operations.

NOMENCLATURE

\( a \) Fluid speed of sound  
\( A \) Rotor disk area, \( \pi R^2 \)  
\( c \) Local rotor blade chord length  
\( C' \) Sectional chord force  
\( c_{\text{tip}} \) Rotor blade tip chord length  
\( C_Q \) Torque coefficient, \( \frac{Q}{(\Omega R)^2 A} \)  
\( C_T \) Thrust coefficient, \( \frac{T}{(\rho \Omega R)^2 A} \)  
\( D \) Rotor diameter  
\( F/M \) Figure of merit, \( \frac{C_{Q}^{3/2}}{\sqrt{2} C_{Q}} \)  
\( M' \) Sectional blade pitching moment  
\( M \) Mach number, \( \frac{V}{a} \)  
\( M_{\text{tip}} \) Mach number at the blade tip, \( \frac{\Omega R}{a} \)  
\( M^2 c_c \) Sectional chord force coefficient, \( \frac{C}{\frac{2}{5} \rho a^2 c} \)  
\( M^2 c_m \) Sectional pitching moment coefficient, \( \frac{M'}{\frac{2}{5} \rho a^2 c^2} \)  
\( M^2 c_n \) Sectional normal force coefficient, \( \frac{N'}{\frac{2}{5} \rho a^2 c} \)  
\( N' \) Sectional blade normal force  
\( NB \) Near-body  
\( OB \) Off-body  
\( r \) Radial position  
\( R \) Rotor radius  
\( Re \) Reynolds number, \( \frac{V L_{\text{ref}}}{\nu} \)  
\( Re_{\text{tip}} \) Reynolds number at the blade tip, \( \frac{\Omega R c_{\text{tip}}}{\nu} \)  
\( V_{\infty} \) Freestream velocity  
\( y^+ \) Non-dimensional viscous wall spacing  
\( \alpha \) Angle of attack, AoA  
\( \delta \) Boundary layer thickness  
\( \Delta \) Grid spacing  
\( \mu \) Advance ratio, \( \frac{V_{\infty} \cos(\alpha)}{\Omega R} \)  
\( \nu \) Fluid kinematic viscosity  
\( \rho \) Fluid density  
\( \Omega \) Rotor rotational speed

INTRODUCTION

Electric multirotor vehicles have grown very popular over the past decade. While originally conceived for military purposes, their simplicity and affordability have made small drones, or Unmanned Aerial Vehicles (UAVs), accessible to the civil market. They are being used for photography, film recording, and also for recreational purposes, to name just a few applications. In addition, the unique ability of multi-rotor vehicles for vertical lift has great potential for human and cargo transportation in urban areas.

NASA and a community of government, industry, and academic partners are working together on this goal, known as Urban Air Mobility (UAM). UAM is a safe and efficient air transportation system where everything from small package delivery drones to passenger-carrying air taxis operate over populated areas, from small towns to the largest cities. Figure 1 shows an artist’s impression of a UAM environment, where different air vehicles move and interact safely in the airspace.

Relying on either fully electric or hybrid propulsion systems, the eVTOL (electrical Vertical Take-Off and Landing) vehicles would transport a small number of passengers from point to point in highly congested cities, avoiding ground traffic and potentially providing a greener means of transportation. Their
capacity to hover and to perform VTOL, together with their
great maneuverability, make multi-rotor vehicles an excellent
choice for UAM aircraft.

Accurate prediction of rotorcraft performance and acoustics
continues to be challenging from a computational point of
view. The flows are inherently unsteady, nonlinear, and com-
plex. The aerodynamic interactions between multiple rotors,
fuselage, and lifting bodies make it even more difficult. High-
fidelity Computational Fluid Dynamics (CFD) approaches of-
fer an advantage over low- and mid-fidelity tools for this type
of flow. In addition, high-fidelity simulations can provide the
information needed to calibrate lower fidelity tools that could
be used for design purposes.

The designs of conventional helicopters and airplanes have
been well established and optimized during the twentieth cen-
tury. However, the design space of eVTOL aircraft remains
wide open, and exploration for improved designs continues.
Previous computational work on small multi-rotor vehicles by
Ventura Diaz et al. (Refs. 1, 2) and Yoon et al. (Refs. 3, 4)
includes studies of the aerodynamic interactions, rotor place-
ment, torque balanced coaxial octo-rotors, and the effect of
wind gusts on a hovering quad-rotor. Johnson et al. (Ref. 5)
and Silva et al. (Ref. 6) have recently conducted a compre-
hensive study for air taxi operations using multidisciplinary
design, analysis, and optimization techniques, coming up with
several multi-rotor air taxi vehicle concepts.

This work focuses on analysis and simulations of one of these
vehicle concepts: the Side-by-Side (SbS) UAM vehicle. The
SbS vehicle has two overlapping intermeshing rotors, as rep-
resented in Figure 2. It is a six-passenger air taxi vehicle, with
a range of 200 nm, thus performing four 50 nm trips without
refueling. A typical UAM mission consists of multiple 50 nm
segments (Ref. 5). The SbS parallel hybrid propulsion sys-

THEORY

The performance of a rotor system can be examined by means
of the one-dimensional momentum theory (Ref. 7). This the-
ory uses the fluid mass, momentum, and energy conservation
equations in an integral form. By using momentum theory, the
ideal power required to hover can be expressed as a function
of the thrust $T$ and the disk area $A$:

$$P_i = T v_i = T \sqrt{\frac{T}{2 \rho A}}$$  \hspace{1cm} (1)

where $v_i$ is the induced velocity at the rotor disk. The disk
loading is defined as $DL = T/A$. The power loading is the
ratio of thrust to power, $PL = T/P$. Therefore, the induced
(ideal) power loading is simply $(PL_i)^{-1} = P_i/T = v_i$.

In helicopter analysis, the rotor thrust coefficient $C_T$, the rotor
power coefficient $C_p$ and the rotor torque coefficient $C_Q$ are
formally defined as:

$$C_T = \frac{T}{\rho A \Omega^2 R^2}$$  \hspace{1cm} (2)

$$C_p = \frac{P}{\rho A \Omega^3 R^3}$$  \hspace{1cm} (3)
\[ C_Q = \frac{Q}{\rho A \Omega^2 R} \]  \hspace{1cm} (4)

where \( \rho \) is the flow density, \( A \) is the rotor disk area (for one rotor), \( \Omega \) is the rotational speed and \( R \) is the radius of the blade. With \( P = \Omega Q \), \( C_P \) and \( C_Q \) are identical.

The figure of merit (FM) is equivalent to a static thrust efficiency and is defined as the ratio of the ideal power required to hover to the actual power required, and it can be expressed as a function of the thrust and torque coefficients using momentum theory:

\[ FM = \frac{\text{Ideal power required to hover}}{\text{Actual power required to hover}} = \frac{C_P}{C_Q} = \frac{C_{T/Q}^{3/2}}{\sqrt{2} C_Q} \hspace{1cm} (5) \]

Accurate prediction of the FM using high-fidelity CFD and hybrid turbulence models has been validated in previous studies (Ref. 8).

The operation of two or more rotors in close proximity will modify the flow field at each, and hence the performance of the rotor system will not be the same as for the isolated rotors. The limit case is the coaxial rotor, with just one-half the disk area of the isolated rotors and therefore twice the disk loading. It follows that by operating the rotors coaxially the induced power required is increased by a factor of \( \sqrt{2} \), a 41\% hover induced power increase.

For hovering side-by-side twin rotors with overlap area \( mA \), the hover performance estimate can be based on the effective disk loading of the system as a whole: \( \frac{P}{P_{iso}} = \left( \frac{T}{T_{iso}} \right)^{3/2} \left( \frac{2}{2 - m} \right)^{1/2} \hspace{1cm} (6) \)

and for the same thrust the interference power is:

\[ \frac{\Delta P}{P} = \left( \frac{2}{2 - m} \right)^{1/2} - 1 \hspace{1cm} (7) \]

In the coaxial limit this gives \( \Delta P/P = 0.41 \) as required.

Note that, with a shaft separation \( d \), as seen in Figure 3, the overlap fraction is:

\[ m = \frac{2}{\pi} \left[ \cos^{-1} \left( \frac{d}{2R} \right) - \frac{d}{2R} \sqrt{1 - \left( \frac{d}{2R} \right)^2} \right] \hspace{1cm} (8) \]

**NUMERICAL APPROACH**

The flow solver used in this study is NASA’s Overflow (Ref. 9) CFD solver. Overflow is a finite-difference, structured overset grid, high-order accurate Navier-Stokes flow solver. NASA’s Chimera Grid Tools (CGT) (Ref. 10) overset grid generation software is used for generating the overset grids of rotors and complete vehicles. Body-fitted curvilinear near-body (NB) grids are generated using CGT. The computational domain is completed with the generation of Cartesian off-body (OB) grids that are automatically generated prior to grid assembly using the domain connectivity framework in Overflow-D mode. The current time-accurate approach consists of an inertial coordinate system where NB curvilinear O-grids for the rotor blades rotate through the fixed OB Cartesian grid system. Overflow is coupled in a loosely manner with the helicopter comprehensive code CAMRAD II (Ref. 11). The CFD provides high-fidelity, nonlinear aerodynamics that replace the comprehensive lifting line aerodynamic analysis from CAMRAD II. The comprehensive code performs the structural dynamics and trim calculations and gives the information to Overflow. The loose coupling allows for a modular approach and communication through input/output. The coupling methodology has been implemented following the approach of Potsdam et al. (Ref. 12).

The numerical approach and the coupling process are described below.

**Overset Grid Generation**

The overset grid generation process using CGT can be divided into the following steps: geometry processing, surface grid generation, volume grid generation, and domain connectivity (Ref. 10).

The geometry is usually obtained from a Computer Aided Design (CAD) model. Figure 4 (a) shows the CAD model of the SbS vehicle. The Boundary Representation (BRep) is an object that holds both the topological entities and the geometric components (Ref. 13). In this work, a pre-processing step converts the analytical BRep solid from a STEP or IGES file into discrete representations for the BRep faces and edges. Access to the model topology and entities is accomplished through EGADS (the Engineering Geometry Aerospace Design System) API which is a foundational component of the Engineering Sketch Pad (Ref. 13). With the egads2srfg tool, discrete representations are generated from each solid. Figure 4 (b) shows the structured surface grid file for the SbS obtained using EGADS. This grid file contains an untrimmed structured patch for each face based on tessellation of the face parameter space. Another file obtained with egads2srfg is a curve grid.
Fig. 4: The side-by-side fuselage. Figure (a) shows the CAD geometry, Figure (b) shows the structured untrimmed patches obtained from the CAD geometry using EGADS. The patches are used as reference surfaces to generate the overset surface grids.

Overset structured surface meshes are typically created using a combination of algebraic and hyperbolic methods, depending on the number of initial curves on each surface domain. The generation of surface grids is the step that requires the most manual effort and experience from the user. Figure 5 shows the surface grids for the SbS vehicle with the rotors. For this study, the landing gear and the strakes have been omitted. Future work will include these elements and their effect when flying near the ground. The grid system consists of the main fuselage, the wing, the engines, the empennage, the rotors, and the rotor hubs.

With sufficient overlap between surface grids, the volume grids can be created easily with hyperbolic marching methods out to a fixed distance from the surface. Such methods provide orthogonal grids with tight clustering characteristics at the wall, which is essential for accurately capturing the boundary layer in viscous flow computations. The distance is chosen such that the outer boundaries of the NB volume grids are well clear off the boundary layer. The NB grids are then embedded inside OB Cartesian grids that extend to the far field. Figure 6 (a) shows the near-body grids for the SbS with rotors, in Figure 6 (b) the NB and the OB Cartesian grids are shown.

Fig. 5: Side-by-side overset surface grids, top view.

Fig. 6: Side-by-side overset volume grids. Figure (a) shows the boundary edges of the near-body volume grids. Figure (b) shows the near- and off-body Cartesian volume grids.
Off-body Cartesian grids with uniform spacing surround the NB grids to resolve the wake region of interest. Coarser Cartesian grids efficiently expand the grid system to the far field, where each successive Cartesian grid is twice as coarse as its previous neighbor. The far field boundary is 20 rotor radii away from the center of the vehicle in all directions. The resolved wake region has a uniform grid spacing of 10% of the tip chord length $c_{tip}$.

While the airframe grids have been generated using the CAD model shown in Figure 4 (a) as reference, the blade grids are generated using the rotor geometry information from Table 1. The profiles used to build the blade are the VR12 airfoils from $r = 0$ to $r = 0.85R$, and the SSCA09 airfoils from $r = 0.95R$ to the tip, $r = R$. The transition between the two different airfoil sections is smooth (linear interpolation with radial station). The blade is tapered and swept near the tip. Figure 7 shows in detail the geometry and the overset grids of the SbS blade. Surface grid resolution on the rotor blades is clustered in the chordwise direction near the airfoils leading and trailing edges to accurately resolve large pressure gradients. Subsequently, the spanwise resolution is clustered near the root and the tip. The normal grid spacing of all grids at the walls maintains $y^+ < 1$.

<table>
<thead>
<tr>
<th>Number of rotors</th>
<th>2</th>
</tr>
</thead>
<tbody>
<tr>
<td>Number of blades/rotor</td>
<td>4</td>
</tr>
<tr>
<td>Radius, $R$</td>
<td>3.203 m</td>
</tr>
<tr>
<td>Linear twist</td>
<td>-16 deg</td>
</tr>
<tr>
<td>Root chord, $c_{root}$</td>
<td>0.217 m</td>
</tr>
<tr>
<td>Tip chord, $c_{tip}$</td>
<td>0.130 m</td>
</tr>
<tr>
<td>Rotor solidity, $\sigma$</td>
<td>0.0832</td>
</tr>
<tr>
<td>Nominal tip speed, $V_{tip}$</td>
<td>167.6 m/s</td>
</tr>
</tbody>
</table>

Table 1: Side-by-side rotor geometry properties.

The multi-rotor system consists of two overlapping rotors, each rotor made up of four blades. The right rotor rotates counter-clockwise (CCW), and the left rotor rotates clockwise (CW), therefore the advancing blade is always outboard. The rotor grids without airframe are shown in Figure 8. The overlapping distance can be changed and its effect on the airloads, performance, and wake geometry in cruise will be the main subject of study in the present work.

The rotor system consists of 30 NB grids with 25 NB million grid points, and 235 to 260 million total (NB + OB) grid points, depending on the overlapping distance. For the complete vehicle, there is a total of 74 NB grids with 40 NB million grid points, the total number of NB and OB grid points is 300 million.

By using a trimmed approach, the domain connectivity step is robust and highly automated: hole cutting is required between
components and with the OB Cartesian grids. In this study, the X-ray hole cutting method is used. An X-ray object is created for every component in the geometry (i.e. the blades, the hubs, the fuselage, the wings, etc.). The user has to supply the list of meshes that each X-ray object is allowed to cut, and an offset distance with which to grow each hole away from the body. The hole cutting process is performed at each time step within the flow solver, allowing for the rotation of the blades relative to the fixed components.

High-Order Accurate Navier-Stokes Solver

Overflow solves the Navier-Stokes equations using finite differences with a variety of numerical algorithms and turbulence models. The time dependent, Reynolds-averaged Navier-Stokes (RANS) equations are solved in strong conservation form.

\[
\frac{\partial \vec{q}}{\partial t} + \frac{\partial (\vec{F} - \vec{F}_v)}{\partial x} + \frac{\partial (\vec{G} - \vec{G}_v)}{\partial y} + \frac{\partial (\vec{H} - \vec{H}_v)}{\partial z} = 0 \tag{9}
\]

where \( \vec{q} = [p, \rho u, \rho v, \rho w, e]^T \) is the vector of conserved variables; \( \vec{F} \), \( \vec{G} \) and \( \vec{H} \) are the inviscid flux vectors; and \( \vec{F}_v \), \( \vec{G}_v \) and \( \vec{H}_v \) are the viscous flux vectors.

In this study, the diagonal central difference algorithm is used with the \( S^{\text{th}} \)-order accurate spatial differencing option with scalar dissipation. The physical time step corresponds to 0.25 degree rotor rotation, together with up to 50 dual-time sub-iterations, to improve time-accuracy by reducing the linearization errors. The numerical approach and time step were previously validated for various rotor flows (Refs. 8, 14, 15).

Hybrid Turbulence Modeling

The Overflow code has a choice of algebraic, one-equation, and two-equation turbulence models, including hybrid Reynolds-Averaged Navier-Stokes / Large Eddy Simulation (RANS/LES) models that close the RANS equations. In this study, the one equation Spalart-Allmaras (Ref. 16) turbulence model is used primarily within the boundary layer.

The intent of the Detached Eddy Simulation (DES) model (Ref. 16) is to be in RANS mode throughout the boundary layer, where the turbulent scales can be very small and need to be modeled, and in LES mode outside the boundary layer where the largest turbulent scales are grid-resolved. In this way, DES is a RANS/LES hybrid approach that mitigates the problem of artificially large eddy viscosity. The turbulence length scale \( \delta \) is replaced by \( \delta_r \), where \( \delta_r \) is the minimum of the distance from the wall, \( \delta \), and the local grid spacing times a coefficient.

The DES approach assumes that the wall-parallel grid spacing \( \Delta_z \) exceeds the thickness of the boundary layer \( \delta \) so that the RANS model remains active near solid surfaces. If the wall-parallel grid spacing is smaller than the boundary layer thickness, \( \Delta_z < \delta \), then the DES Reynolds stresses can become under-resolved within the boundary layer; this may lead to non-physical results, including grid-induced separation. Using Delayed Detached Eddy Simulation (DDES) (Ref. 17), the RANS mode is prolonged and is fully active within the boundary layer. The wall-parallel grid spacing used in this study does not violate the hybrid-LES validity condition; thus DES and DDES should give similar results. Nevertheless, all computations have been performed using the DDES model for both NB and OB grids.

Comprehensive Analysis

Structural dynamics and rotor trim for the coupled calculations are performed using the comprehensive rotorcraft analysis code CAMRAD II (Ref. 11). CAMRAD II is an aeromechanics analysis of rotorcraft that incorporates a combination of advanced technologies, including multibody dynamics, nonlinear finite elements, and rotorcraft aerodynamics. The trim task finds the equilibrium solution for a steady state operating condition, and produces the solution for performance, loads, and vibration. The aerodynamic model for the rotor blade is based on lifting-line theory, using two-dimensional airfoil characteristics and a vortex wake model. CAMRAD II has undergone extensive correlation with performance and loads measurements on rotorcraft.

Loose Coupling Overflow – CAMRAD II

A loose coupling approach between Overflow and CAMRAD II based on a trimmed periodic rotor solution is implemented. The comprehensive code provides the trim solution and blade motions. The high-fidelity CFD calculates the airloads. Figure 9 shows the flow diagram of the loose coupling strategy. In summary, the CFD airloads replace the comprehensive airloads while using lifting line aerodynamics to trim and computational structural dynamics to account for blade deformations. In this work, the SbS blades have flap hinges, but are considered rigid, and consequently there will not be elastic motions in the blades. The comprehensive analysis is still needed to get the trim angles and rigid blade motions for the CFD.

![Flow diagram for CFD/Comprehensive Analysis loose coupling methodology.](Image)
The iterative loose coupling process is summarized next. The simulation is initialized with a comprehensive analysis resulting in a trimmed rotor solution obtained with lifting line aerodynamics. This analysis creates initial quarter chord motions as a function of the radius $r$ and the azimuth $\psi$, for each rotor. In addition, the aircraft attitude is also obtained from CAMRAD II. The motions and aircraft pitch angle are given to the CFD. The CFD analysis accounts for the entire flow field, and therefore it only requires the structural motion. The CFD is run with the prescribed motions and angles, for two to three full rotor revolutions for the first coupling step. Overflow outputs the normal force $N$, pitching moment $M'$, and chord force $C'$ as a function of radius and azimuth.

Then, the aerodynamic force and moment coefficients increments ($\Delta c$) that are used in the comprehensive code at the next iteration $n+1$ are calculated. The increments are the difference between the CFD loads and the comprehensive lifting line solution required to trim from the previous step $n$, plus the load increments from the previous step:

$$\Delta c^{n+1} = \Delta c^n + (c_{CFD} - c_{total})$$

For the initial step the increments are the difference between CFD and the total loads from the $0^{th}$ run in CAMRAD II:

$$\Delta c^1 = (c_{CFD} - c_{total})$$

The sectional pitching moment $M^2 c_m$, normal force $M^2 c_n$ and chord force $M^2 c_c$ coefficients are defined as:

$$M^2 c_m = \frac{M'}{\frac{1}{2} \rho c^2}$$

$$M^2 c_n = \frac{N'}{\frac{1}{2} \rho c^2}$$

$$M^2 c_c = \frac{C'}{\frac{1}{2} \rho c^2}$$

With the new quarter chord motions of the retrimmed rotor and the new aircraft attitude, the CFD is rerun. The previous CFD flow solution is used as restart condition. The coupling is performed every half or quarter rotor revolution. The coupling solution is considered to be converged when collective and cyclic control angles and the CFD aerodynamic forces do not change between iterations. The CFD flow solution is usually converged after 10 to 20 rotor revolutions.

The coupling procedure is valid as long as the rotor loads are periodic. This approach is still good if there is some aperiodicity in the vortex wake, which is often the case in a high-resolution turbulent simulations.

The coupled Overflow–CAMRAD II simulations are completely automatized using a Python code and Fortran post-processing functions. All simulations were run on NASA’s supercomputer Electra, using from 1000 to 2000 cores on Broadwell nodes.

RESULTS

The comprehensive rotorcraft code CAMRAD II and the high-fidelity CFD code Overflow are used throughout this study to simulate the side-by-side urban air taxi in edgewise forward flight. The main flow features and parameters used in forward flight are presented in Figure 10, for a rotor rotating CCW. For a CW rotor, the direction of the azimuth $\psi$ is reversed.

First, the side-by-side vehicle performances are obtained using CAMRAD II, with its free wake model. Then, the loosely coupled Overflow - CAMRAD II simulations for the intermeshing rotors without airframe are run, for different rotor overlaps. Finally, the complete vehicle simulated with CAMRAD II - Overflow will be presented.

CAMRAD II provides the trim solution and the rigid motion of the blades to Overflow. Overflow solves the Navier-Stokes flow equations with the rotor information. The loosely coupled methodology provides a more accurate strategy to calculate the rotor loads.

CAMRAD II Results: SbS in Edgewise Flight

The first part of this work consists in the analysis of different rotor overlaps and advancing ratios $\mu$, using the comprehensive code CAMRAD II. CAMRAD II provides a fast low-fidelity solution, and it only needs to be run on a single core. CAMRAD II can be used in the early stages of design for different rotorcraft configurations. Its fast turnaround is very useful when the design is changed multiple times, until a better or optimized configuration is obtained.

The side-by-side vehicle is analyzed for edgewise forward flight. The SbS helicopter model within CAMRAD II is set up to have zero flapping using cyclic pitch, the rotor collective trims the thrust to the weight, and the aircraft pitch angle
trims the propulsive force to the drag. Lateral $\theta_{lc}$ and longitudinal $\theta_{ls}$ cyclic pitches are set up to trim the $1/\text{rev}$ flapping ($\beta_{ls}$ and $\beta_{lc}$) to zero. The wake model is based on tandem helicopter models.

The twin rotor performance results for different overlaps and advancing ratios are shown in Figure 11. The rotor lift-to-drag $L/D_e$ ratio is a measure of rotor lifting efficiency, appropriate for cruise flow conditions, and is defined as follows:

$$\frac{L}{D_e} = \frac{WV}{P_i + P_0} = \frac{LV}{P + XV} \quad (14)$$

where $L$ and $X$ are the sum of the wind-axis lift and drag of both rotors, and $P$ is the sum of the shaft power of the two rotors.

Twin rotor performance is calculated for spans of 0.75$D$, 0.85$D$, 0.95$D$, 1.05$D$, and 10$D$. The 10$D$ case is representative of two isolated rotors. The most efficient overlap is found for a span of 0.85$D$: overlapped side-by-side rotor performance is about 15% better than separate, non-overlapped twin rotors. This favorable interference is due to each rotor operating in the upwash and downwash flow fields of the other rotor. The combined wake system approximates that of a single wing with a large span, which reduces the induced power. With no overlap, the rotors function as individual wings, hence larger induced power. With too much overlap, the system has small span, hence larger induced power.

Figure 12 shows the geometry of the vortex wake calculated with CAMRAD II for a span of 0.85$D$. There is wake rollup, and the outboard advancing blades generate two large super-vortices outboard and a weak interaction inboard.

**Overflow - CAMRAD II Results: SbS Rotors in Edgewise Flight**

The side-by-side rotors without airframe are simulated in forward flight using the Overflow - CAMRAD II loosely coupled strategy described in the Numerical Approach section, using the overset grids shown in Figure 8. Four different overlap distances are considered in this work: 25%, 15%, 5% and 0% overlaps, or equivalently, 0.75$D$, 0.85$D$, 0.95$D$ and 1$D$ spans. From the comprehensive analysis study described in the previous paragraphs, the rotors with 15% overlap are the most efficient in cruise, and thus the nominal case for study in the coupled simulations will be the side-by-side rotors with 15% overlap distance.

The flight condition simulated is the best-range cruise, at 5000 ft ISA +20$^\circ$C, with an advance ratio of $\mu = 0.356$. Table 2 summarizes the flow conditions.

<table>
<thead>
<tr>
<th>Condition</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>$h$</td>
<td>5000 ft</td>
</tr>
<tr>
<td>$T_\infty$</td>
<td>ISA +20$^\circ$C</td>
</tr>
<tr>
<td>$V_\infty$</td>
<td>115 kts</td>
</tr>
<tr>
<td>$N$</td>
<td>499.7 rpm</td>
</tr>
<tr>
<td>$M_{tip}$</td>
<td>0.484</td>
</tr>
<tr>
<td>$M_{aoa}$</td>
<td>0.172</td>
</tr>
<tr>
<td>$\mu$</td>
<td>0.356</td>
</tr>
<tr>
<td>$Re$</td>
<td>$1.9 \cdot 10^6$</td>
</tr>
</tbody>
</table>

**Table 2: Best range cruise flow conditions for the side-by-side urban air taxi CFD simulations.**

The Overflow - CAMRAD II coupling approach is considered to be converged when the trim angles do not change from one coupling step to the next, or when this change is relatively small. Figure 13 shows the trim angles history, for the right and left rotors (rotor 1 and rotor 2, respectively) in the 15% overlap case. Loose coupling convergence is reached after 15 to 20 coupling steps. As mentioned earlier, the cyclic pitch trims the $1/\text{rev}$ flapping to zero. The final values of collective pitch $\theta_0$, the first harmonics $\theta_{lc}$ and $\theta_{ls}$, and the coning angle $\beta_0$, for different rotor overlaps, for the right and left rotors, are shown in Table 3.
Fig. 13: Loose coupling angles convergence history, for intermeshing rotors with 15% overlap. There is no $1/rev$ flapping or lag motion.

The converged values of the trim angles are very similar for the different overlaps. There is a small dissymmetry between the right and left rotors for the collective and cyclic pitch angles.

The blade motions are transferred to Overflow using Fortran post-processing functions. In addition, the aircraft’s attitude (yaw, pitch, and roll angles) should also be updated in the CFD simulations at each coupling step. Previous loosely coupled Overflow - CAMRAD II computational work (Refs. 12, 18) does not update the aircraft’s attitude in the CFD. In this work, the aircraft pitch angle $\theta_A$, calculated with the comprehensive code, is given to Overflow for more accurate calculations. Figure 14 shows the history of the aircraft pitch angle for the different rotor overlaps. The initial value of the aircraft pitch angle is $(\theta_A)_0 = -0.96\,^\circ$, while the final values are close to $(\theta_A)_n = 0.55\,^\circ$. This difference is important, and updating the aircraft pitch angle at every coupling step has shown improved convergence of the CAMRAD II - Overflow loose coupling approach. The reader should note that while the converged aircraft pitch angles are positive, the rotor shafts have $3^\circ$ forward tilt relative to the airframe axes, and therefore the thrust is leaning forward, its $x$ wind-axis component producing propulsive force $X$. Table 4 shows the converged rotor shaft pitch angles for different overlaps.

<table>
<thead>
<tr>
<th>Overlap</th>
<th>25%</th>
<th>15%</th>
<th>5%</th>
<th>0%</th>
</tr>
</thead>
<tbody>
<tr>
<td>Rotor #</td>
<td>1 2</td>
<td>1 2</td>
<td>1 2</td>
<td>1 2</td>
</tr>
<tr>
<td>$\beta_0$ [(^\circ)]</td>
<td>3.1</td>
<td>3.1</td>
<td>3.1</td>
<td>3.1</td>
</tr>
<tr>
<td>$\theta_0$ [(^\circ)]</td>
<td>9.2</td>
<td>9.1</td>
<td>9.4</td>
<td>9.1</td>
</tr>
<tr>
<td>$-\theta_1$ [(^\circ)]</td>
<td>8.2</td>
<td>8.1</td>
<td>8.5</td>
<td>8.0</td>
</tr>
<tr>
<td>$\theta_2$ [(^\circ)]</td>
<td>3.5</td>
<td>3.5</td>
<td>3.6</td>
<td>3.5</td>
</tr>
</tbody>
</table>

Table 3: Converged trim angles for different rotor overlaps. Values are in degrees. The $1/rev$ flapping is trim to zero with the lateral $(\theta_1)$ and longitudinal $(\theta_2)$ cyclic pitch. Rotor 1 corresponds to the right rotor (CCW), rotor 2 is the left rotor (CW).

<table>
<thead>
<tr>
<th>Overlap</th>
<th>25%</th>
<th>15%</th>
<th>5%</th>
<th>0%</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\theta_R$ [(^\circ)]</td>
<td>-2.46</td>
<td>-2.44</td>
<td>-2.46</td>
<td>-2.52</td>
</tr>
</tbody>
</table>

Table 4: Converged rotor shaft pitch angles for different rotor overlaps. Values are in degrees.

**Airloads**

The airloads for the side-by-side rotors with 15% overlap are shown in Figures 15 and 16. Figure 15 shows the mean normal force coefficient $c_n$, chord force coefficient $c_c$, and pitch moment coefficient $c_m$, as a functions of the radial distance $r/R$, for the right and left rotors. The sectional coefficients are multiplied by the square of the sectional Mach number $M$. The mean is obtained by averaging over a complete rotor revolution, $0 \leq \psi < 360$, where $\psi$ is the azimuth. There are some asymmetries between the right and left rotor airloads, and the differences slightly increase near the tip. They are probably due to the turbulent character of the vortex wake. The mean sectional airload values for the other overlap distances (25%, 5% and 0%) are very similar to the 15% case.

The azimuth variations of the computed sectional loads at four radial stations are shown in Figure 16. The mean values have been removed. The radial stations considered in this study are

---

1 The positive direction for the pitch angles is aft.
Fig. 15: Mean sectional airloads for two side-by-side rotors with an overlap distance of 15%, $\mu = 0.356$.

(a) Normal force coefficient (mean removed), $c_n M^2$ - mean, at four different radial stations.

(b) Chord force coefficient (mean removed), $c_c M^2$ - mean, at four different radial stations.

(c) Pitch moment coefficient (mean removed), $c_m M^2$ - mean, at four different radial stations.

Fig. 16: Azimuthal distribution of airloads for two side-by-side rotors with an overlap distance of 15%.
Fig. 17: Airloads distribution for two side-by-side rotors for three overlap distances, 25%, 15% and 0%, at four radial stations. The first row of each overlap distance shows the normal force coefficient, and the second row shows the chord force coefficient.
$r/R = 0.5, r/R = 0.75, r/R = 0.85,$ and $r/R = 0.965$. Figure 16 (a) shows the normal force coefficient $c_n M^2$. The normal force coefficient varies significantly with the azimuth at each radial station. For the most outboard location, $r/R = 0.965$, the peaks observed near $\psi = 90^\circ$ and $\psi = 270^\circ$ are the consequence of blade-vortex interaction (BVI). The peaks before $\psi = 360^\circ$ are due to the separated wake from the hubs, and they can be observed at the four radial stations.

The nomenclature from Figure 10 for a rotor in forward flight rotating CCW will be used in the following paragraphs. From the airloads distribution we can identify three BVI events happening on the rotor:

1. BVI near $\psi = 90^\circ$ (advancing side), seen by the peak in normal force near the tip of the blade. The BVI is from the vortex trailed by the previous blade of the same rotor.

2. BVI before $\psi = 270^\circ$ (retreating side). In this case, the vortex is coming from the blade of the opposite rotor. This is a strong BVI event.

3. BVI right after $\psi = 270^\circ$ (retreating side), see the smaller peak right after the peak from BVI from the opposite rotor. This time again, the vortex is trailed by the previous blade of the same rotor.

In single rotors, the BVI observed are the BVI events (1) and (3) described previously. BVI event (2) is characteristic of ShS overlapping rotors.

Figure 16 (b) shows the chord force coefficient. The peaks in chord force appear near the beginning and end of the overlap region; that also agrees well with the BVI events. The pitch moment coefficient is shown in Figure 16 (c). The variations of the pitching moment along the azimuth are relatively smaller.

In all the cases, for all the radial stations, there are fluctuations near $\psi = 360^\circ$, where the blades are downwind and immersed in the separated wake of the hub.

The effect of the overlap distance on the airloads (mean removed) can be observed on Figure 17. The first two rows show the normal and chord coefficients $(M^2 c_n - \text{mean}, M^2 c_c - \text{mean})$ for four different radial stations for two side-by-side rotors with an overlap of 25%. The third and fourth row are for an overlap distance of 5%. The last two rows represent the case with 0% overlap.

For the 25% overlap distance (rows one and two of Figure 17), the three BVI events can be seen even more clearly than for the 15% overlap case. This time, with a 25% overlap, there is only 50% of the radius that is out of the overlap zone at $\psi = 270^\circ$. The BVI events happening on the retreating side appear to be equally strong near the tip; see the last column of rows one and two (25% overlap: $M^2 c_n - \text{mean}, M^2 c_c - \text{mean}$) of Figure 17, $r/R = 0.965$. The BVI events at the retreating side can also be observed further from the tip, at $r/R = 0.85$. With a rotor separation of 0.85D, the double BVI event near $\psi = 270^\circ$ was clearly depicted in the outermost radial station ($r/R = 0.965$), but not in the $r/R = 0.85$ station. The second BVI event appears to be stronger at this radial station. Traveling further inboard, $r/R = 0.75$, the third BVI event is no longer observable on the force coefficients. There is a relatively stronger peak on the chord force for the second BVI event.

The third and fourth rows of Figure 17 represent the case with 5% overlap. The BVI events near $\psi = 270^\circ$ happen very close to each other due to the reduced overlap distance. In addition, there seem to be some discrepancies between the left and right rotor airloads near this region; this is probably caused by the turbulent character of the vortex wake.

When looking at the chord force coefficient distribution, there is a single large peak at $r/R = 0.85$, instead of the two peaks observed with 25% overlap, which suggests a combination of both BVI events in one single stronger event for certain radial stations.

The last two rows of Figure 17 show the normal and chord force coefficient for two side-by-side rotors with 0% overlap. There is still interference between the two rotors because each rotor is operating in the upwash flow field of the other rotor. In this case, there is only BVI from vortices trailed by the blades from the same rotor.

When comparing the airloads for the different overlaps at the innermost radial station represented here, $r/R = 0.5$, the distributions seem very similar to each other; see the first column of Figures 16 and 17. That is, outside of the overlap zone, the airloads do not change for different overlaps.

Overall, there are important interactions between the rotors that affect the airload distributions. The BVI event from the vortices originating from the opposite rotor changes with the overlap distance. The right and left rotors have similar airload distributions, and the small dissymmetries are probably caused by the turbulent character of the vortex wake.

**Rotor Disk Forces**

Figure 18 shows the rotor disk normal force coefficient $M^2 c_n$ for the side-by-side rotors for 25%, 15%, 5% and 0% overlap distances. The grey dashed lines represent the overlap region. For a reminder of the flow features and parameters in forward flight on the rotor disk, the reader should refer to Figure 10.

The effect on the normal force from the interactions between the rotor blades can be clearly seen for the 25% and 15% overlap distances, around the azimuth $\psi = 270^\circ$; see the red regions. In the 5% case, the affected region is smaller but the magnitude of the normal force is relatively higher. The regions of higher normal force downstream near $\psi = 360^\circ$ are due to the separated wake of the hub.

When comparing the right rotor (CCW) and the left rotor (CW), there is symmetry between the force coefficients, and the slight differences are not unexpected. Overflow is solving the unsteady Navier-Stokes equations, and a small asymmetry between rotors is normal.

The chord force coefficients on the rotor disk for 25%, 15%, 5% and 0% overlaps are shown in Figure 19. The main interest of these images is the clear effect of the overlap distance.
Fig. 18: Rotor disk normal force coefficient for side-by-side rotors with different overlap distances. The grey dashed lines show the overlap zone.
Fig. 19: Rotor disk chord force coefficient for side-by-side rotors with different overlap distances. The grey dashed lines show the overlap zone.
Fig. 20: Vortex wake for two side-by-side rotors with 25% overlap.

Fig. 21: Vortex wake for two side-by-side rotors with 15% overlap.
Fig. 22: Vortex wake for two side-by-side rotors with 5% overlap.

Fig. 23: Vortex wake for two side-by-side rotors with 0% overlap.
on the polar plots. These images agree well with the linear plots from Figures 16 and 17.

Wake Geometry
The vortex wakes for 25%, 15%, 5% and 0% overlap distances are shown in Figures 20, 21, 22, and 23, respectively. The vortices are visualized by using iso-surfaces of the Q-criterion vorticity, and they are colored with the vorticity magnitude. All figures use the same values of Q-criterion and vorticity magnitude.

Some common flow features that appear in the vortex wakes of all four cases are summarized here:

- Blade-vortex interactions.
- Blade tip-vortex rollup downstream at the exterior (advancing side) of both rotors.
- Interaction in the overlap region (retreating side) of one vortex wake with another.
- Development of worm-like vortical structures (Ref. 18) in the downstream wake.
- Separated wake flow from the hubs.

The rotor-rotor interactions are stronger when the rotors are closer. An interesting flow feature can be observed for the 0% case: the vortices from each rotor do not interact in the beginning; however, further downstream there are clear interactions and vortex wrapping and rolling can be seen.

The three BVI events can be easily found in Figure 20. Using the nomenclature from Figure 10, where blade 1 is the blade in the first quadrant (0\(^\circ\) < \(\psi\) < 90\(^\circ\)), blade 2 is the blade positioned in the second quadrant (90\(^\circ\) < \(\psi\) < 180\(^\circ\)), blade 3 is the blade positioned in the third quadrant (180\(^\circ\) < \(\psi\) < 270\(^\circ\)), and blade 4 is the blade positioned in the fourth quadrant (270\(^\circ\) < \(\psi\) < 360\(^\circ\)). The first BVI event happens near \(\psi = 90^\circ\), see blade 2 (\(\psi = 90^\circ\)) from the right rotor: the blade encounters the vortex trailed by blade 3 from the same rotor. The second BVI event happens before \(\psi = 270^\circ\); see blade 4 from the right rotor: this blade meets with the vortex shed by blade 4 from the left rotor (located at \(\psi = 315^\circ\)). As we decrease the overlap distance, the second BVI event is delayed in azimuth, see for example Figure 22 where the BVI event happens at \(\psi = 270^\circ\). The third BVI event is harder to see, because it happens for an azimuth located between 270\(^\circ\) < \(\psi\) < 315\(^\circ\). It can still be discerned for blade 4 of the left rotor; see Figure 23 for the case with 0%, where the blade has already encountered the vortex from the previous blade (blade 1), from the same rotor. In the cases with some overlap, blade 4 will again encounter the vortex from blade 4 of the opposite rotor. The multiple BVI events will likely lead to increased noise.

In conclusion, the wake in the overlap zone leads to unsteadiness in the flow, with significant blade-vortex interactions and therefore likely increased noise.

Performances
To finish the study of different overlap distances on side-by-side rotors for urban air mobility vehicles, the performances of the rotor systems are presented. The efficiency of the rotor system is measured using the rotor-to-lift drag ratio, defined as in equation 11.

\[
\frac{(L/D_e)/(L/D_{e,ref})}{\text{Overlap}} = \begin{array}{l}
25\% \quad 1.077 \\
15\% \quad 1.095 \\
5\% \quad 1.036 \\
0\% \quad 1
\end{array}
\]

Table 5: Side-by-side loosely coupled Overflow - CAMRAD II performance calculations of the rotor system (right rotor + left rotor). The relative rotor efficiency in cruise, \((L/D_e)/(L/D_{e,ref})\), is shown. The reference \((L/D_{e,ref})\) is for the case with 0% rotor overlap.

The calculated performances for each overlap are shown in Table 5. The trend obtained from the comprehensive analysis is confirmed here with the high-fidelity CFD results: the rotors with 15% overlap are the most efficient in cruise.

The study of two isolated side-by-side rotors, where the separation between rotors is high enough that there are no rotor-rotor interactions, remains to be done. The comparison of different overlapping distances with the two isolated rotors will be left for a future study.

Overflow - CAMRAD II Results: SbS Urban Air Taxi in Edgewise Flight

This final section presents NASA’s urban air taxi concept with side-by-side rotors, for a span of 0.85\(D\). The complete vehicle is a complicated geometry as it has multiple components: the rotors, the fuselage, the engine, the empennage, etc. The mesh generation process is very long and requires considerable expertise from the user in order to produce high quality meshes. In addition, the convergence of the flow solver is more difficult than the cases with only rotors. A complete study of the full vehicle and an exhaustive comparison with the results of rotors only are left for a future paper. In this section we briefly show some results of the full side-by-side UAM vehicle in order to illustrate the main effects of adding the airframe.

![Fig. 24: Rotor disk normal force coefficient \(M^2c_n\) for the side-by-side urban air taxi with 0.85\(D\) span, or 15\% overlap. The dashed grey lines show the overlap zone.](Image)
The rotor disk normal force coefficient is shown in Figure 24. It seems to follow a similar pattern as the side-by-side rotors without airframe. However, there are some differences; for example, the separated flow coming from the hubs seems to affect the rotor disk less when the airframe is present.

The vortex wake for the full vehicle composed of overlapping rotors and airframe is shown in Figures 25 and 26. Figure 25 shows the top view of the Q-criterion vorticity iso-surfaces colored by the vorticity magnitude. Some similarities in the wake geometry with the case of the side-by-side rotors without the airframe (Figure 21) can be observed, i.e., the BVI events seem to happen at comparable times. However, many more turbulent structures coming from the airframe are present: there is separated flow from the engines, junction vortices, and other complicated turbulent features.

Figure 26 shows an oblique view of the vortex wake colored by the pressure. The side-by-side urban air taxi offers some advantages compared to other UAM concepts: a more compact design with an increased efficiency in cruise might be one key factor enabling its feasibility. Although there are no validation data, the computational results may assist in future validation testing.

**SUMMARY**

The effects of the rotor overlap distance of NASA's side-by-side urban air taxi concept have been analyzed. A loose coupled approach to simulate the flow of the twin rotor system in cruise conditions has been followed. The rotorcraft comprehensive code CAMRAD II and the overset finite-difference Navier-Stokes high-order accurate CFD solver Overflow have been coupled. CAMRAD II provides the blade motions and aircraft attitude to Overflow and Overflow accurately solves the flow, and the aerodynamic loads obtained with Overflow are sent as input to CAMRAD II, which recalculates the new trim solution and updates the blade motions and aircraft angles.

The airloads for different overlaps have been studied. The intermeshing region causes multiple blade-vortex interaction events, with vortices shed from the same and the opposite rotor. The opposite rotor BVI's event varies when the overlap distance is changed. The right and left rotors have similar airload distributions, and the small dissymmetries are probably caused by the turbulent character of the vortex wake.

The main features of the wake geometry with different overlap distances have been analyzed.

The performances of the rotors have been presented. First, the results with CAMRAD II have shown that overlapping rotors are more efficient in cruise than two isolated rotors, and the best configuration is for 15% rotor overlap. Then, for the best range cruise, the coupled CAMRAD II - Overflow simulations have confirmed this trend.

The computational results for the complete side-by-side vehicle have been introduced. Future work will include a more extensive analysis of the full vehicle, for different flight conditions.

The side-by-side urban air taxi is one of the concept vehicles expected to focus and guide NASA research activities in support of aircraft development for emerging aviation markets, in particular VTOL air taxi operations.
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REFERENCES


